



PARTICLES BREAKING OUT OF DEBRIS FLOWS



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Debris flows

- Fast moving, subaerial gravitational flows of water, sediments and coarse material (rocks, trees boulders)
- A general term encompassing lahars, landslides, jökulhlaups.

e.g. Vargas, Venezuela 1999





- With our debris flow experiments we want to
 - Understand the effect of various flow variables, e.g. surface roughness, particle size
 - Use large particle sizes to (try and) achieve Froude and particle Reynolds number similarity
 - Simultaneously measure velocity profiles, pore pressure and basal shear and normal stress

Design Criterion - Similarity

- Always difficult in particle laden flows!

Parameter	Name	Force Balance	Notts Chute	USGS Chute	1982 Oddstad
$N_{\text{Bag}} = \frac{\phi_s \rho_s d^2 \dot{\gamma}}{(1 - \phi_s) \mu}$	Bagnold number	Inertial grain stress to viscous shear stress	2	400	4
$N_{\text{Sav}} = \frac{\rho_s d^2 \dot{\gamma}^2}{(\rho_s - \rho_f) g h \tan \theta}$	Savage number	Inertial grain stress to friction	0.2	0.2	2×10^{-4}
$N_{\text{fric}} = \frac{N_{\text{Bag}}}{N_{\text{Sav}}}$		Friction to viscous shear stress	9	2×10^3	2×10^4
$N_{\text{mass}} = \frac{\phi_s}{(1 - \phi_s)} \frac{\rho_s}{\rho_f}$	Mass number c.f. Stokes number	Solid to fluid inertia	1	4	4
$N_{\text{Rey}} = \frac{N_{\text{Bag}}}{N_{\text{mass}}}$	c.f. Reynolds number	Fluid inertial stress to viscous shear stress	2.5	100	1
$\text{Fr} = \frac{u}{\sqrt{gh}}$	Froude number	Inertial to gravitational	0.6	10	3

Data from

Iverson Richard M., 1997, Physics of debris flows, *Rev. Geophys* 35, 3, 245-296

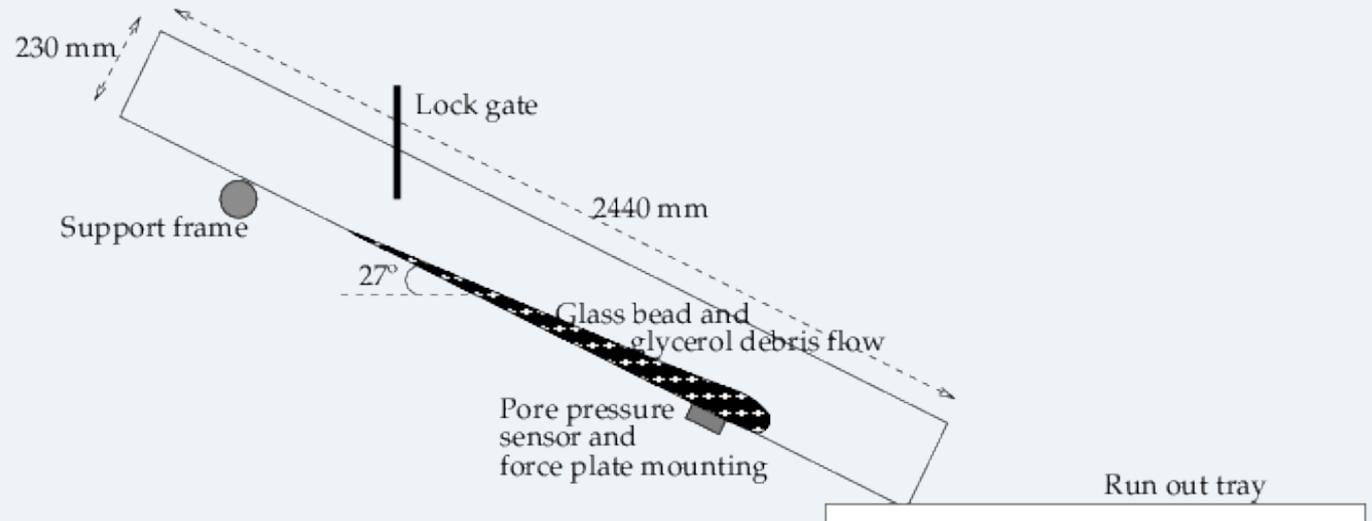
Experiment Design

2D CHUTE

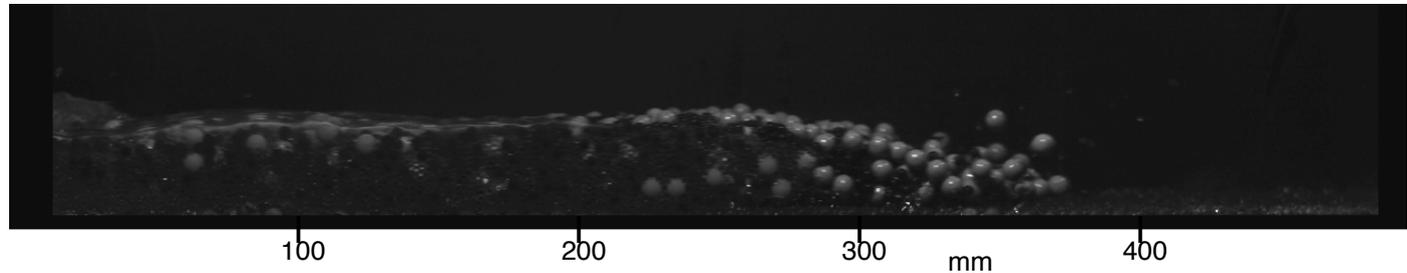
Lock release

Variable	Notation	Values	(units)
Solids volume fraction	ϕ_s	0.6	
Volume of solids	$\phi_s V$	1	litre
Roughness length	$[d_{r1}, d_{r2}, d_{r3}]$	[2, 4, 8]	$\times 10^{-3}$ m
Angle of inclination	θ	27°	
Solids: glass beads			
Density	ρ_s	2600	kg m^{-3}
Diameter	$[d_1, d_2, d_3]$	[2, 4, 8]	$\times 10^{-3}$ m
Fluids: water, glycerol			
Density	$[\rho_{f1}, \rho_{f2}]$	[1000, 1260]	kg m^{-3}
Viscosity	$[\mu_1, \mu_2]$	[1.41, 0.8]	Pas

Experiment



What happens?

