

# Stability of Alternate Bars and Oblique Dunes



Alternate bars in the Tokachi River, Japan  
Image courtesy of V. Langlois



# Stability of River Bed Forms

## Part III



- Part I – linear stability of roll-waves, dunes & antidunes;
- Part II – linear stability of ripples and dunes;





# Stability of River Bed Forms

## Part III



- Part I – linear stability of roll-waves, dunes & antidunes;
- Part II – linear stability of ripples and dunes;
- Part III – linear stability of alternate bars & oblique dunes;





# Stability of River Bed Forms

## Part III

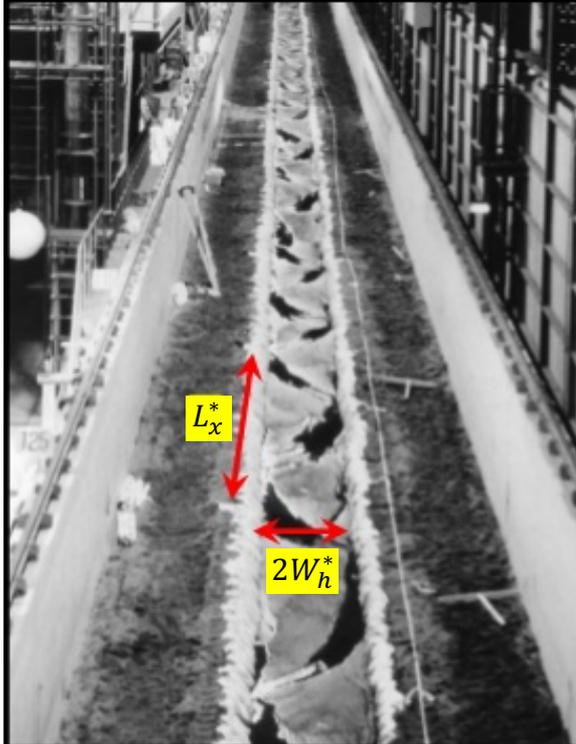


- Part I – linear stability of roll-waves, dunes & antidunes;
- Part II – linear stability of ripples and dunes;
- Part III – linear stability of alternate bars & oblique dunes;
- Part IV – weakly nonlinear analysis





## Alternate bars in straight channels

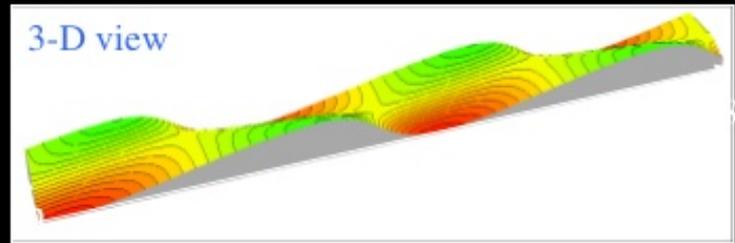


Lisle et al. (1997)

Naka River, Parker's e-book



$$\beta = \frac{W_h^*}{D^*}$$



Federici and Colombini (2003)

Ingeniería de ríos, Jorge Abad 2011





The name game: alternate bars



Super fun group game!



Two-Phase Continuum Models for Geophysical Particle-Fluid Flows

Max Planck Institute for the Physics of Complex Systems

14 March - 15 April 2016



The name game: alternate bars

## Alternate bars in a flume.

Image courtesy of S. Ikeda.



Super fun group game!



Two-Phase Continuum Models for Geophysical Particle-Fluid Flows

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The name game: alternate bars



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"A small dune sings only the few days  
in which there is no wind and no  
clouds"

(Andreotti, B. - LiveScience, 2005)



Super fun group game!

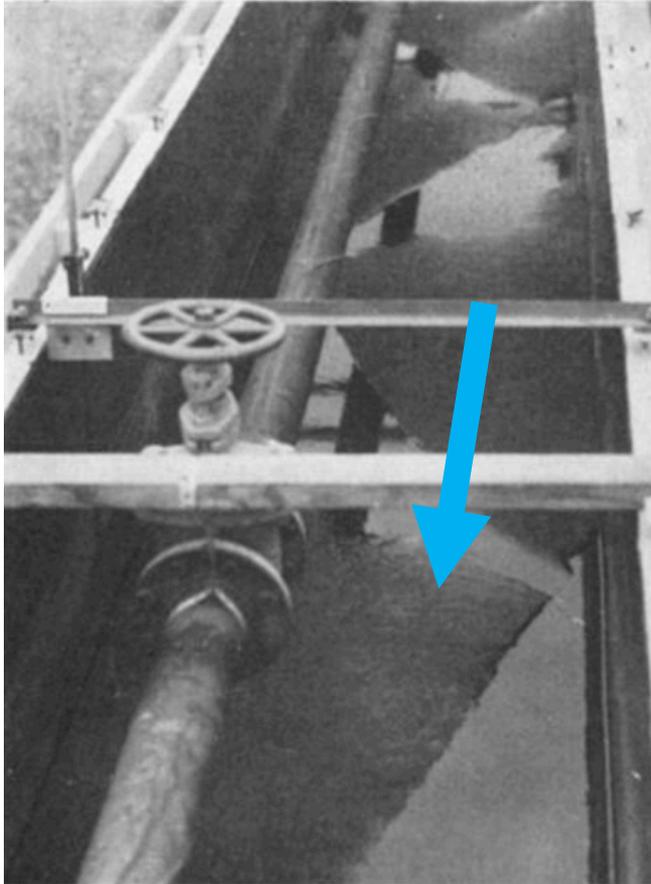




The name game: alternate bars

## Diagonal bars in a flume.

From Einstein & Shen, JGR 1964



Super fun group game!



Two-Phase Continuum Models for Geophysical Particle-Fluid Flows

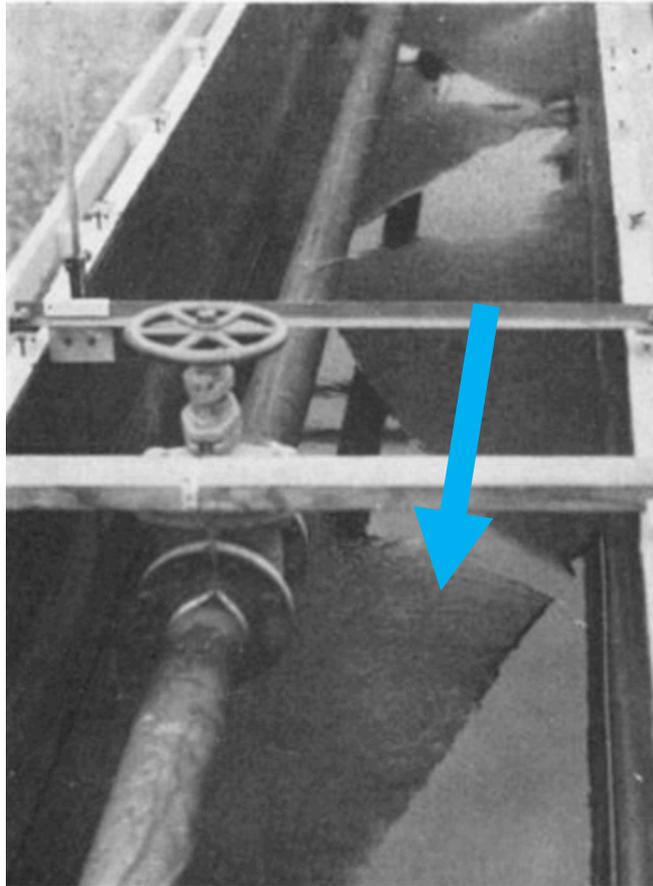
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## Diagonal bars in a flume.

From Einstein & Shen, JGR 1964



## Alternate bars in a flume.

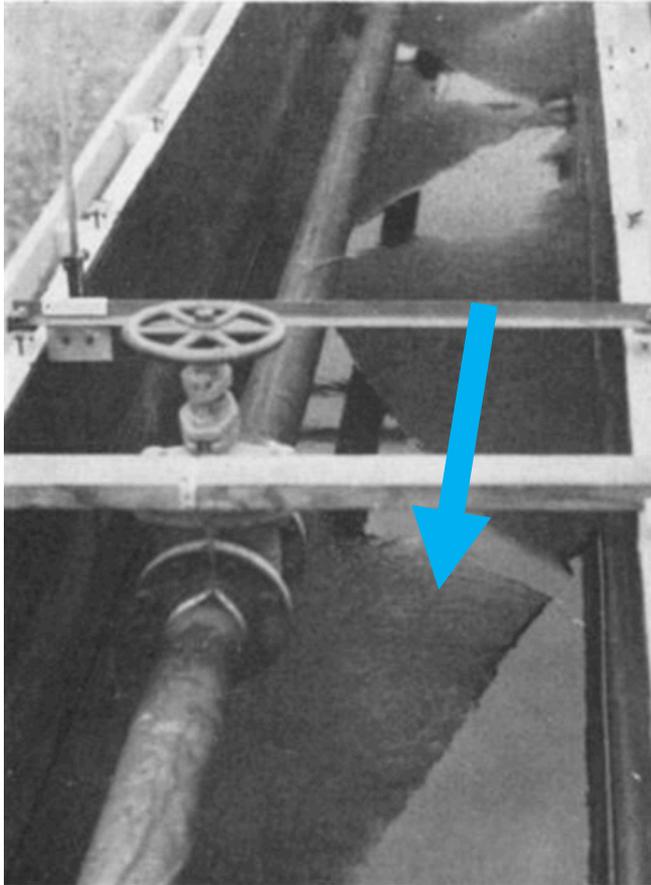
Image courtesy of S. Ikeda.





## Diagonal bars in a flume.

From Einstein & Shen, JGR 1964



- “ ... it is probable that diagonal bars have sometimes been classified as alternate bars by some authors.”
- “ ... experiments seem to indicate that a grouping of three-dimensional mesoforms [i.e. scaling with flow depth], in which the fronts of the mesoforms were diagonally aligned over the channel width, was responsible for these features.”

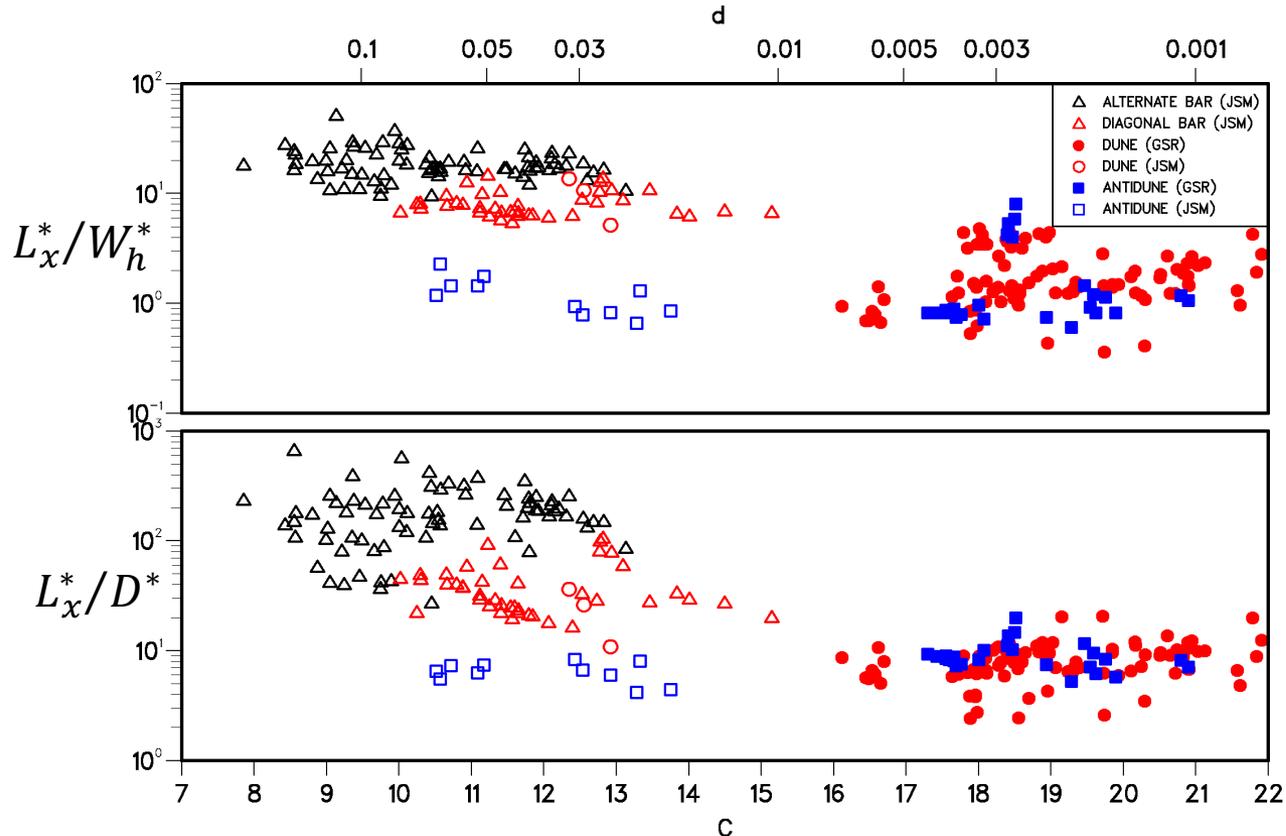
From Jaeggi, JHE 1984





JSM: Jaeggi (1984), Sukegawa (1971), Muramoto & Fujita (1978)

GSR: Guy, Simons & Richardson (1966)



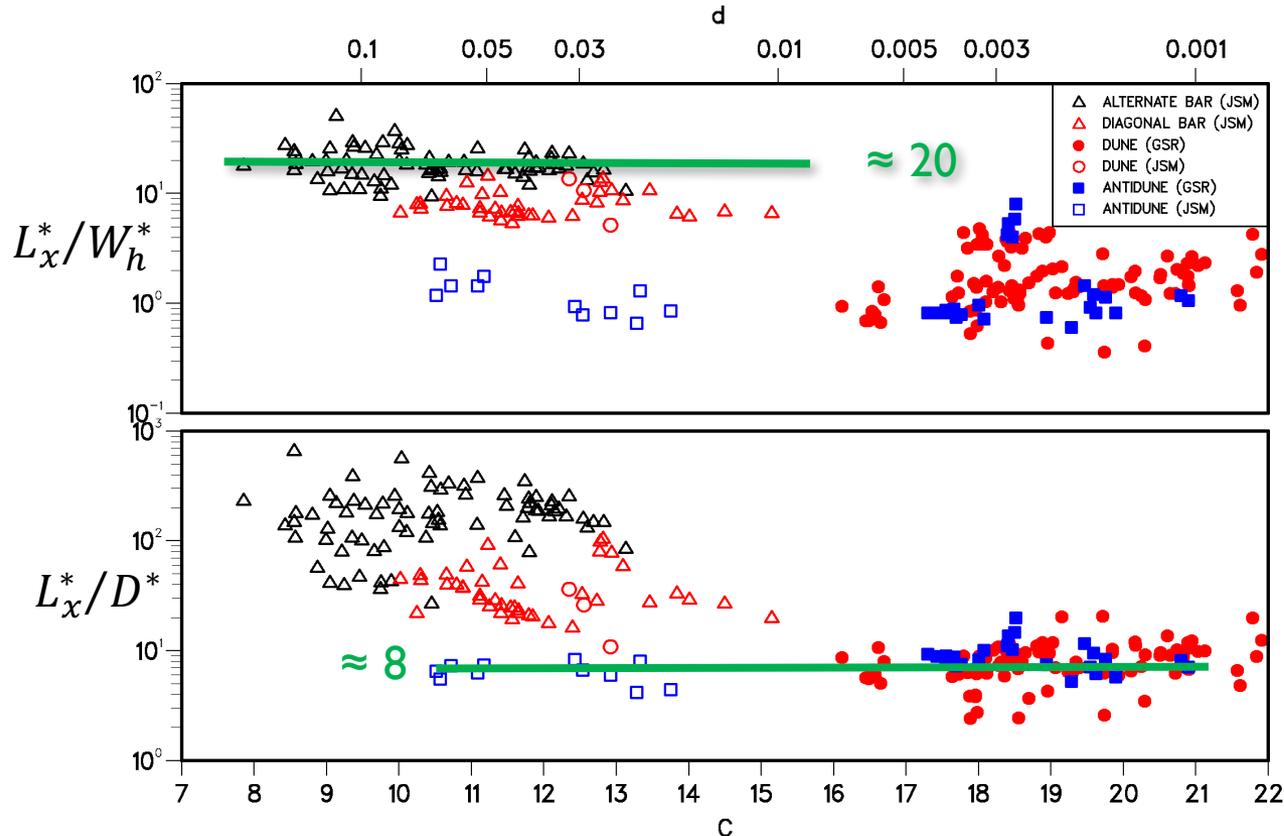
$$C = \frac{U^*}{u_f^*} = \sqrt{\frac{8}{f}} = \frac{1}{\kappa} \ln \left( \frac{11.09}{2.5d} \right) \quad (\text{Rough regime})$$





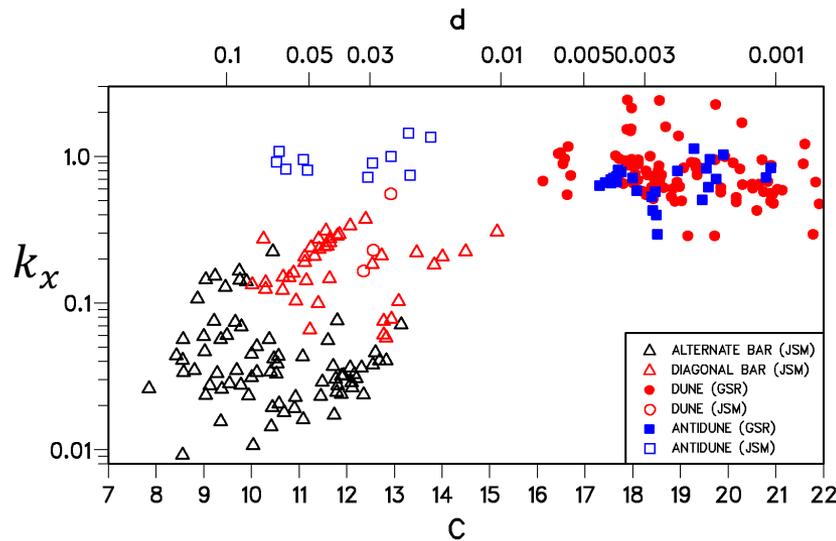
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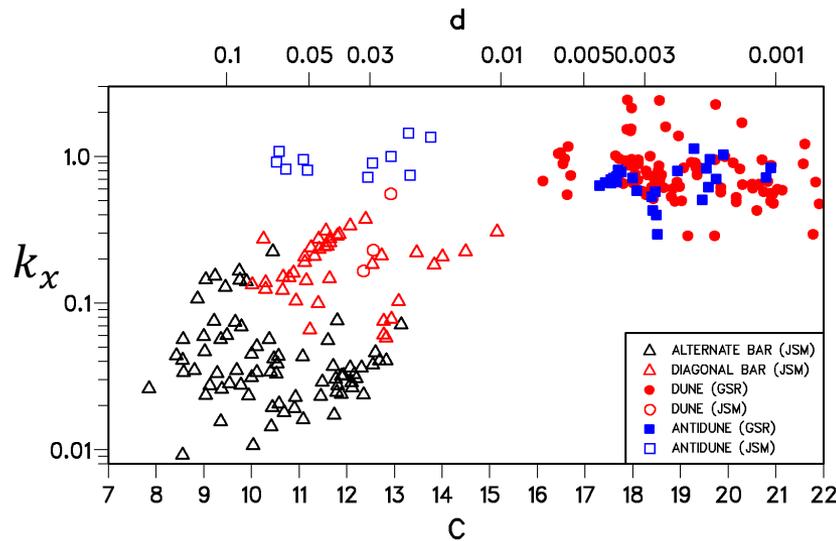




