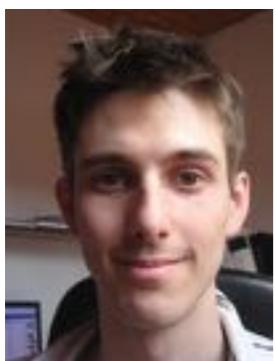


Interaction-induced transparency in the strong coupling regime of polaritons in photonic crystal waveguides

Francesco Piazza

(Max-Planck Institute for the Physics of Complex Systems)



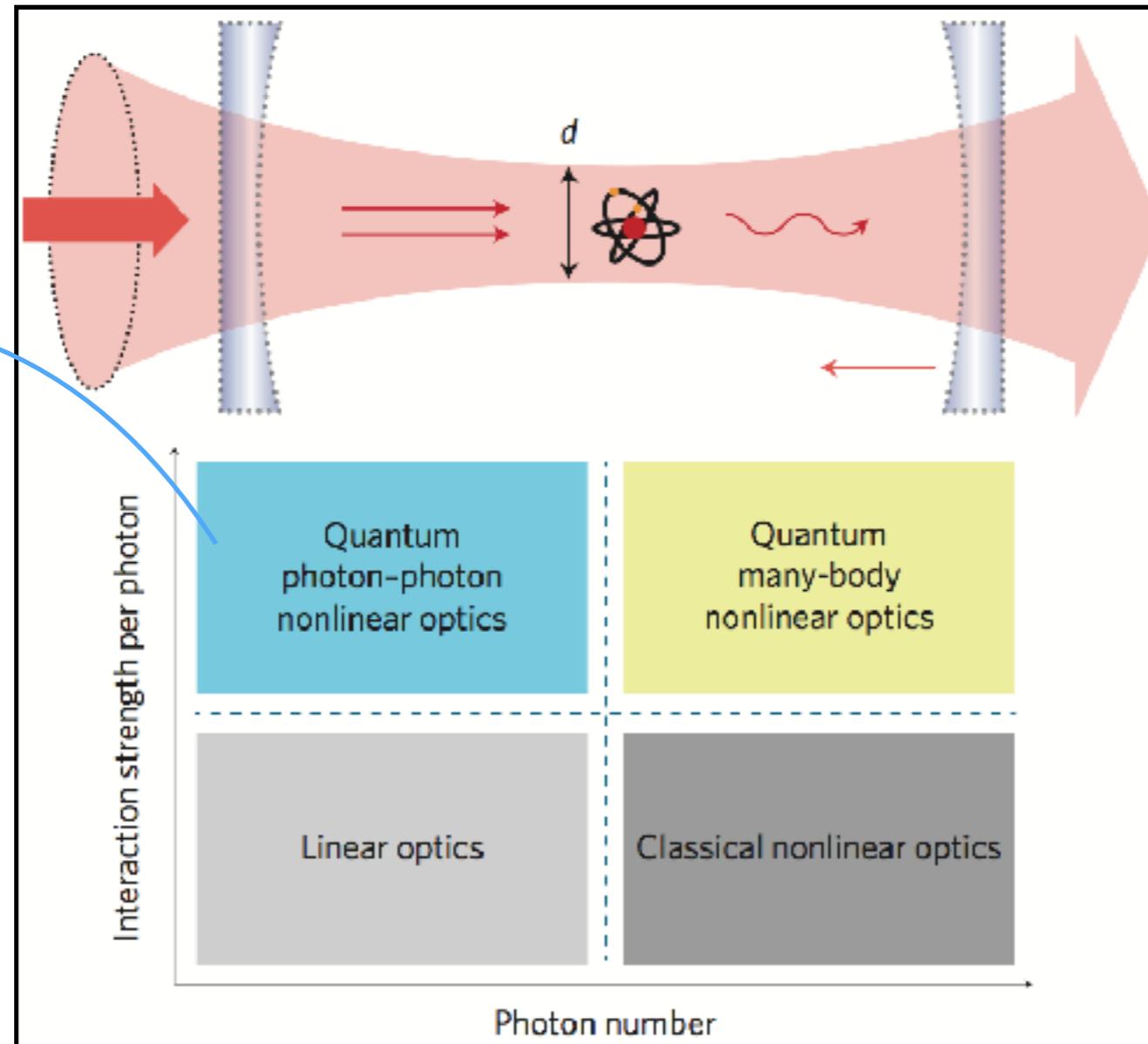
Johannes Lang
(TU Munich)



Darrick Chang
(ICFO Barcelona)

Quantum nonlinear optics

Review: Chang, Vuletic, Lukin, Nat. Phot. 8, 685 (2014).

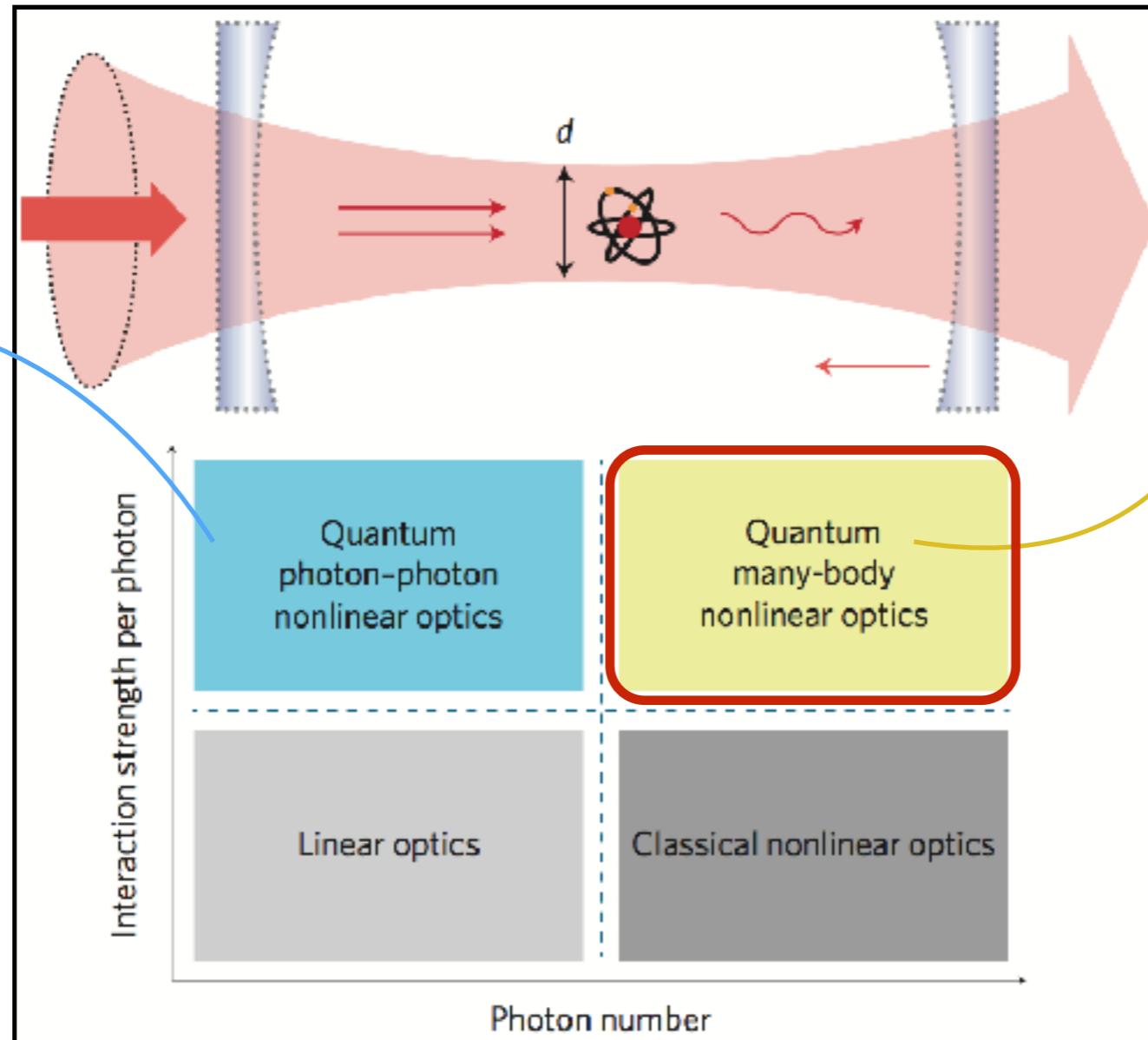


Progress has been made with several devices, coupling the electromagnetic field to:

- Trapped ions (Innsbruck)
- Molecules (Erlangen)
- Superconducting circuits (Paris, Yale, Santa Barbara, ETH, New York)
- Neutral Atoms

The challenge of quantum nonlinear optics

Review: Chang, Vuletic, Lukin, Nat. Phot. 8, 685 (2014).

