

Pathway and free energy barriers of rupturing the hemi-fission intermediate under tension

A Wang-Landau sampling study

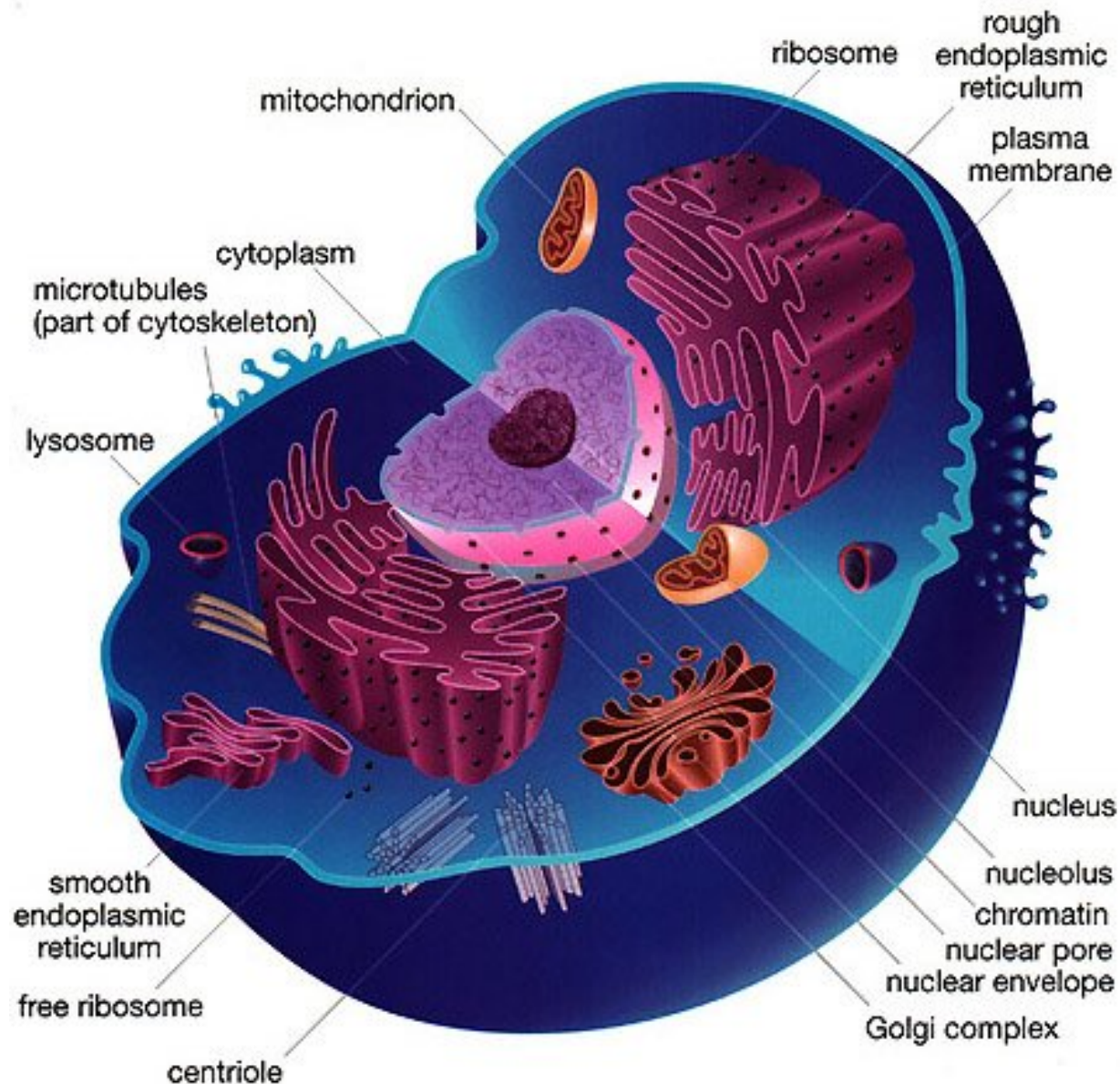
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Georg-August-University Göttingen*

Workshop of “Bridging time-scale techniques and their
applications in atomistic computational science”
at Max-Planck-Institute for the Physics of Complex Systems

12-15 Sep. 2016, Dresden

Cell membranes

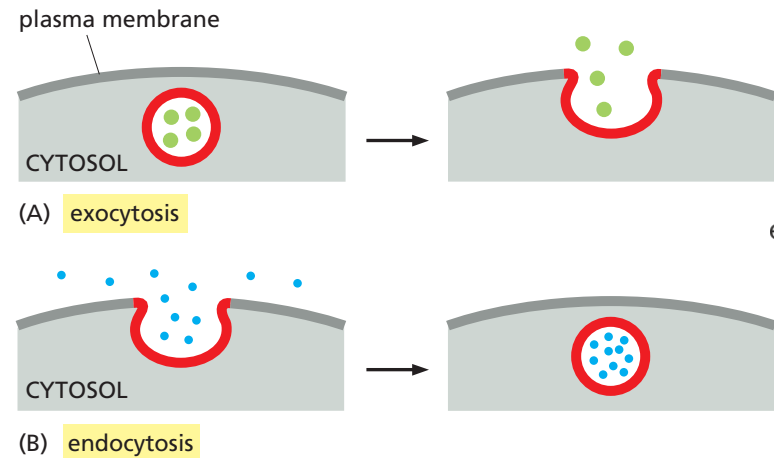


Everything you see here is membrane!

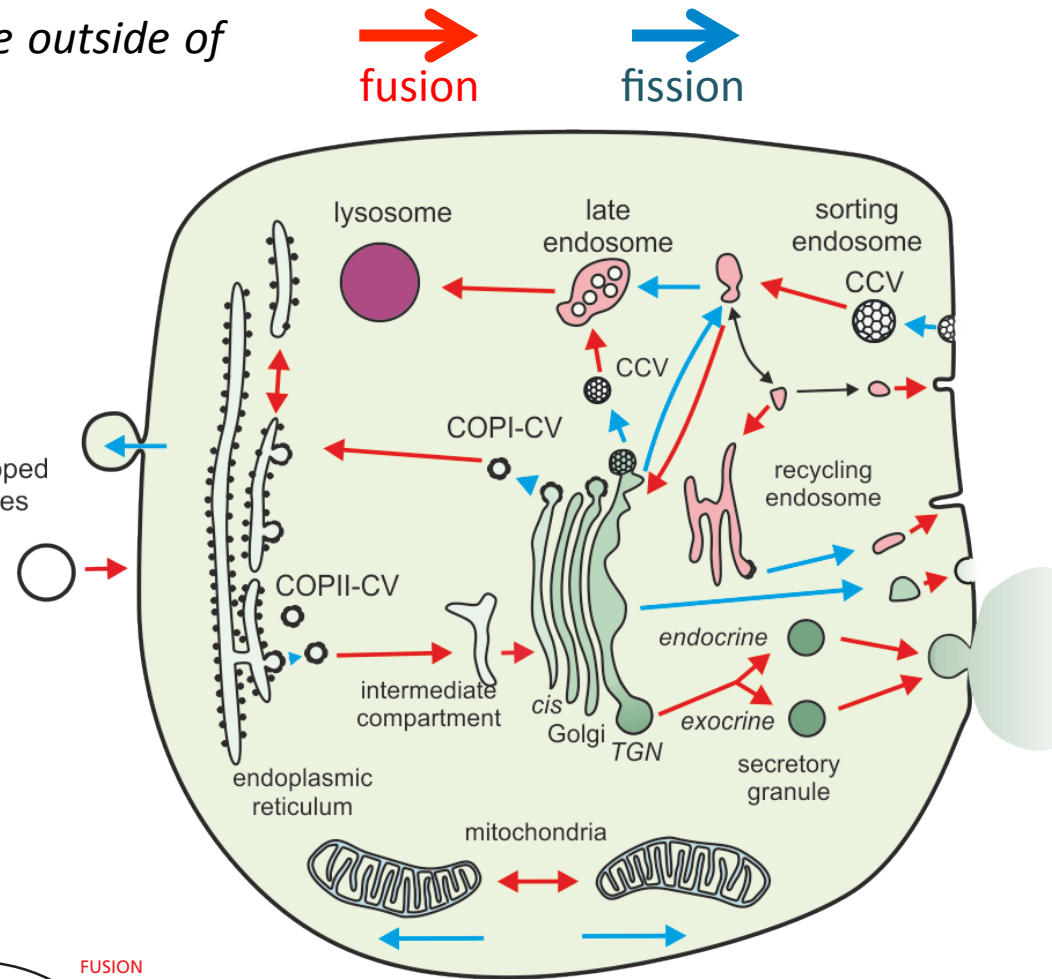
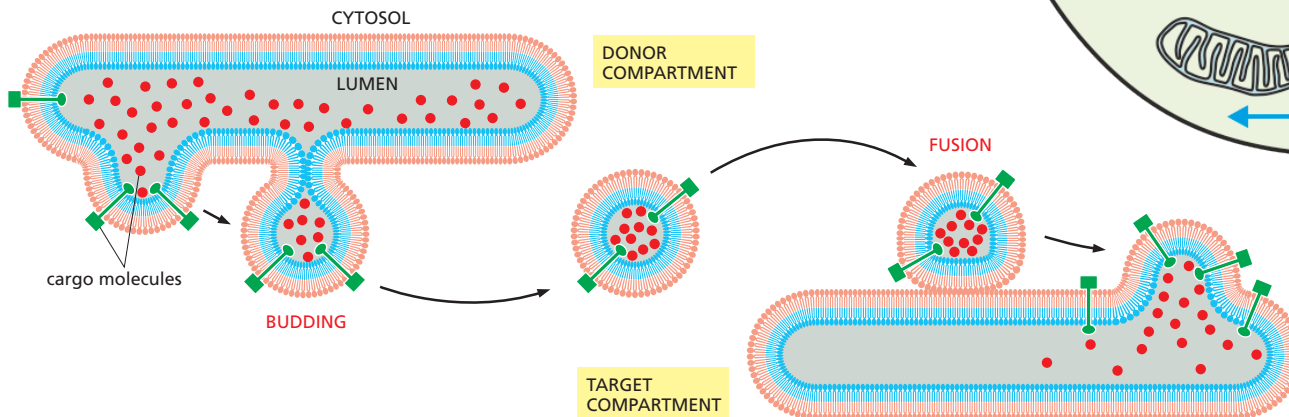
Cell must eat via fusion and fission

To make itself alive, cell must *exchange materials* between intracellular compartments or the outside of the cell.

✧ Exocytosis and endocytosis

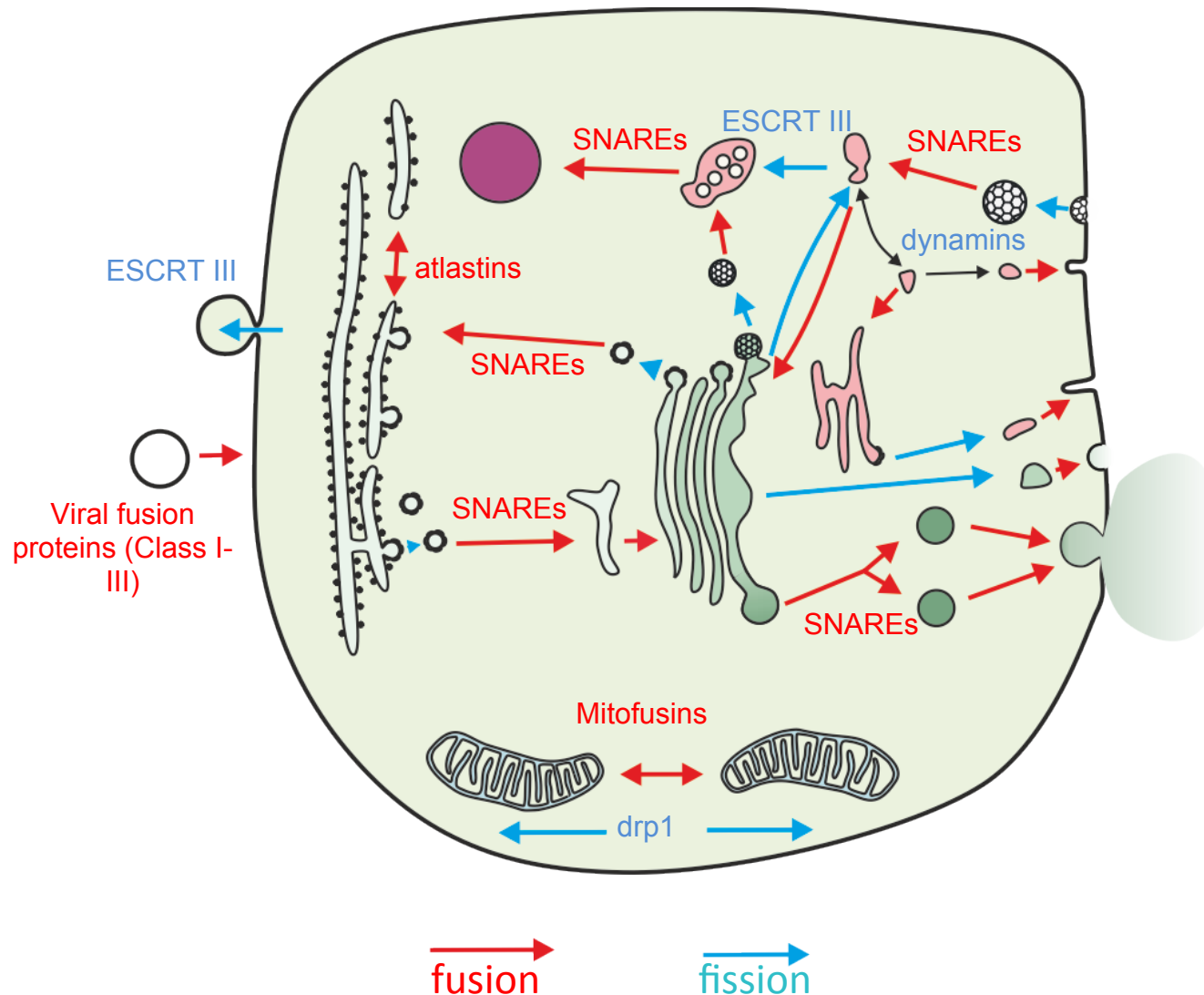


✧ Fission and fusion in vesicle transport



Alberts, et al.
Molecular biology of the cell, 6th Ed.
Garland Science (2015).