➤ Longitudinal wavenumbers for Diagonal Bars are larger than for Alternate Bars

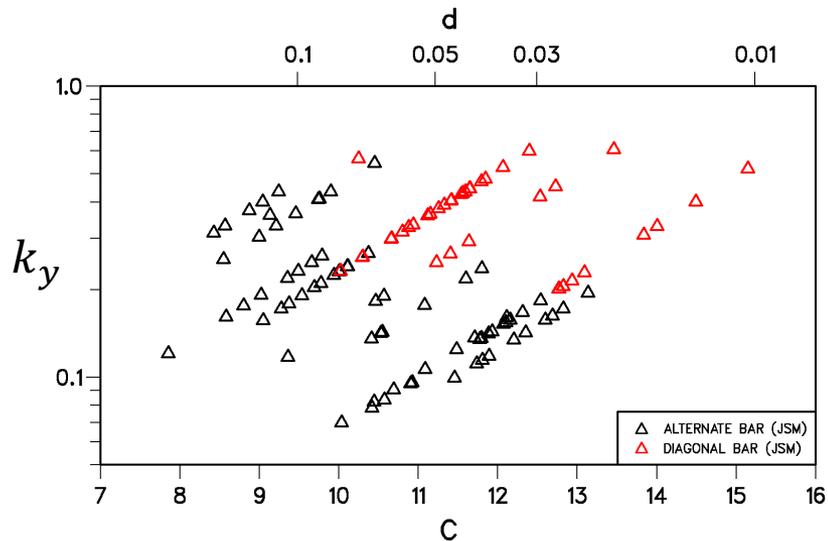
$$k_x = \frac{2\pi D^*}{L_x}$$





- Longitudinal wavenumbers for Diagonal Bars are larger than for Alternate Bars

$$k_x = \frac{2\pi D^*}{L_x}$$



- Transverse wavenumbers for Diagonal Bars are larger than for Alternate Bars

$$k_y = \frac{2\pi D^*}{L_y} = \frac{2\pi D^*}{4W_h^*} = \frac{\pi}{2\beta}$$





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Instability and river channels  
R. A. Callander

[Journal of Fluid Mechanics, Volume 36, Issue 03, May 1969, pp 465-480](#)

doi: 10.1017/S0022112069001765, Published online by Cambridge University Press 29 Mar 2006

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- Linear
- 2D Flow model
- Free Bars

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A unified bar–bend theory of river meanders  
P. Blondeaux and G. Seminara

[Journal of Fluid Mechanics, Volume 157, August 1985, pp 449-470](#)

doi: 10.1017/S0022112085002440, Published online by Cambridge University Press 20 Apr 2006

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- Linear
- 2D Flow model
- Free and Forced Bars

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Finite-amplitude alternate bars  
M. Colombini and G. Seminara and M. Tubino

[Journal of Fluid Mechanics, Volume 181, September 1987, pp 213-232](#)

doi: 10.1017/S0022112087002064, Published online by Cambridge University Press 21 Apr 2006

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- Weakly Non Linear
- 2D Flow model
- Free Bars

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Three-dimensional river bed forms  
M. Colombini and A. Stocchino

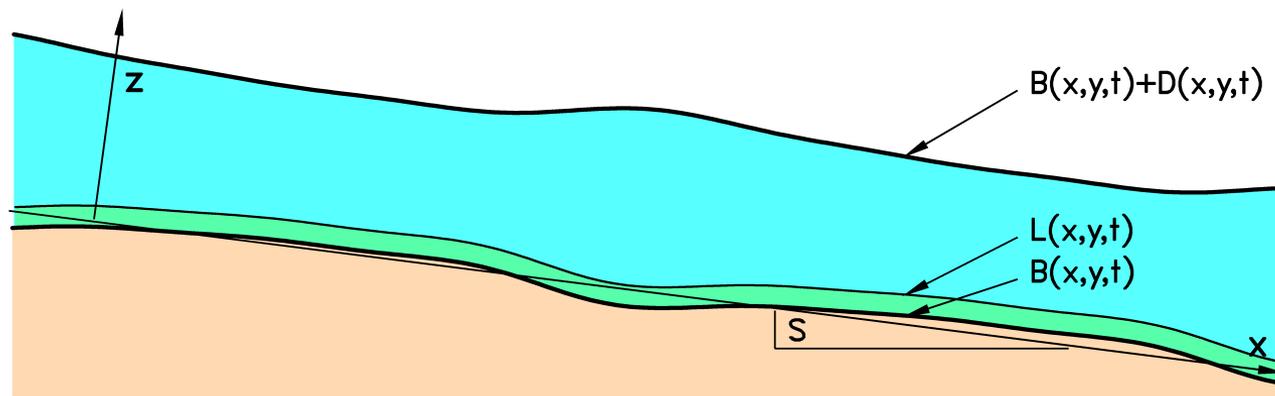
[Journal of Fluid Mechanics, Volume 695, March 2012, pp 63-80](#)

doi: 10.1017/jfm.2011.556, Published online by Cambridge University Press 07 Feb 2012

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- Linear
- 3D Flow model
- Free Bars





## FLOW MODEL

- 3D ROTATIONAL FLOW MODEL (infinitely wide channel)
- BOUSSINESQ'S CLOSURE (algebraic mixing length)
- COORDINATE TRANSFORMATION (rectangular domain)

## SEDIMENT TRANSPORT MODEL

- EQUILIBRIUM MODEL (Exner)
- BEDLOAD ONLY (MPM bedload function)
- CORRECTIONS FOR SEDIMENT WEIGHT (x – Fredsøe, y – Engelund)
- CORRECTION FOR BEDLOAD LAYER THICKNESS





$$G(x, y, z, t) = G_0(z) + \varepsilon g_1(z) \exp[ik_x(x - \omega t) + ik_y y + \Omega t] + c.c.$$

- **LINEAR LEVEL: differential eigenvalue problem**




$$G(x, y, z, t) = G_0(z) + \varepsilon g_1(z) \exp[ik_x(x - \omega t) + ik_y y + \Omega t] + c.c.$$

- **LINEAR LEVEL: differential eigenvalue problem**

$$\Omega = \Omega(k_x, k_y; Fr, C)$$

**DUNE FLAVOUR**




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**DUNE FLAVOUR**

$$\Omega = \Omega(\lambda, \beta; \vartheta, d)$$

**BAR FLAVOUR**





$$G(x, y, z, t) = G_0(z) + \varepsilon g_1(z) \exp[ik_x(x - \omega t) + ik_y y + \Omega t] + c.c.$$

- **LINEAR LEVEL: differential eigenvalue problem**

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**DUNE FLAVOUR**

**BAR FLAVOUR**

$$\lambda = k_x \beta$$

$$\beta = \frac{\pi}{2k_y}$$

$$\vartheta \cong 0.14 \frac{Fr^2 e^{\kappa C}}{C^2}$$

$$C = \frac{1}{\kappa} \ln \left( \frac{11.09}{2.5d} \right)$$

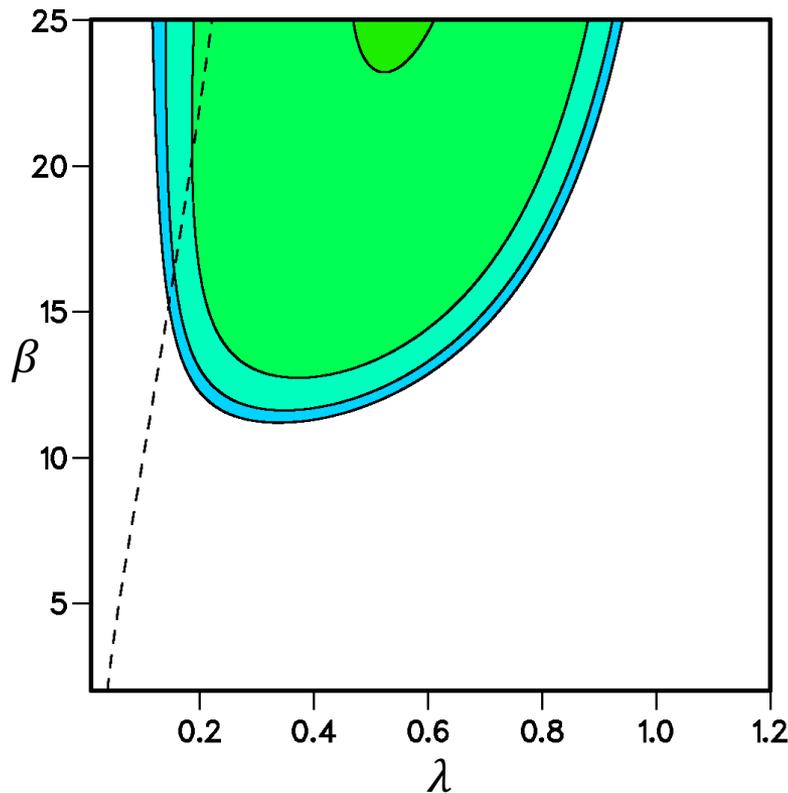




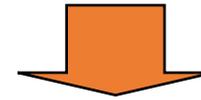
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- **LINEAR LEVEL: differential eigenvalue problem**

$$\Omega = \Omega(\lambda, \beta; \vartheta, d) \quad \lambda = k_x \beta \quad \beta = \frac{\pi}{2k_y}$$



$$\frac{\vartheta}{\vartheta_c} = 3 \quad d = 0.025$$



$$Fr = 1 \quad C = 13$$

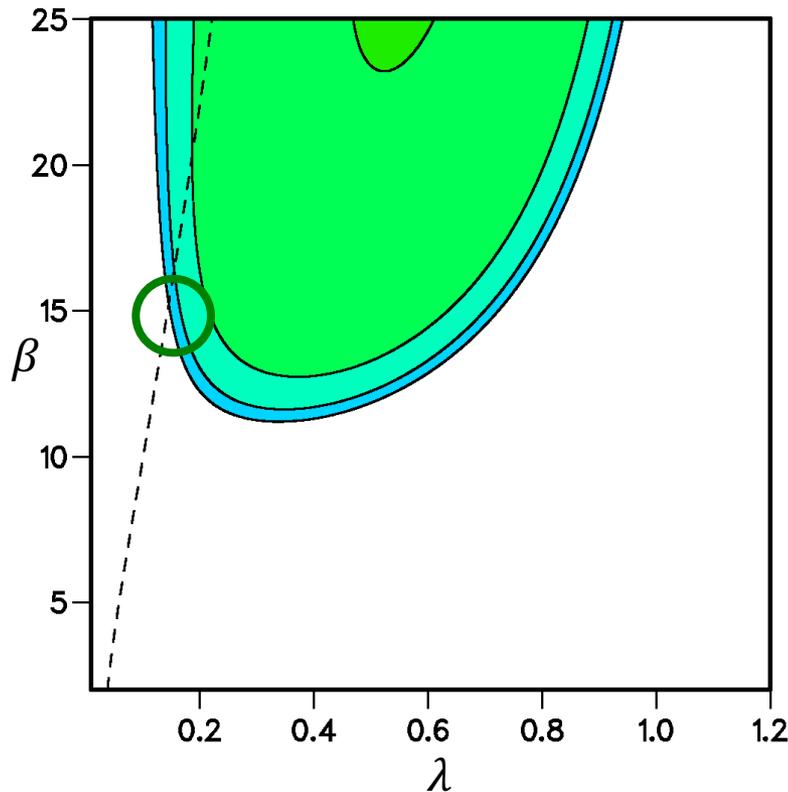




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○ Resonant conditions -  $(\lambda_R, \beta_R)$

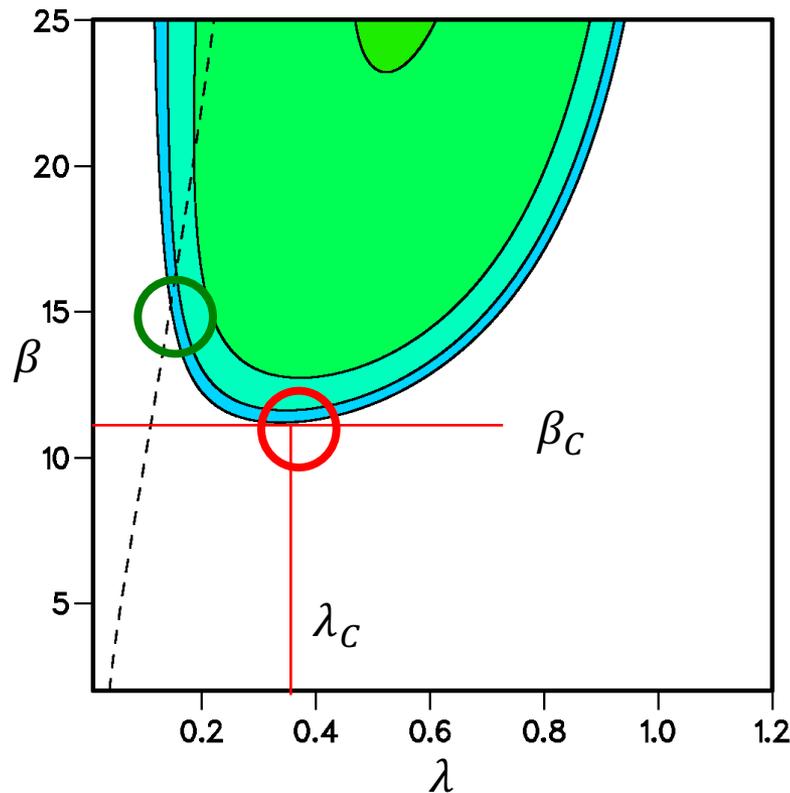




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- Resonant conditions -  $(\lambda_R, \beta_R)$
- Critical conditions -  $(\lambda_C, \beta_C)$

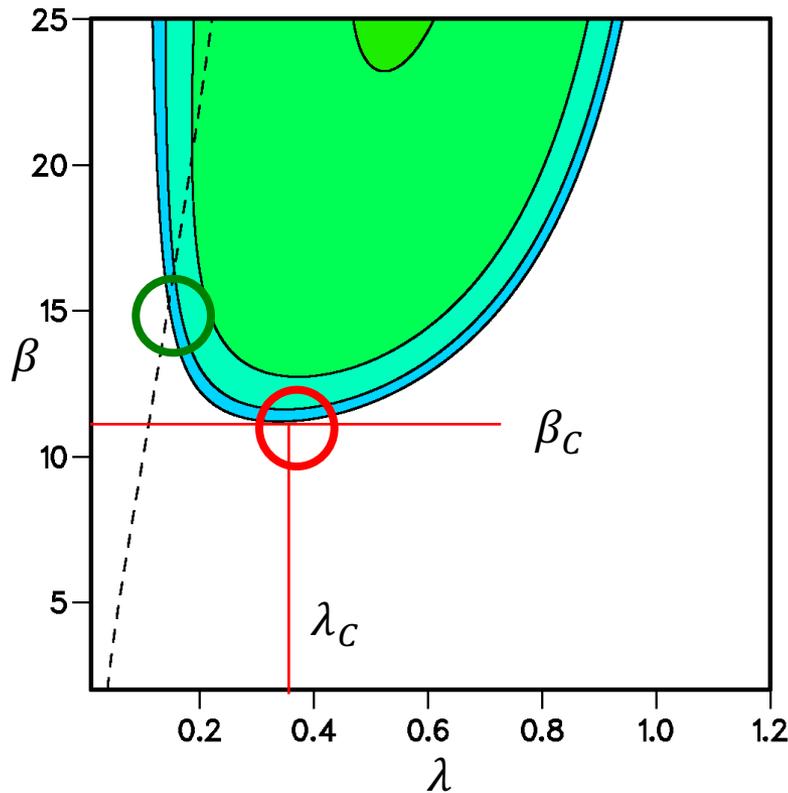




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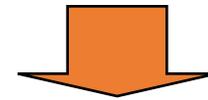
- **LINEAR LEVEL: differential eigenvalue problem**