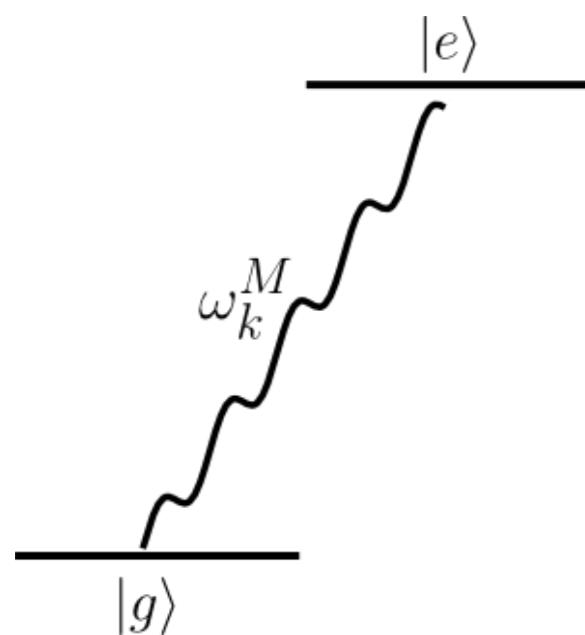


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- **Neutral Atoms**

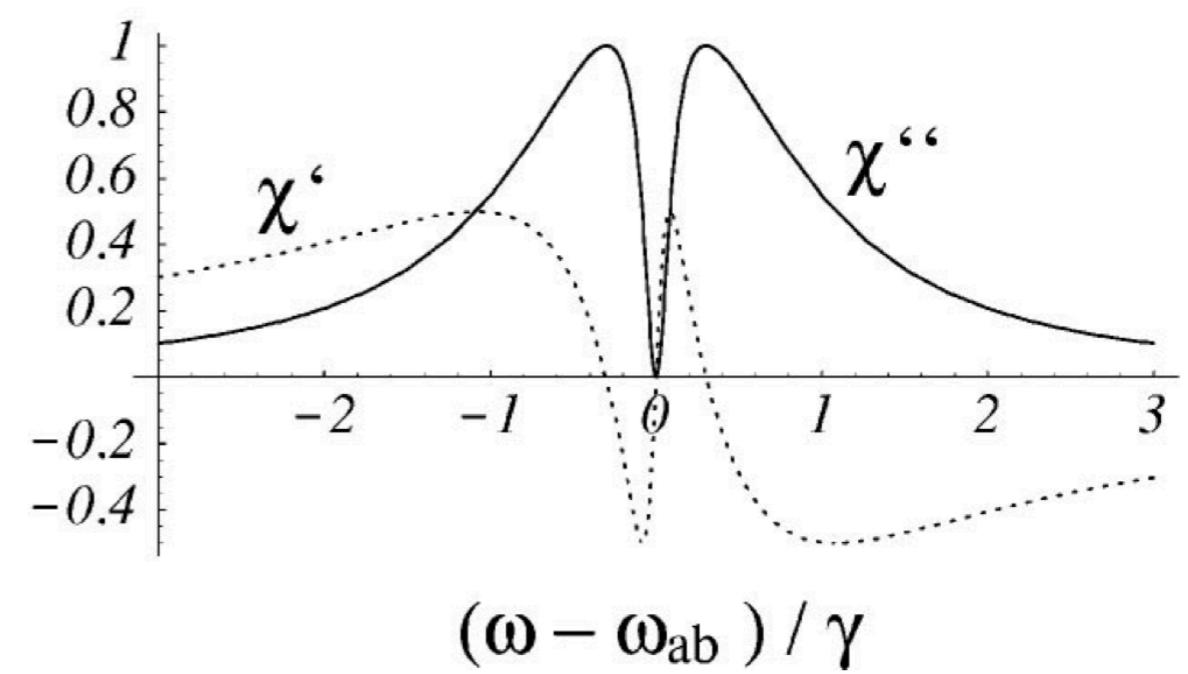
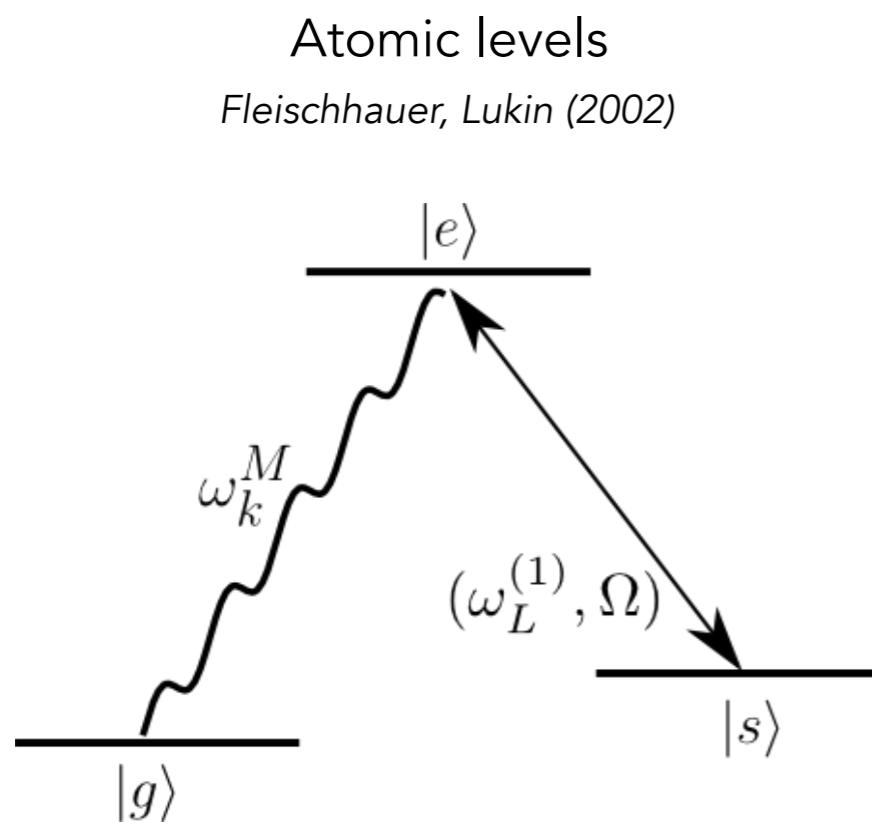
Electromagnetically-Induced Transparency (EIT)

Increase light/atom interaction probability by
slowing down the light without absorbing it



Electromagnetically-Induced Transparency (EIT)

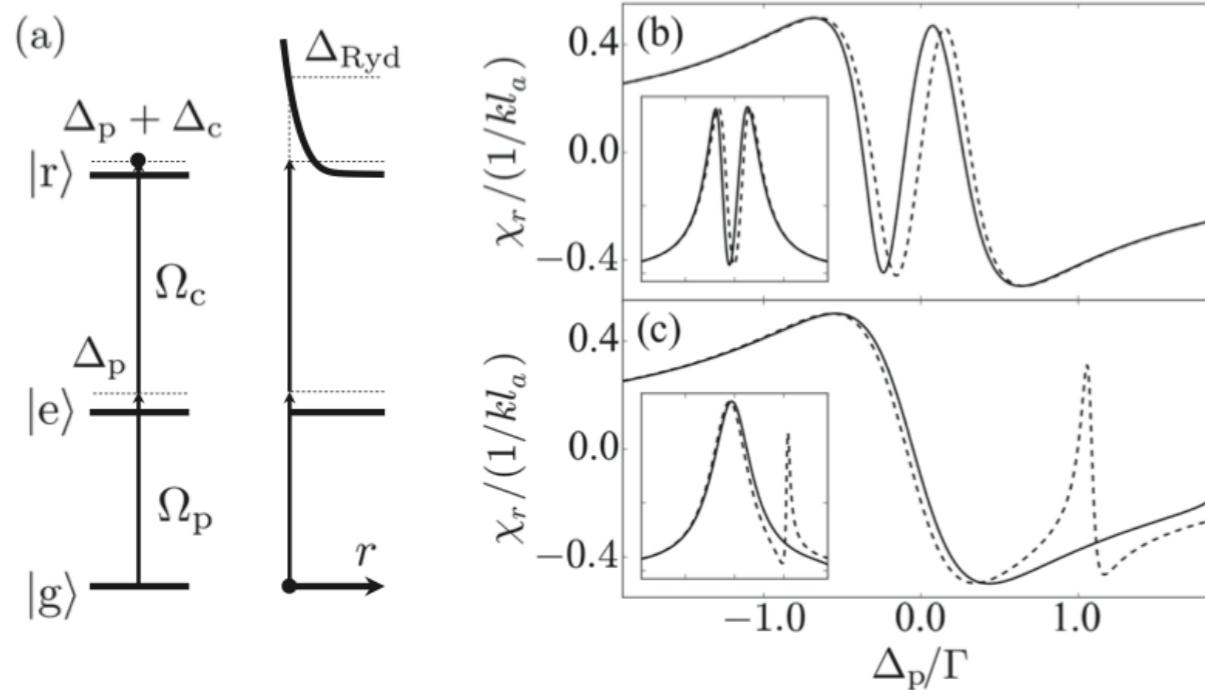
Increase light/atom interaction probability by
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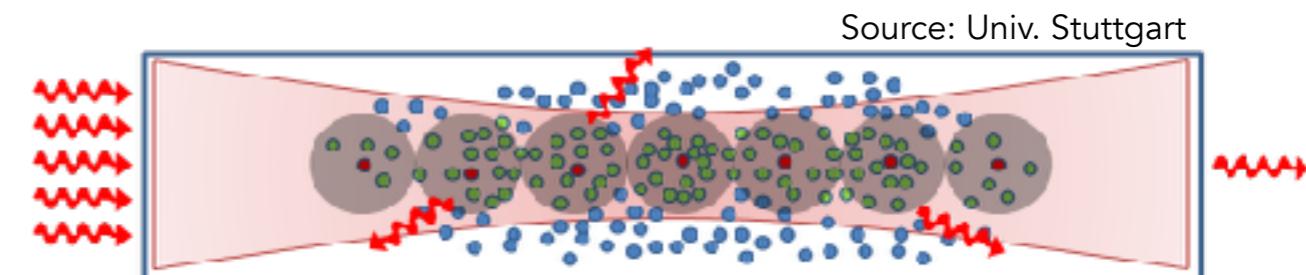
EIT with Rydberg interactions

J. Phys. B: At. Mol. Opt. Phys. 49 (2016) 152003

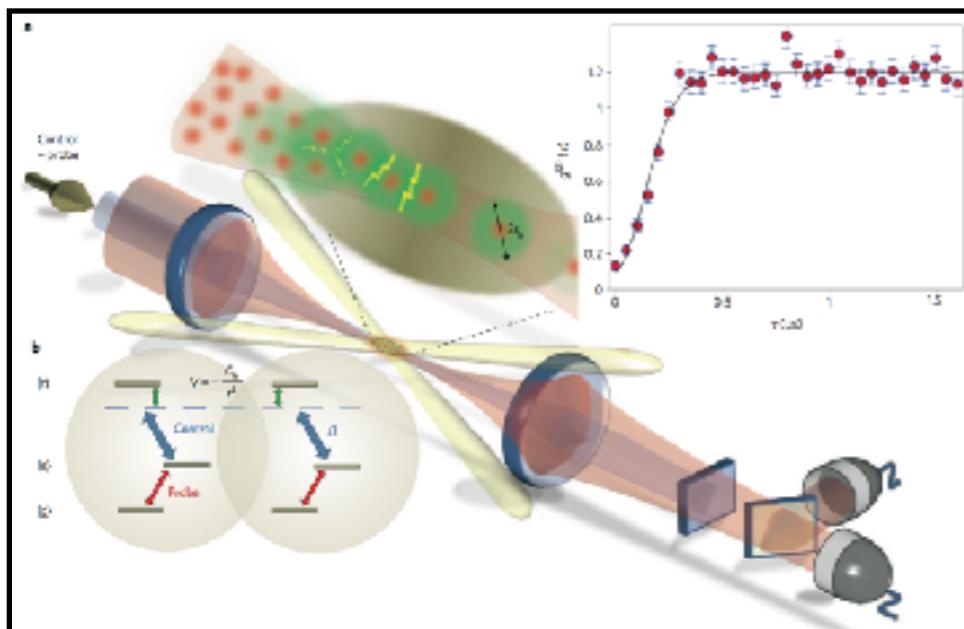
To



Rydberg atom interact strongly over a **blockade radius**
Only one atom can be excited within radius