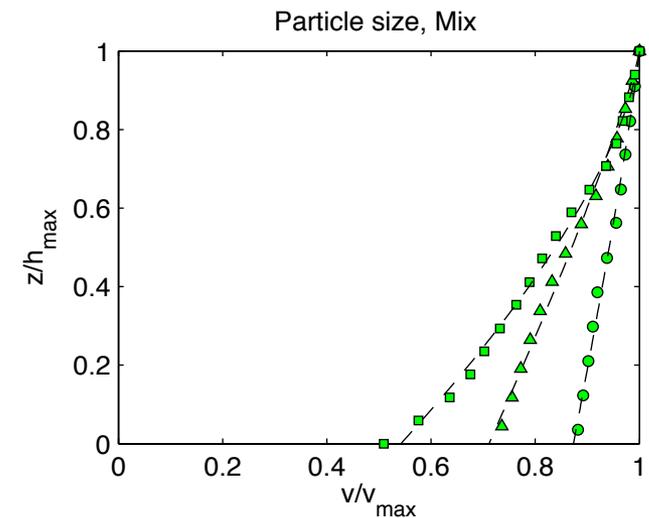
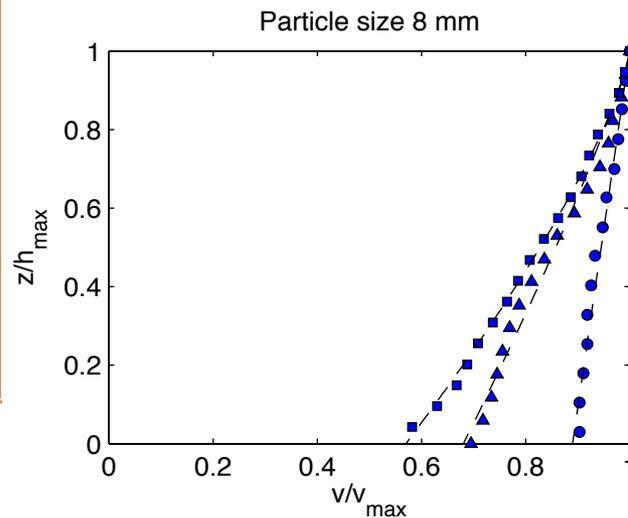
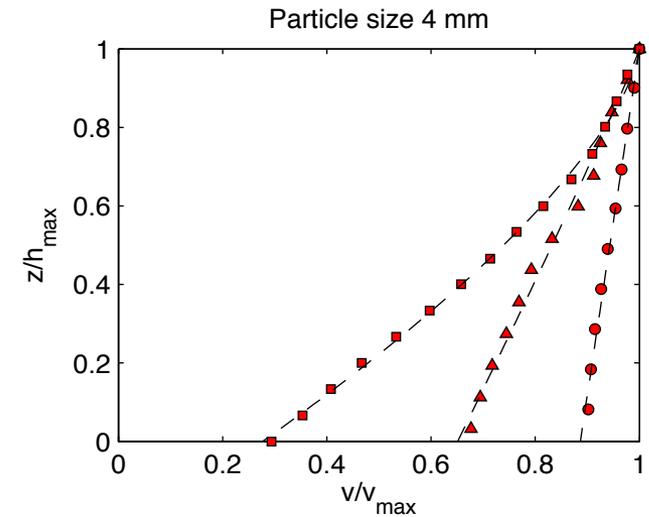
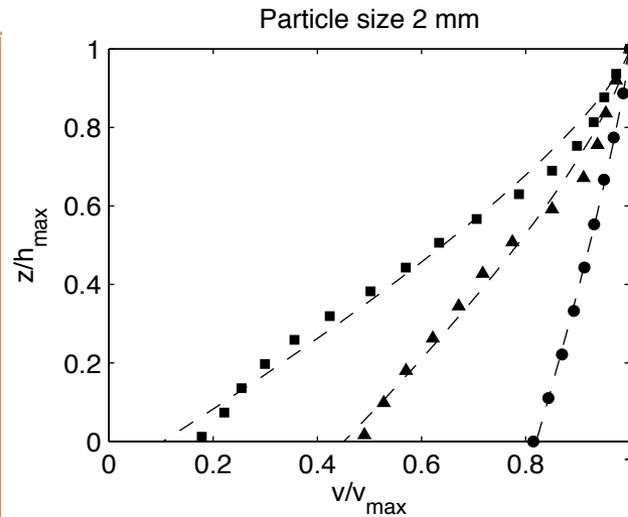


- Snout formation
- Longitudinal and vertical particle size and volume fraction variation
- Distinct collisional and continuum regions

Power law profiles

By particle size

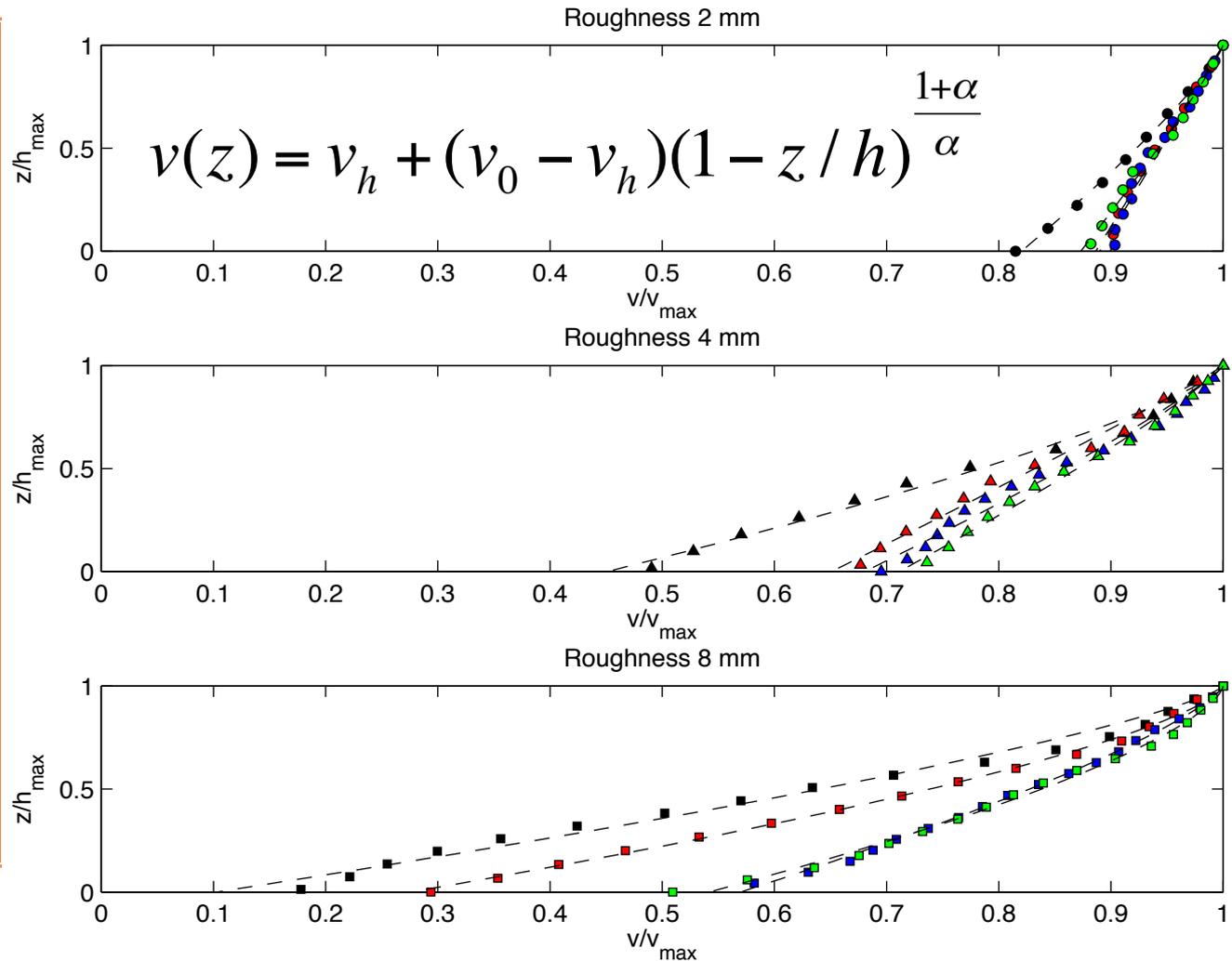
Roughness increases in the order circles, triangles, squares