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- Resonant conditions -  $(\lambda_R, \beta_R)$
- Critical conditions -  $(\lambda_C, \beta_C)$

$$\beta > \beta_C(\vartheta, d)$$



Alternate Bars do not form in a narrow channel



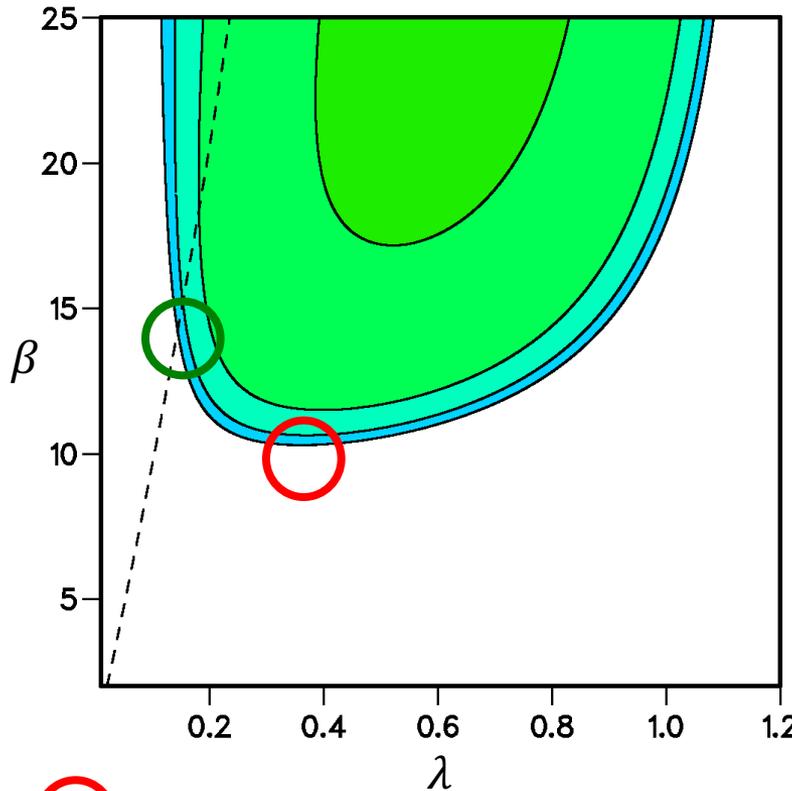


$$\frac{\vartheta}{\vartheta_C} = 3 \quad C = 13$$

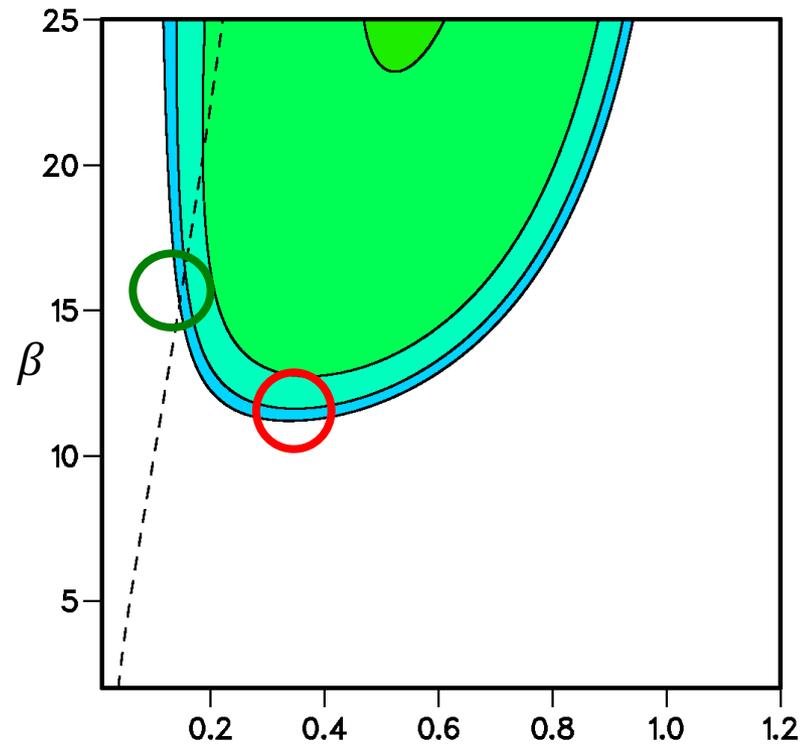


$$Fr = 1 \quad d = 0.025$$

SW MODEL



3D MODEL



 Critical conditions -  $(\lambda_C, \beta_C)$

 Resonant conditions -  $(\lambda_R, \beta_R)$



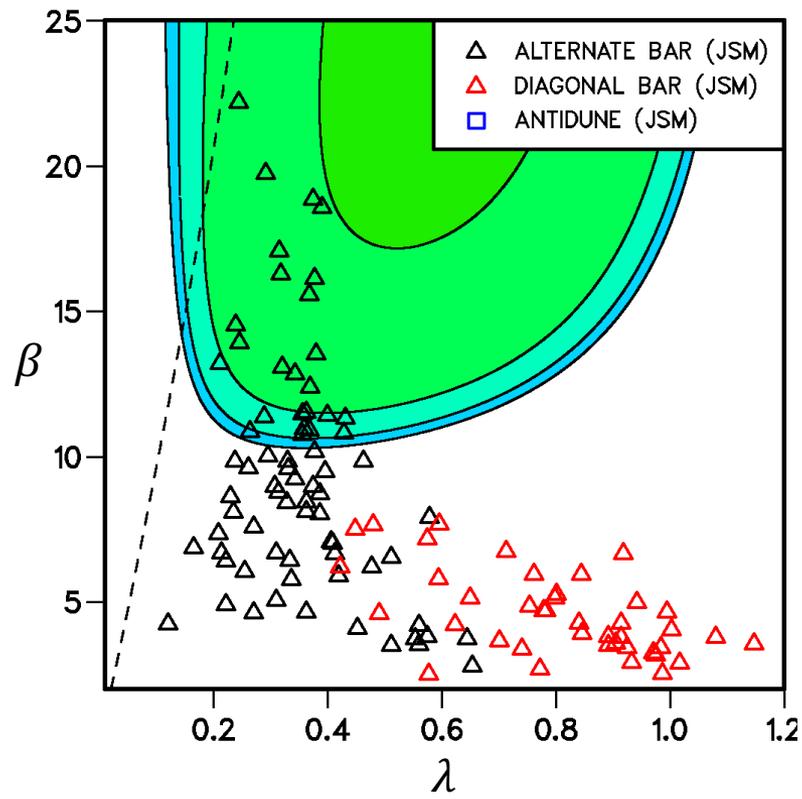


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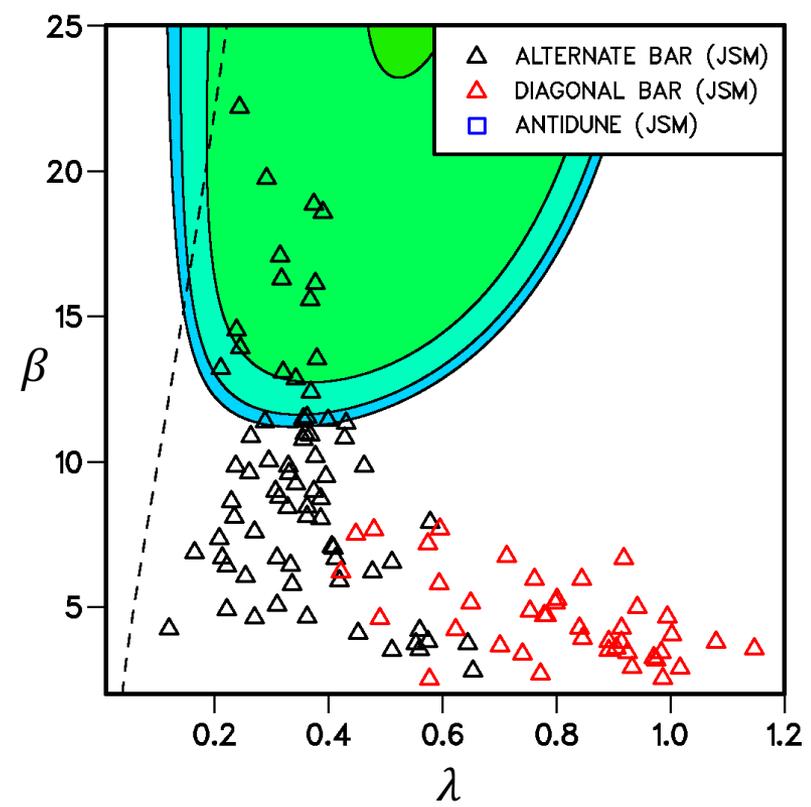


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SW MODEL



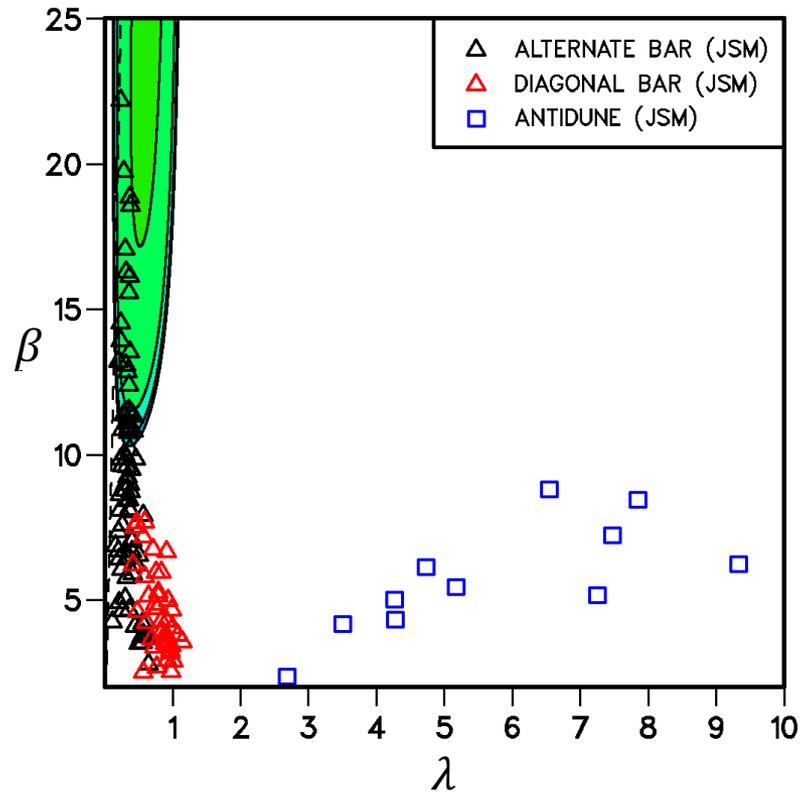
3D MODEL



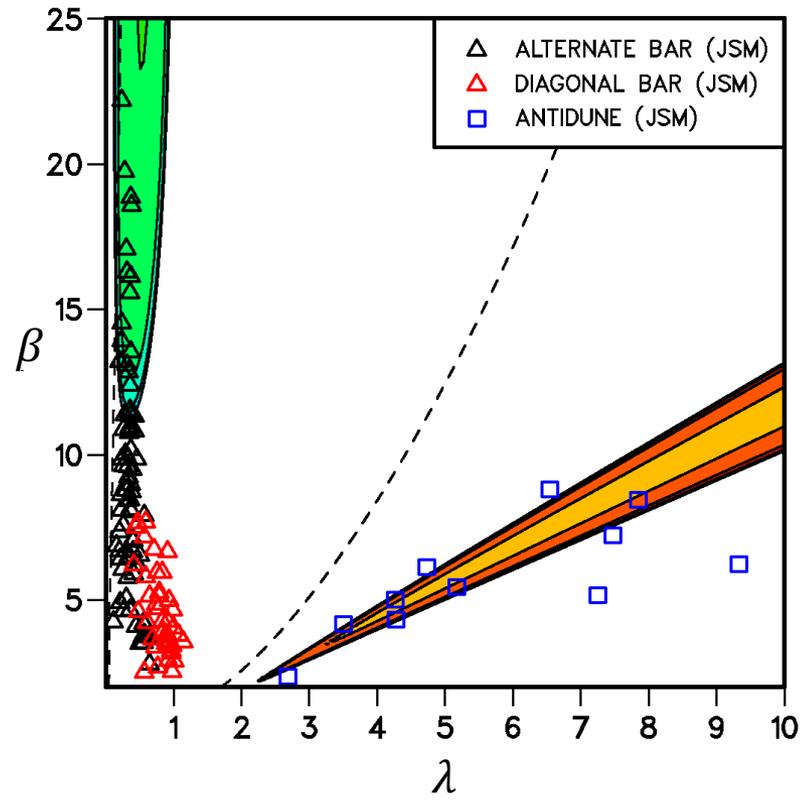


$$\frac{\vartheta}{\vartheta_C} = 3 \quad C = 13 \quad \longleftrightarrow \quad Fr = 1 \quad d = 0.025$$

SW MODEL



3D MODEL

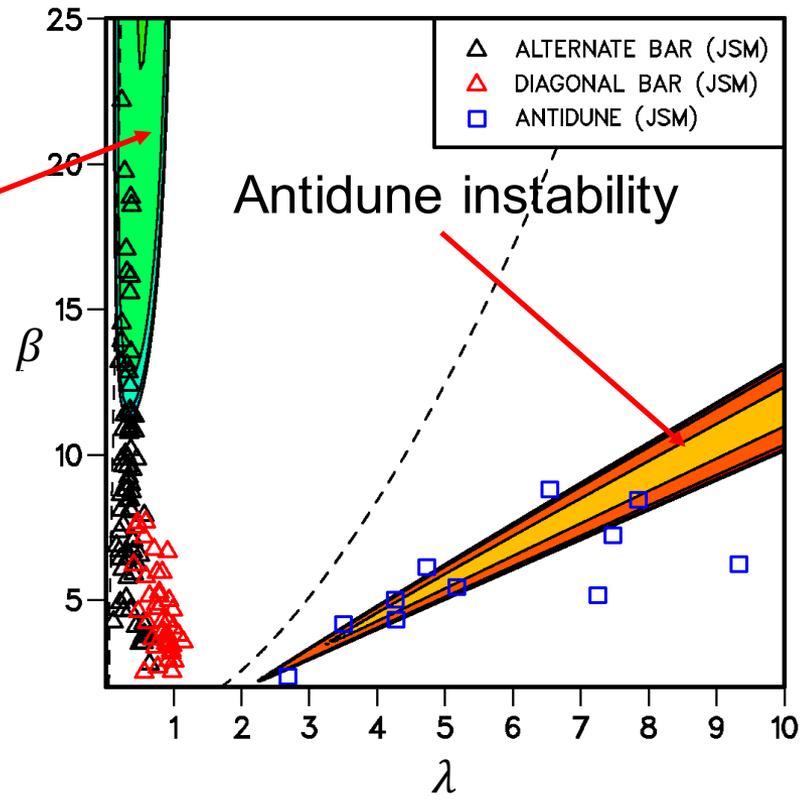
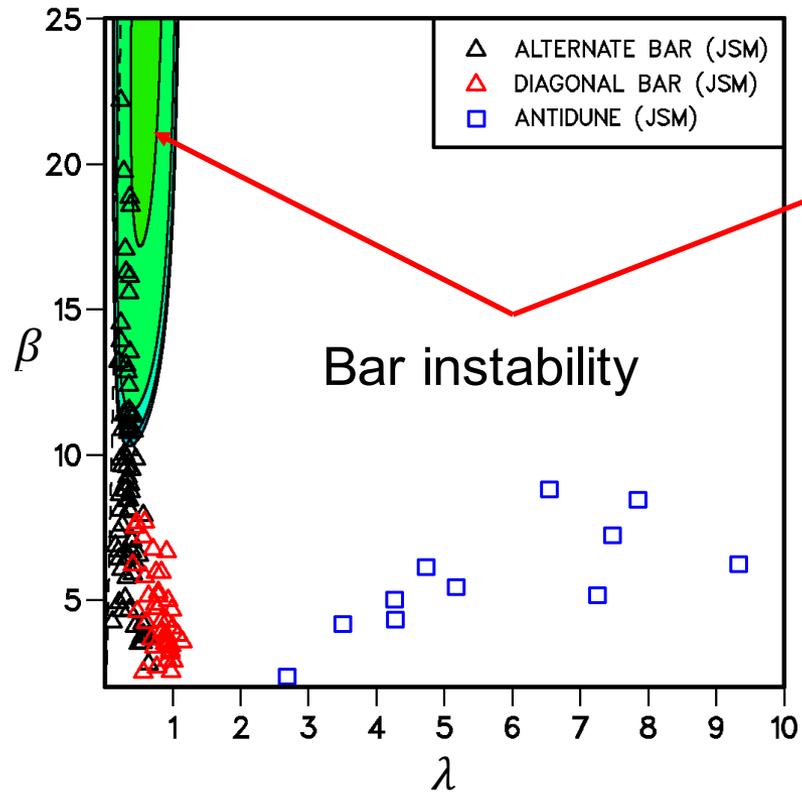




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SW MODEL

3D MODEL

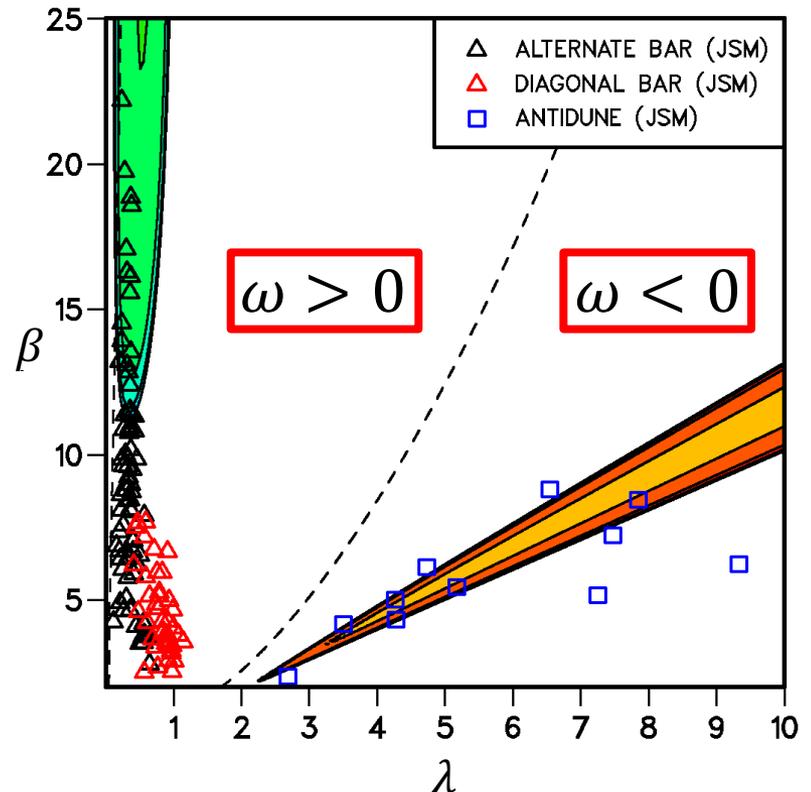
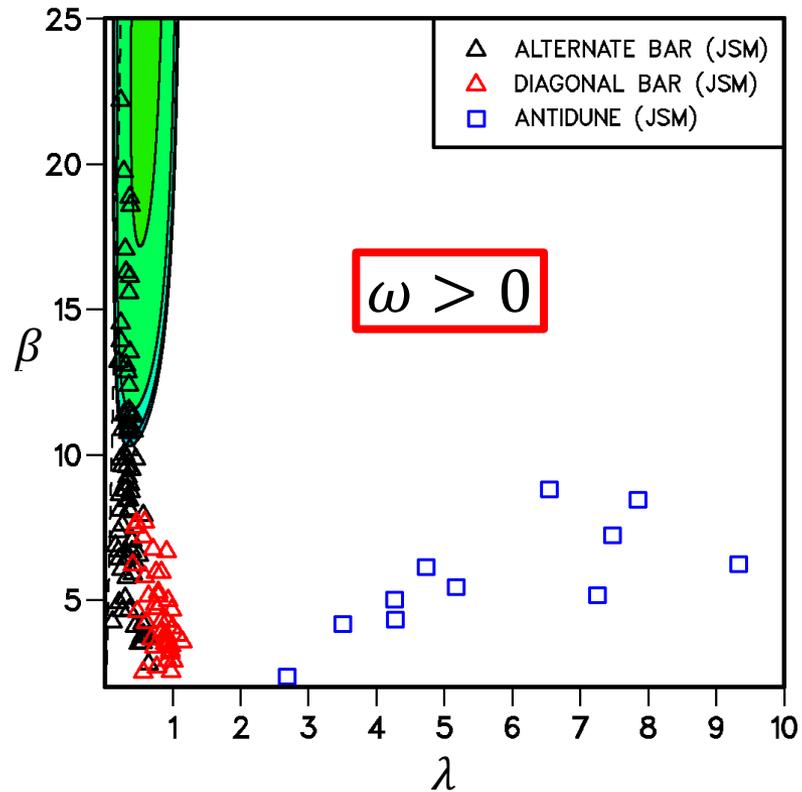