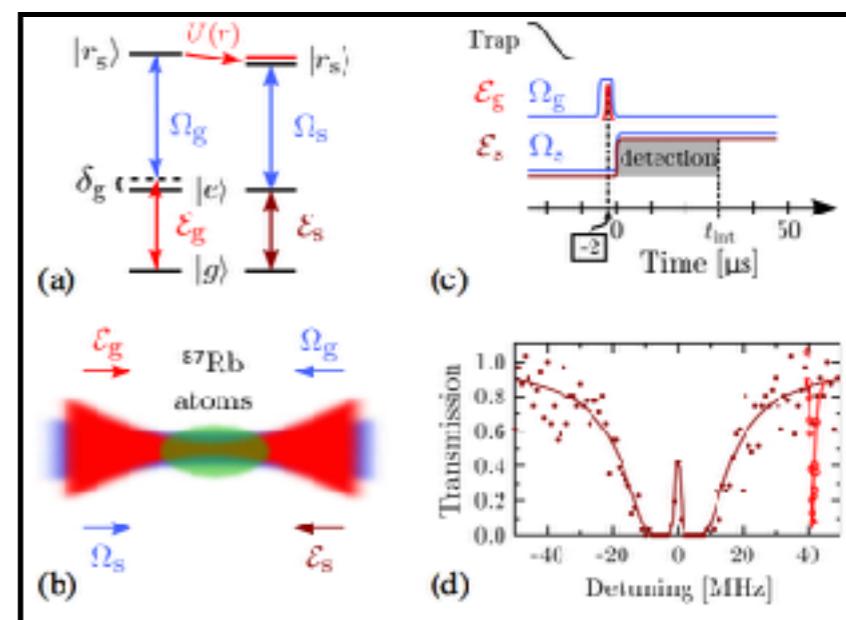


Application: Photon Blockade



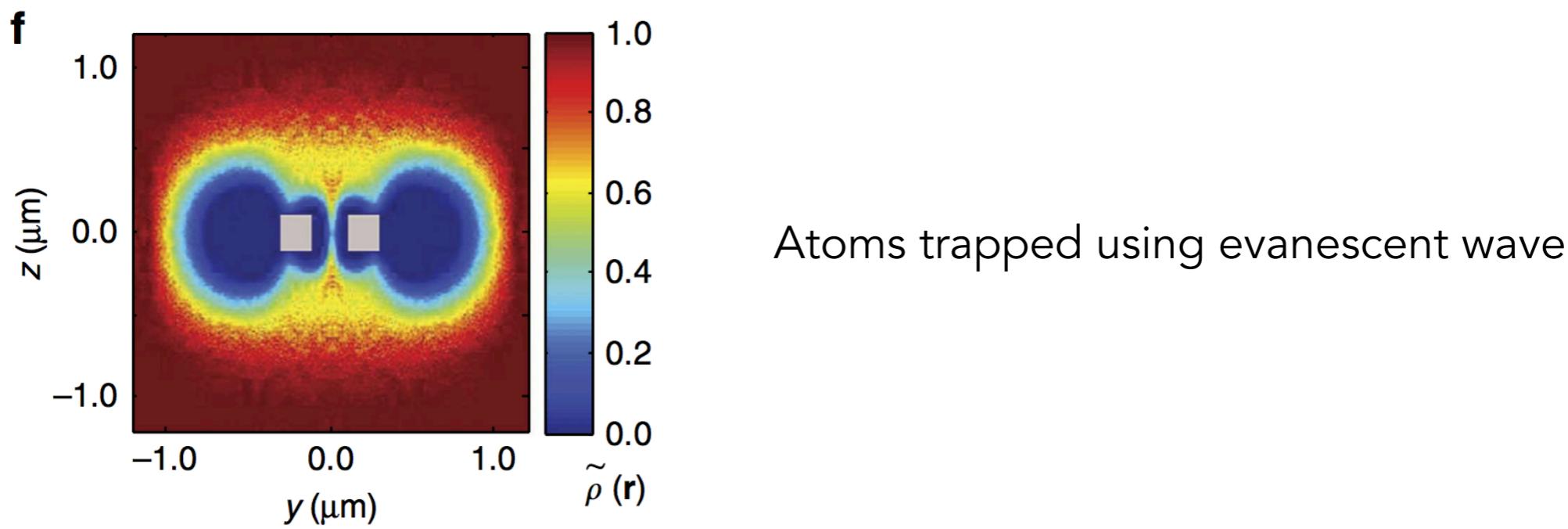
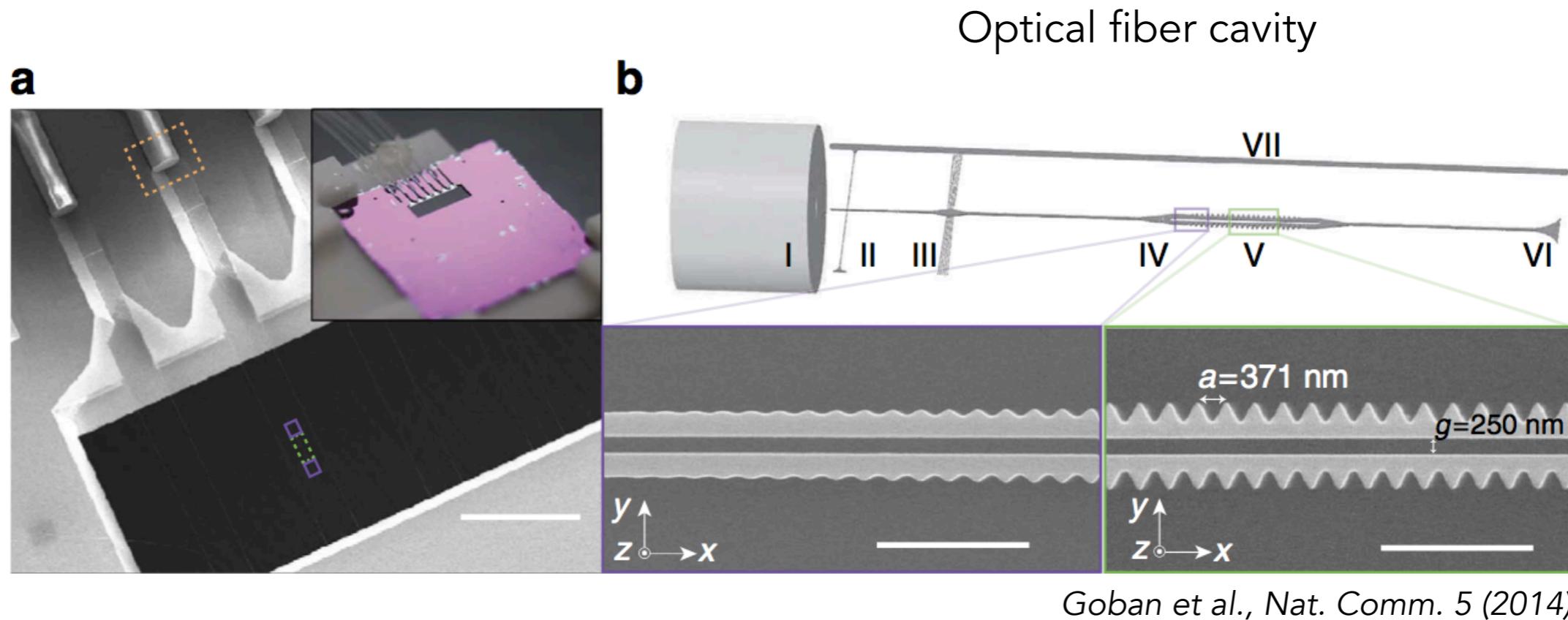
Exp.: Peyronel, et al., Nature 488, 57 (2012).