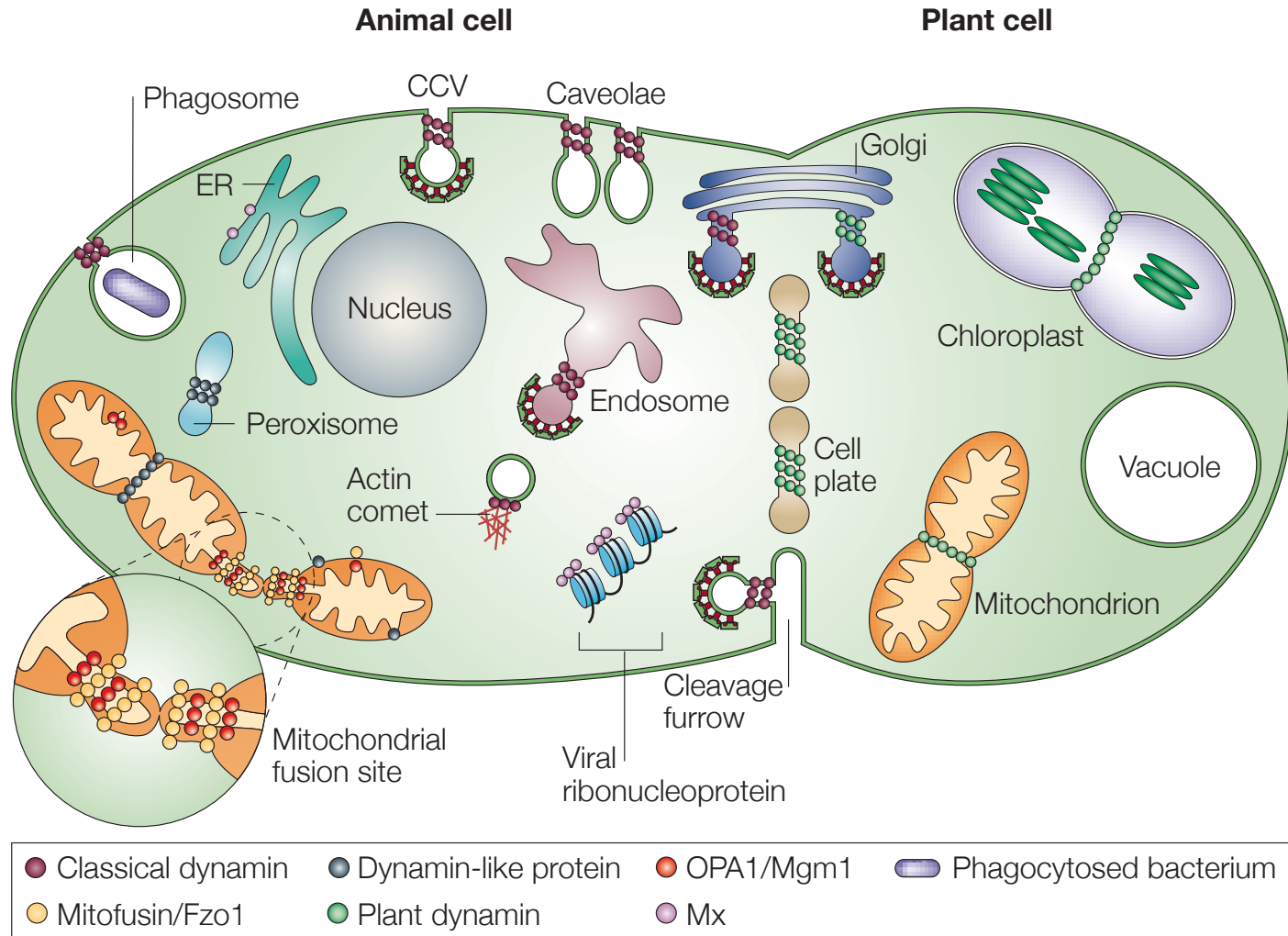
Membrane fusion and fission



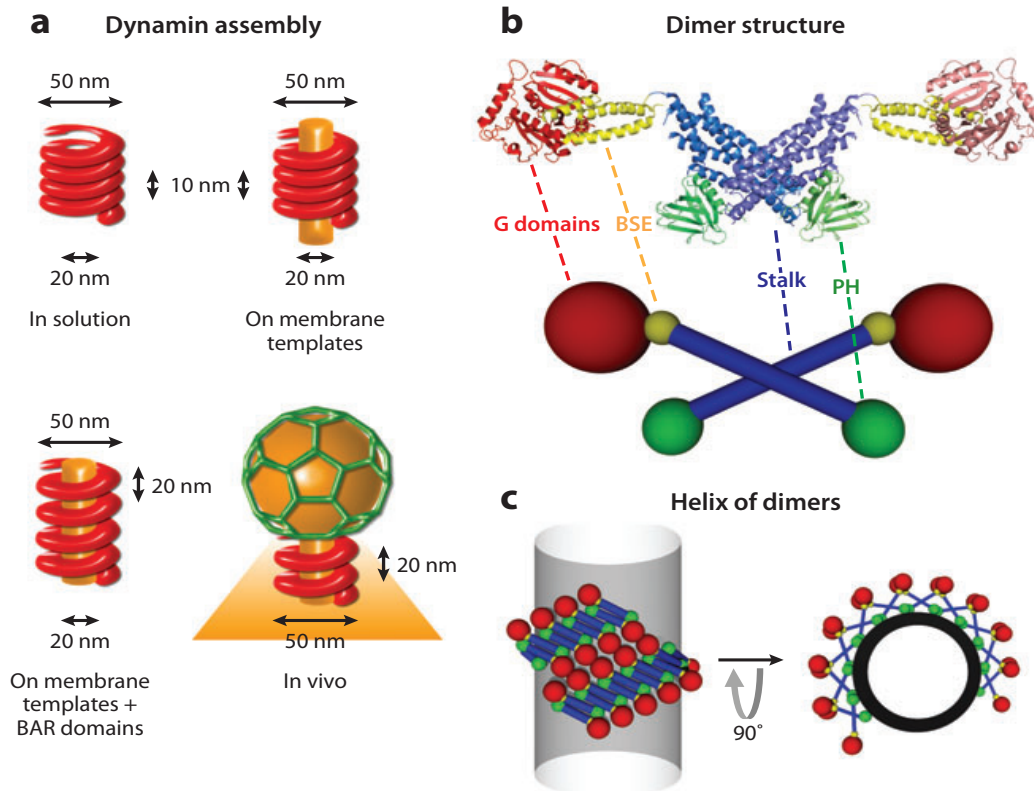
Membrane fusion and fission are *mediated by proteins*, e.g. SNAREs for fusion, and dynamin for fission.

Protein-mediated membrane fission

Membrane fission: A process by which a bud separates from a lipid membrane, and which is mostly mediated by (**dynammin-family**) proteins.



Self-assembly of dynamin around the neck of a budding vesicle



What we know:

- ✧ Existence as tetramer in solution
- ✧ Self-assembly (into rings/helices)
- ✧ Inner diameter of helices independent from membrane template.
- ✧ Membrane tube is constricted by PH domain insertion of dynamin rings.
- ✧ Constriction is necessary, but insufficient for completing fission.

Nat. Rev. Mol. Cell Biol. 5, 133 (2004)
Annu. Rev. Cell Dev. Biol. 27, 79 (2011)
Annu. Rev. Biochem. 81, 407 (2012)
Annu. Rev. Biophys. 42, 629 (2013)

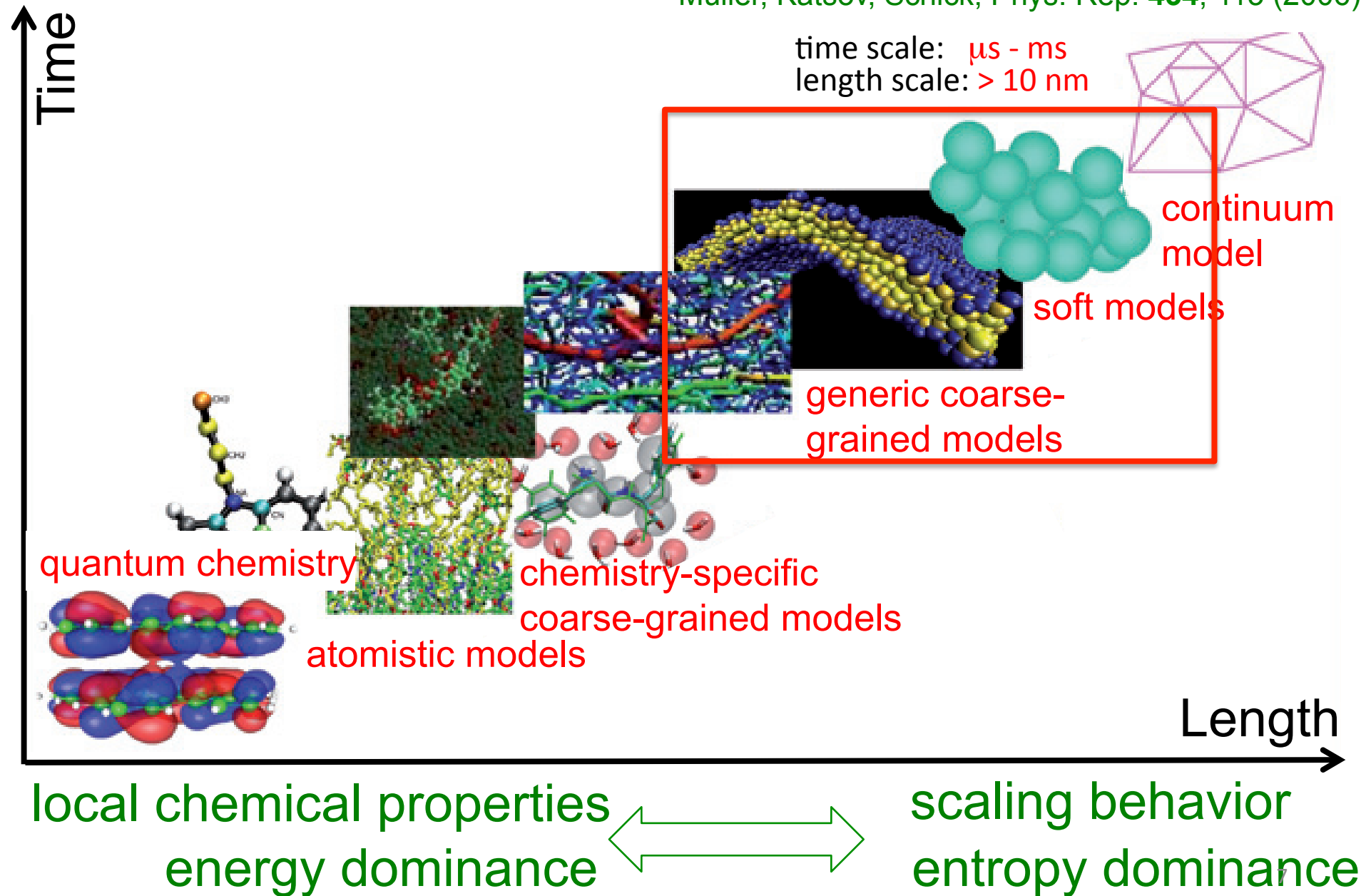
Questions: 1. *Role* of dynamin's *constriction*.