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SW MODEL

3D MODEL



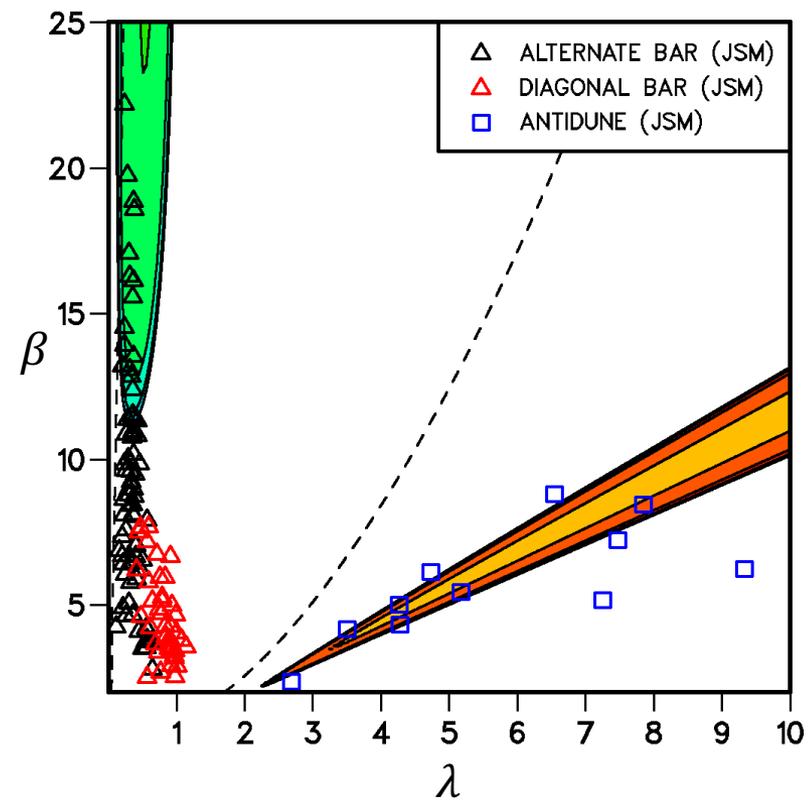
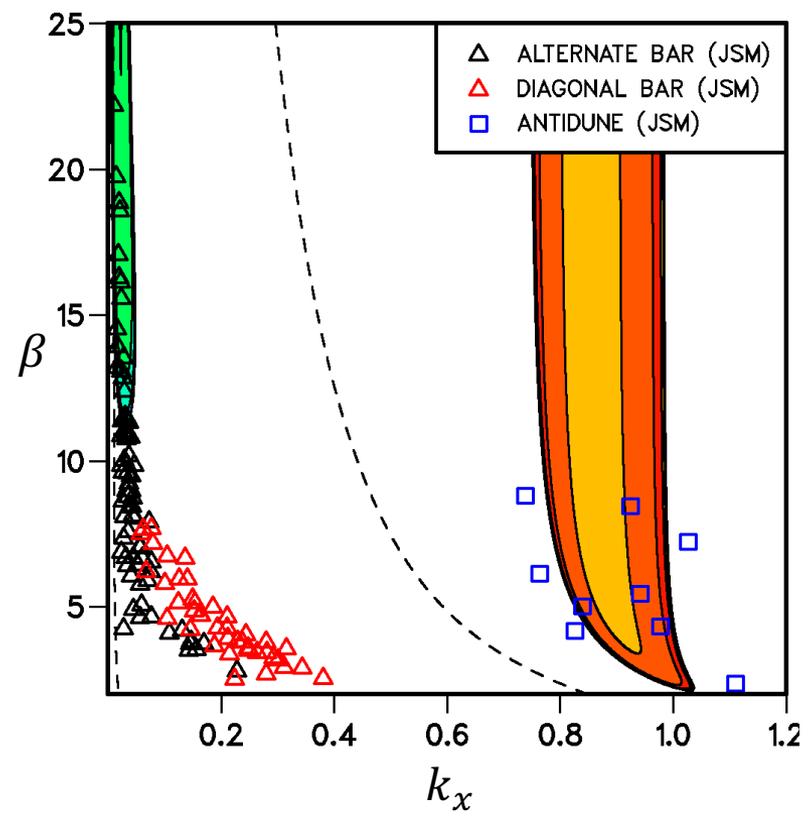
$$\beta = \frac{\lambda}{k_x}$$





# Alternate Bars: 3D linear stability

$$\frac{\vartheta}{\vartheta_c} = 3 \quad C = 13 \quad \longleftrightarrow \quad Fr = 1 \quad d = 0.025$$



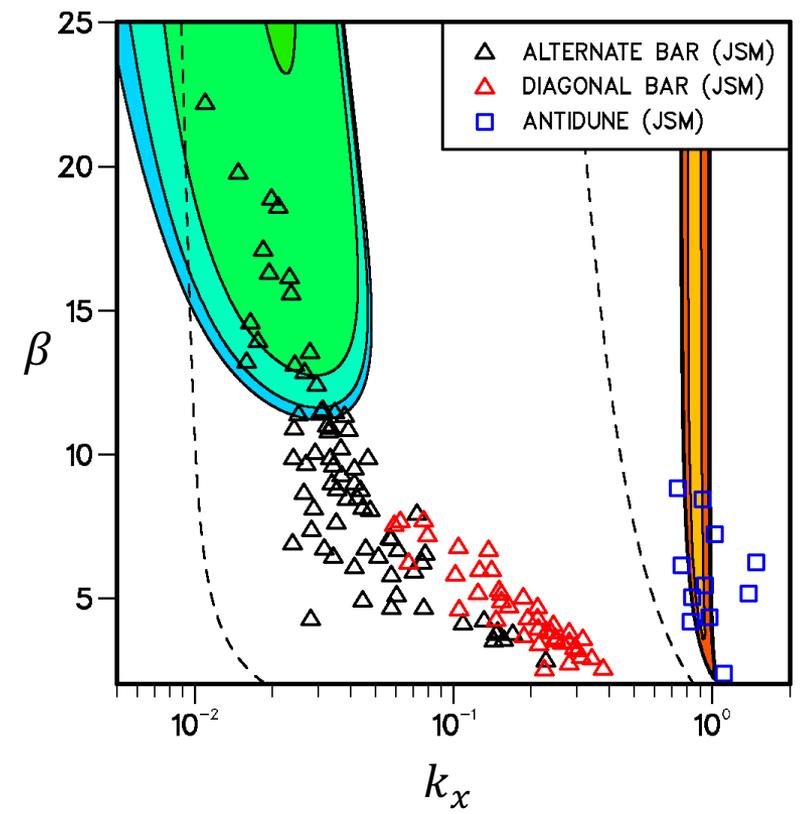
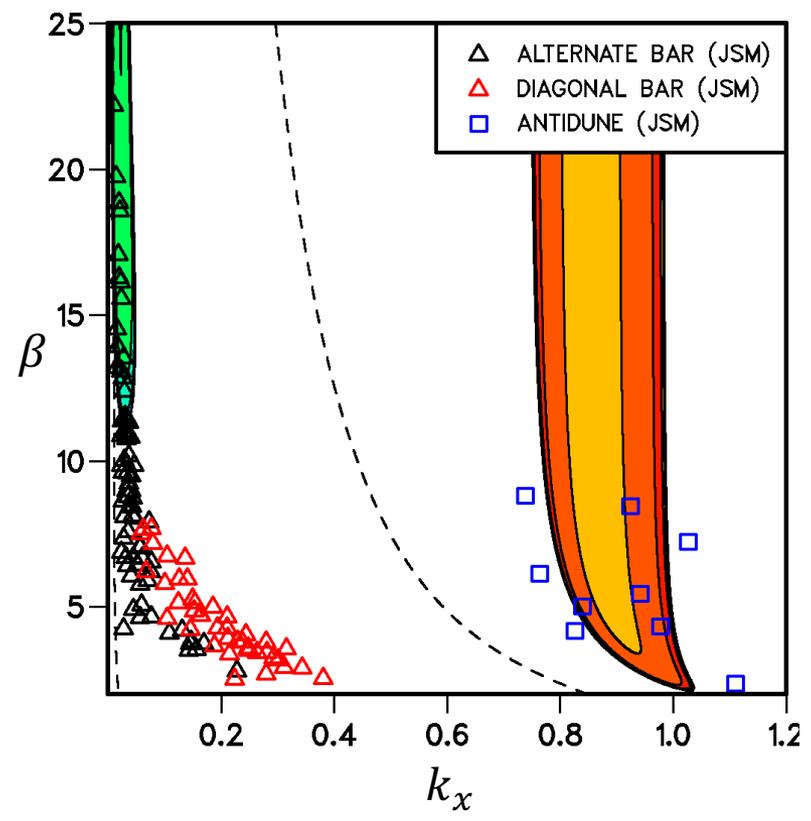
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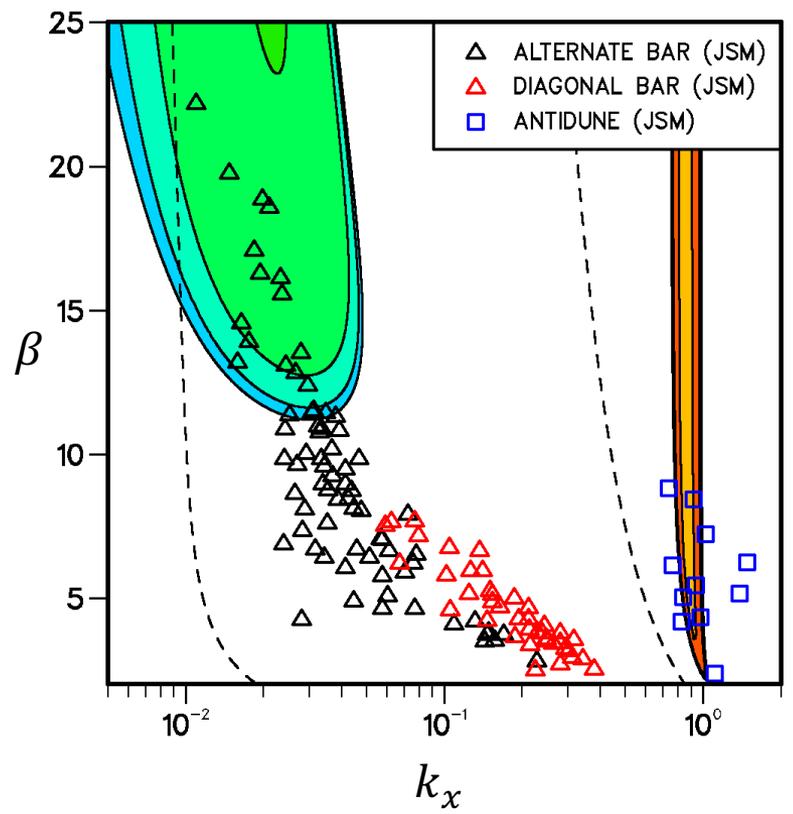
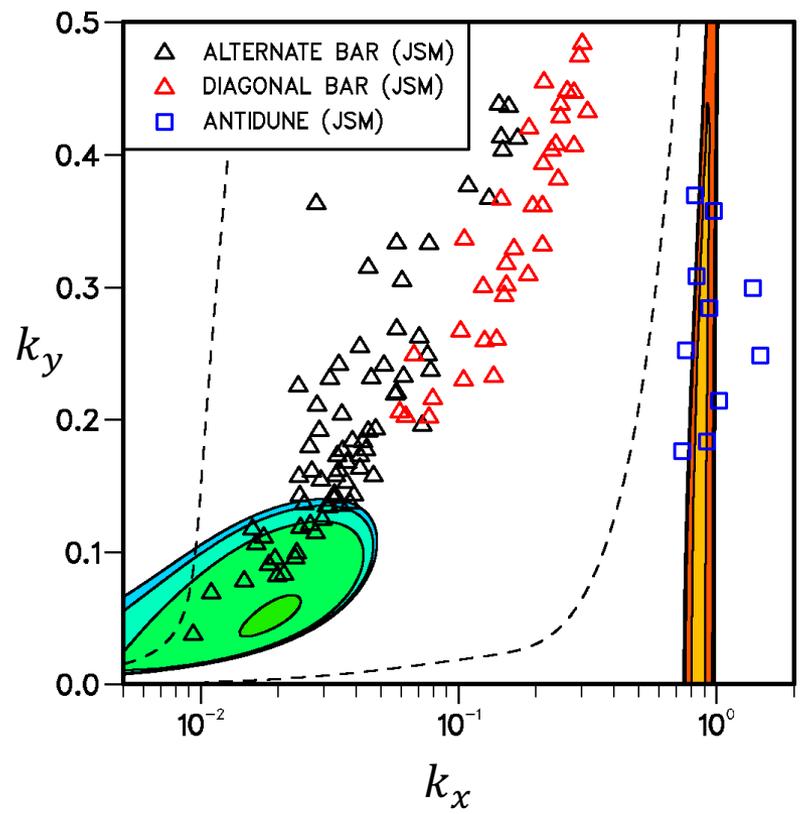


# Alternate Bars: 3D linear stability

$$\frac{\vartheta}{\vartheta_C} = 3 \quad C = 13$$



$$Fr = 1 \quad d = 0.025$$



$$k_y = \frac{\pi}{2\beta}$$

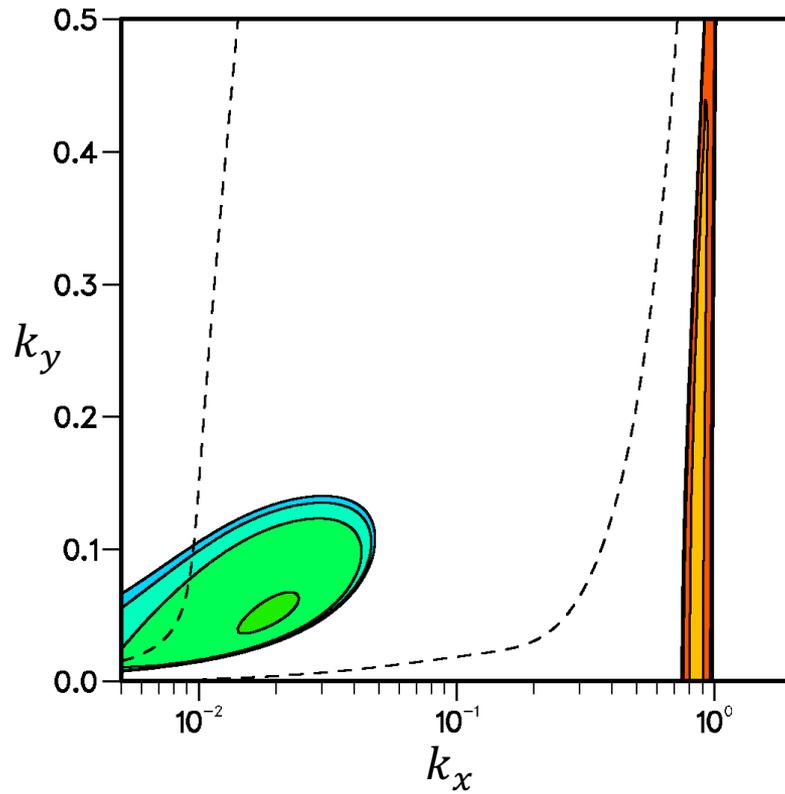




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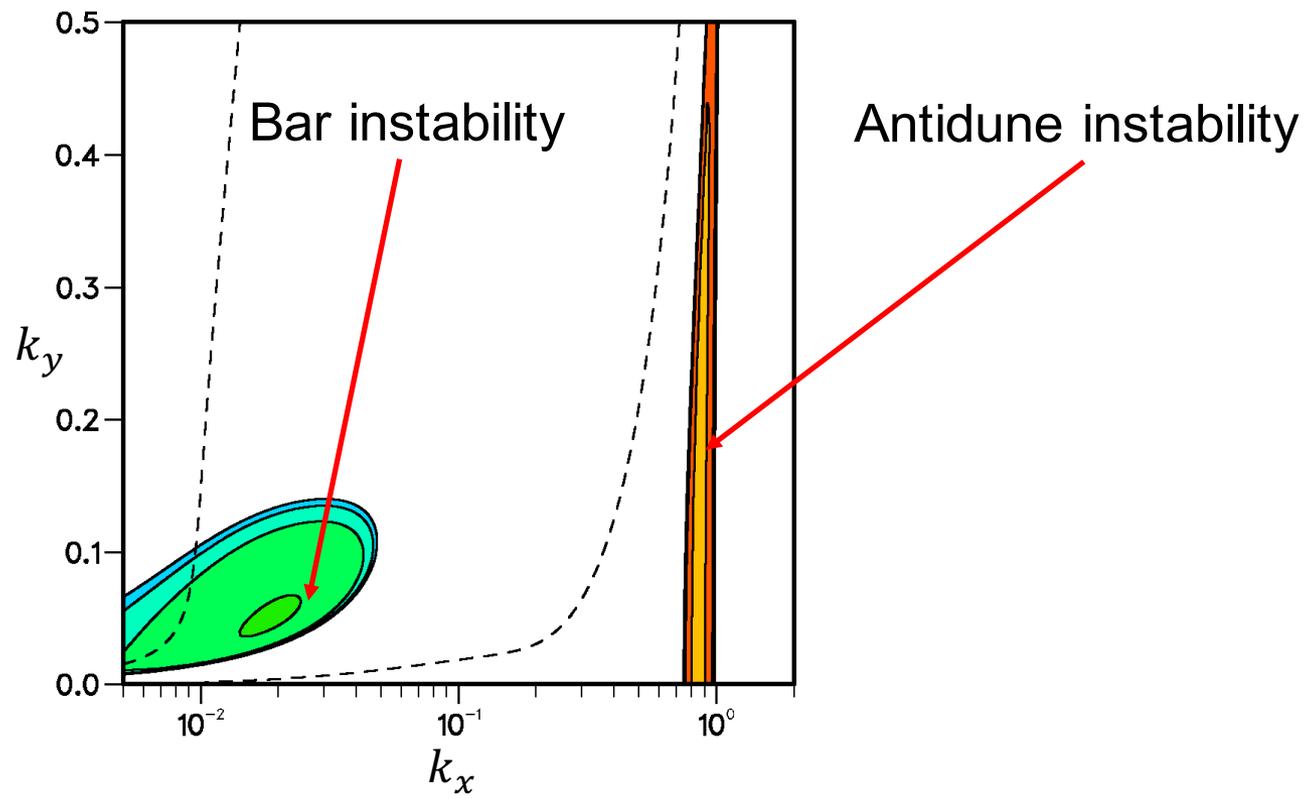