Power law profiles

By roughness length

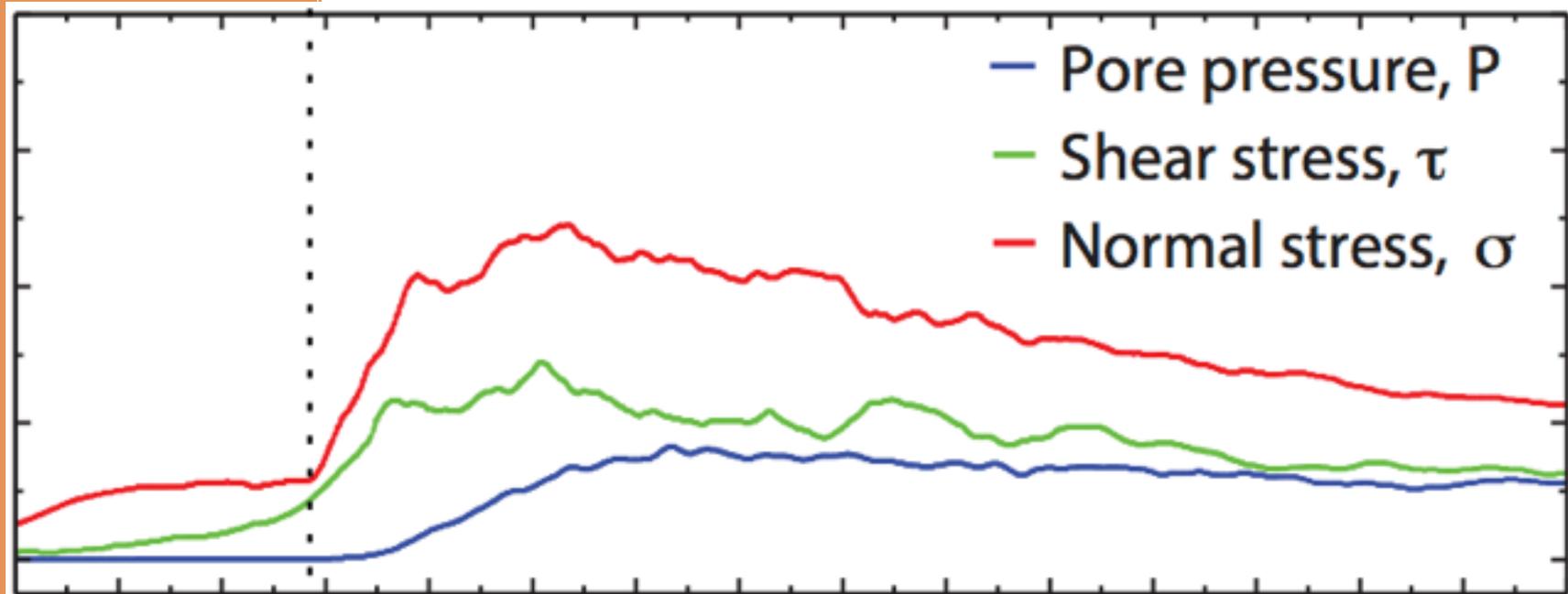
Particle size increases with black, red, blue. Green is a mixture

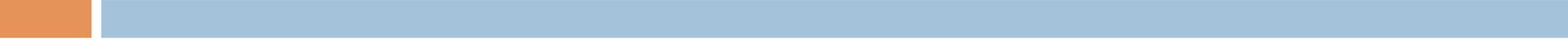
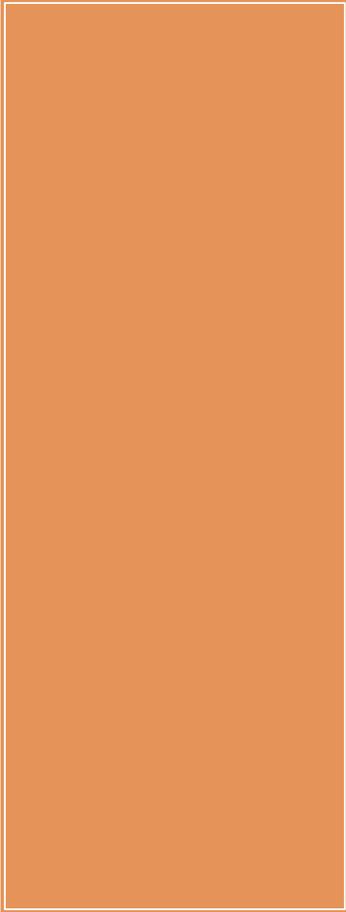


Pressures

4mm particles

8mm roughness

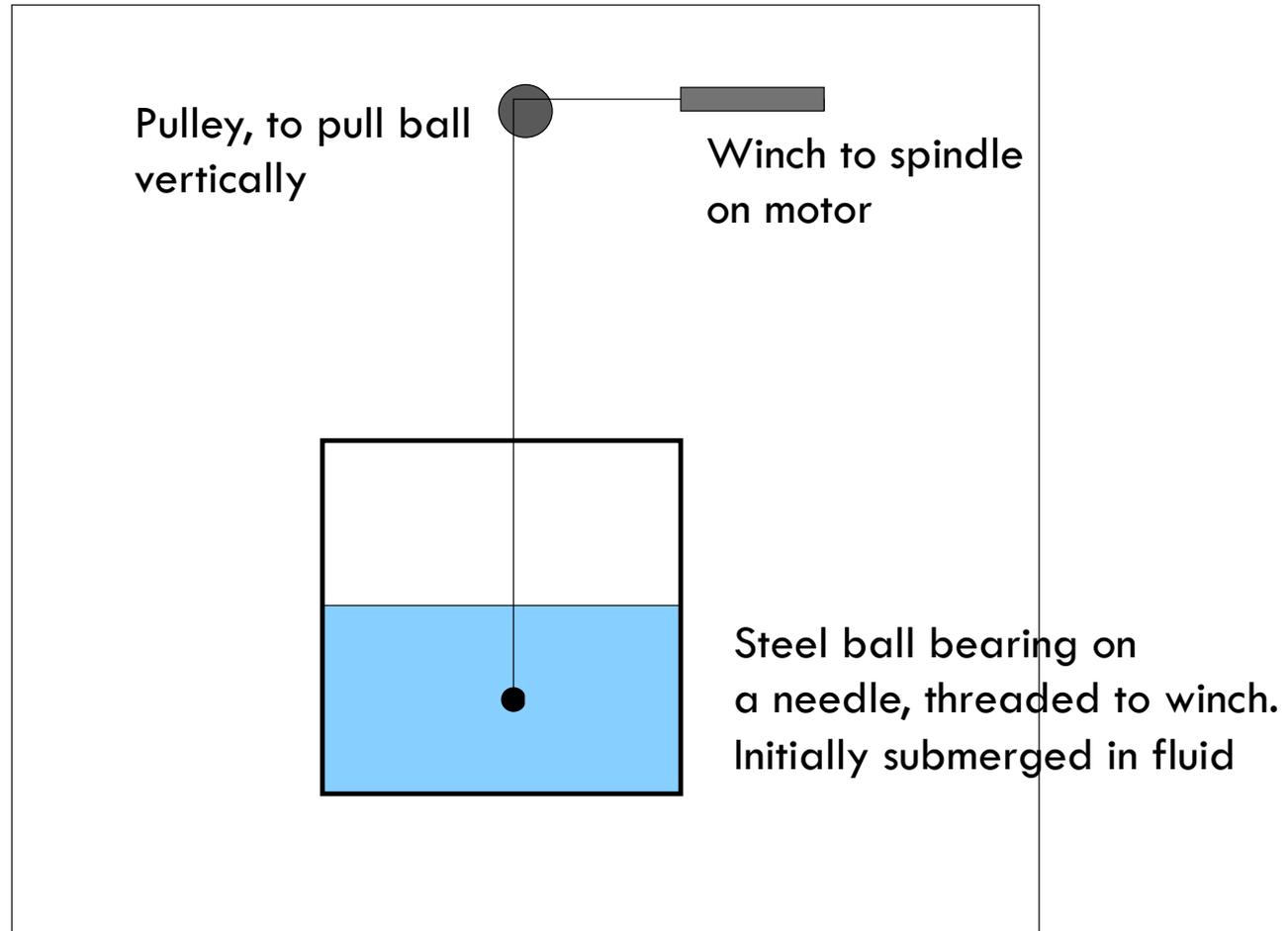
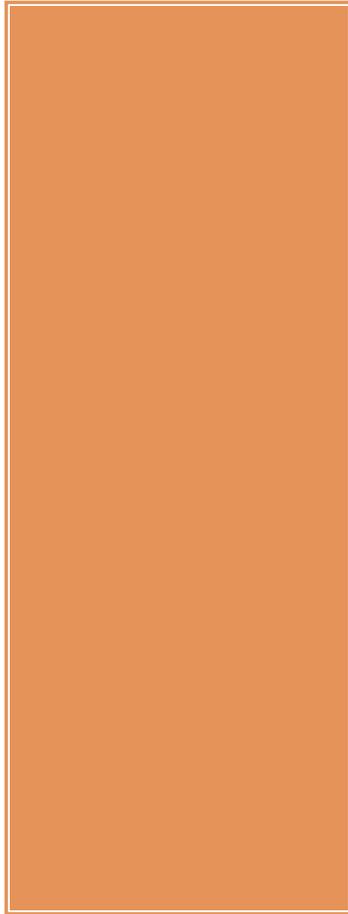