Application: All optical single-photon transistor

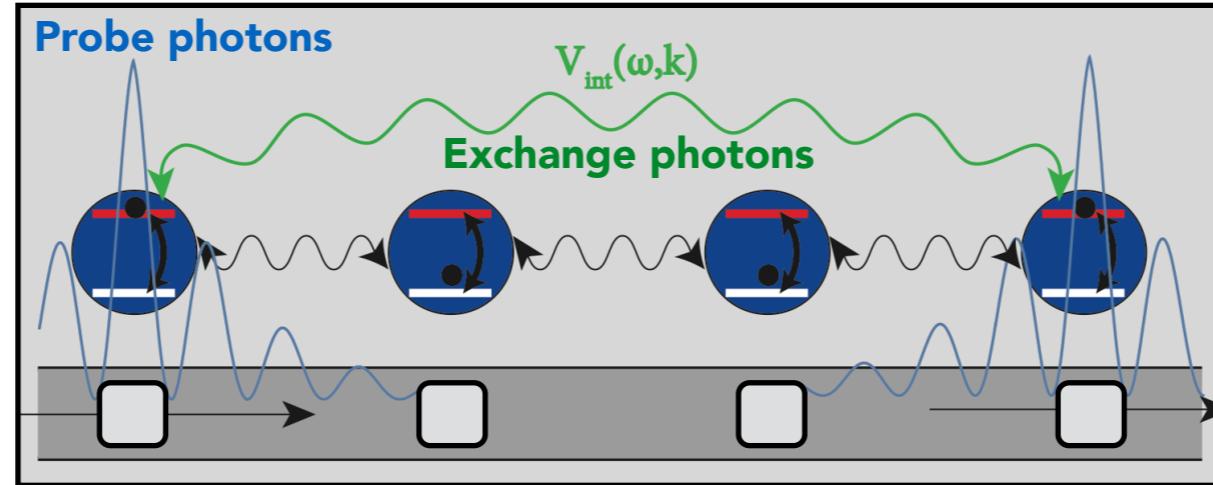


Exp.: Gorniaczyc, et al., Phys. Rev. Lett. 113, 053601

Experiment: atoms coupled to photonic crystal waveguides

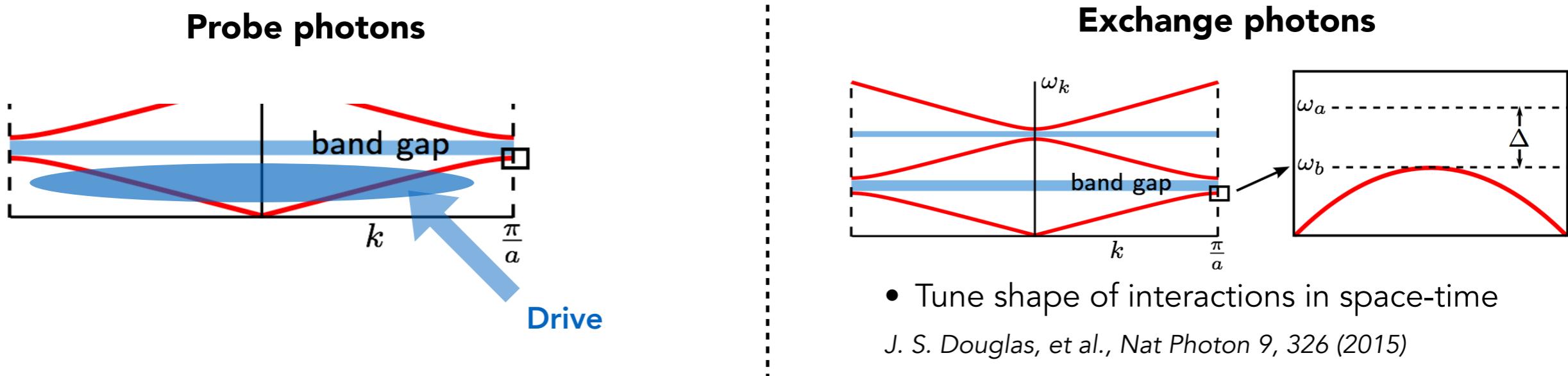


Interacting polaritons in photonic crystal waveguides

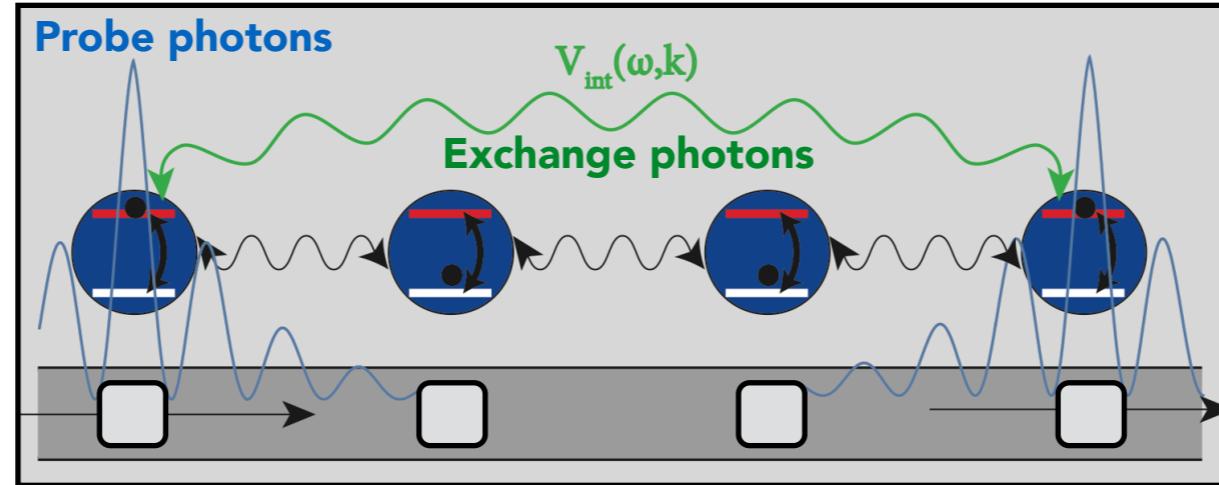


Hybrid atom-photon architecture

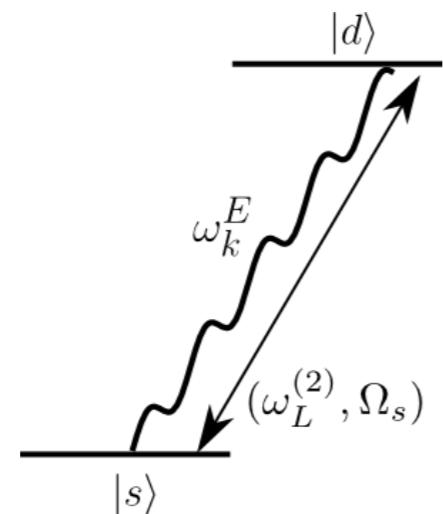
- Multi-level atoms fixed positions
- Probe photons within propagating band of lattice
- Exchange photons mediating interactions



Shaping the light-mediated interactions



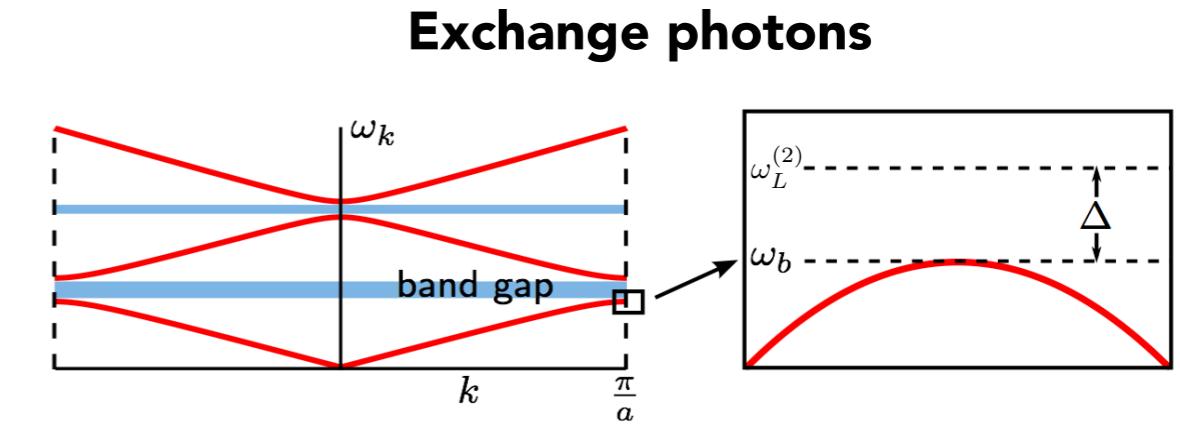
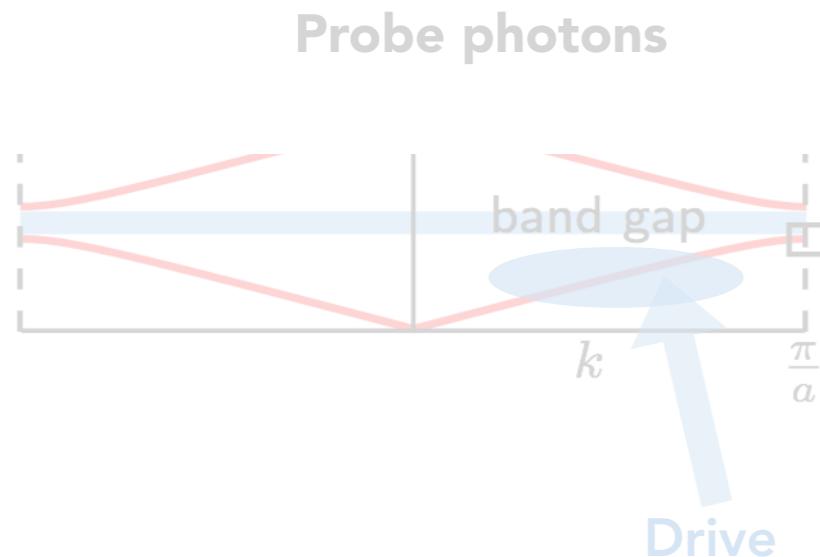
Atomic levels



$$\hat{H}_{\text{int}} = \sum_z \left[\Omega_s e^{-i\omega_L^{(2)} t} \hat{a}_d^\dagger(z) \hat{a}_s(z) + g_E \hat{a}_E(k) e^{ikz} u_k^E(z) \hat{a}_d^\dagger(z) \hat{a}_s(z) + h.c. \right]$$

$$\hat{H}_{\text{int}}^{\text{eff}} \sim \sum_{zz'} \hat{a}_s(z)^\dagger \hat{a}_s(z) e^{-|z-z'|/L_E} \hat{a}_s^\dagger(z') \hat{a}_s(z') \quad (\text{adiabatic elimination})$$

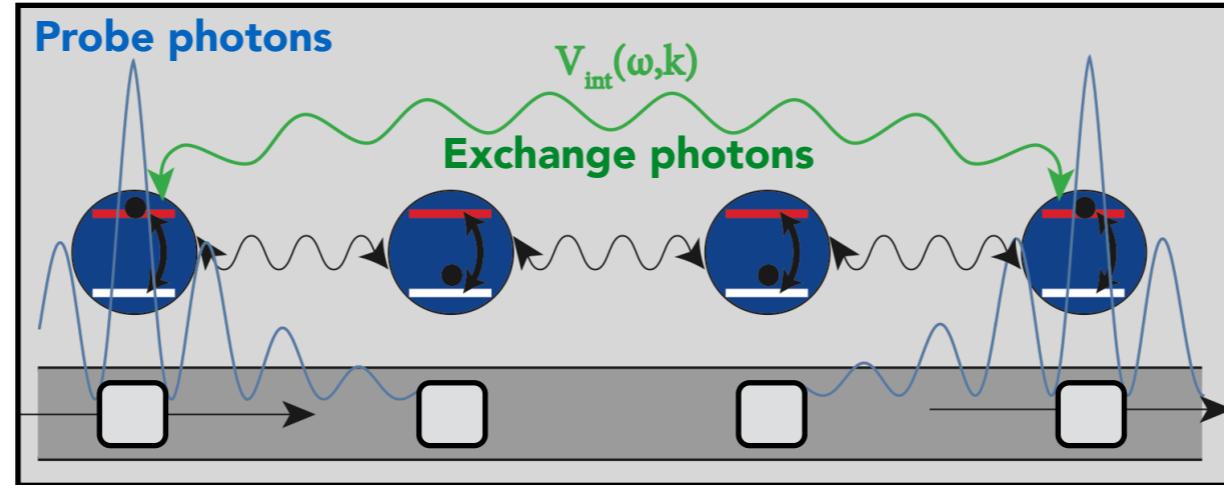
Exponential interaction range: $L_E \simeq \sqrt{\frac{\alpha_E}{\Delta}}$



Dispersion:
$$\int \frac{dk}{2\pi} (\alpha_E(k - k_0)^2) \hat{a}_E^\dagger(k) \hat{a}_E(k)$$

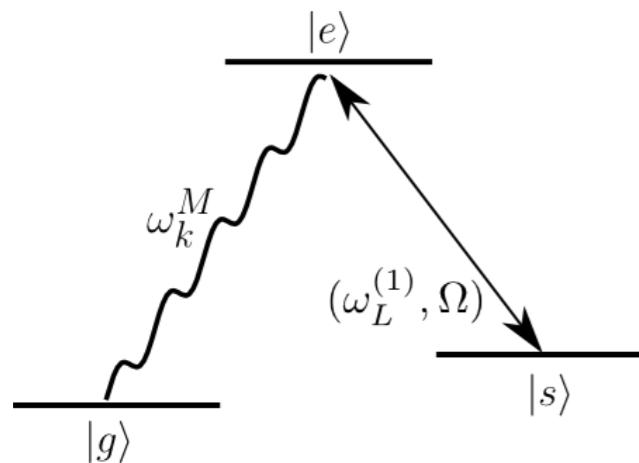
Losses: $\kappa_E \mathcal{D}[\hat{a}_E(k)]$

Probe photons propagation: Electromagnetically-Induced Transparency



Atomic levels

Fleischhauer, Lukin (2002)