Alternate Bars: 3D linear stability

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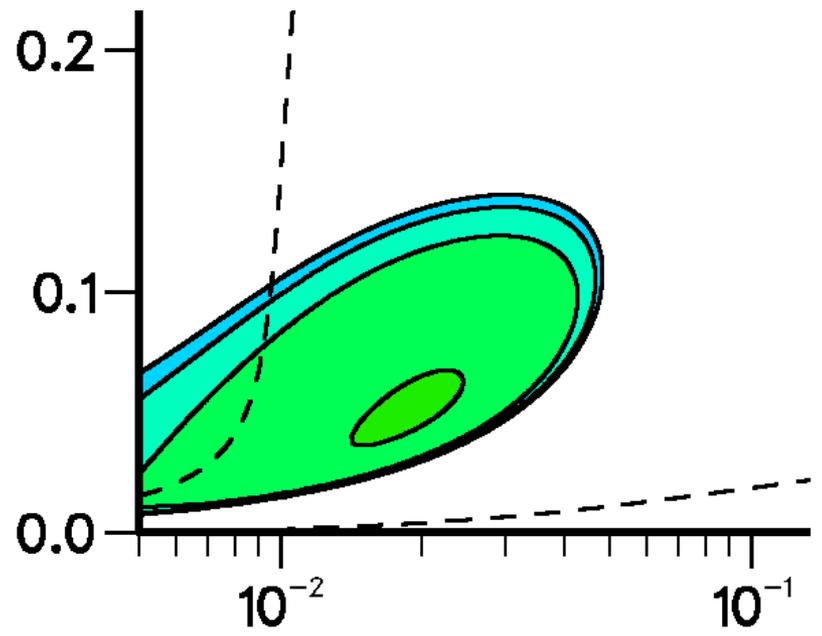
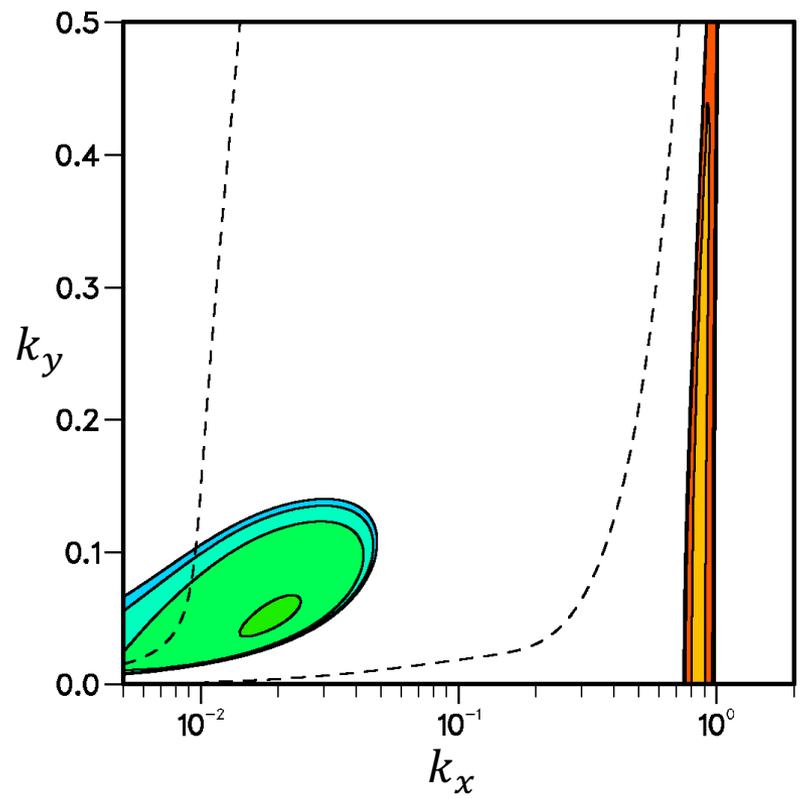


Alternate Bars: 3D linear stability

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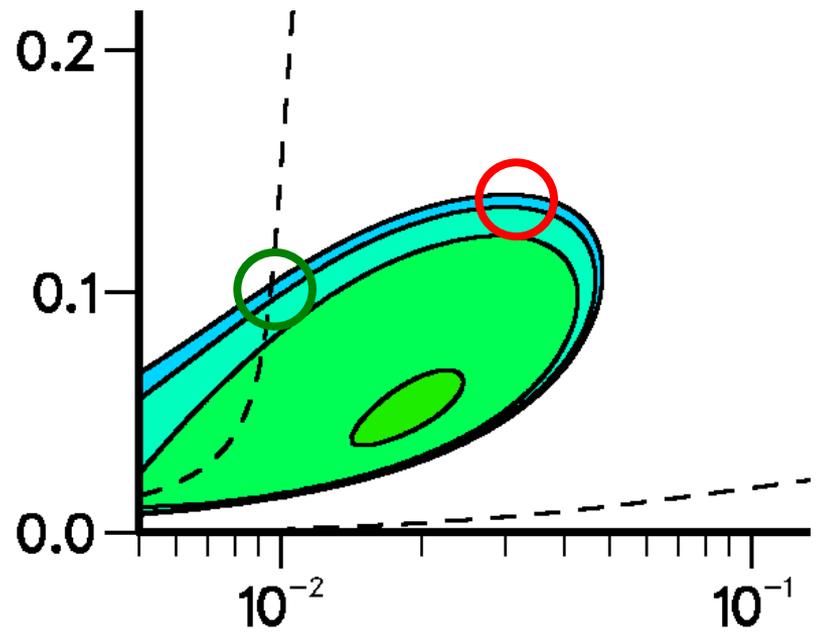
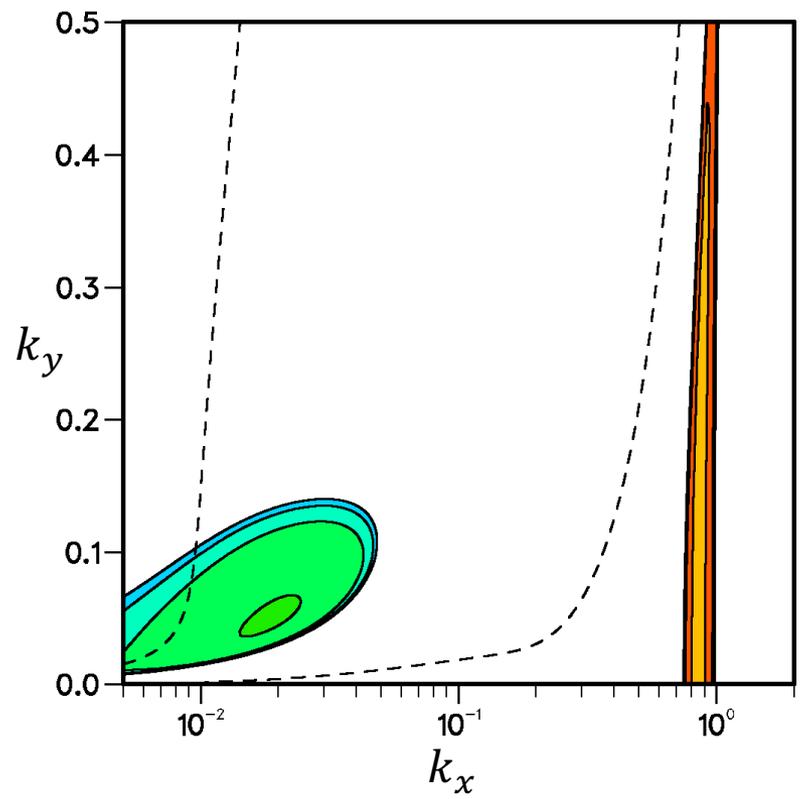


Alternate Bars: 3D linear stability

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-  Critical conditions -  $(\lambda_C, \beta_C)$
-  Resonant conditions -  $(\lambda_R, \beta_R)$



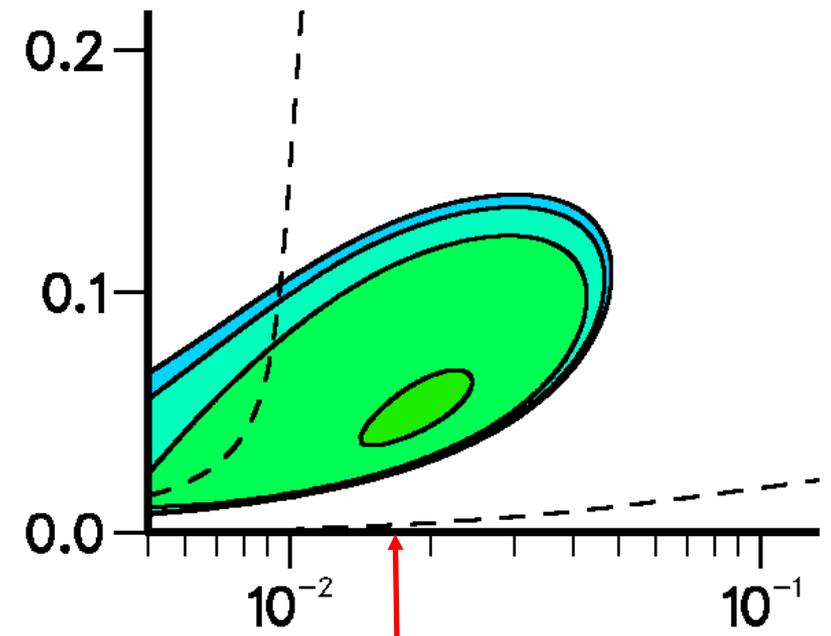
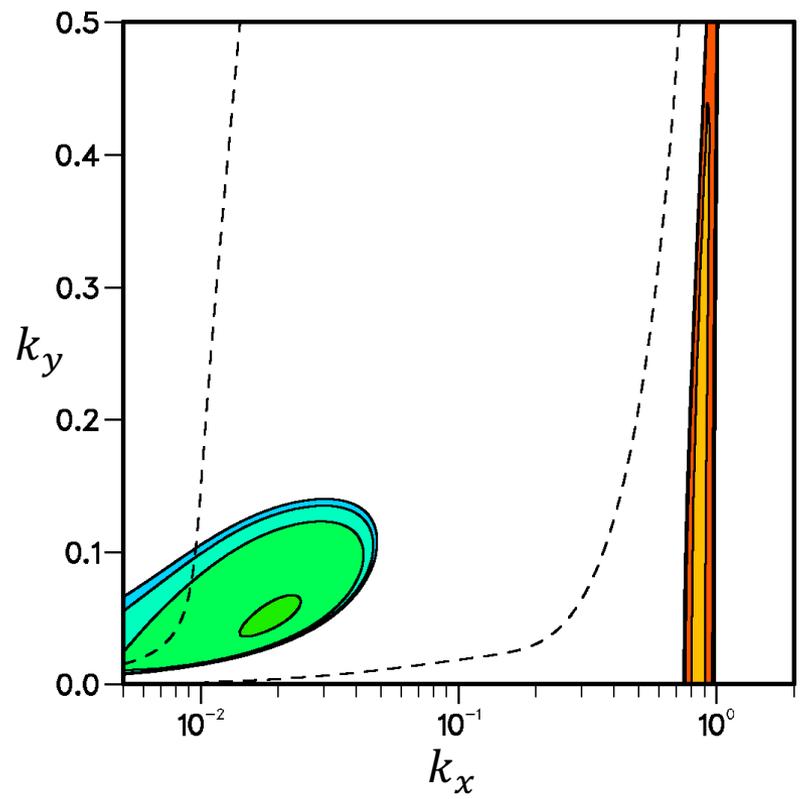


Alternate Bars: 3D linear stability

$$\frac{\vartheta}{\vartheta_C} = 3 \quad C = 13$$



$$Fr = 1 \quad d = 0.025$$



2D disturbances are stable!

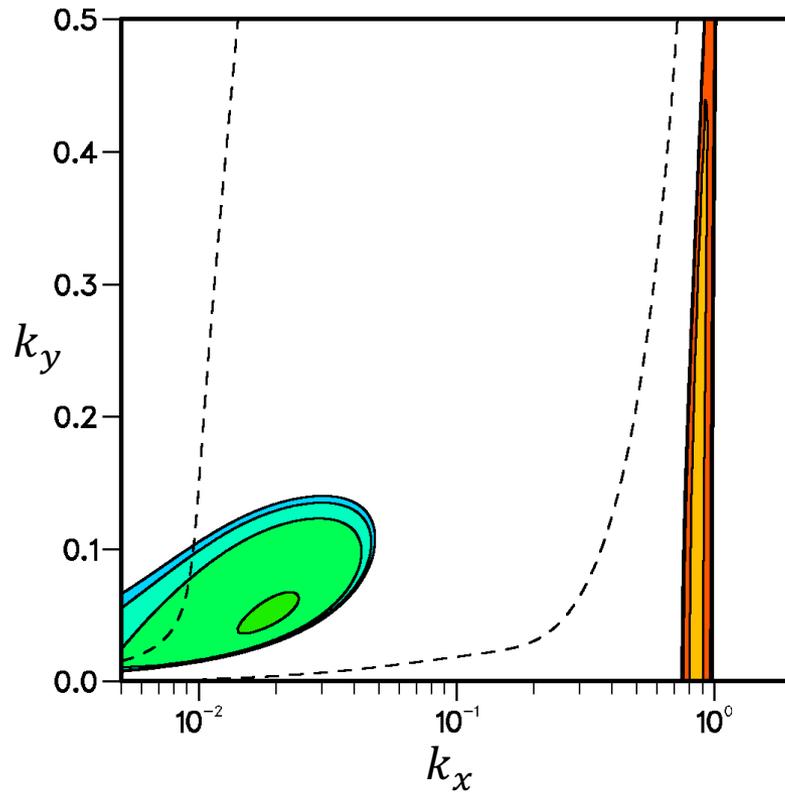




$$\frac{\vartheta}{\vartheta_C} = 3 \quad C = 13$$

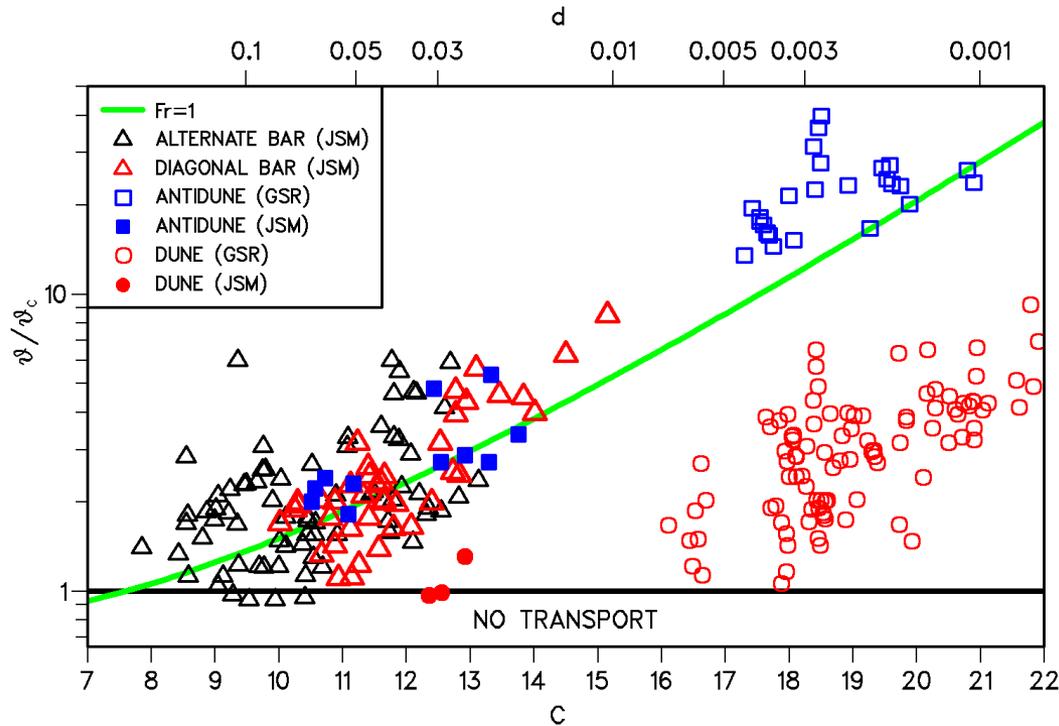


$$Fr = 1 \quad d = 0.025$$



2D disturbances are the most unstable!