2. Pathway and *free energy landscapes* of fission process.

Goals: Study and understand dynamin's role in membrane fission via *computer simulation*.

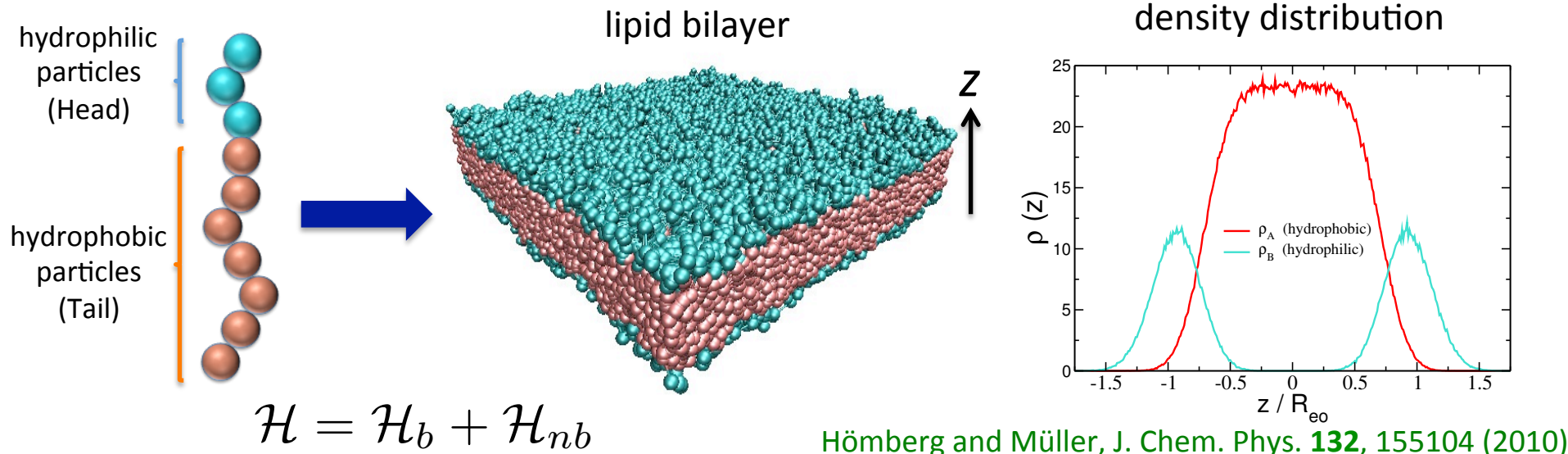
Characteristic length and time scales in soft matter

Müller, Katsov, Schick, Phys. Rep. **434**, 113 (2006)



A soft, solvent-free model for lipid membrane

Coarse-grained model for lipid membrane



- Harmonic springs and bending angle potentials

$$\frac{\mathcal{H}_b}{k_B T} = \sum_i \frac{k_s}{2} (\mathbf{r}_{i+1} - \mathbf{r}_i)^2 + \sum_i k_b (1 - \cos \theta_i)$$

- “density-functional theory”-based description for \mathcal{H}_{nb}

$$\frac{\mathcal{H}_{nb}}{k_B T} = \int \frac{d\mathbf{r}}{R_{eo}^3} \hat{\rho}_\alpha(\mathbf{r}) \left[\frac{v_{\alpha\beta}}{2} \bar{\rho}_{2\beta}(\mathbf{r}) + \frac{\omega_{\alpha\beta\gamma}}{3} \bar{\rho}_{3\beta}(\mathbf{r}) \bar{\rho}_{3\gamma}(\mathbf{r}) \right]$$

$$\hat{\rho}_\alpha(\mathbf{r}) = \frac{R_{eo}^3}{N} \sum_i \delta(\mathbf{r}_i - \mathbf{r}) \delta_{\alpha t(i)} \quad v_{\alpha\beta}, \omega_{\alpha\beta\gamma} : \text{thermodynamic properties}$$

$$\bar{\rho}_{m\alpha}(\mathbf{r}) = \frac{R_{eo}^3}{N} \sum_i w_m(|\mathbf{r}_i - \mathbf{r}|) \delta_{\alpha t(i)} \quad w_m : \text{local structures (packing)}$$

- Capabilities of the model

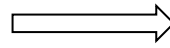
- ✧ Self-assembly of lipids
- ✧ Mechanical properties
- ✧ Main phase transitions

- Particle-based simulation Implementation: MD, MC, DPD, ...

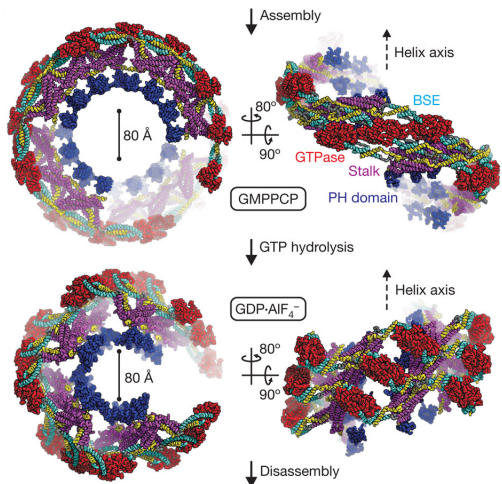
Two coarse-grained descriptions for dynamin

□ A very crude model for dynamin constriction

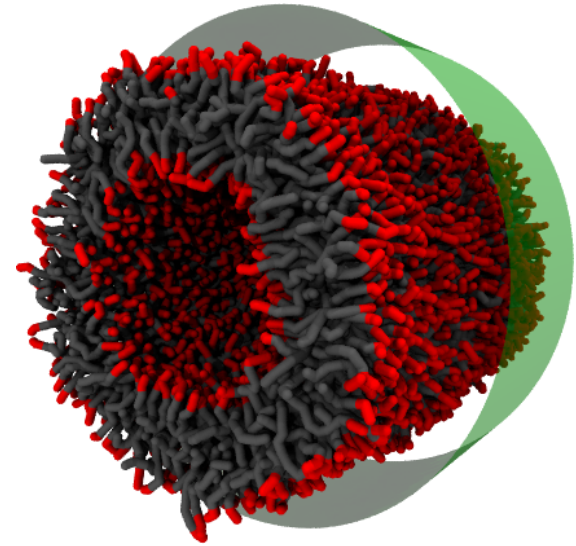
constriction by *dynamin coating*



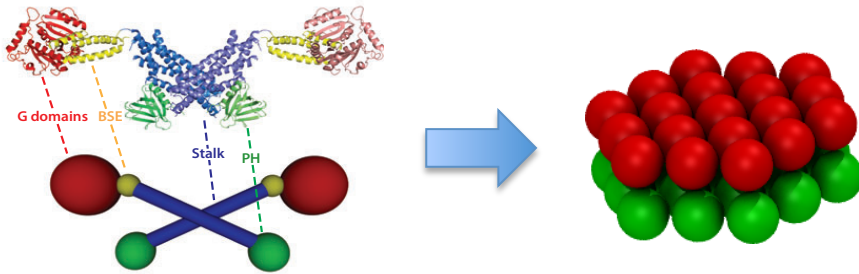
constriction by *a cuff potential*



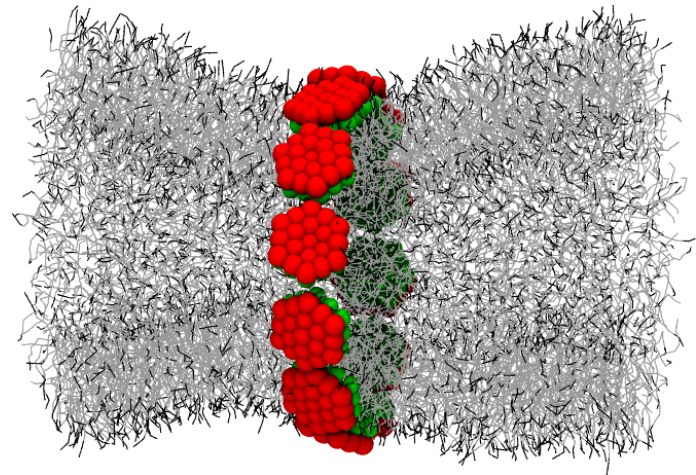
Ford et al., *Nature*, **477**, 561 (2011)



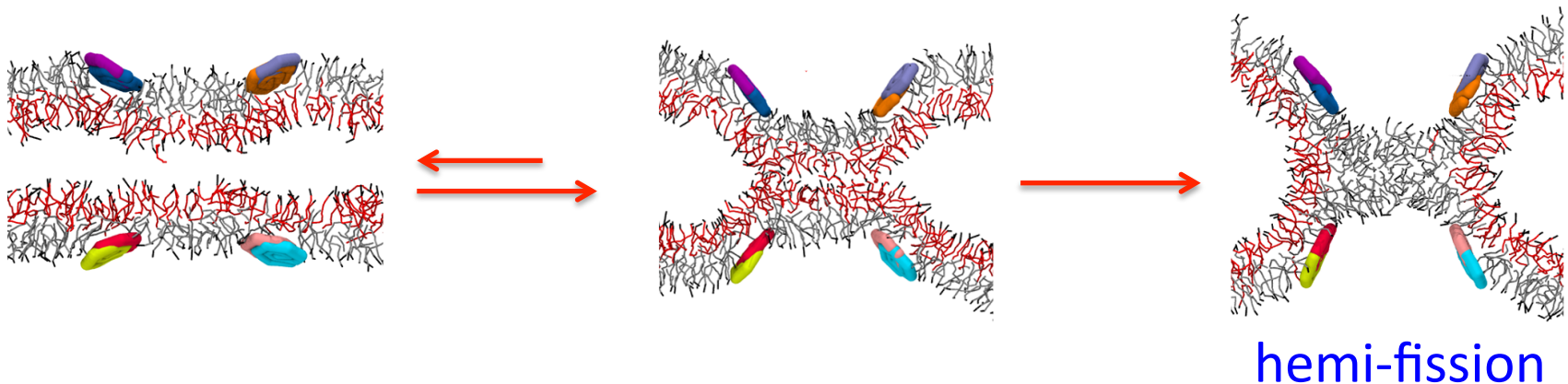
□ A more realistic model of dynamin (*constriction + PH domain insertion*)