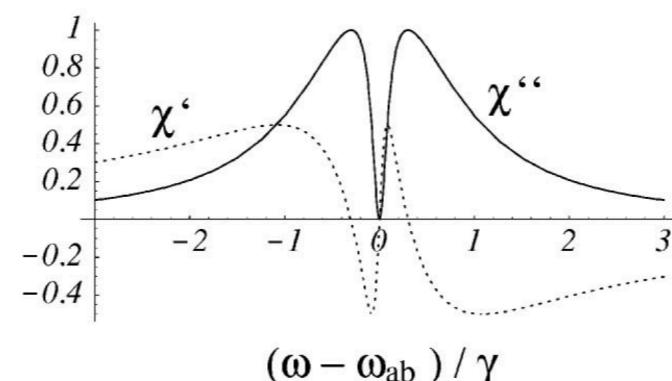
$$\hat{H}_{\text{int}} = \sum_z \left[\Omega e^{-i\omega_L^{(1)} t} \hat{a}_e^\dagger(z) \hat{a}_s(z) + g_M \hat{a}_M(k) e^{ikz} u_k^M(z) \hat{a}_e^\dagger(z) \hat{a}_g(z) + h.c. \right]$$

$$\omega_M(k) = \omega_{sg} - \omega_L^{(1)}$$

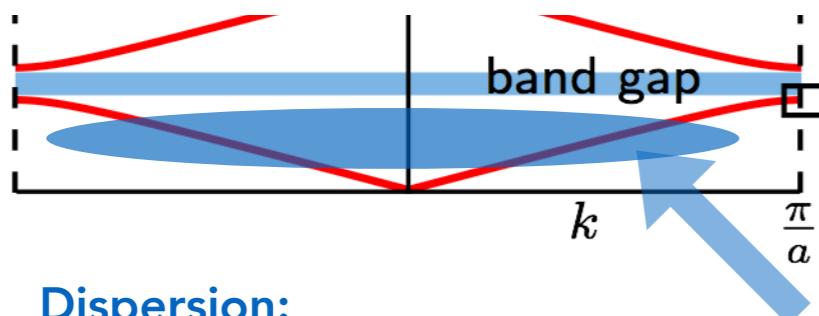
Large dispersion without absorption

EIT propagation range:

$$L_M \simeq \frac{v_M}{\kappa_M}$$



Probe photons



Dispersion:

$$\int \frac{dk}{2\pi} J_M (1 - \cos(ka)) \hat{a}_M^\dagger(k) \hat{a}_M(k)$$

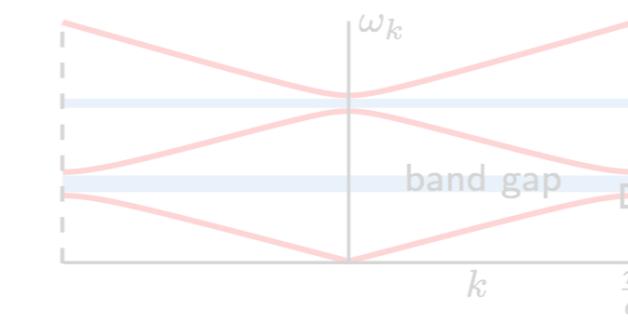
Losses:

$$\kappa_M \mathcal{D}[\hat{a}_M(k)]$$

Incoh. Drive:

$$\kappa_s \mathcal{P}[\hat{a}_M(k)]$$

Exchange photons



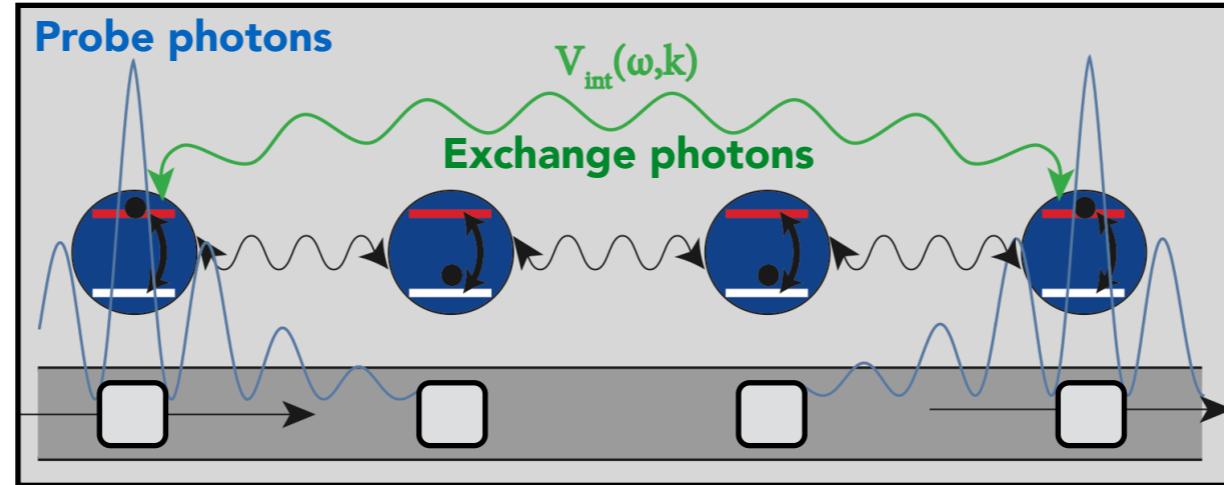
Dispersion:

$$\int \frac{dk}{2\pi} (\alpha_E(k - k_0)^2) \hat{a}_E^\dagger(k) \hat{a}_E(k)$$

Losses:

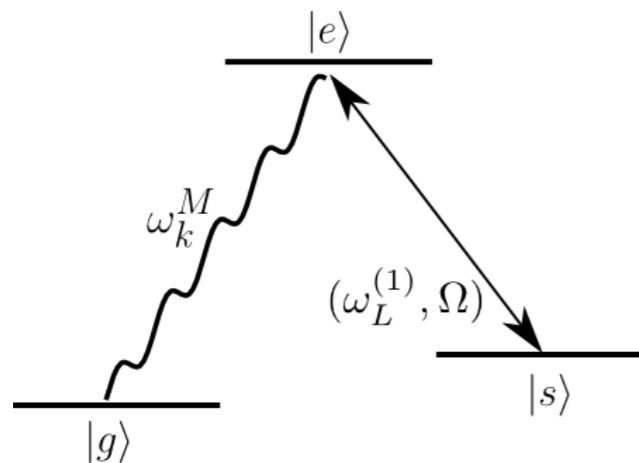
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Probe photons propagation: Electromagnetically-Induced Transparency



Atomic levels

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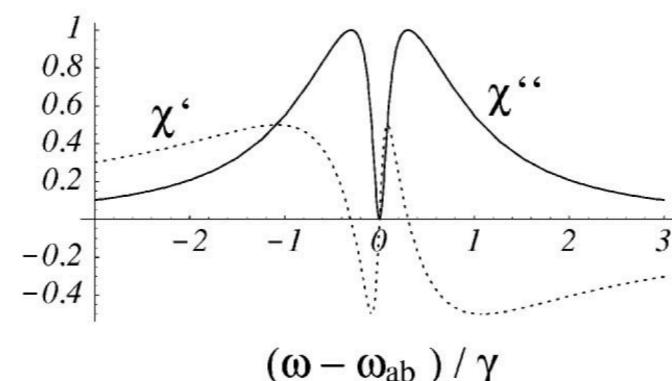
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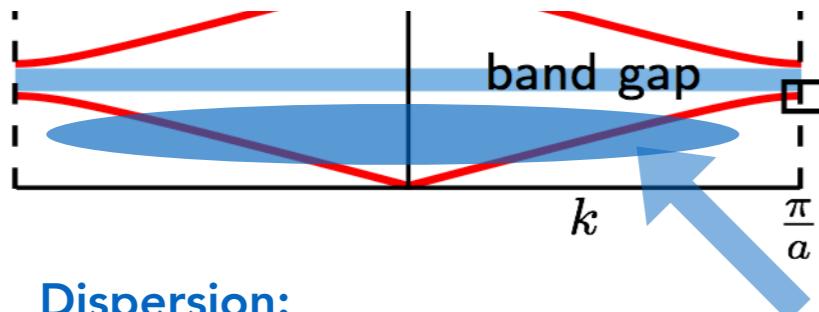
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Probe photons



Dispersion:

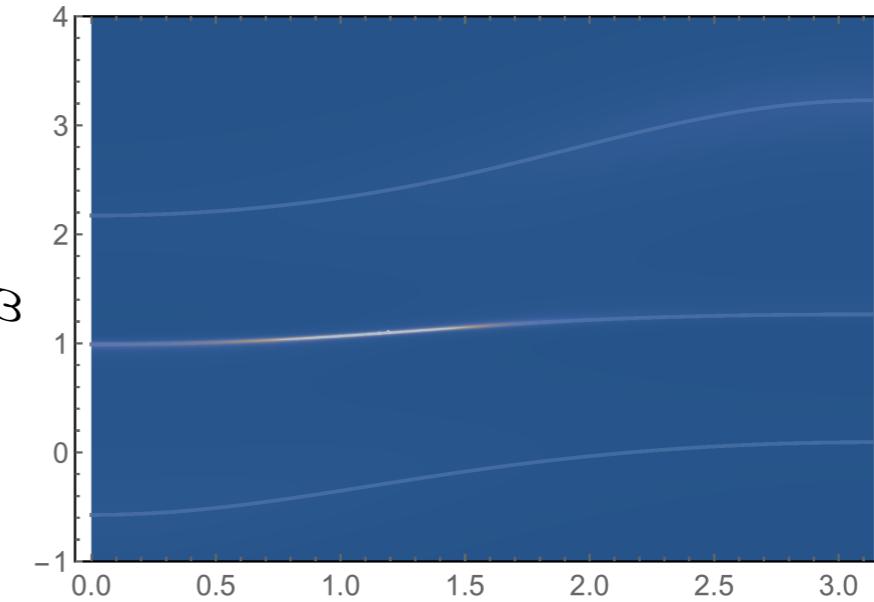
$$\int \frac{dk}{2\pi} J_M (1 - \cos(ka)) \hat{a}_M^\dagger(k) \hat{a}_M(k)$$

Losses:

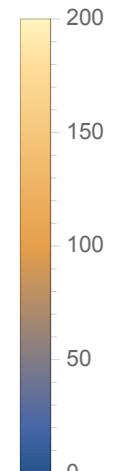
$$\kappa_M \mathcal{D}[\hat{a}_M(k)]$$

Incoh. Drive:

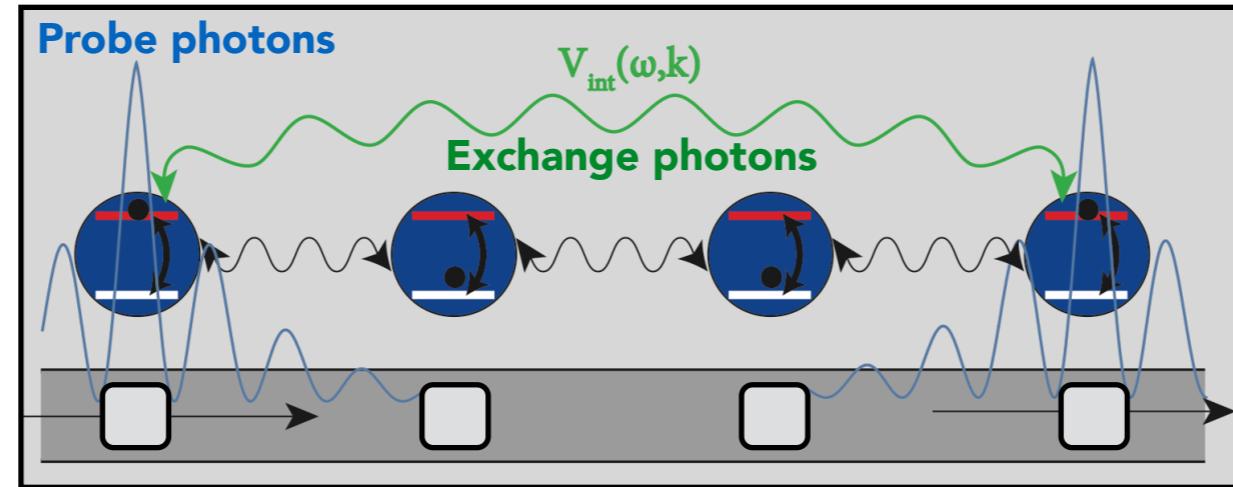
$$\kappa_s \mathcal{P}[\hat{a}_M(k)]$$



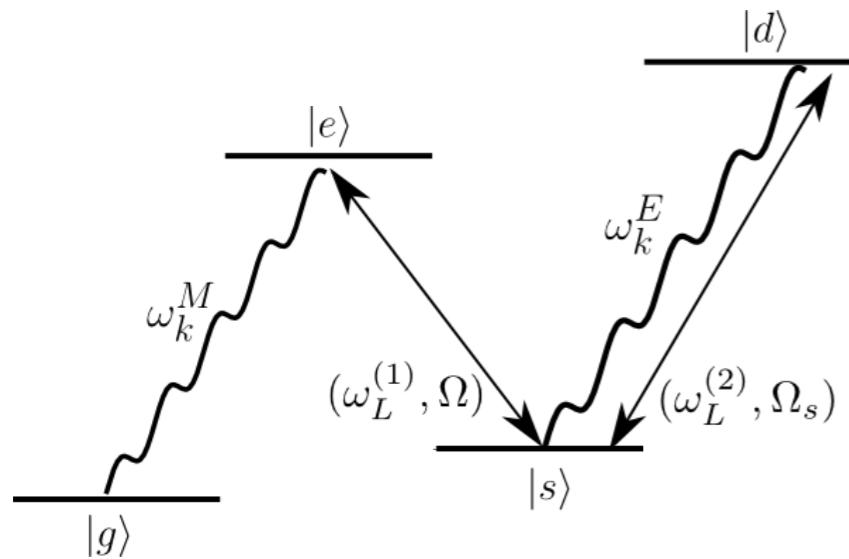
$$n_M(\omega, k)$$



Interacting EIT polaritons



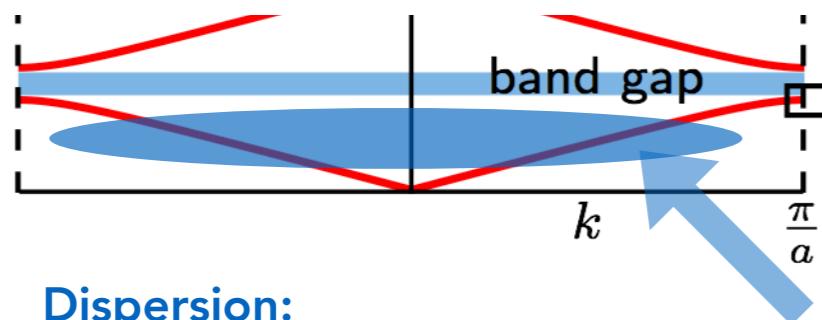
Atomic levels



$$\text{EIT propagation range: } L_M \simeq \frac{v_M}{\kappa_M}$$

$$\text{Exponential interaction range: } L_E \simeq \sqrt{\frac{\alpha_E}{\Delta}}$$

Probe photons



Dispersion:

$$\int \frac{dk}{2\pi} J_M (1 - \cos(ka)) \hat{a}_M^\dagger(k) \hat{a}_M(k)$$

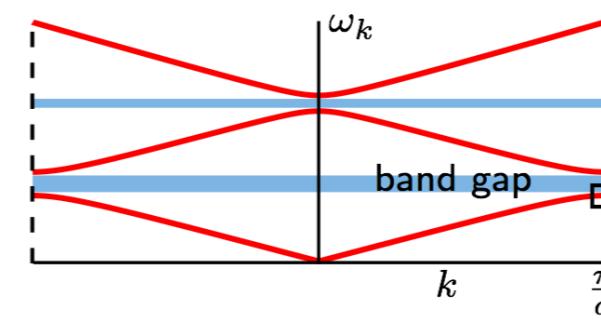
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Exchange photons



Dispersion:

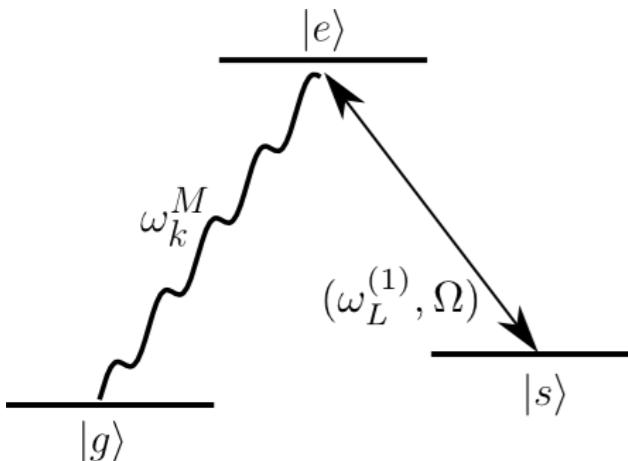
$$\int \frac{dk}{2\pi} (\alpha_E(k - k_0)^2) \hat{a}_E^\dagger(k) \hat{a}_E(k)$$

Losses:

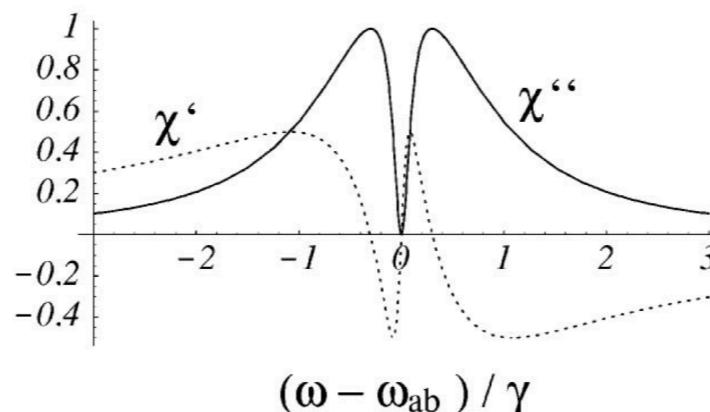
$$\kappa_E \mathcal{D}[\hat{a}_E(k)]$$

Non-equilibrium diagrammatic approach to EIT

Atomic levels



$$\hat{H}_{\text{int}} = \sum_z \left[\Omega e^{-i\omega_L^{(1)}t} \hat{a}_e^\dagger(z) \hat{a}_s(z) + g_M \hat{a}_M(k) e^{ikz} u_k^M(z) \hat{a}_e^\dagger(z) \hat{a}_g(z) + h.c. \right]$$



Large dispersion without absorption

EIT propagation range:

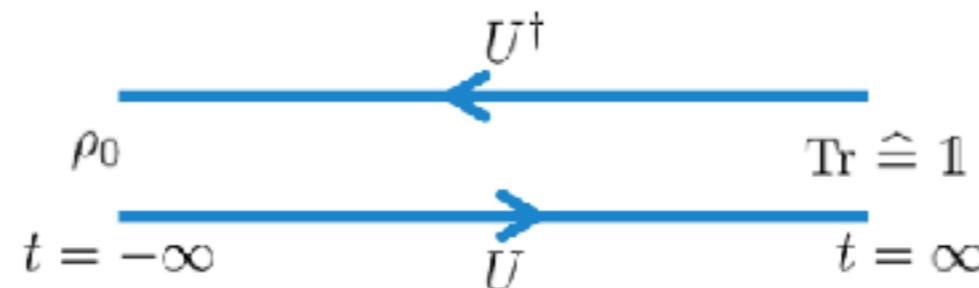
Dyson equations for Green's function:

A diagrammatic equation for the Dyson equation for the Green's function. It shows a bare propagator ω_k^M (represented by a wavy line) equal to a bare propagator ω_k^M plus a loop correction. The loop correction consists of two parts: one where the loop is closed by a $|e\rangle$ vertex and another where it is closed by a $|s\rangle$ vertex. A blue arrow points from this equation to the formula below.

$$\chi(\omega, k) = \frac{g_M^2 |u_k^M(0)|^2}{\omega - \omega_e - \frac{\Omega^2}{\omega - \omega_s - \omega_L^{(1)} + i\gamma_e/2}}$$

destructive interference = EIT

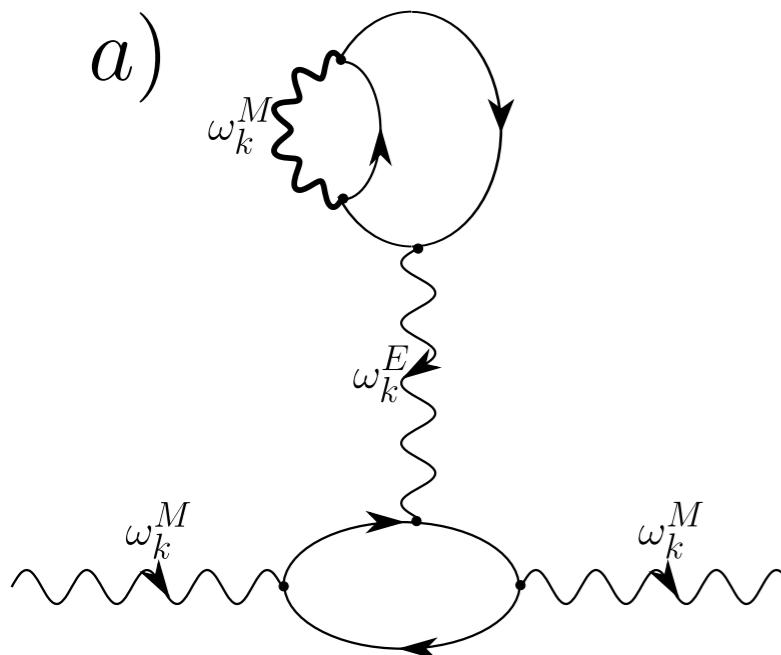
**Keldysh
real-time contour:**



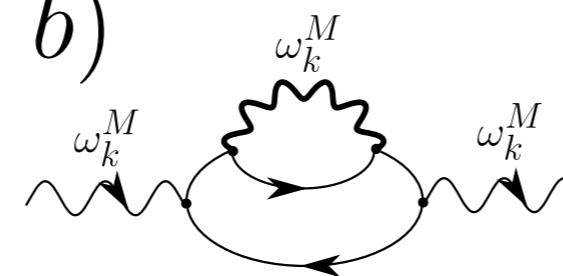
Forward-backward branch:
double # of GF's