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- This is a method for systematically determining the extent of continuum versus intermittent/collisional behaviour within a laboratory debris flow
 - Roughness is only important in the when the roughness length is greater than or equal to the mean particle size
 - But, we see that snout-body architecture formation does not require mixtures of particle sizes

What about saturation?

- So if all that snout and body architecture needs is for the may-or-may-not be larger material at the top to be moving more quickly than the rest of the flow, then the level of saturation should be important
- While rebuilding the chute to investigate this, I was wondering over what range a particle leaving a fluid surface remained 'in touch' with the surface

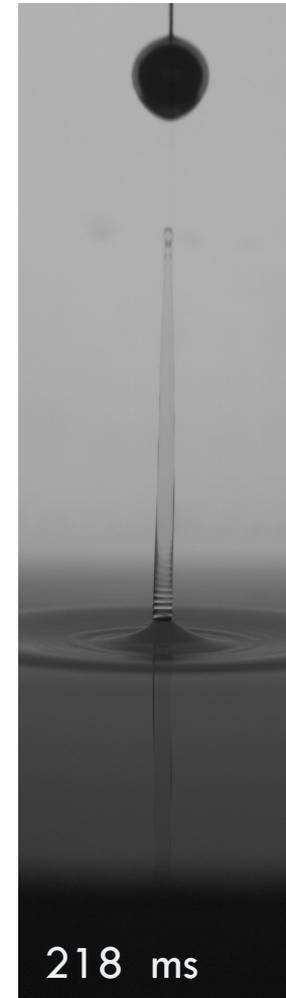
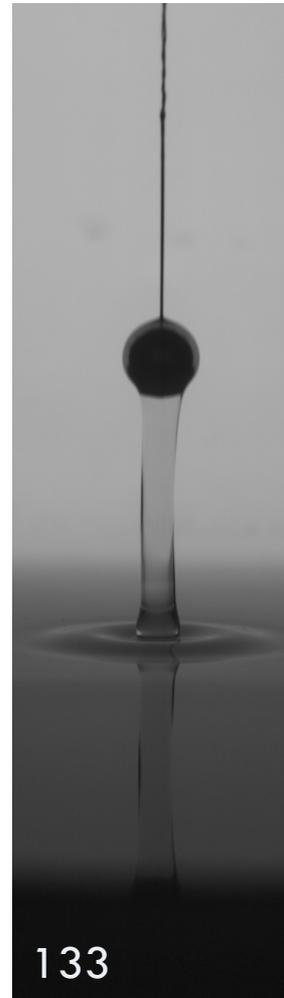
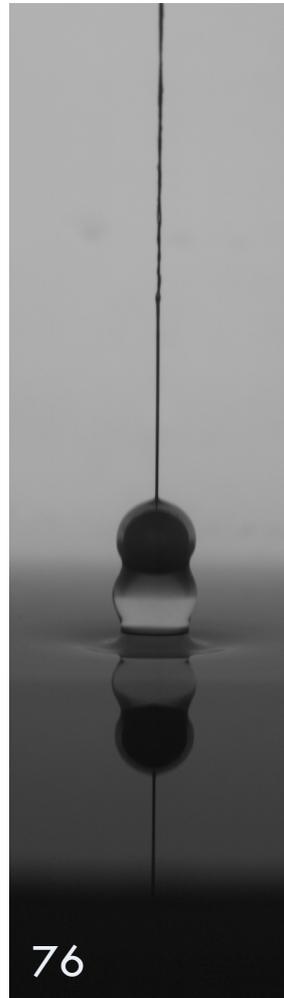
To start with, consider one....



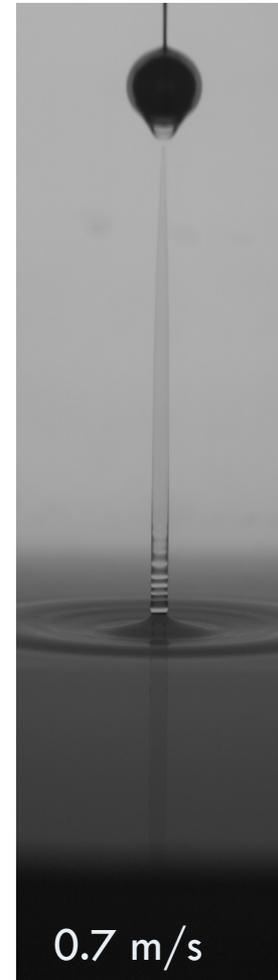
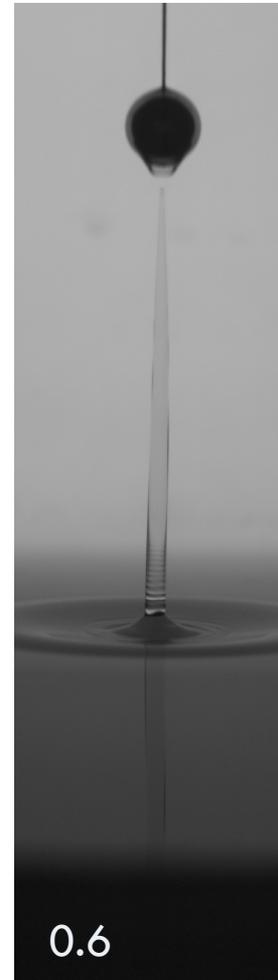
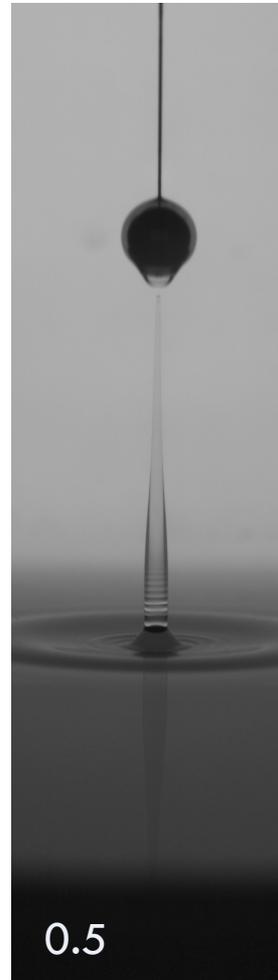
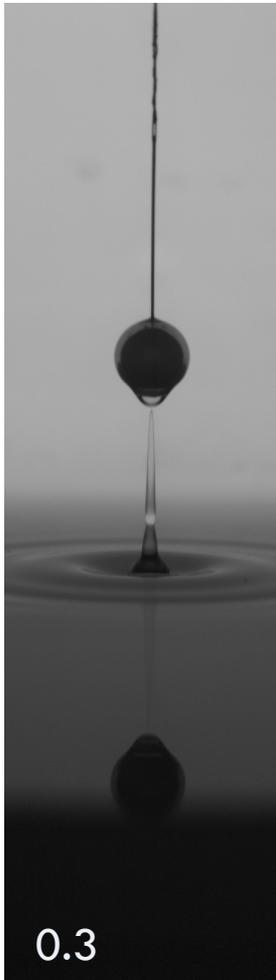
Experiments