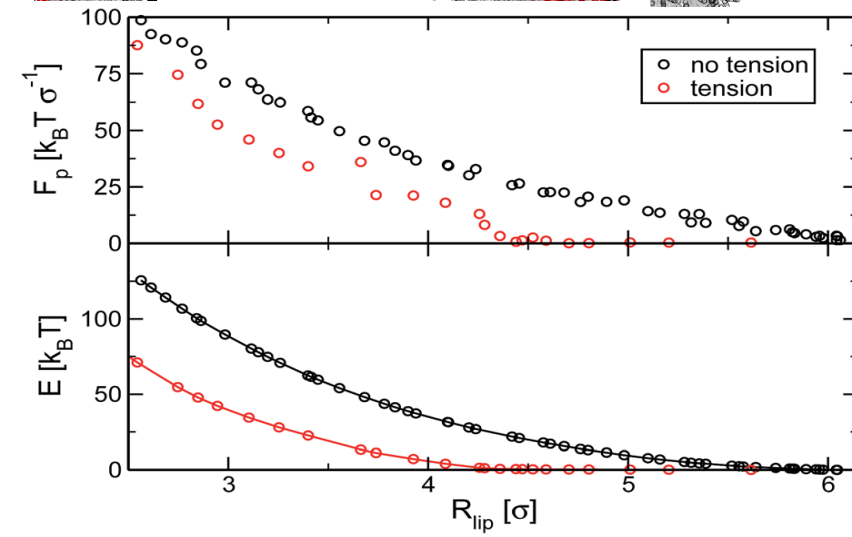
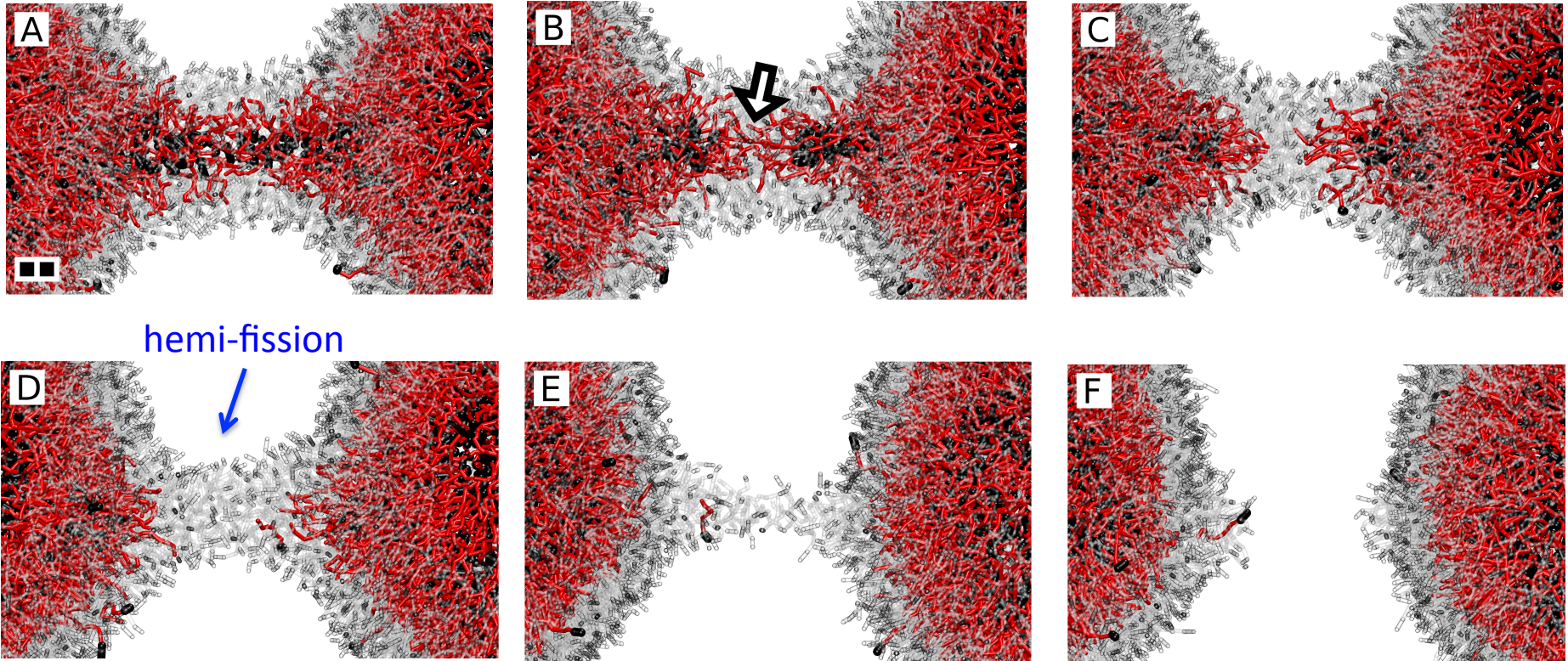
Fuhrmans and Müller, *Soft Matter* **11**, 1464 (2015)



Part I: Role of dynamin's constriction



Dynamin's constriction as a cuff potential



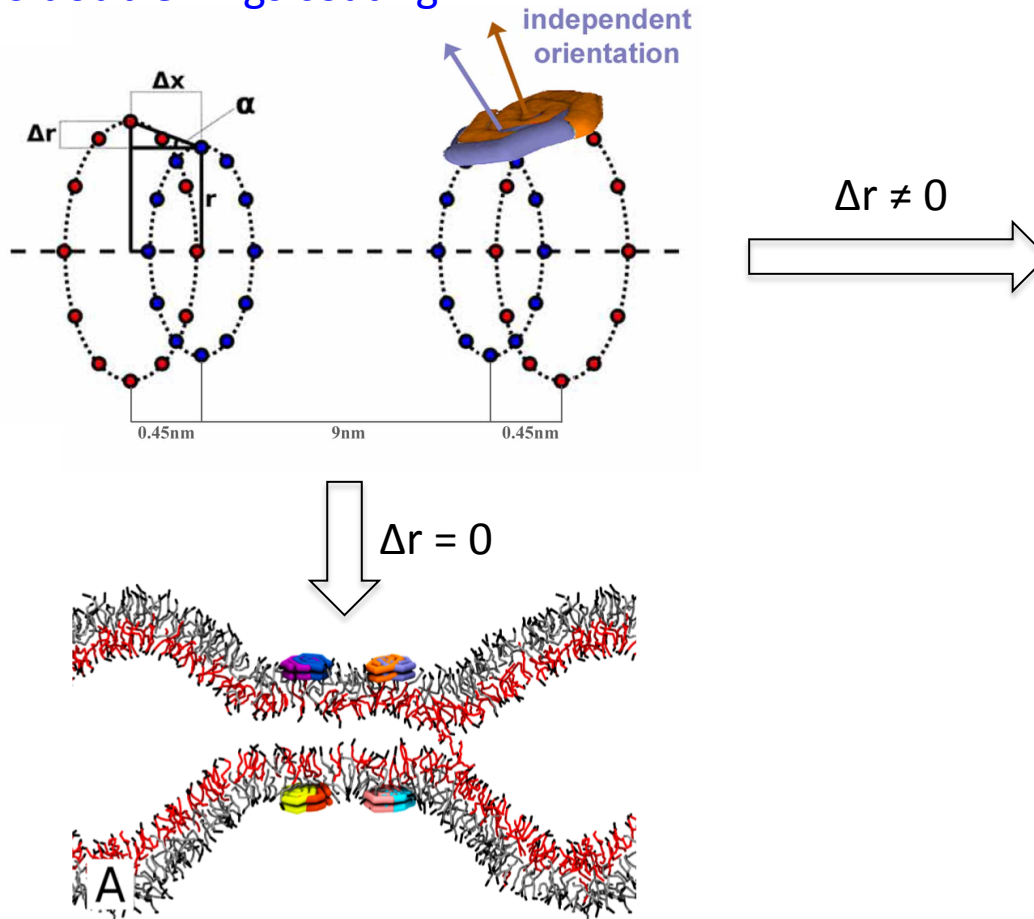
- ✧ non-leaky fission through formation of a worm-like micelle (*hemi-fission*) intermediate
- ✧ very *high* free-energy cost (about $150 k_B T$) for the constriction.
- ✧ *two* free-energy barriers involved: (i) formation of the worm-like micelle; (ii) rupture of the micelle.

Hemi-fission from a model with two double-dynamamin rings

Model of dynamin

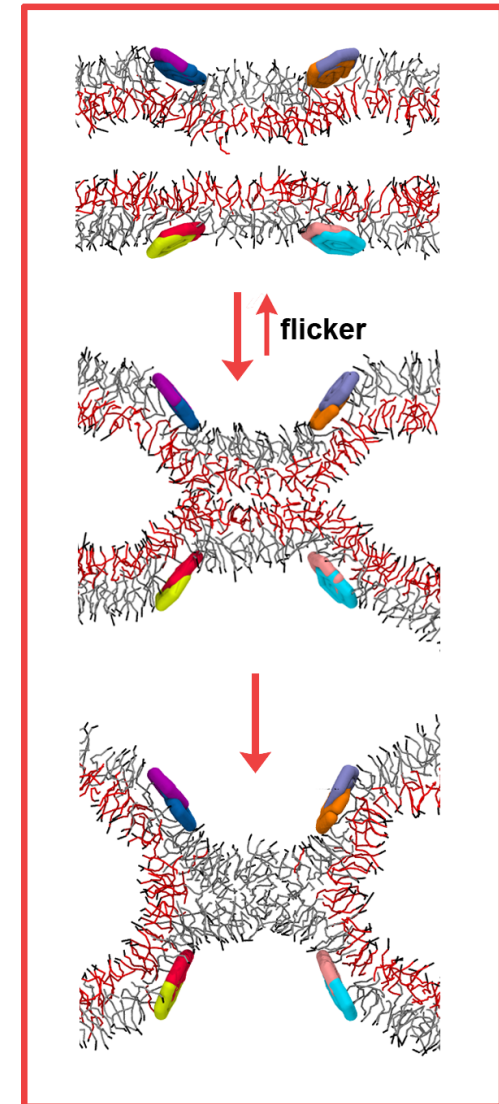


Two double-rings coating

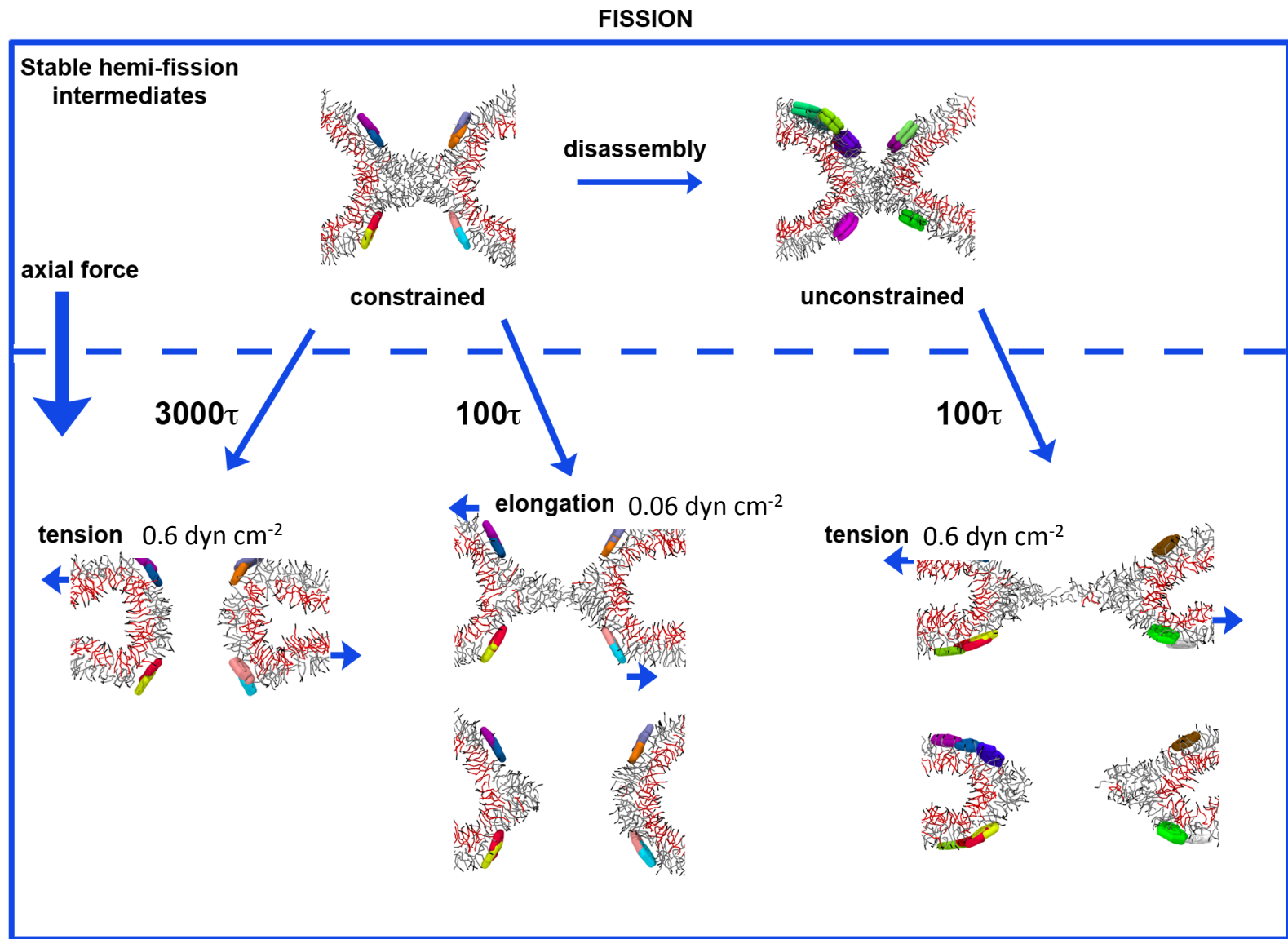


no hemi-fission intermediate formation
only thinning of membrane tube

HEMI-FISSION

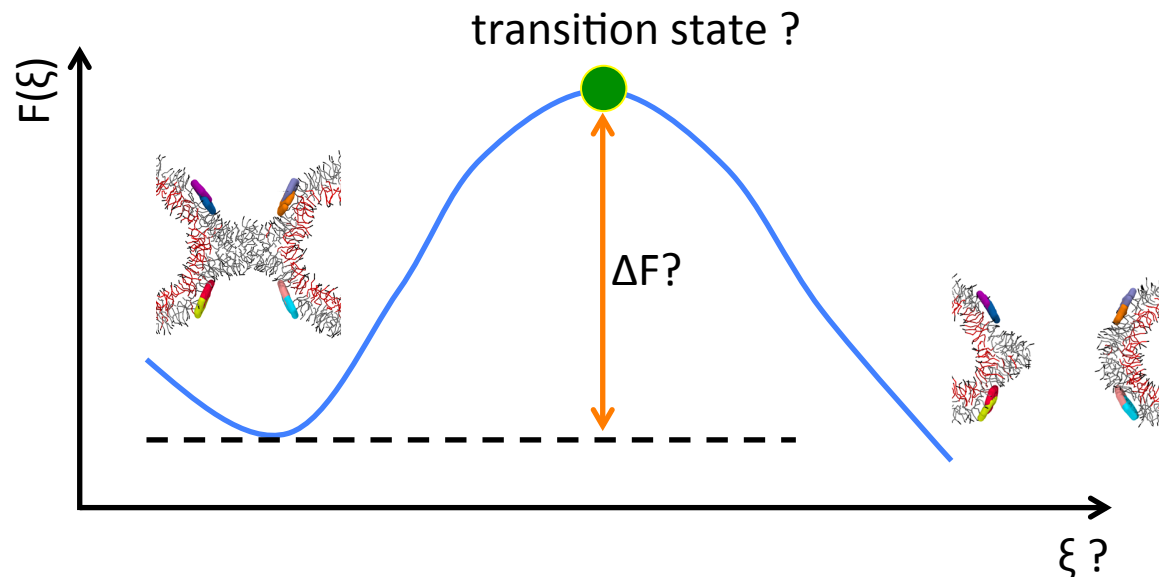


Completion of fission via rupture of the hemi-fission intermediate



The hemi-fission state is *(meta)stable*, i.e., rupture of hemi-fission state requires external effects.