$$C = \frac{1}{\kappa} \ln \left( \frac{11.09}{2.5d} \right)$$

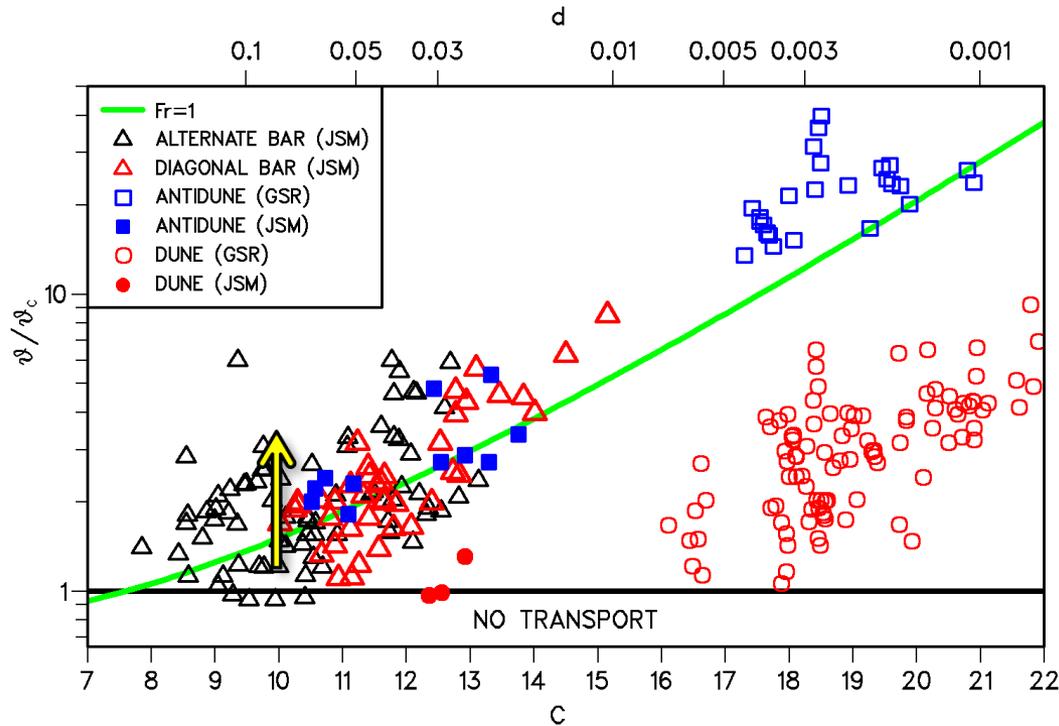
$$\vartheta \cong 0.14 \frac{Fr^2 e^{\kappa C}}{C^2}$$

## Experimental data sets

JSM: Jaeggi (1984), Sukegawa (1971), Muramoto & Fujita (1978)

GSR: Guy, Simons & Richardson (1966)





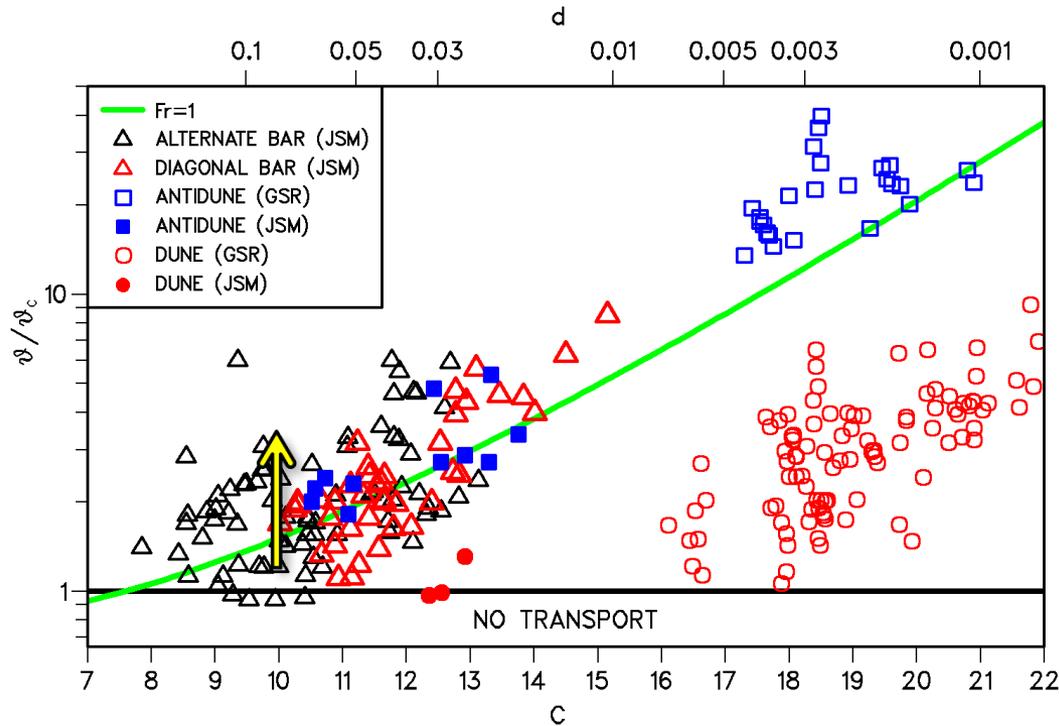
$$C = \frac{1}{\kappa} \ln \left( \frac{11.09}{2.5d} \right)$$

$$\vartheta \cong 0.14 \frac{Fr^2 e^{\kappa C}}{C^2}$$

Moving up along a vertical line means (same sediment):

- $C$  constant  $\Rightarrow d$  constant  $\Rightarrow D^*$  constant
- $\theta$  increases  $\Rightarrow Fr$  increases  $\Rightarrow U^*$  increases





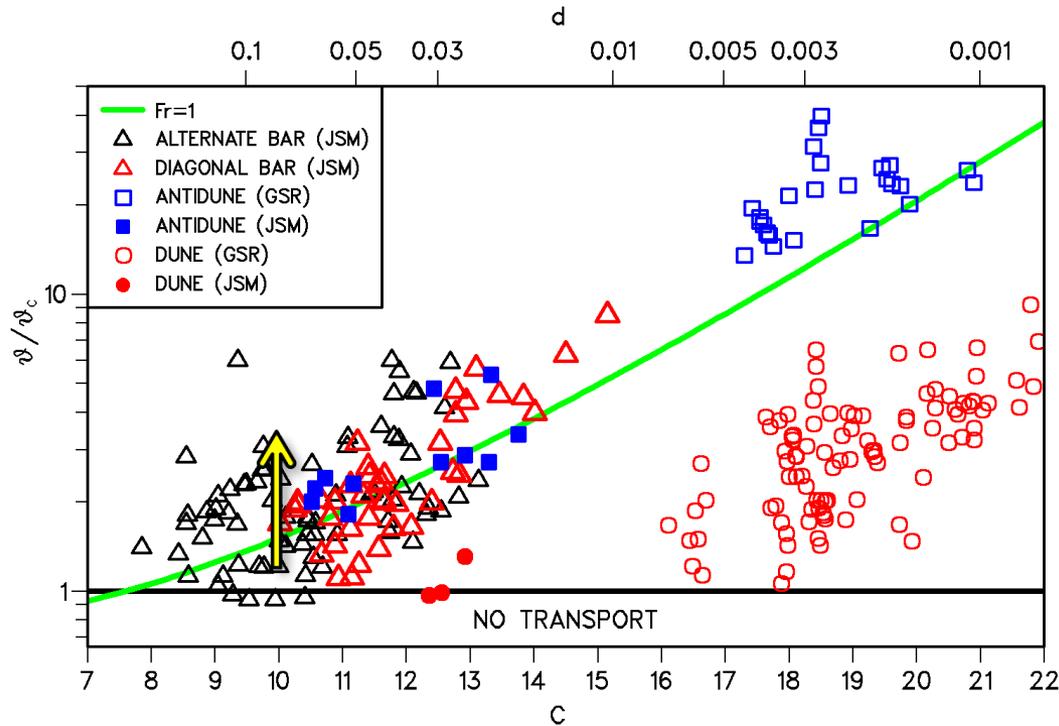
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- }  $\Rightarrow S \propto Fr^2$  increases





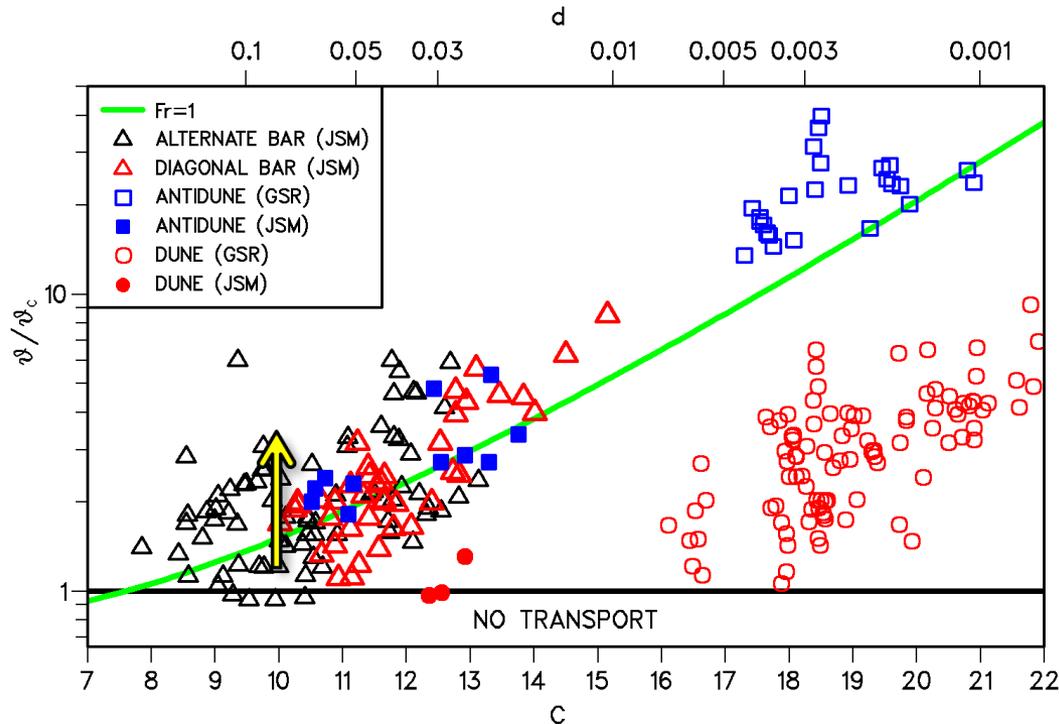
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  - Flow rate increases with constant flow depth
- }  $\Rightarrow S \propto Fr^2$  increases





$$C = \frac{1}{\kappa} \ln \left( \frac{11.09}{2.5d} \right)$$

$$\vartheta \cong 0.14 \frac{Fr^2 e^{\kappa C}}{C^2}$$

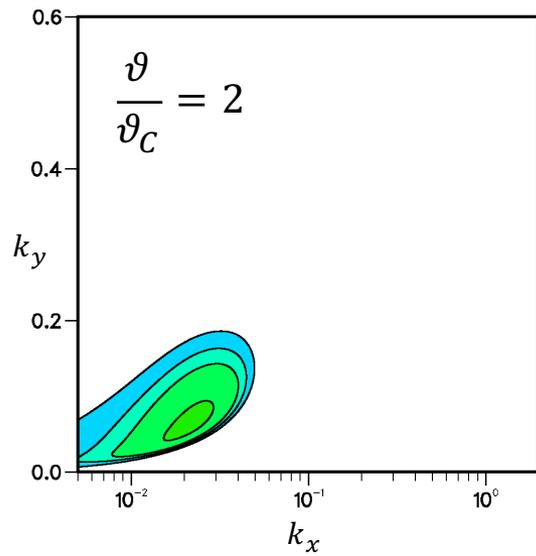
Moving up along a vertical line means (same sediment):

- $C$  constant  $\Rightarrow d$  constant  $\Rightarrow D^*$  constant
- $\theta$  increases  $\Rightarrow Fr$  increases  $\Rightarrow U^*$  increases
- Flow rate increases with constant flow depth  $\Rightarrow$  Slope increases
- **From Alternate Bars to Alternate Bars & Antidunes**



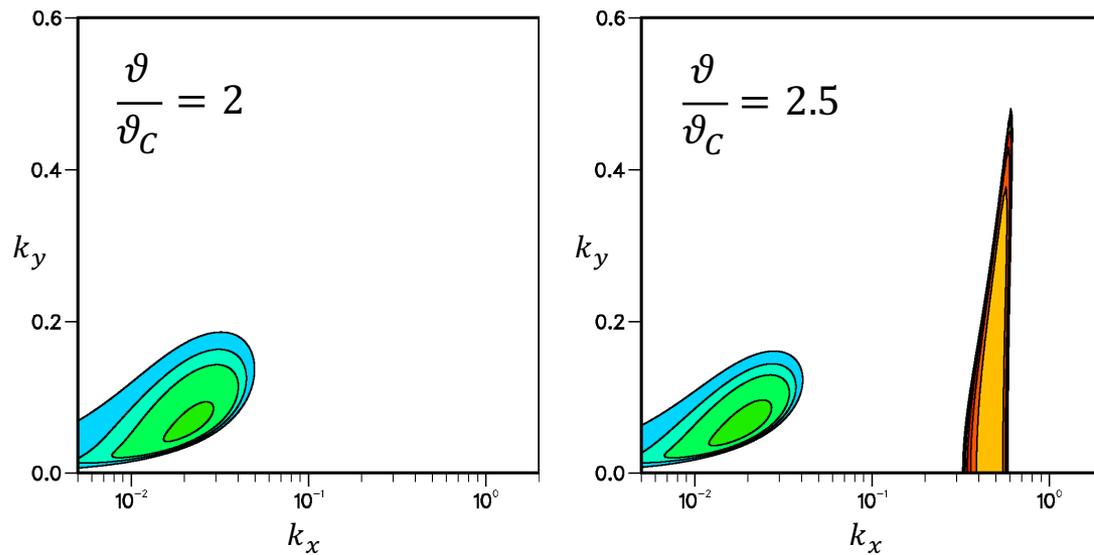


Moving up:  $C = 10$



As  $\vartheta$  increases:

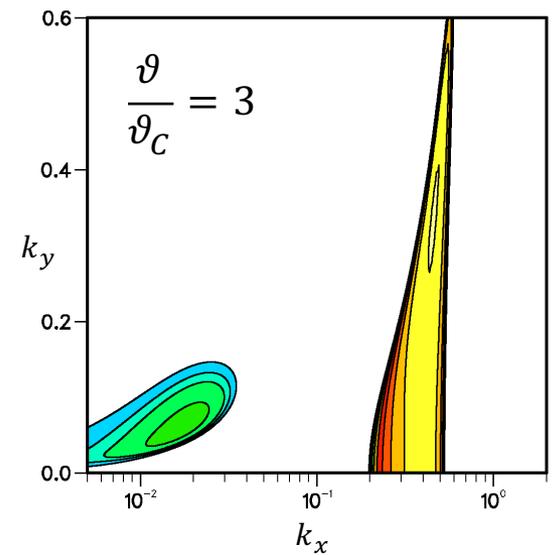
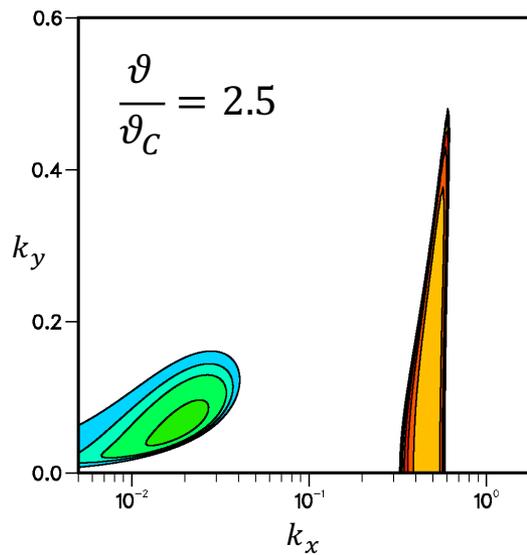
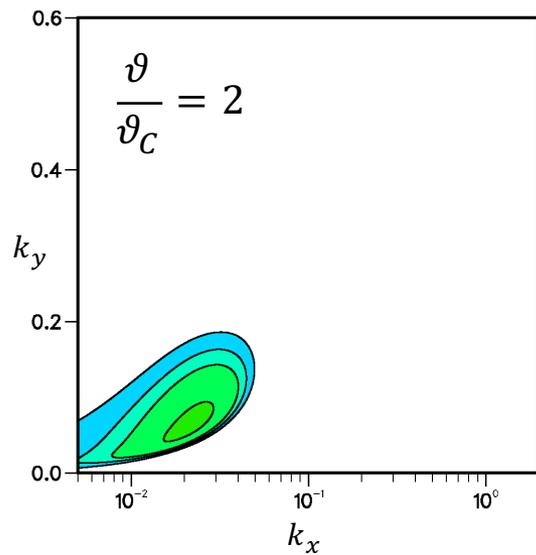




As  $\vartheta$  increases:

- A new region of instability appears: 2D antidunes;
- The regions of instability for bars and antidunes are distinct;
- Bars and antidunes linearly coexist: bars are less unstable;

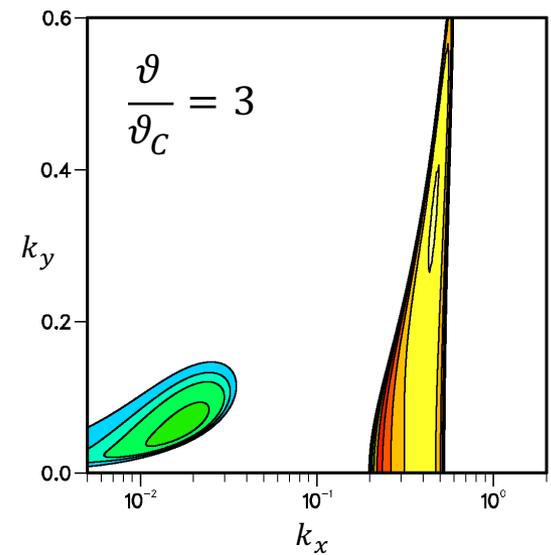
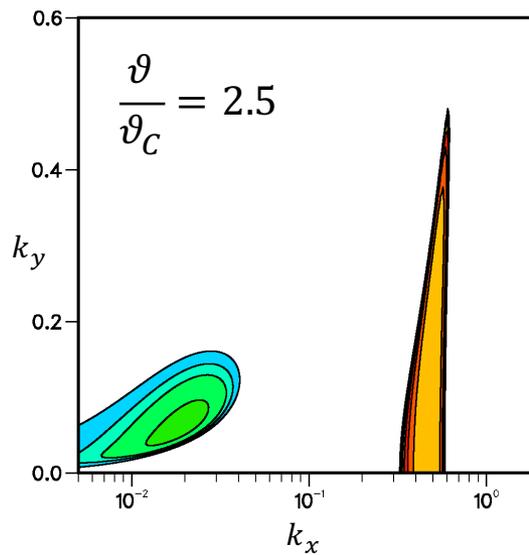
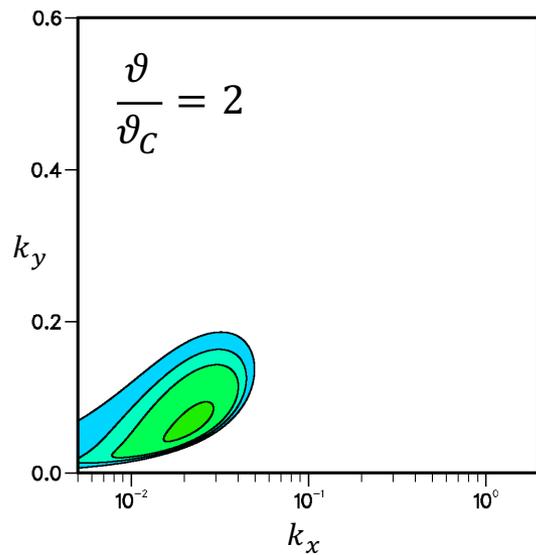