Controlled 1/range expansion for the interaction diagrams

"Alien"



"M-king"

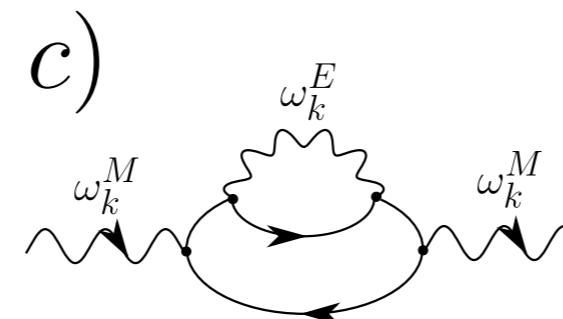


- Expansion for large propagation range

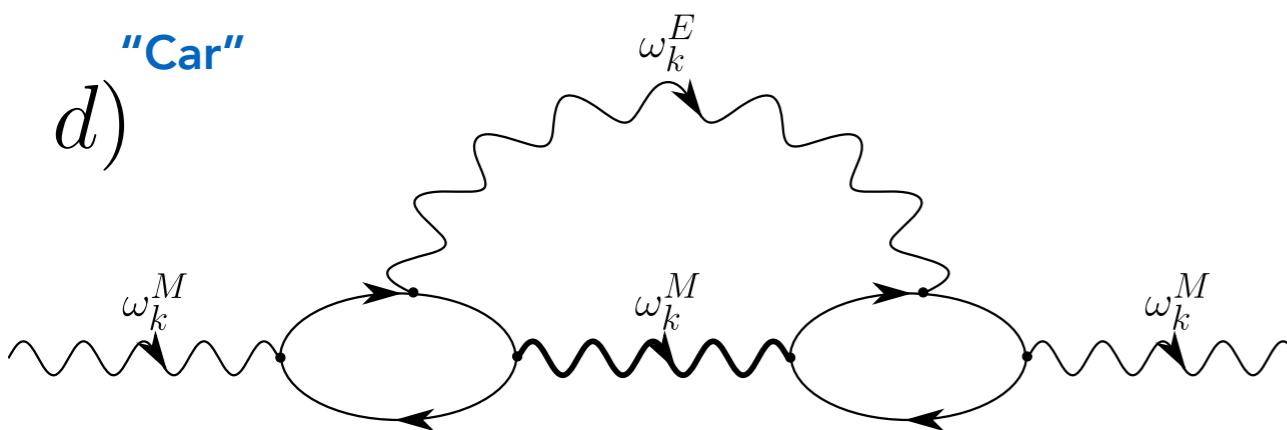
Effective coupling scales like $1/\sqrt{L}$:
multiple scattering suppressed

$$\left(\frac{a}{L}\right)^{\# \text{loops} - \# \text{atom-loops}}$$

"E-king"



"Car"



- EIT is leading order $\mathbf{O}(1)$:

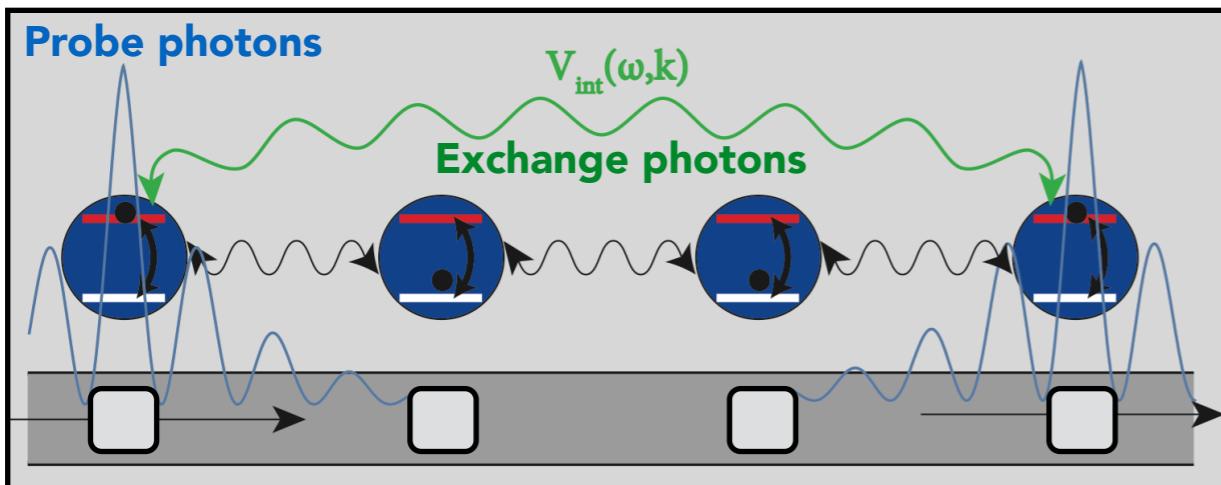


All next-to-leading order $\mathbf{O}(1/L)$ diagrams

for probe photons

(Keldysh structure and atom internal structure not shown)

Parameter regimes for selected interaction processes



Tuneability:

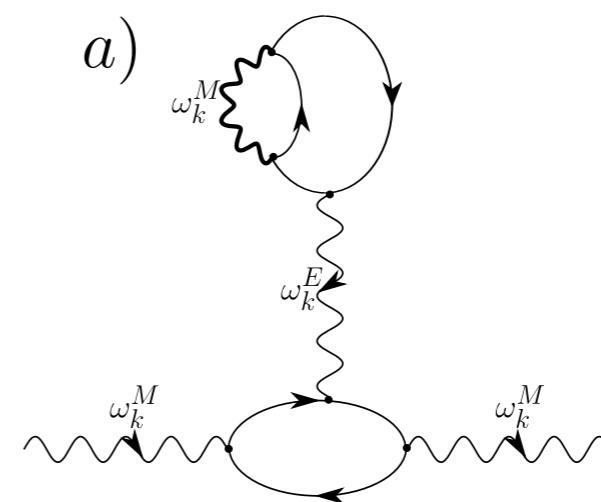
$$\text{EIT propagation range: } L_M \simeq \frac{v_M}{\kappa_M}$$

$$\text{Exponential interaction range: } L_E \simeq \sqrt{\frac{\alpha_E}{\Delta}}$$

Still very complex theory at order a/L :

- Three-loop (leading interaction contribution)
- 4 atom + 2 photon degrees of freedom
- Keldysh structure

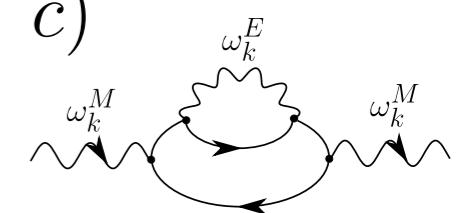
"Alien"



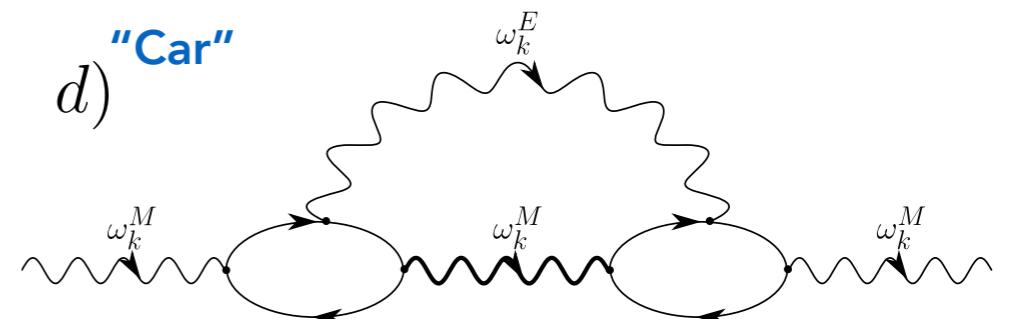
"M-king"



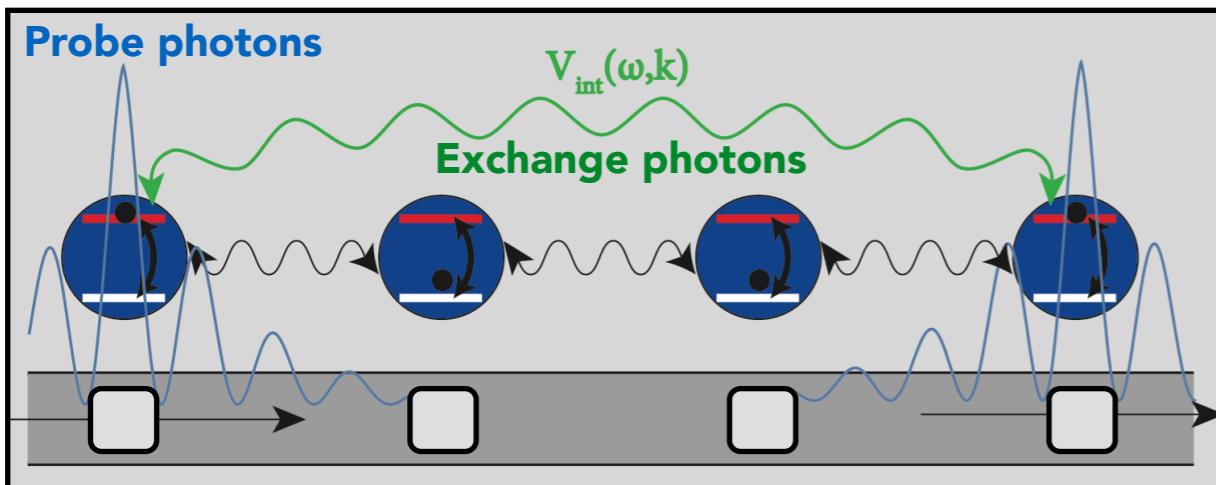
"E-king"



"Car"