Part II: Pathway and free-energy landscape of rupturing the hemi-fission state (Stability of the hemi-fission state)



Free energy calculations: with explicit reaction coordinates

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$$F(\xi) = -k_{\text{B}}T \ln \int d\mathbf{r}^n e^{-H(\{\mathbf{r}^n\})/k_{\text{B}}T} \delta[\hat{\xi} - \xi] = -k_{\text{B}}T \ln P(\xi) - k_{\text{B}}T \ln Z(n, V, T)$$

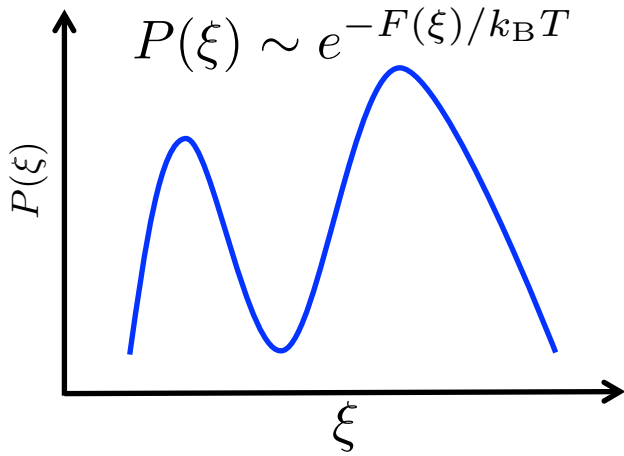
ξ : reaction coordinate

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□ Direct sampling in a *conventional* ensemble

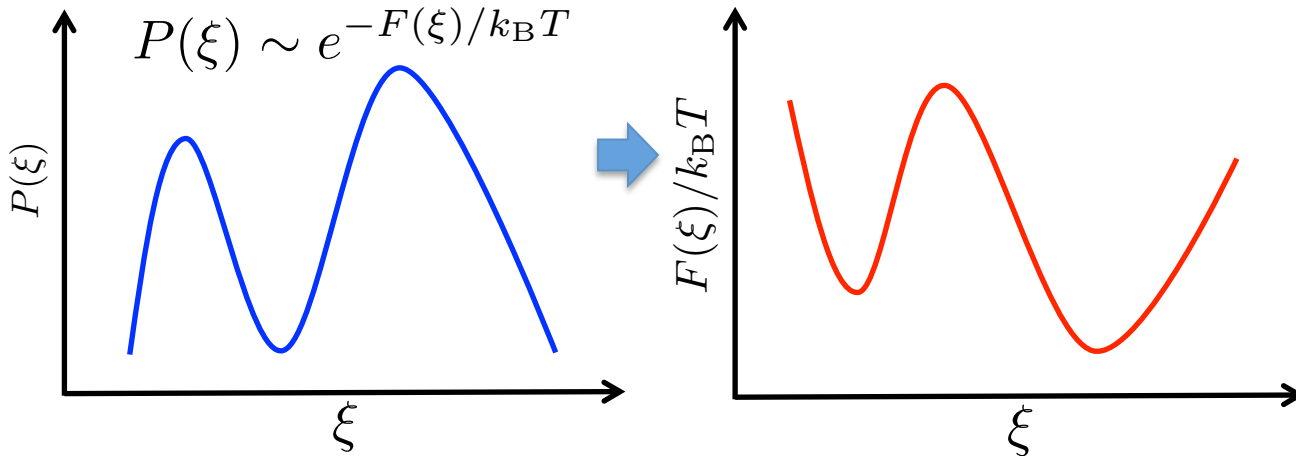


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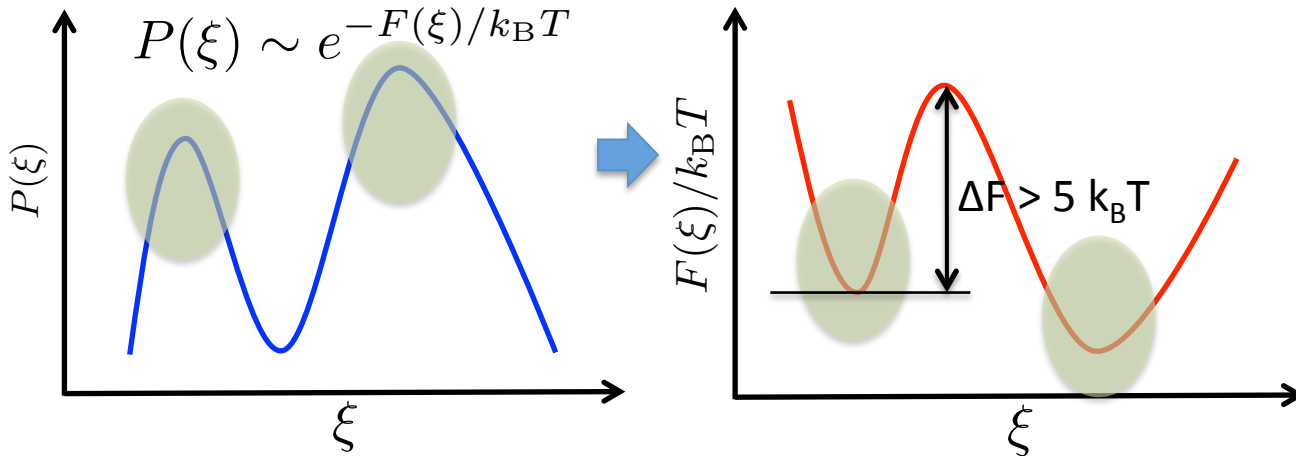


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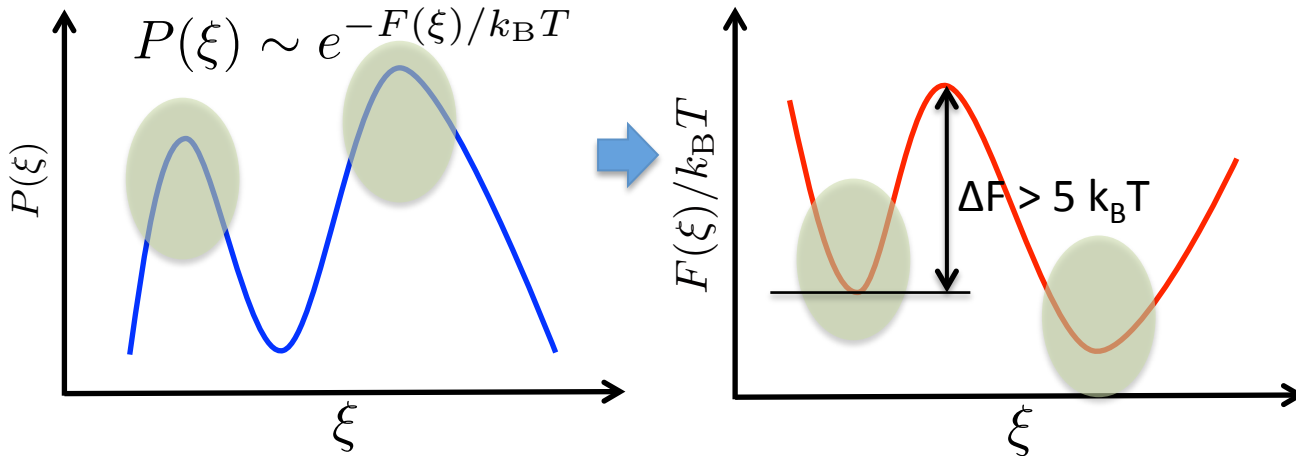


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Sampling problem in rare event

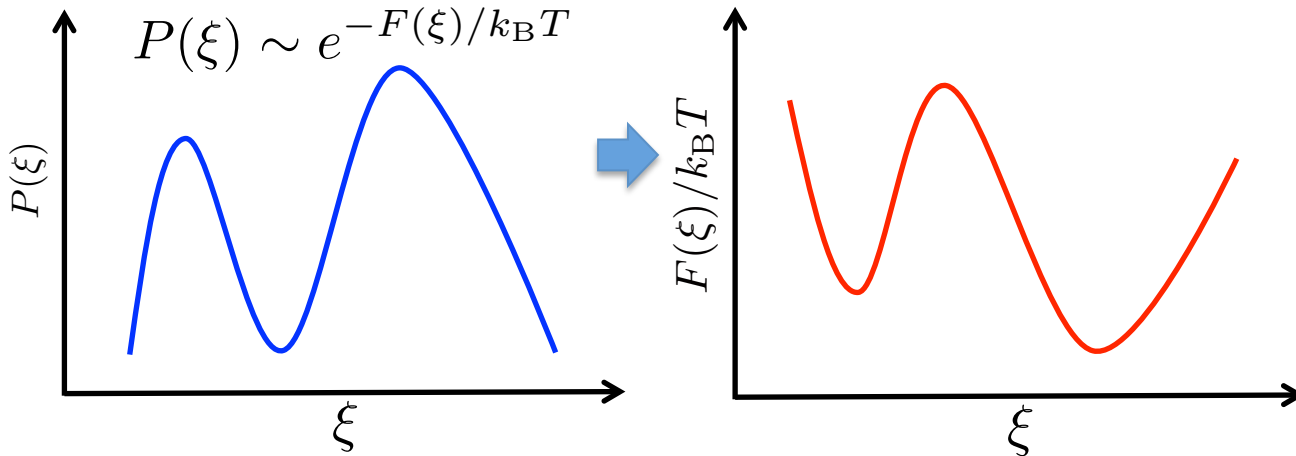
non-ergodic in a case with *large* free energy barrier ($> 5 k_B T$).

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❑ Biased sampling in an *extended* ensemble

➤ *Umbrella sampling*: enhancing sampling over the reaction coordinate by *biasing* the Hamiltonian