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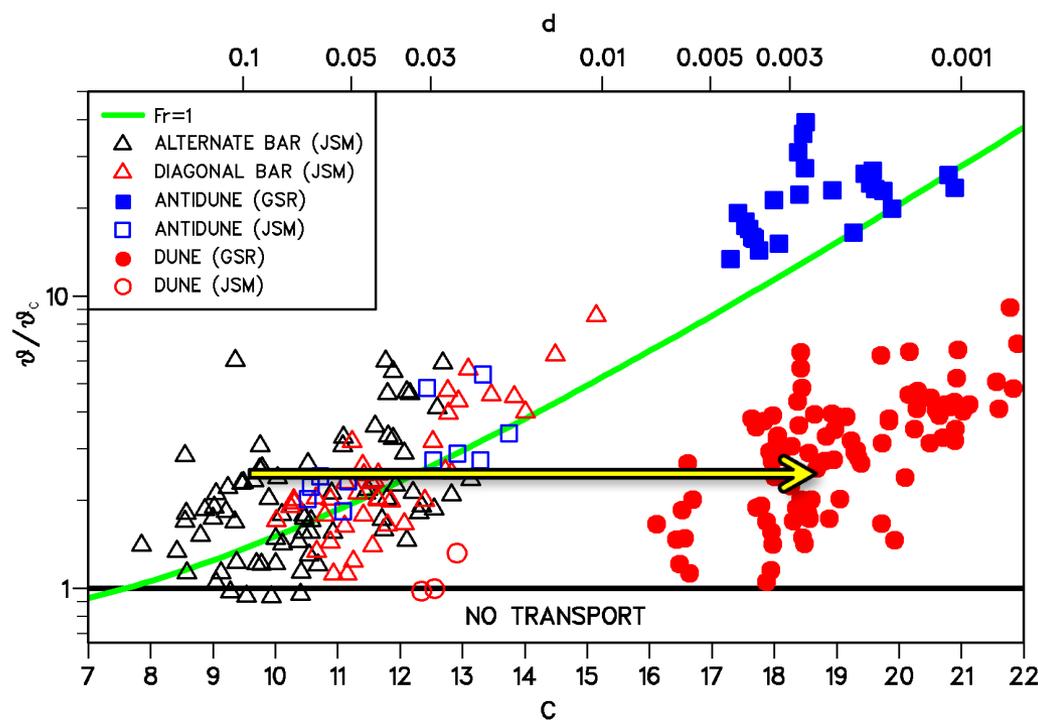




As  $\vartheta$  increases:

- A new region of instability appears: 2D antidunes;
- The regions of instability for bars and antidunes are distinct;
- Bars and antidunes linearly coexist: bars are less unstable;
- 3D antidunes become the most unstable;
- This results in a transition from 2D to 3D antidunes;





$$C = \frac{1}{\kappa} \ln \left( \frac{11.09}{2.5d} \right)$$

$$\vartheta \cong 0.14 \frac{Fr^2 e^{\kappa C}}{C^2}$$

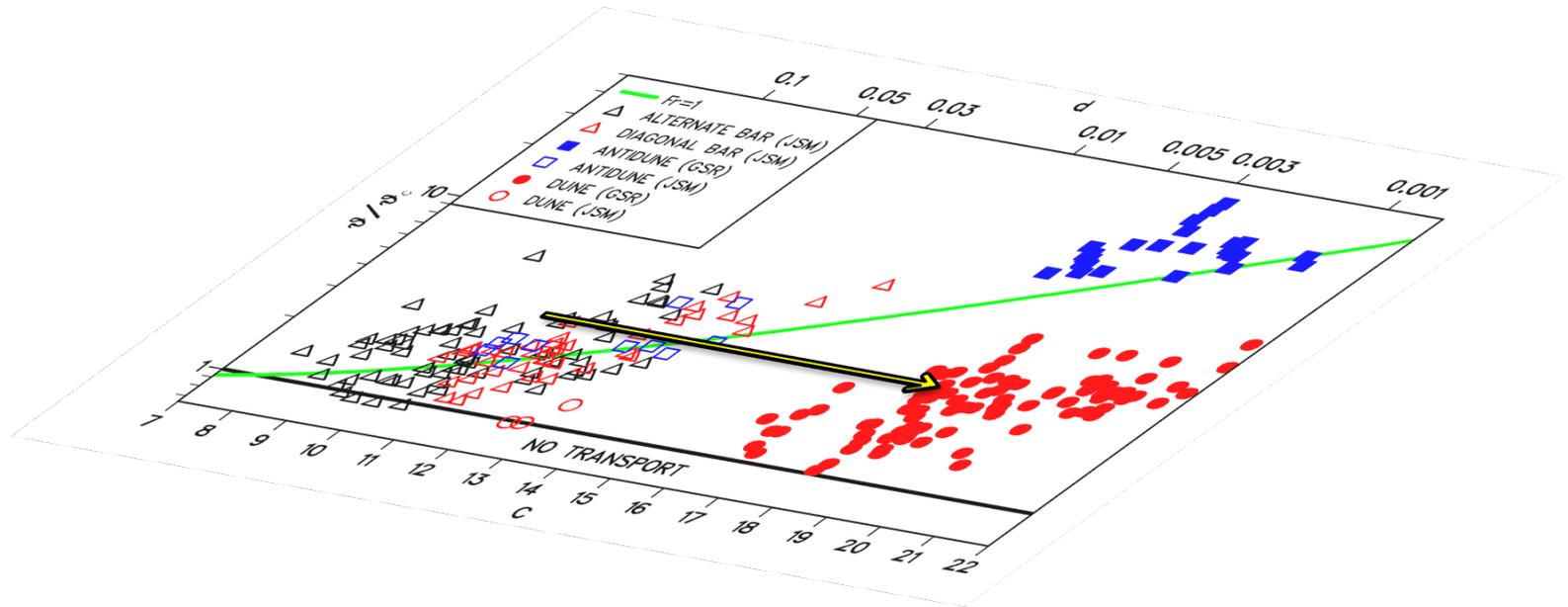
Moving right along a horizontal line means (same sediment):

- $C$  increases  $\Rightarrow d$  decreases  $\Rightarrow D^*$  increases
  - $\theta$  constant  $\Rightarrow U^* \propto C$  increases
  - Flow rate increases but  $Fr$  decreases
  - From Alternate Bars to Diagonal Bars to 2D Dunes
- }  $\Rightarrow S \propto 1/D^*$  decreases