- Particle diameters: 4, 6, 8, 10, 12 mm
- Winch speeds: 0.2, 0.3, 0.4, 0.5, 0.6, 0.7 m/s
- Fluids: Water, Kaolin solution (10% vol)

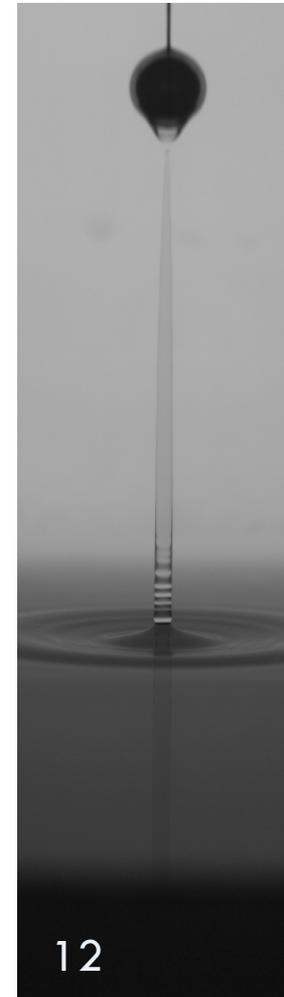
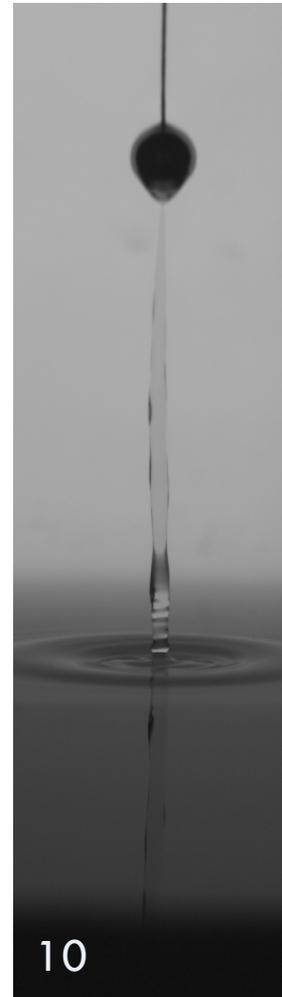
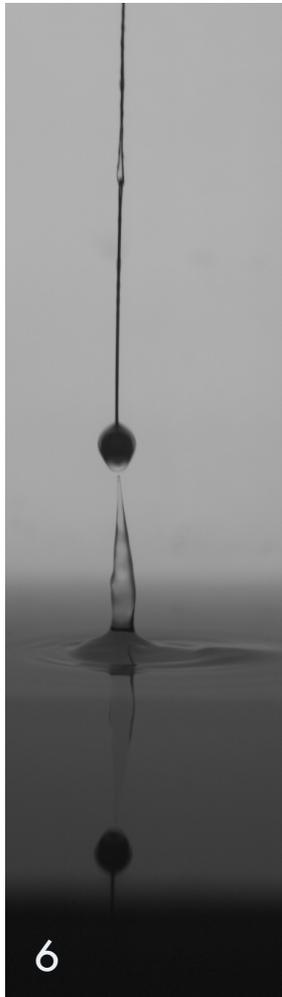
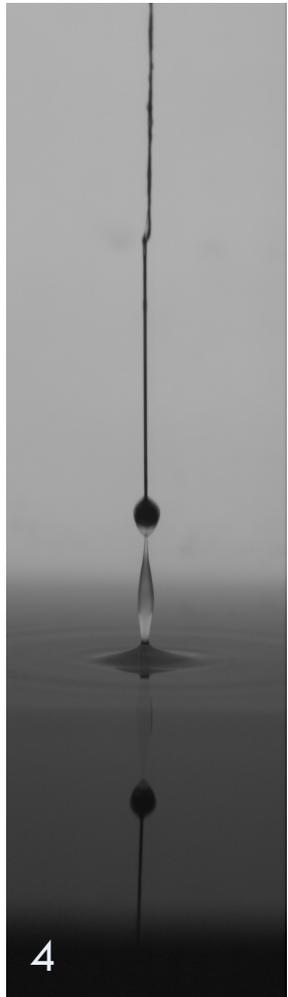
What happens?



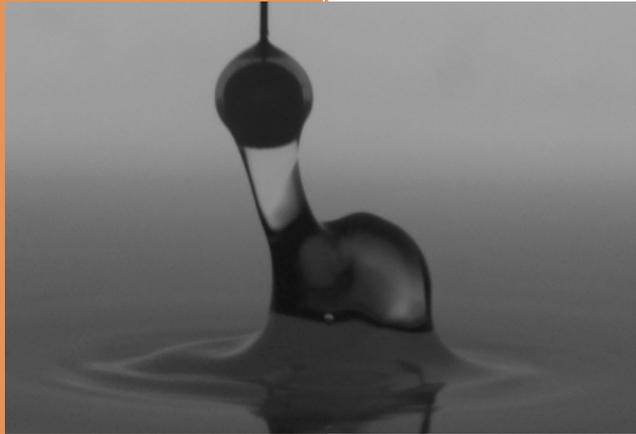
Variation with winch speed, 12 mm



Variation with particle size (0.7 m/s)

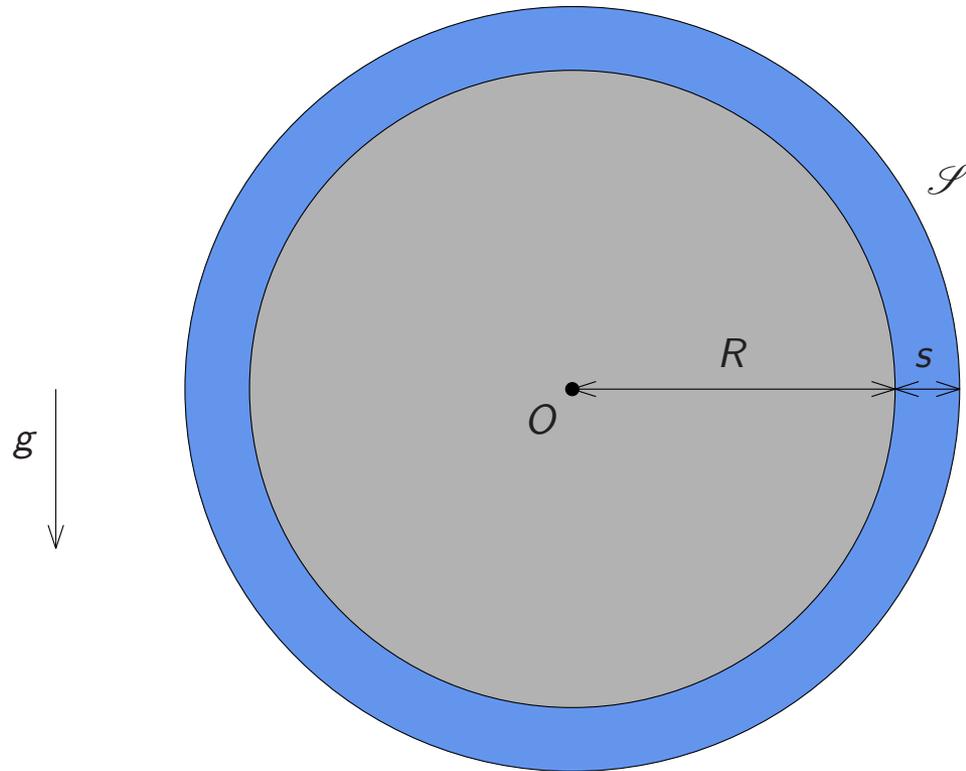
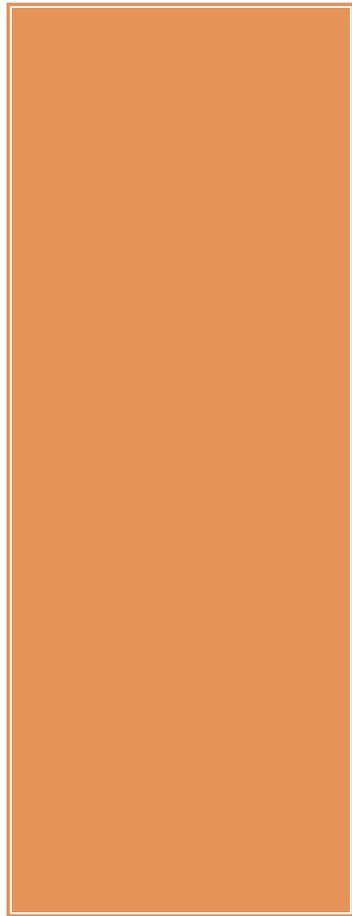