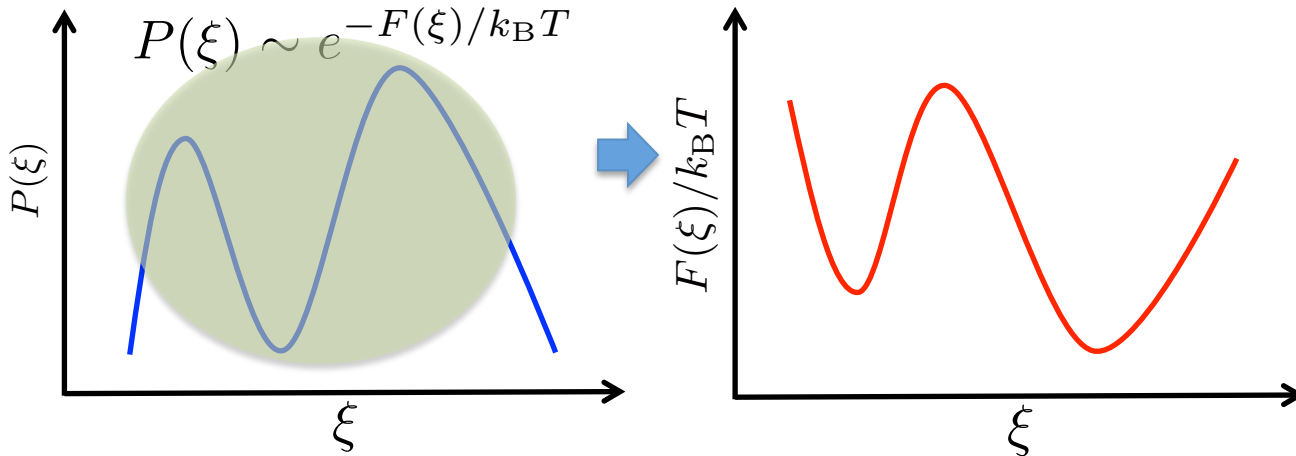
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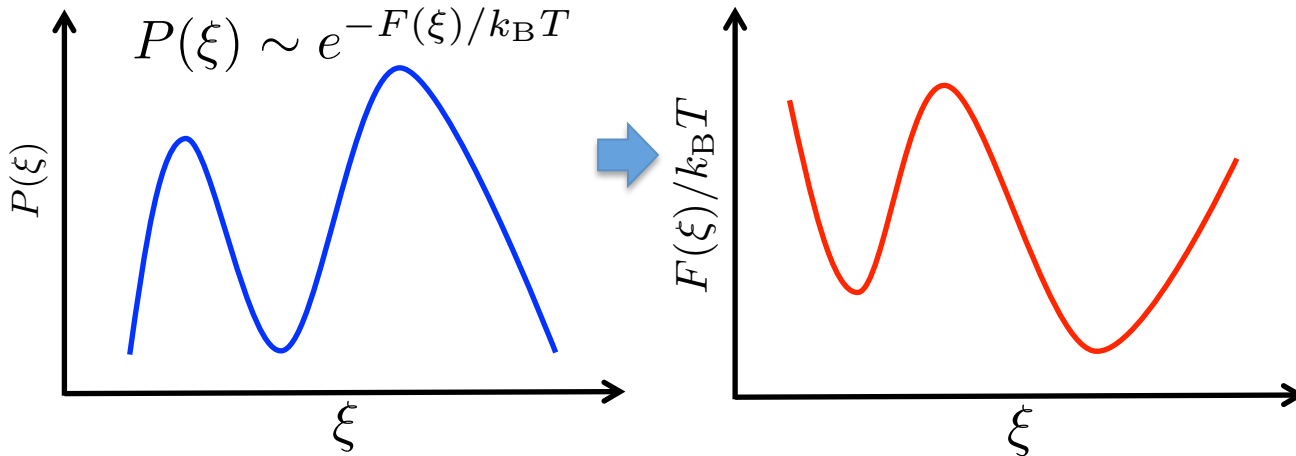
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➤ *Flat-histogram method*: when the biased potential is the *free energy*

$$P_{\text{fh}}(\xi) \sim \int d\mathbf{r}^n e^{[-H(\mathbf{r}^n) + F(\hat{\xi})]/k_B T} \delta[\hat{\xi} - \xi]$$

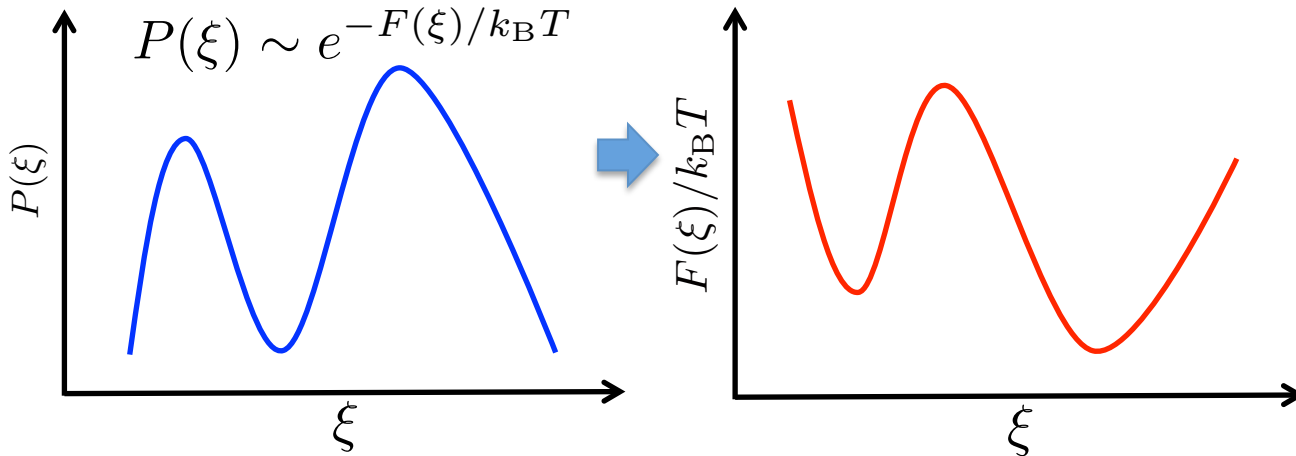
$$\sim e^{-F(\xi)/k_B T} e^{F(\xi)/k_B T} = O(\xi^0)$$

Free energy calculations: with explicit reaction coordinates

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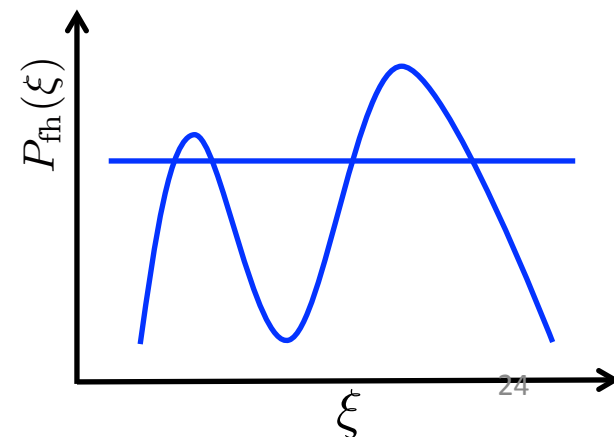
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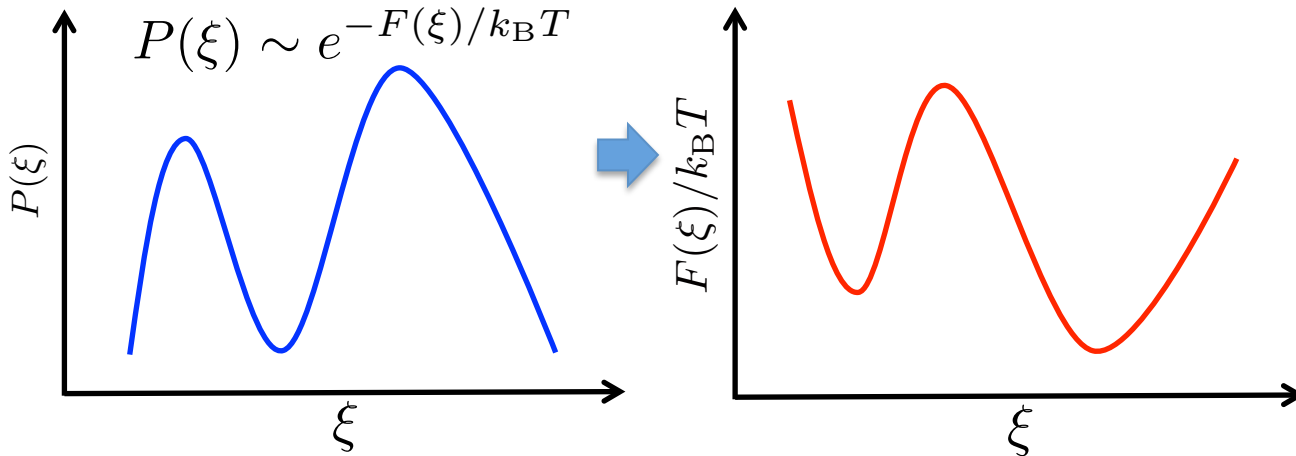


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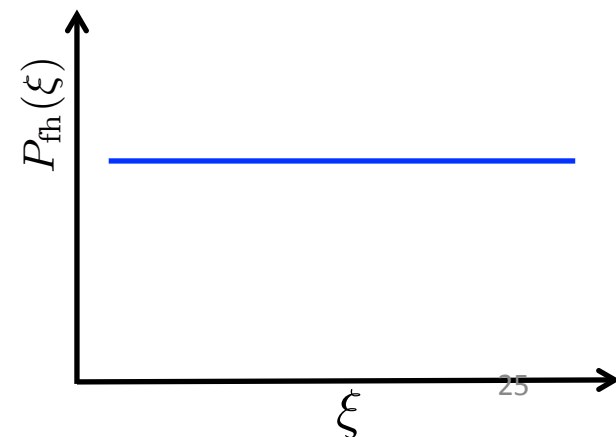
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✧ Goal of the method is to obtain $F(\xi)$ by realizing a flat-histogram of $P(\xi)$.

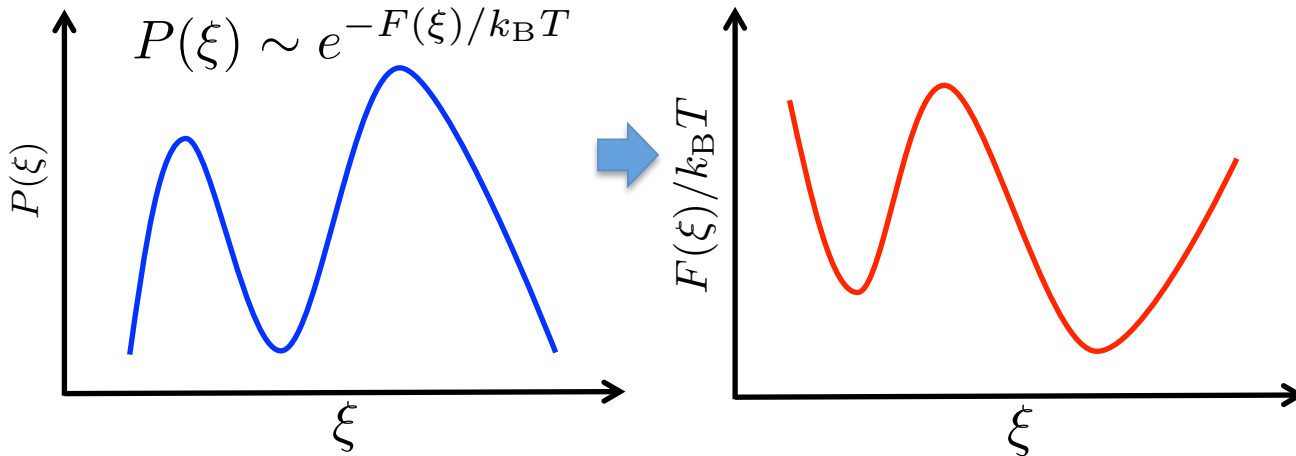


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❑ Direct sampling in a **conventional** ensemble



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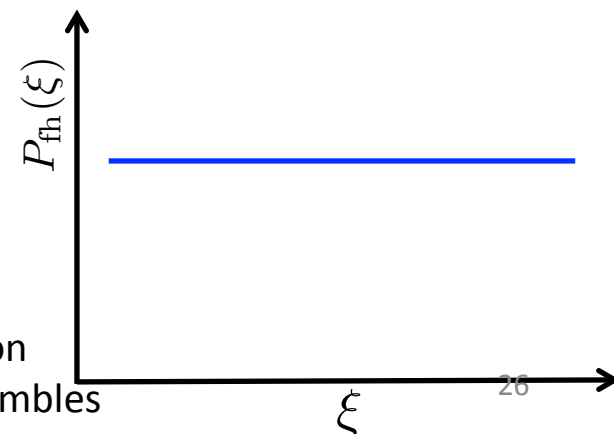
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Wang-Landau (WL) sampling

Monte Carlo simulation
convenient for any ensembles

Metadynamics

molecular dynamics simulation
not easy for grand-canonical ensembles

A practical recipe for Wang-Landau sampling

Goal: to derive a Metropolis acceptance criterion in Wang-Landau sampling in a general sense.