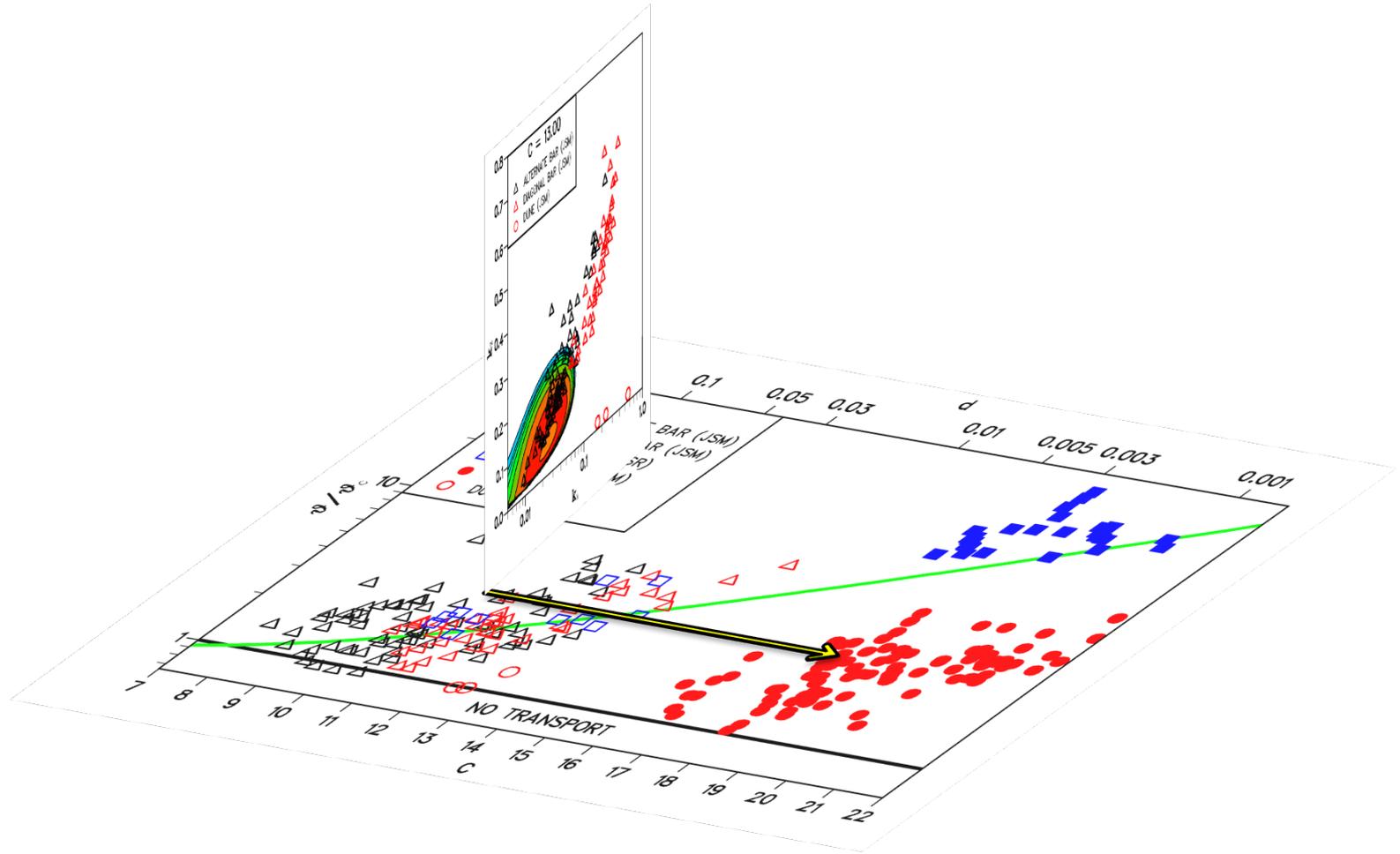


# Regime diagram: bars & dunes



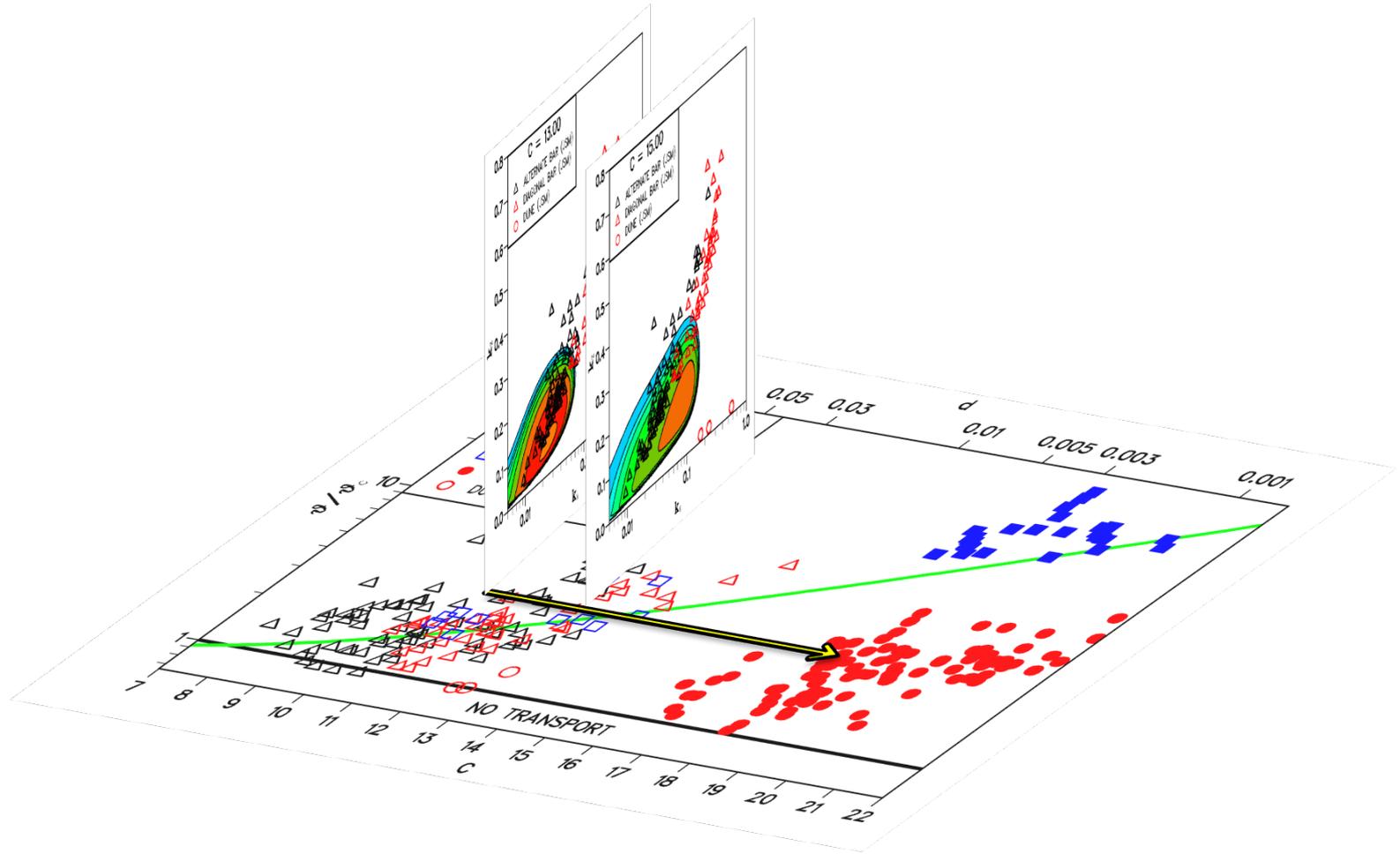


# Regime diagram: bars & dunes



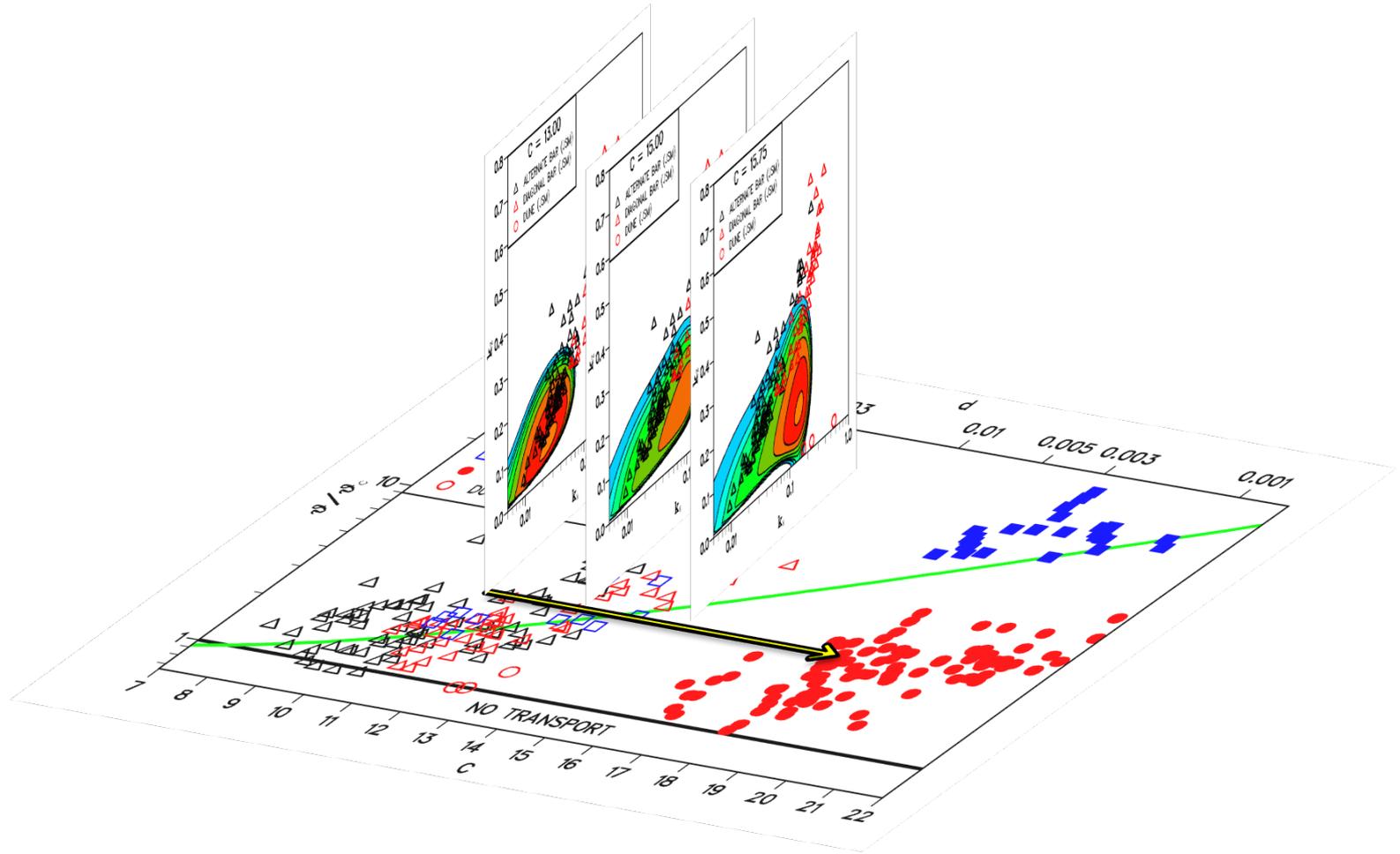


# Regime diagram: bars & dunes



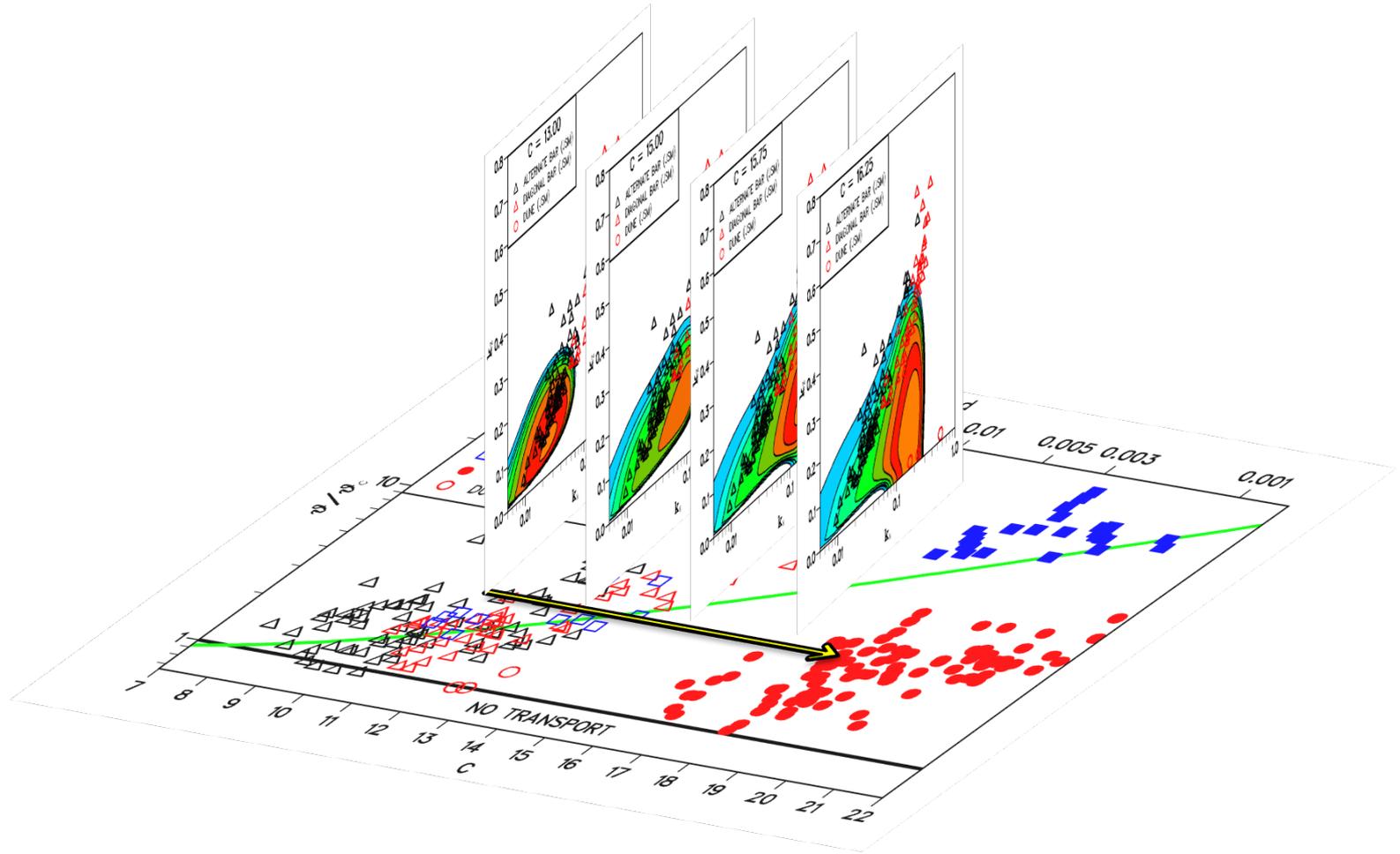


# Regime diagram: bars & dunes



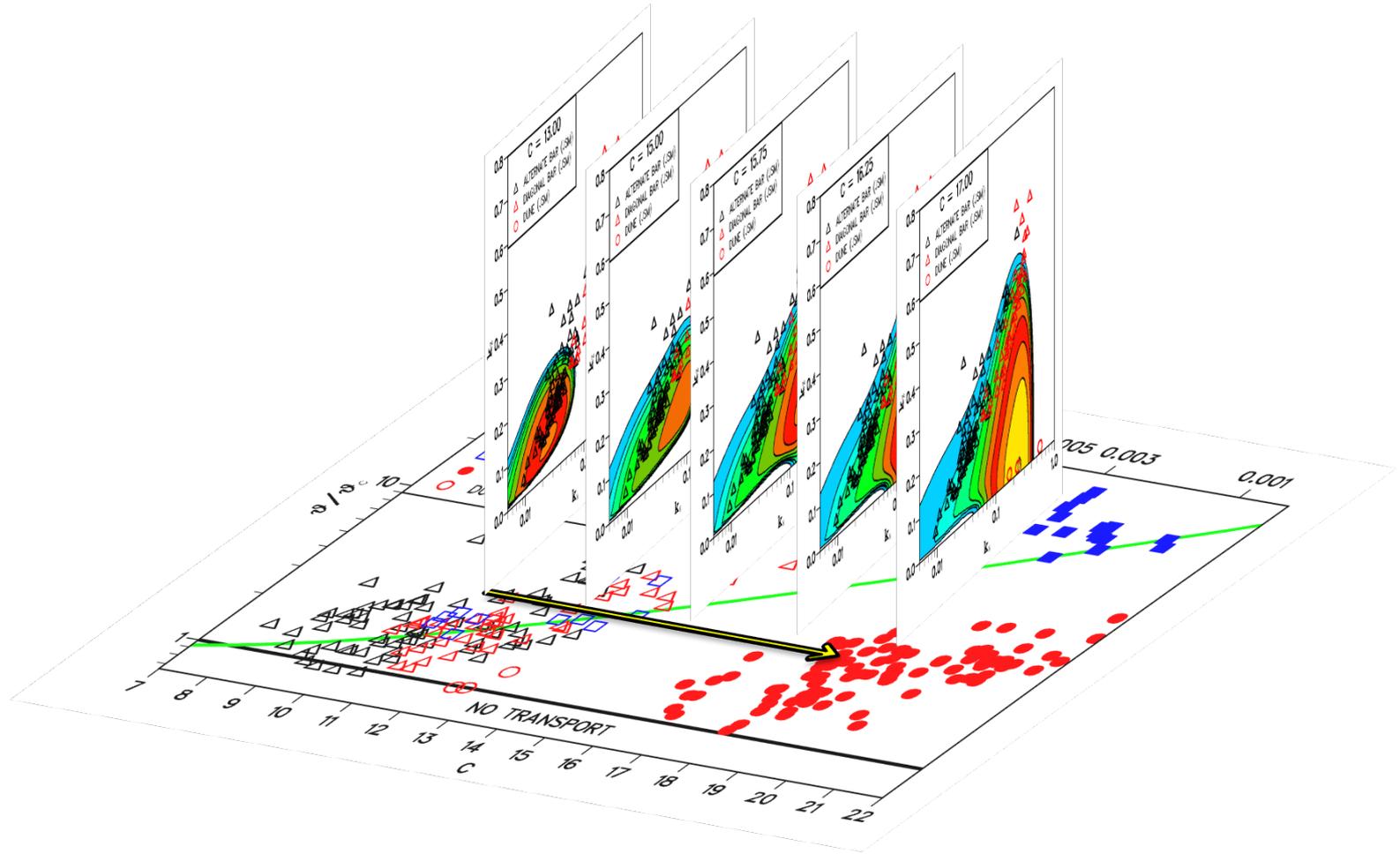


# Regime diagram: bars & dunes



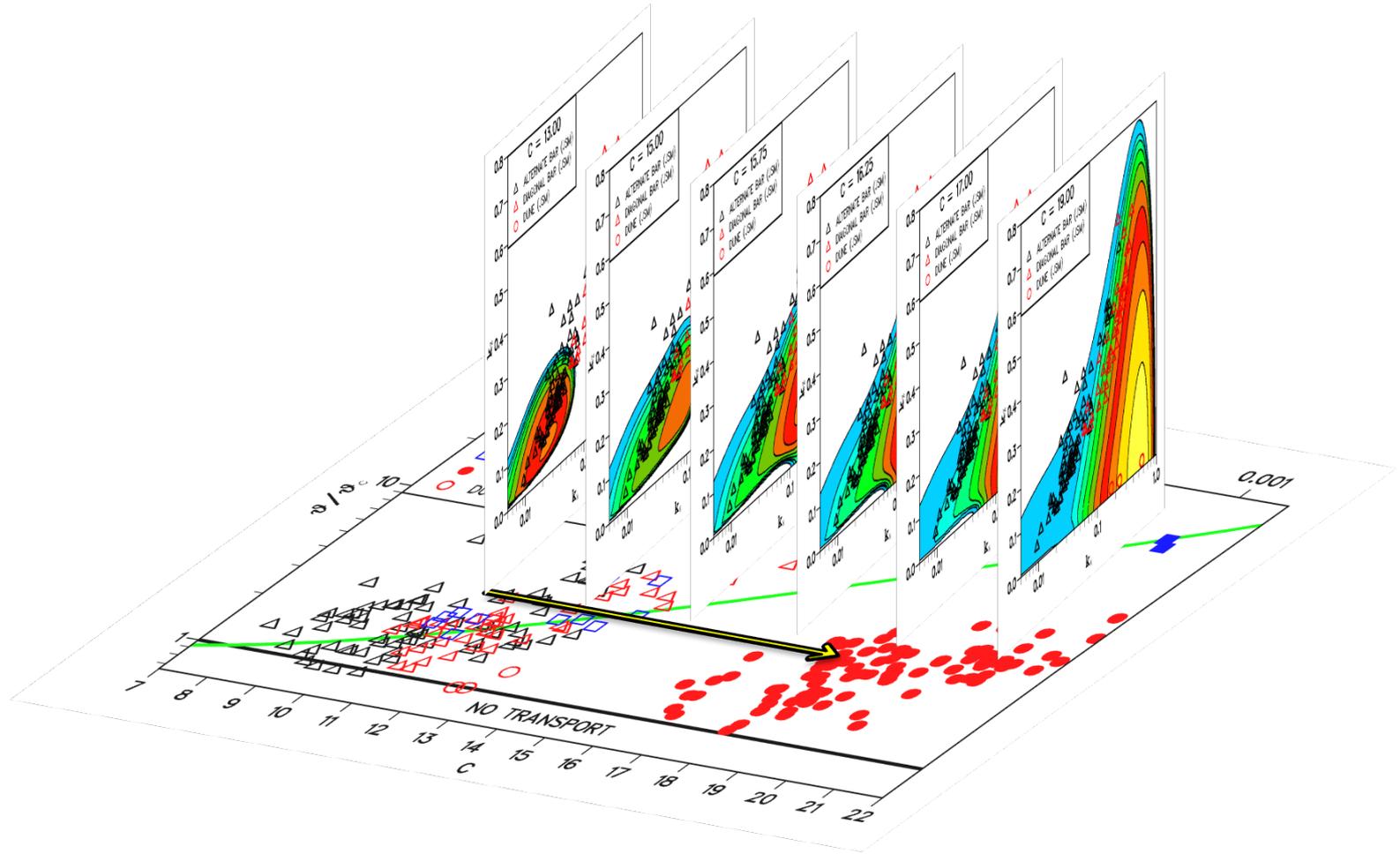


# Regime diagram: bars & dunes



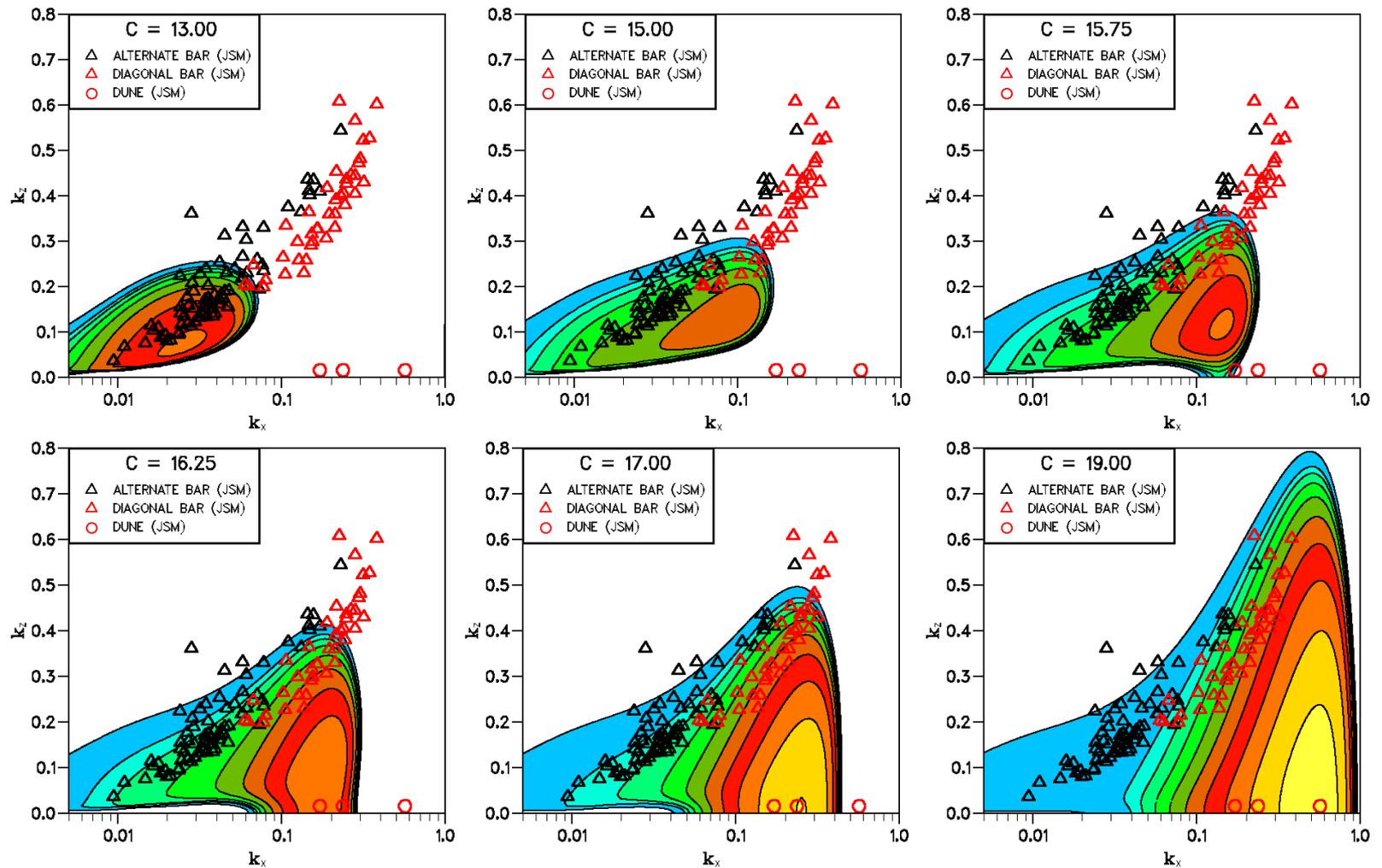


# Regime diagram: bars & dunes





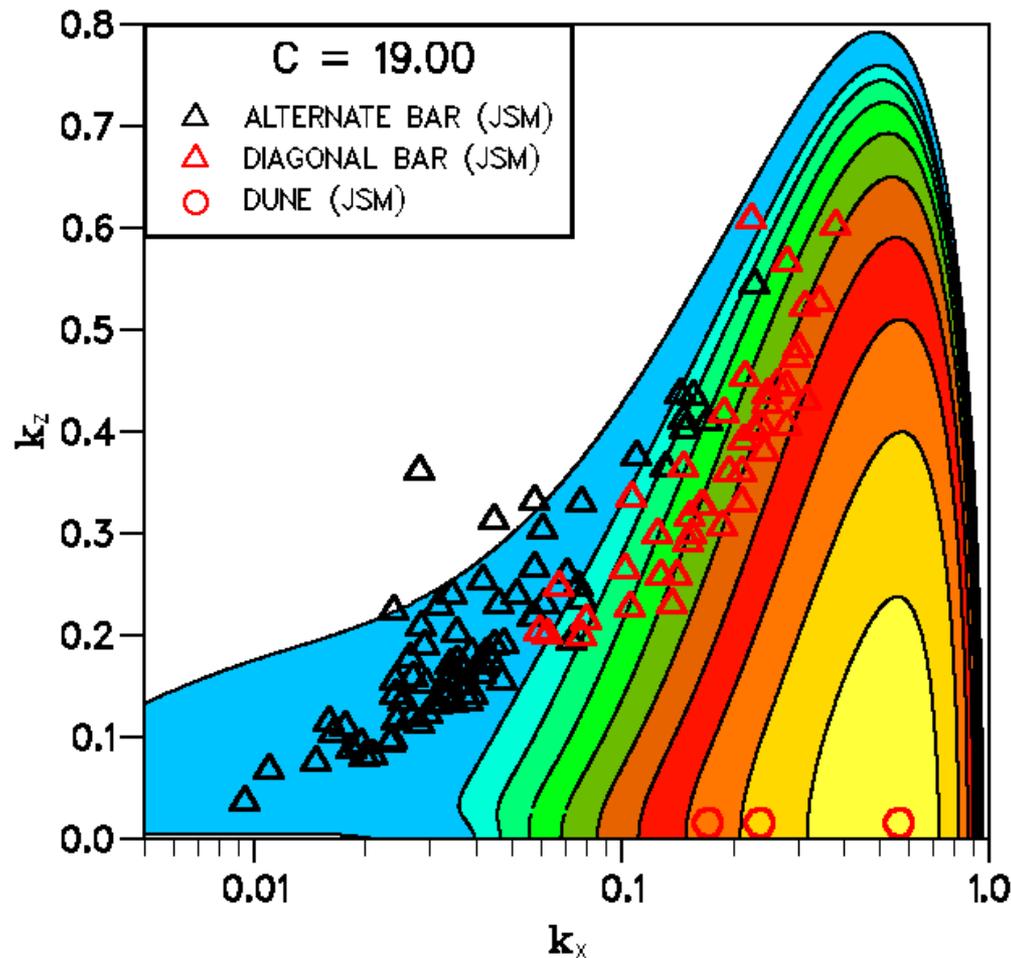
# Dune-ripple transition



As  $C$  increases:

- the longitudinal wavenumber of maximum growth rate increases: bars become shorter;
- 2D disturbances become unstable but are less unstable than 3D disturbances;
- 2D disturbances become the most unstable;
- This results in a transition from 3D bars to 2D dunes via diagonal bars (3D dunes);





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## Two-Phase Continuum Models for Geophysical Particle-Fluid Flows

Max Planck Institute for the Physics of Complex Systems

14 March - 15 April 2016