Ejection jet



- As the particle exits the surface, it leaves behind a little vortex ring
- This continues to accelerate fluid through it after the particle has left
- It may even *accelerate* the particle
- A force balance allows to estimate a timescale $T_j = \frac{u}{g}$

Draining flow



Stokes flow in a thin layer, invoke lubrication assumptions. Short time asymptotics...

Draining timescale

This draining flow has a layer Reynolds number that looks like

$$\text{Re}_l = \frac{g(\varepsilon R)^3}{\nu^2}$$

and, without surface tension, that leads to a timescale

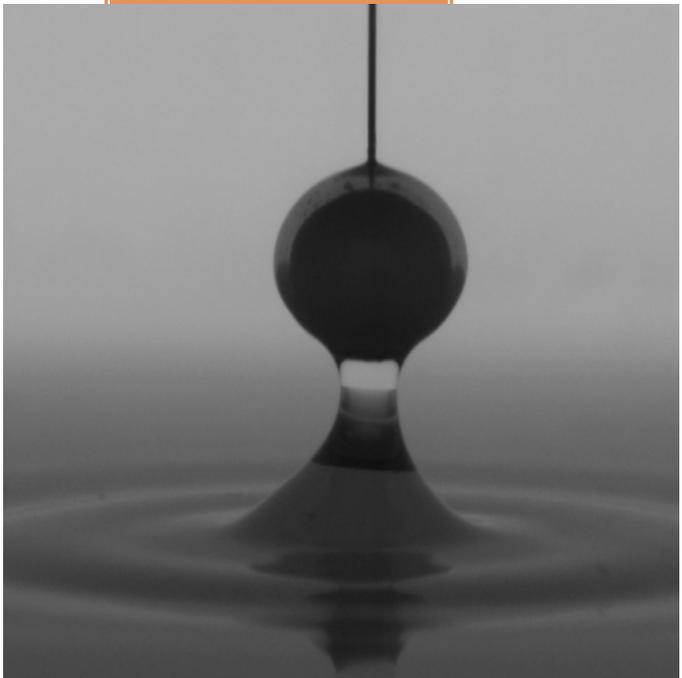
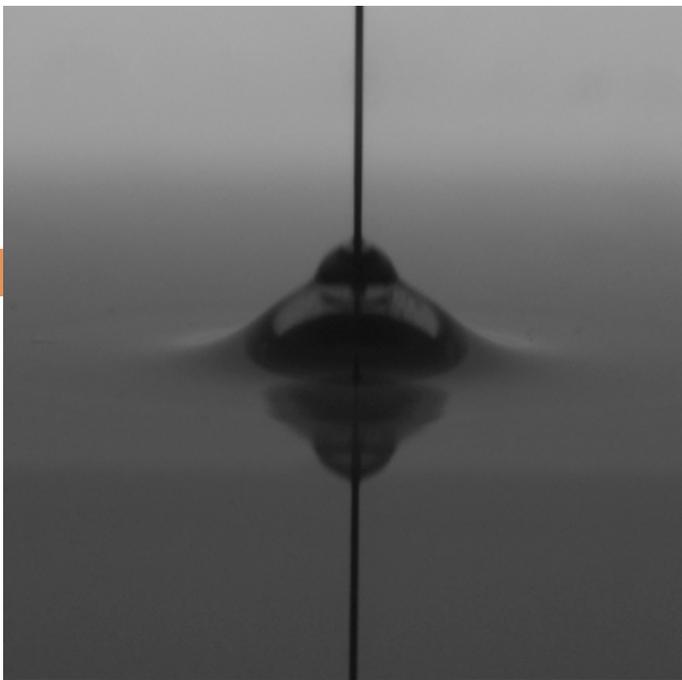
$$T_d = \frac{\nu}{\varepsilon R g}$$

Clearly, this depends on the original layer thickness

Coating

- So how thick is the initial layer in this draining problem? i.e. ε
- 1. Inertial?
- 2. Capillary?
- If inertial, then I can imagine that the fluid that gets dragged with the particle is the added mass
- If capillary, ($Ca \ll 1$, ours are $\sim 10^{-2}$), the coat thickness varies with $Ca^{2/3}$

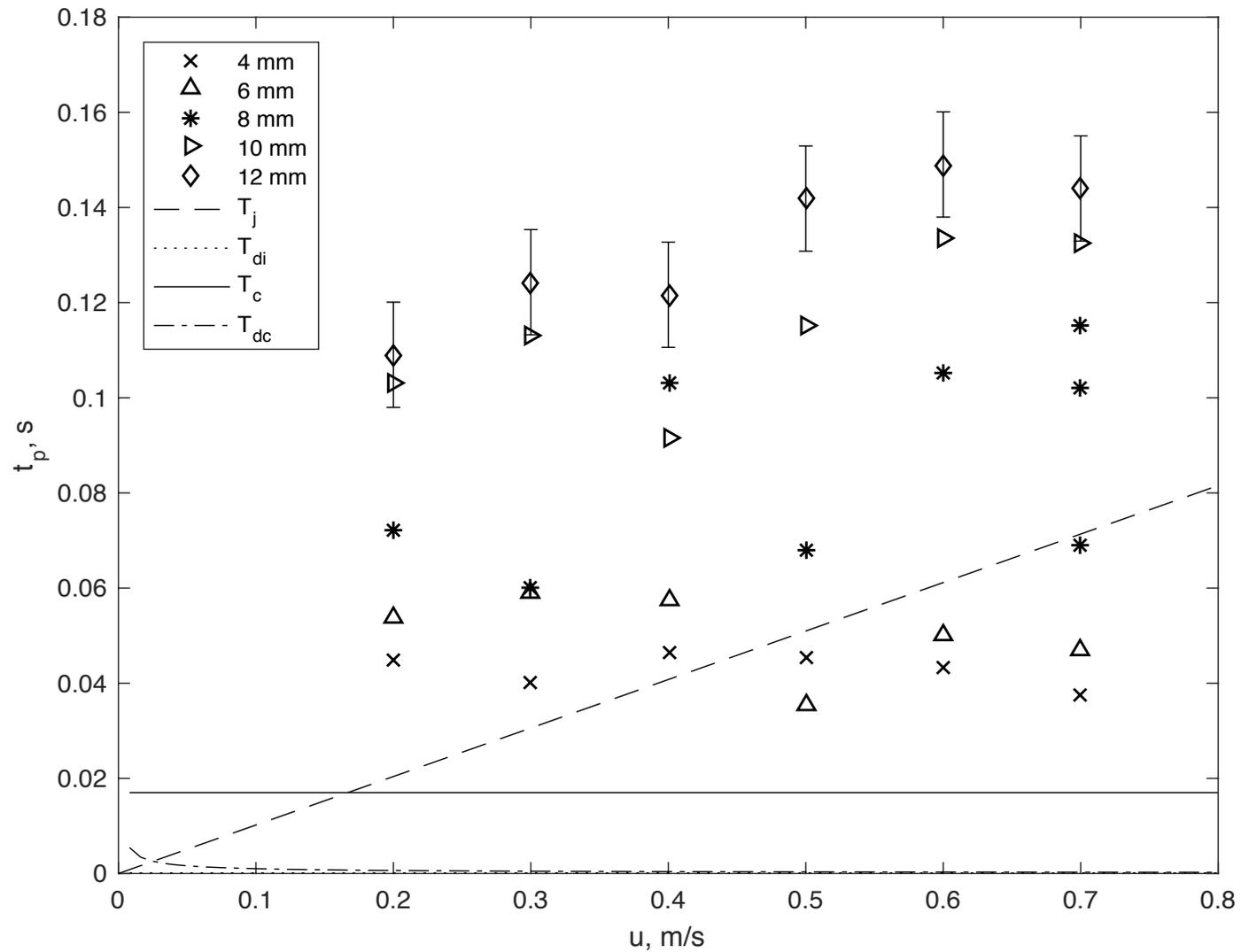
(Landau Levitch 1942) $Ca = \frac{w\nu\rho}{\gamma}$