Wang and Landau, Phys. Rev. Lett. 86, 2050 (2001)

Ganzenmueller and Camp, J. Chem. Phys. 127, 154594 (2007)

Methodology:

(1) The microscopic state distribution in Wang-Landau sampling

$$1 = p_{WL}(\xi) = \int \underline{p_{WL}(\Gamma, \xi)} d\Gamma \quad \Gamma: \text{microstate} \quad \xi: \text{reaction coordinate}$$

$$\boxed{p_{WL}(\Gamma, \xi) = \frac{p_0(\Gamma, \xi)}{p(\xi)}} \quad \longrightarrow \quad p_{WL}(\xi) = \int p_{WL}(\Gamma, \xi) d\Gamma = \int \frac{p_0(\Gamma, \xi)}{p(\xi)} d\Gamma = 1$$

(2) The acceptance criterion from detailed balance

$$p_{WL}(\Gamma_o, \xi_o) \text{acc}(o \rightarrow n) = p_{WL}(\Gamma_n, \xi_n) \text{acc}(n \rightarrow o)$$

$$\text{acc}(o \rightarrow n) = \min \left[1, \frac{p_{WL}(\Gamma_n, \xi_n)}{p_{WL}(\Gamma_o, \xi_o)} \right]$$

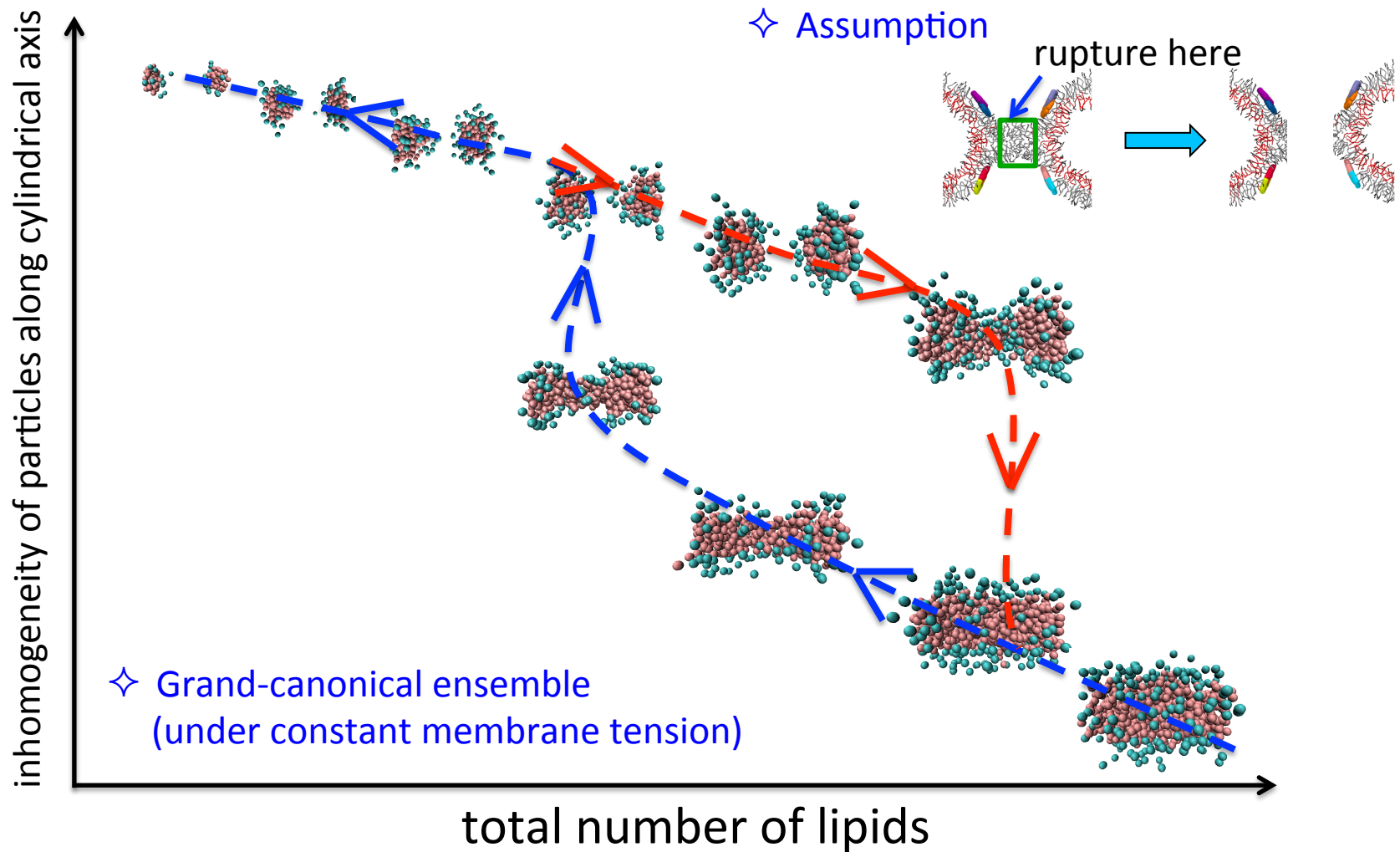
Example: $\xi = E$ in canonical ensemble

$$\left. \begin{aligned} p_0(\Gamma_o, E_o) &= \frac{V^N e^{-\beta E_o}}{N! \Lambda^{3N} Q_{NVT}} \\ p(E_o) &= \frac{V^N e^{-\beta E_o} \Omega(N, V, E_o)}{N! \Lambda^{3N} Q_{VNT}} \end{aligned} \right\} \longrightarrow p_{WL}(\Gamma_o, E_o) = \frac{p_0(\Gamma_o, E_o)}{p(E_o)} = \frac{1}{\Omega(N, V, E_o)}$$

$$\text{acc}(o \rightarrow n) = \min \left[1, \frac{\Omega(N, V, E_o)}{\Omega(N, V, E_n)} \right]$$

Microcanonical ensemble partition function is obtained from WL sampling in *canonical* ensemble.

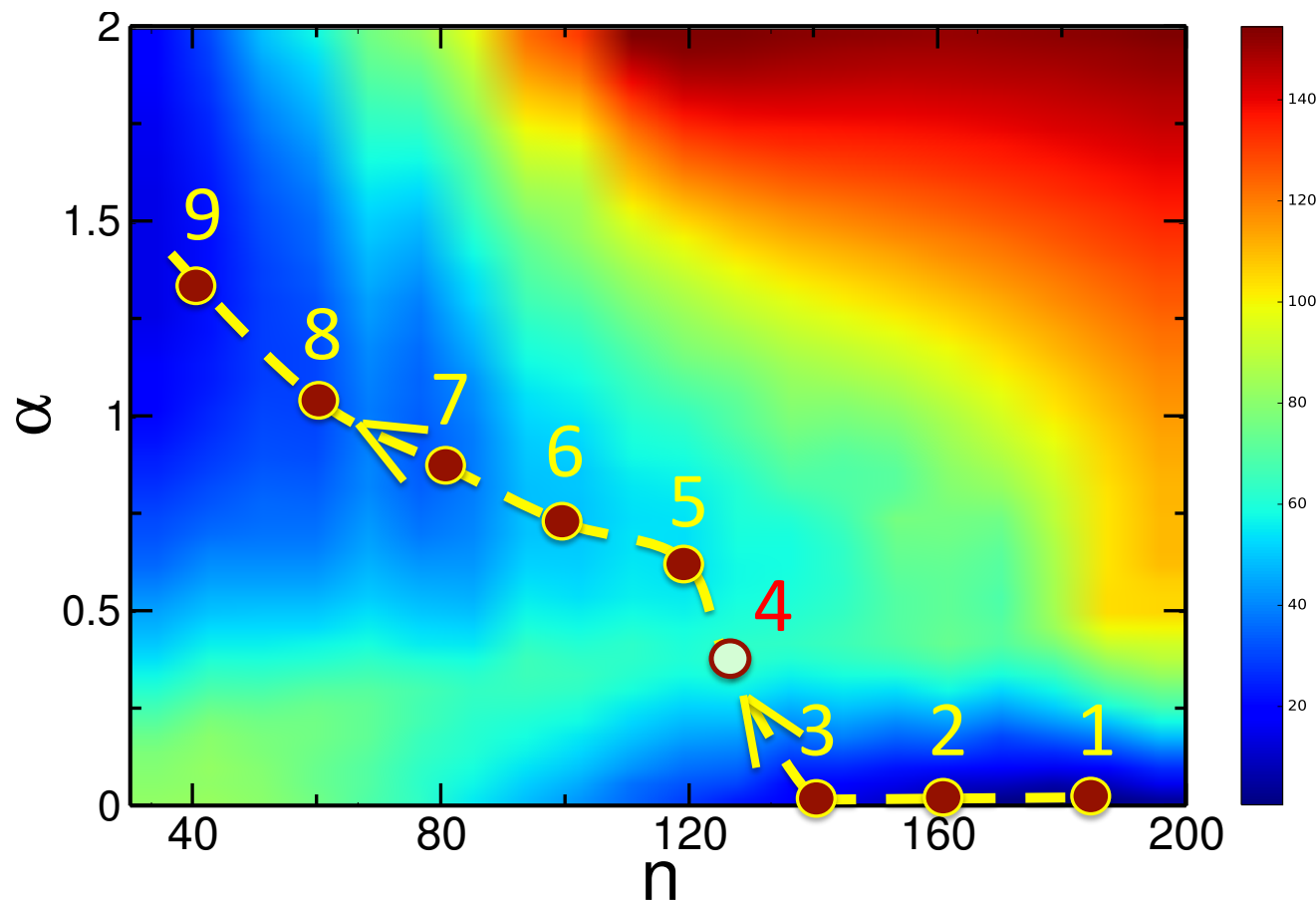
Identification of suitable reaction coordinates



- *Hysteresis* occurs if only one order parameter is used.
- One order parameter *isn't* sufficient.
- A *second* order parameter is identified as the inhomogeneity of particles along the cylindrical axis.

Free-energy landscape and transition path

□ Projection of an intrinsic *high-dimensional* free-energy landscape to a *two-dimensional* one.

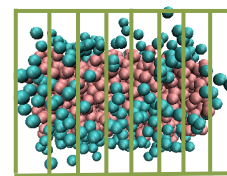


Reaction coordinates

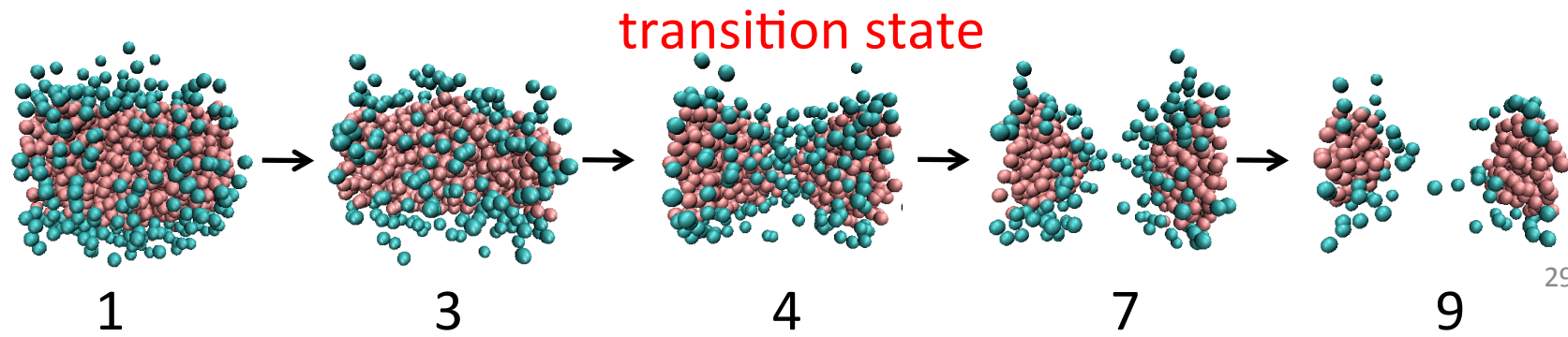
n : total number of lipids

α : inhomogeneity of particles along cylinder axis, which is defined as:

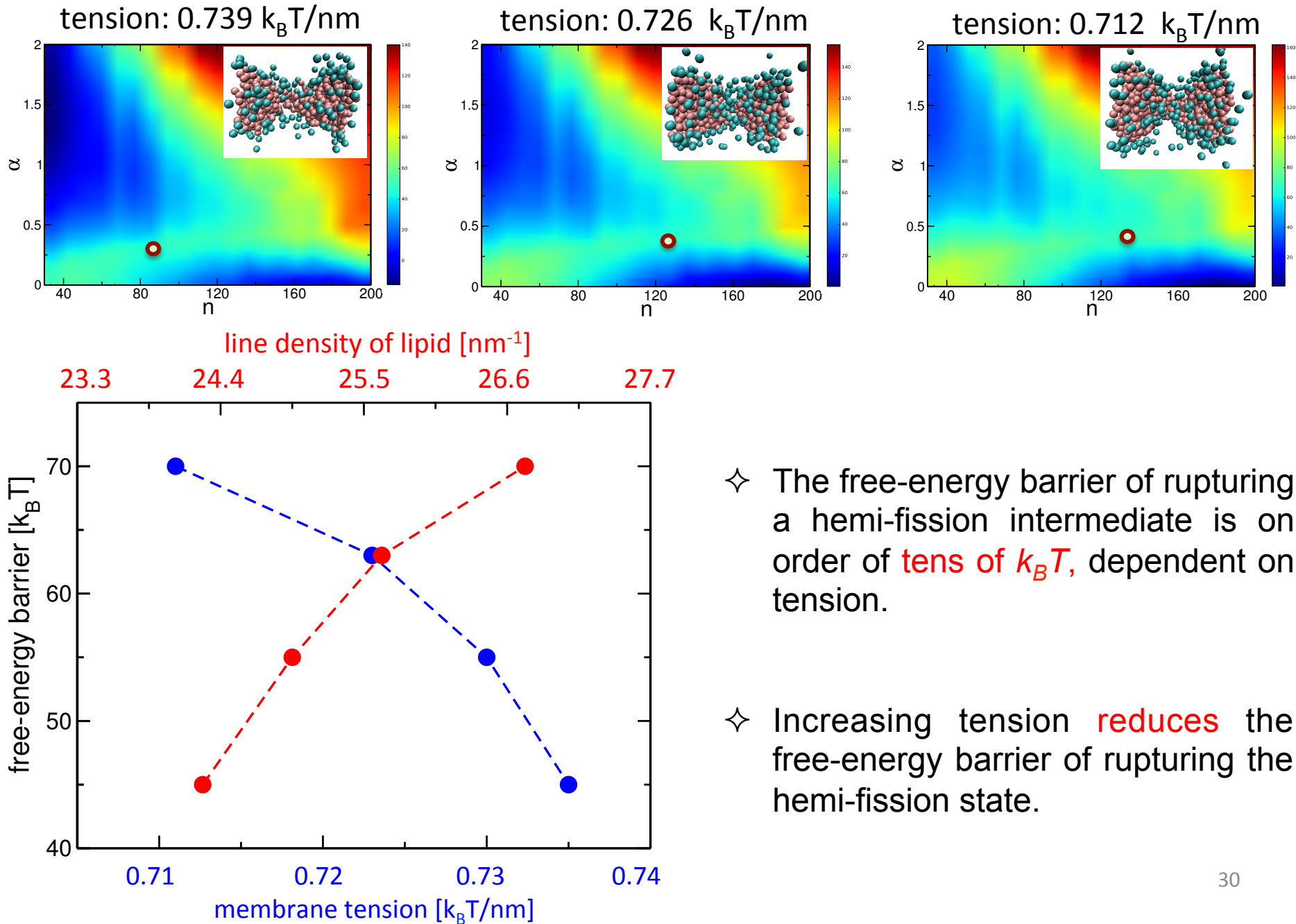
$$\alpha = \frac{\langle N_i^2 \rangle - \langle N_i \rangle^2}{\langle N_i \rangle^2}$$



N_i particles in the i -th slab



Effect of tension on free-energy barriers



Free energy calculations: without explicit reaction coordinates

*In general, identifying suitable reaction coordinates is very **challenging**.*

□ String method for minimum free-energy path (MFEP)

(1) Theory

- ✧ The MFEP is defined by condition that the derivative **perpendicular** to the path **vanishes**.

$$\nabla_{\perp} \mathcal{F}[m] = \frac{\delta \mathcal{F}}{\delta m_s(\mathbf{r})} - \underbrace{\frac{dm_s(\mathbf{r})}{ds} \frac{\int d\mathbf{r} \frac{\delta \mathcal{F}}{\delta m_s(\mathbf{r})} \frac{dm_s(\mathbf{r})}{ds}}{\int d\mathbf{r} \left(\frac{dm_s(\mathbf{r})}{ds} \right)^2}}_{\text{tangent term along the path}} \stackrel{!}{=} 0$$

- ✧ A path: $m_s(\mathbf{r}), 0 < s < 1$.

- ✧ Free energy: $\mathcal{F}[m_s] = \int_0^s ds \int d\mathbf{r} \frac{dm_s(\mathbf{r})}{ds} \frac{\partial \mathcal{F}[m_s]}{\partial m_s(\mathbf{r})}$

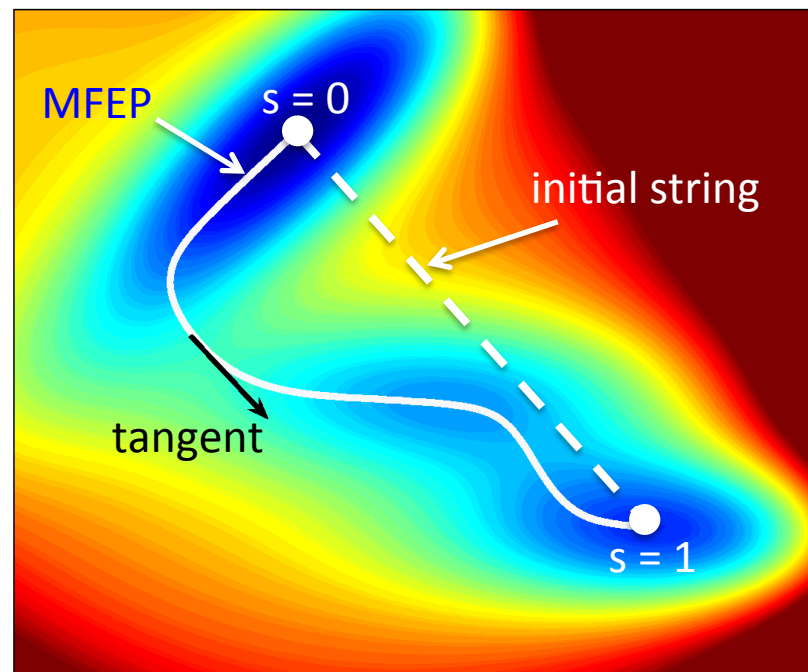
- ✧ Transition state: $d\mathcal{F}[m_s]/ds = 0$

(2) Numerical implementation

- ✧ **Step 1:** $\Delta m_s(\mathbf{r}) = -\mu(\mathbf{r}|m_s)\Delta\varepsilon$

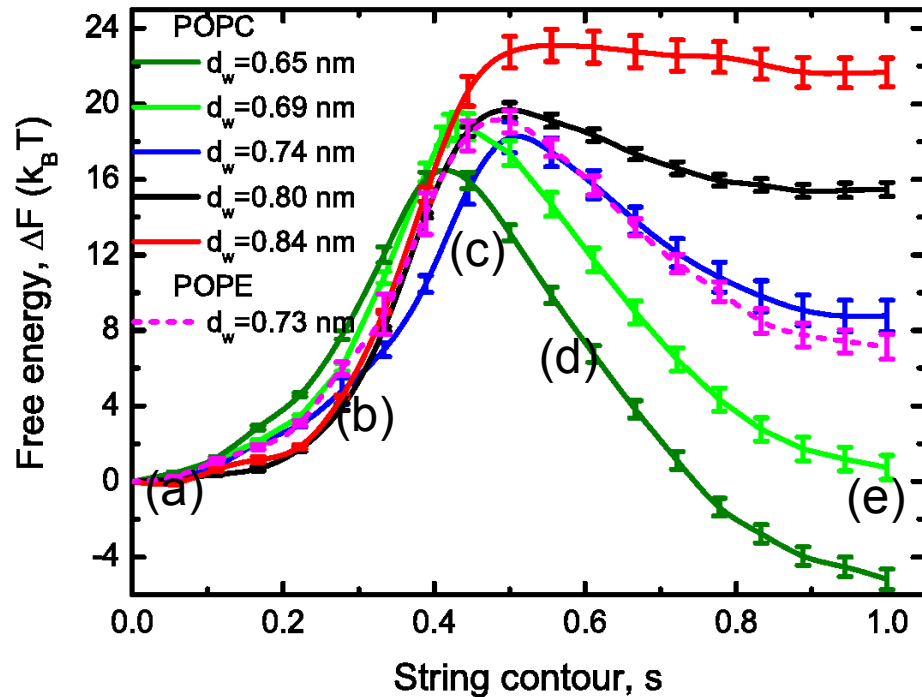
$$\mu(\mathbf{r}) = \frac{\delta \mathcal{F}}{\delta m_s(\mathbf{r})} \approx k_B T \lambda \left\langle m(\mathbf{r}) - \hat{m}(\mathbf{r}|\{\mathbf{r}\}) \right\rangle_{\text{umbrella}}$$

- ✧ **Step 2:** re-parameterization of morphologies such that they are distributed **evenly** along the string.



Application of string method to membrane fusion

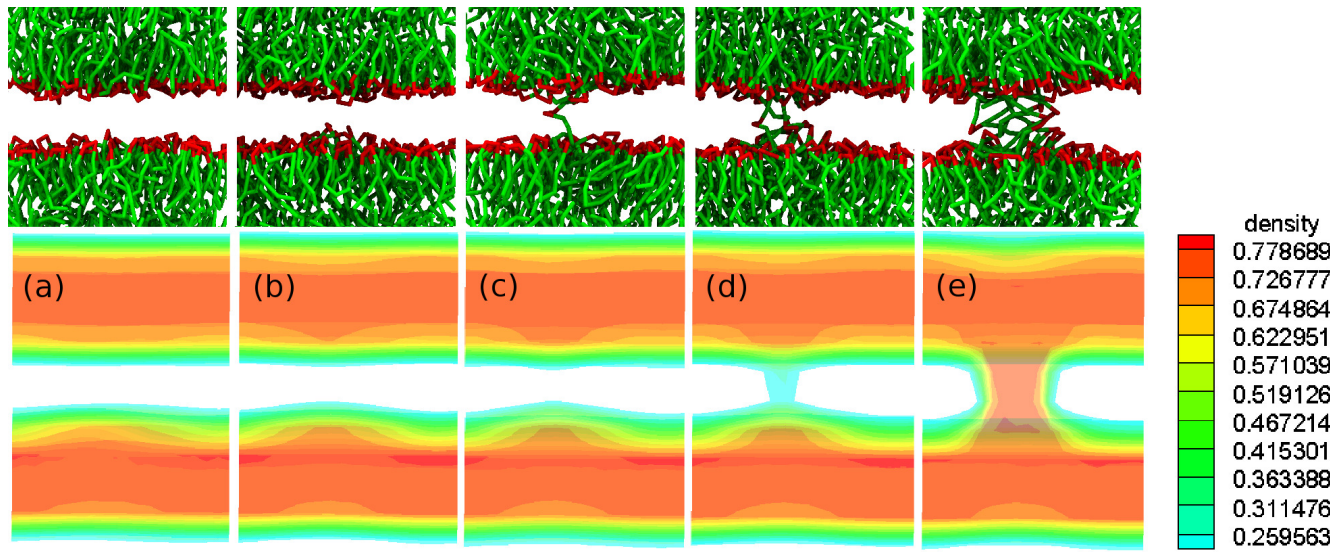
□ Model for lipid membrane: MARTINI model



✧ Degree of dehydration (thickness of bilayers) slightly modifies the free-energy barriers.

✧ It significantly affects the stability of the stalk morphology, higher dehydration favors formation of the stalk state.

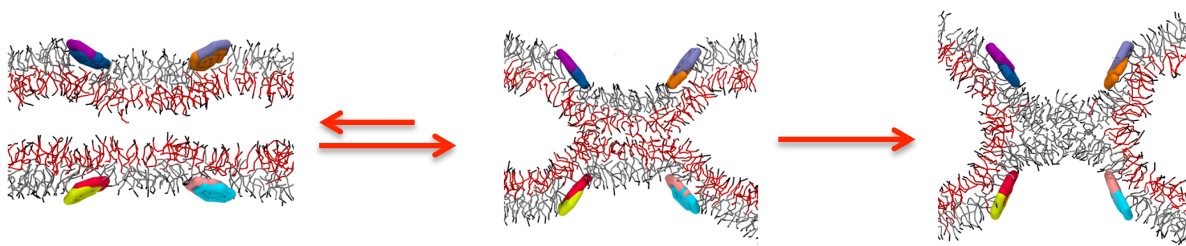
Smirnova, Müller, manuscript in preparation.



Summary

□ Role of dynamin's constriction in fission

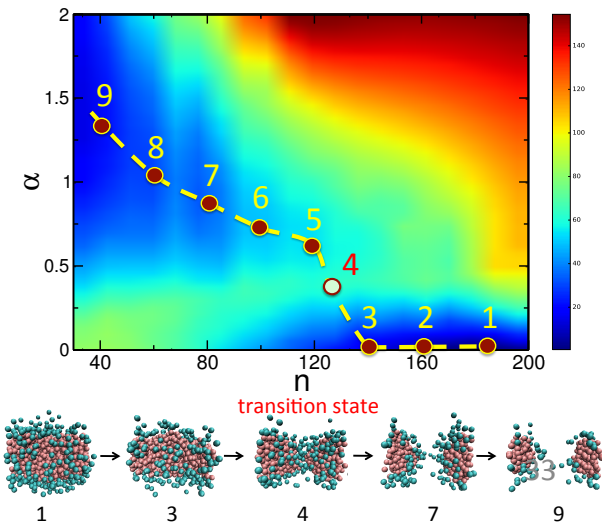
- ✓ The constriction model using a *cuff potential* gives rise to quite *large free energy cost* for constriction even before a topological change occurs.
- ✓ The more “realistic” model using *double-dynamin rings* leads to a formation of the (meta)stable *hemi-fission intermediate*.



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□ Free energy landscape of rupturing a hemi-fission state

- ✓ Two suitable *order parameters* has been identified.
- ✓ Two-dimensional free-energy landscape is calculated.
- ✓ Increasing membrane tension *lowers* the free-energy barrier.



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Thanks for your attention!