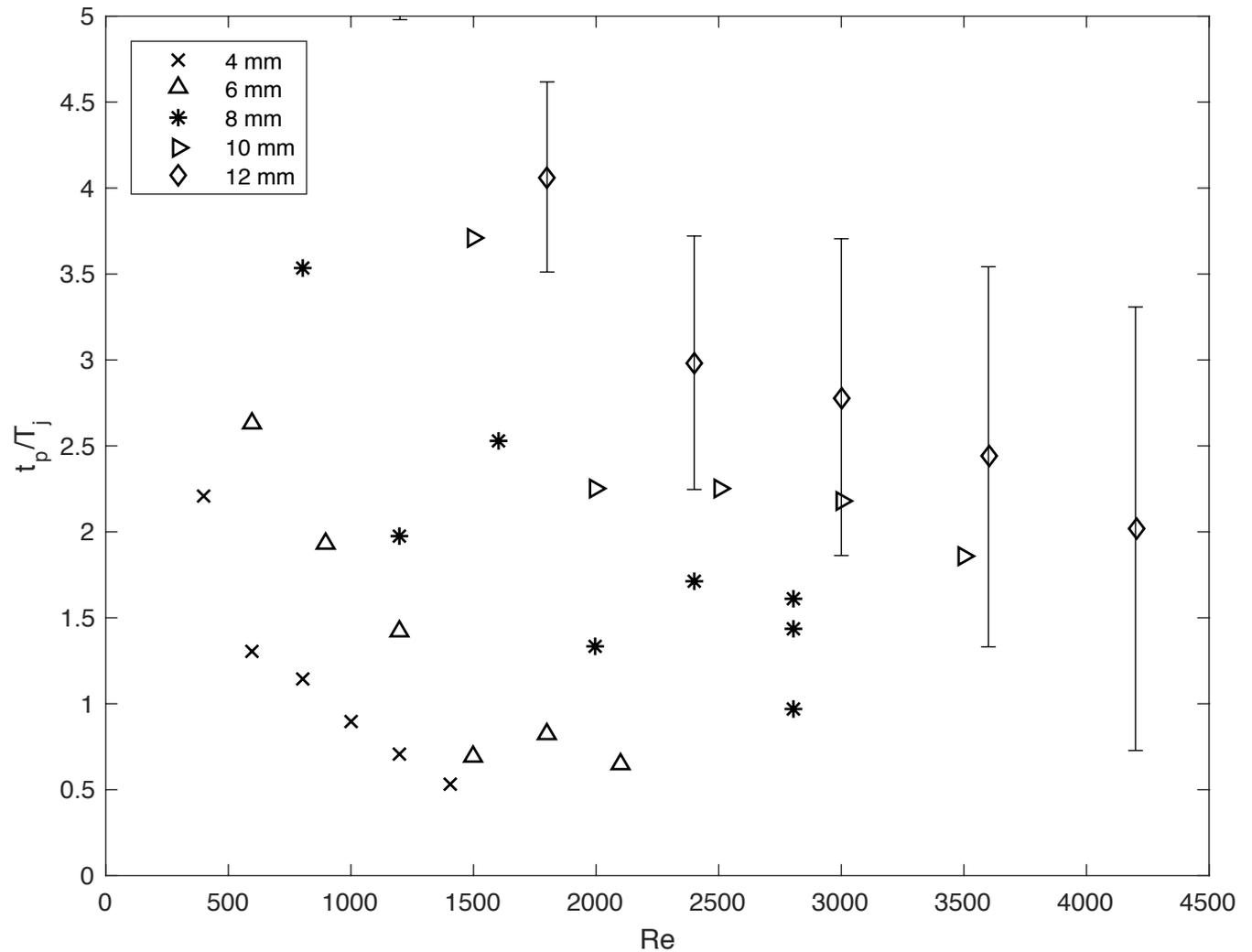
Timescales

- Ejection jet $T_j = \frac{u}{g}$
- Draining $T_d = \frac{\nu}{\varepsilon R g}$
 - where $\varepsilon = \text{Ca}^{2/3}$
 - or $\varepsilon = 0.14$

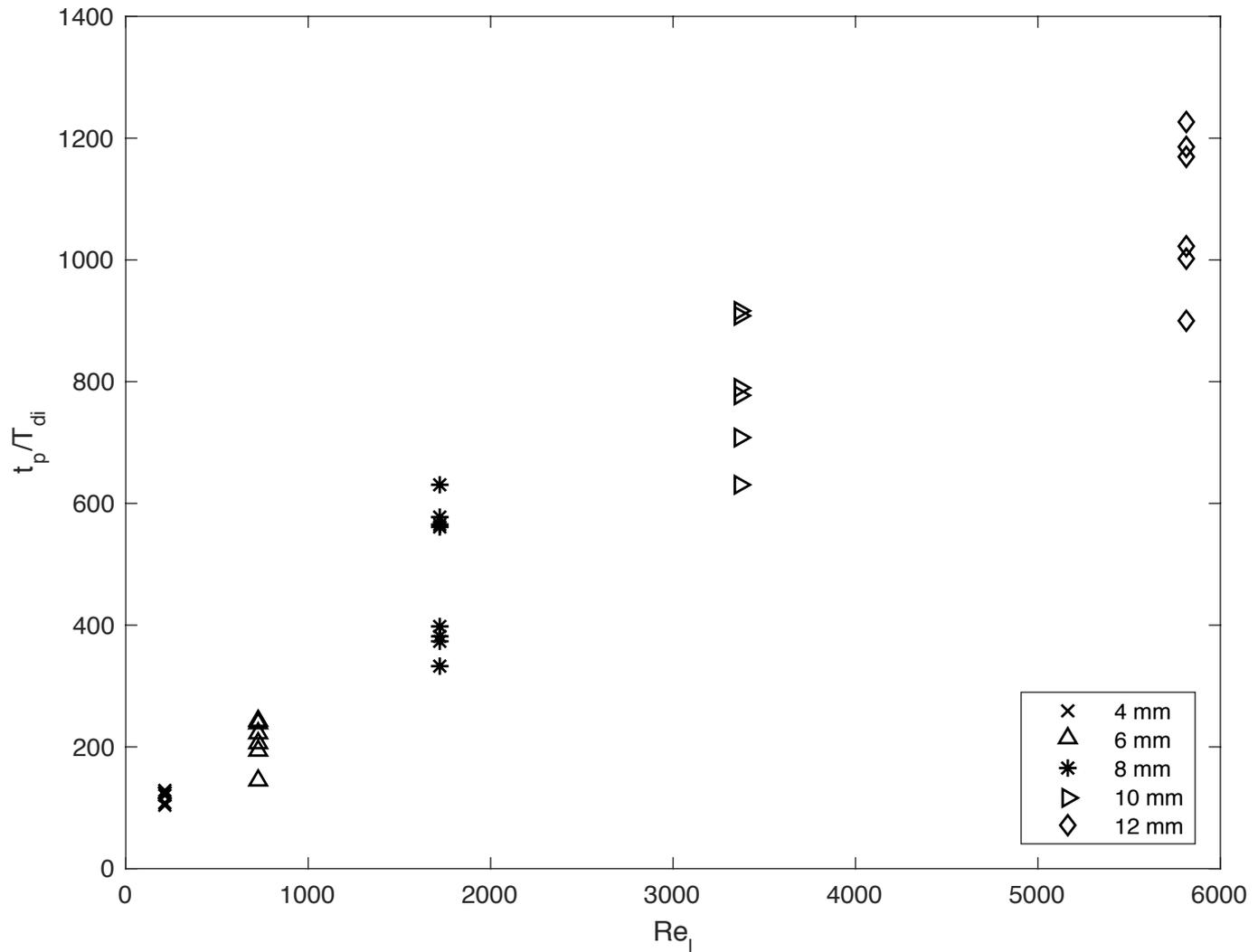
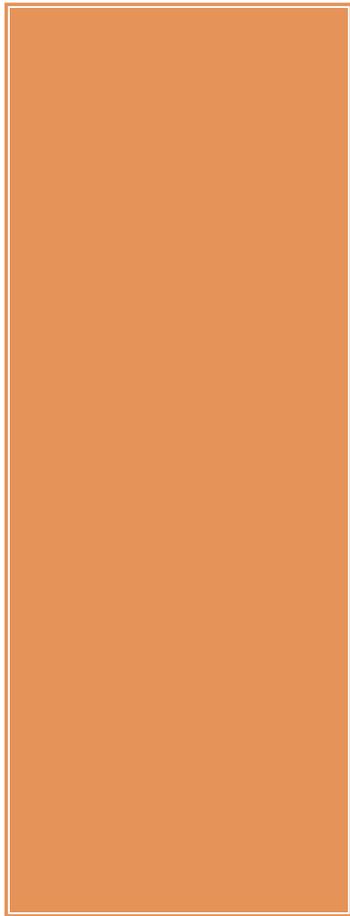
Raw data



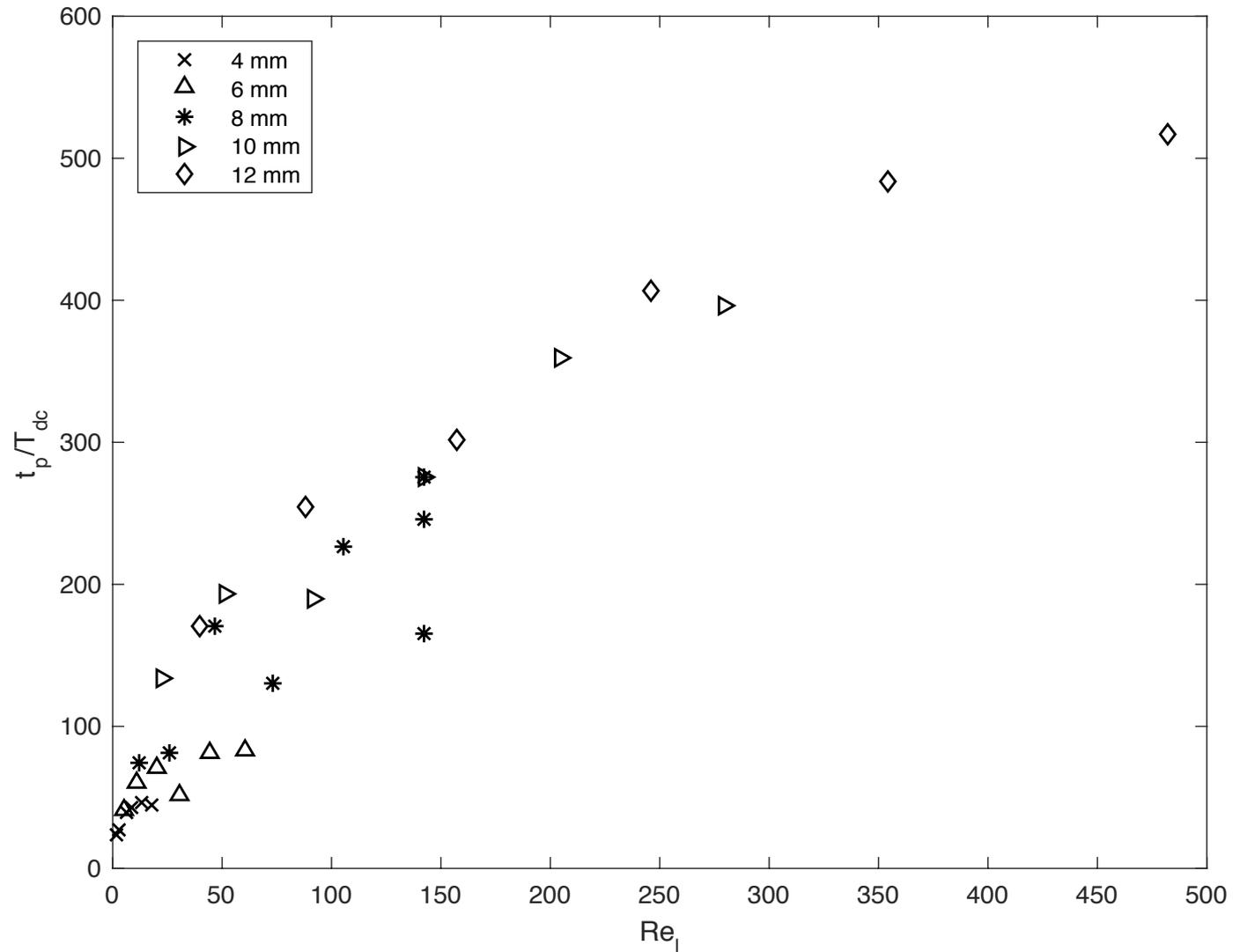
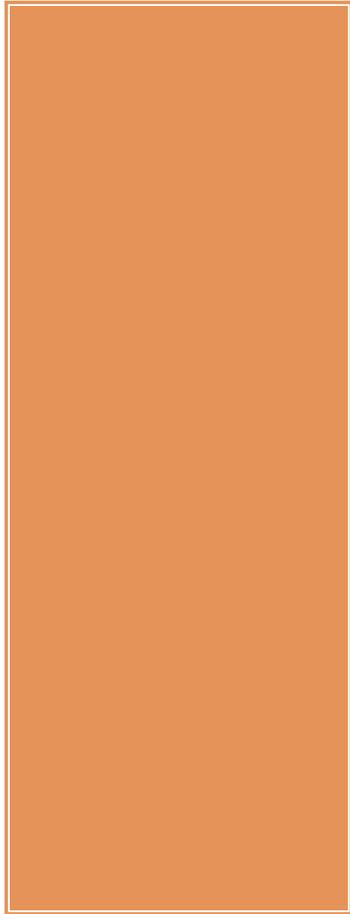
Ejection jet scaling



Draining scaling, 'inertial coating'



Draining timescale, 'capillary coating'



So...

- The sustaining of the tendril scales more convincingly with time scales inferred from analysis of the draining flow
- But this itself depends on the way in which the particle became coated
- Include surface tension in the draining flow
- I would have done the experiments a bit differently if I had known this!

In reality?

- Rarely have real debris flows where a particle could eject in this way
- But this might give a handle on the extent of the 'intermittent' zone
- Or why particles stay *in* the flow
- More significant for lab flows