Parameter regimes for selected interaction processes



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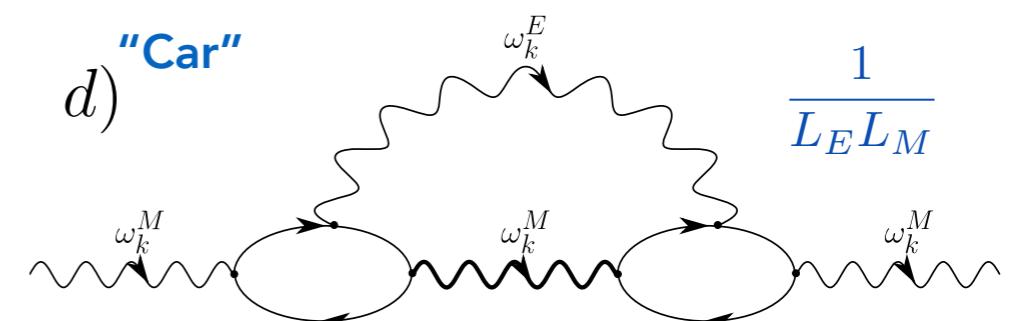
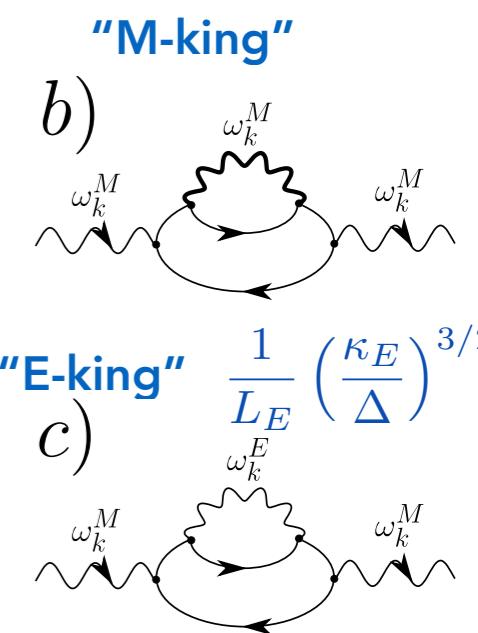
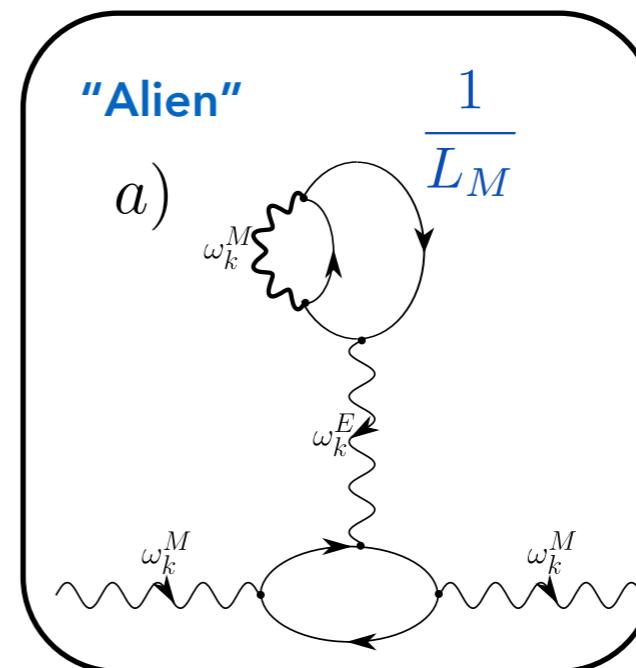
$$\text{Large interaction range: } \frac{L_E}{L_M} \gg \left| \frac{\kappa_E}{\Delta} \right|^{3/2}$$

Interaction-induced shifts dominate over scattering

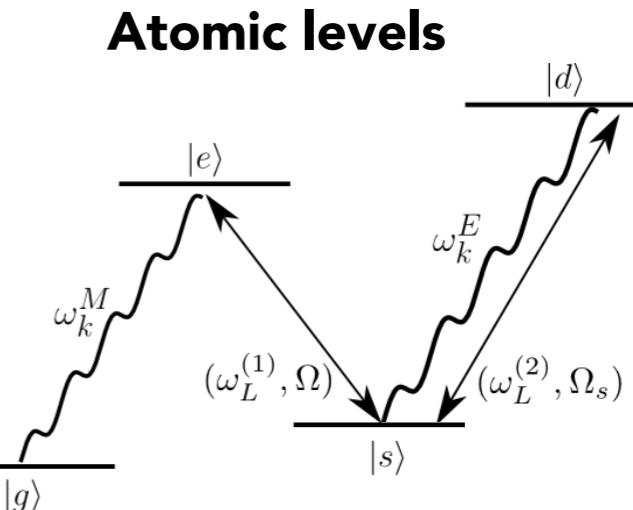
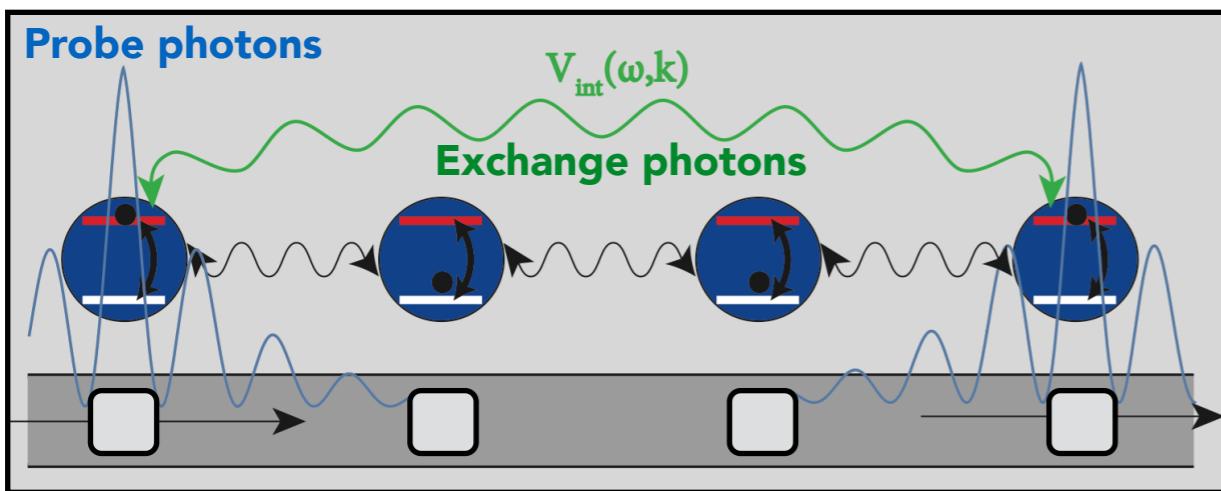
"Alien" (Hartree) diagram dominates

Still very complex theory at $1/L$:

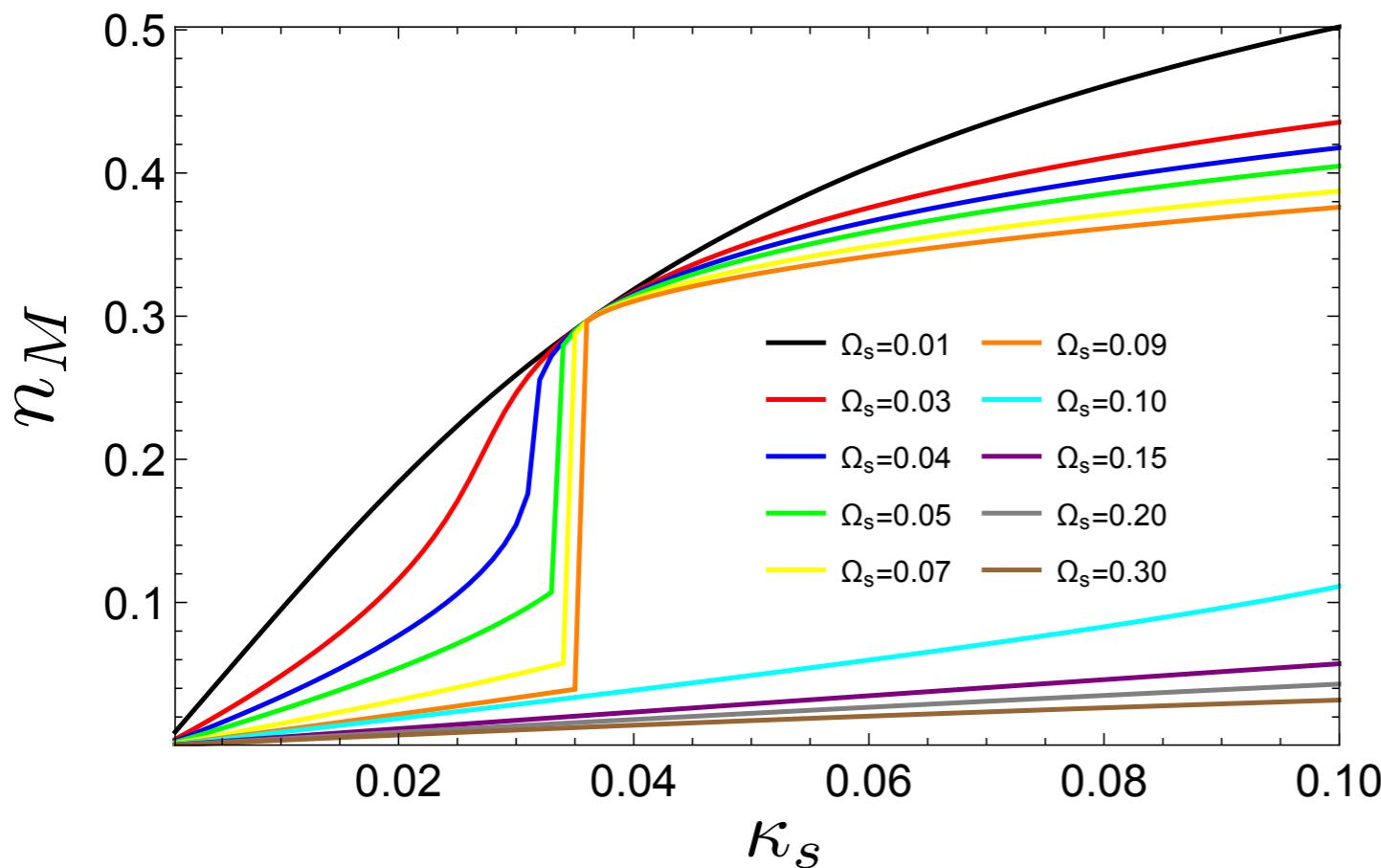
- Three-loop (leading interaction contribution)
- 4 atom + 2 photon degrees of freedom
- Keldysh structure



Non-equilibrium phase transition



Number of probe photons in the steady state

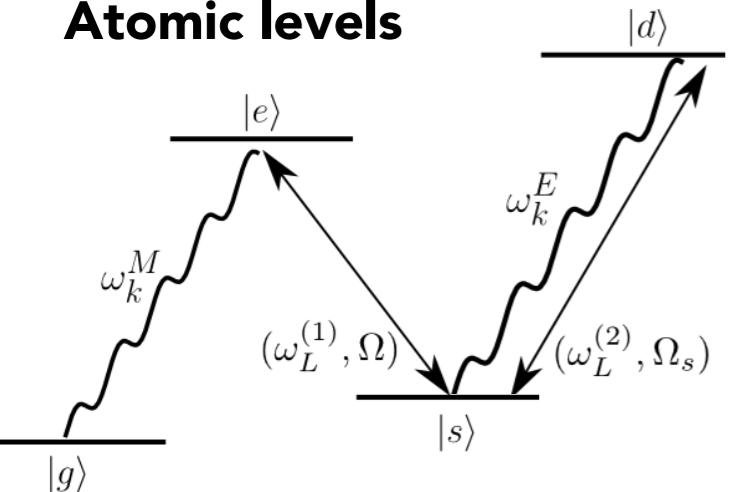


- **Transition between bright and dark phase**
First order, culminates in bi-critical point

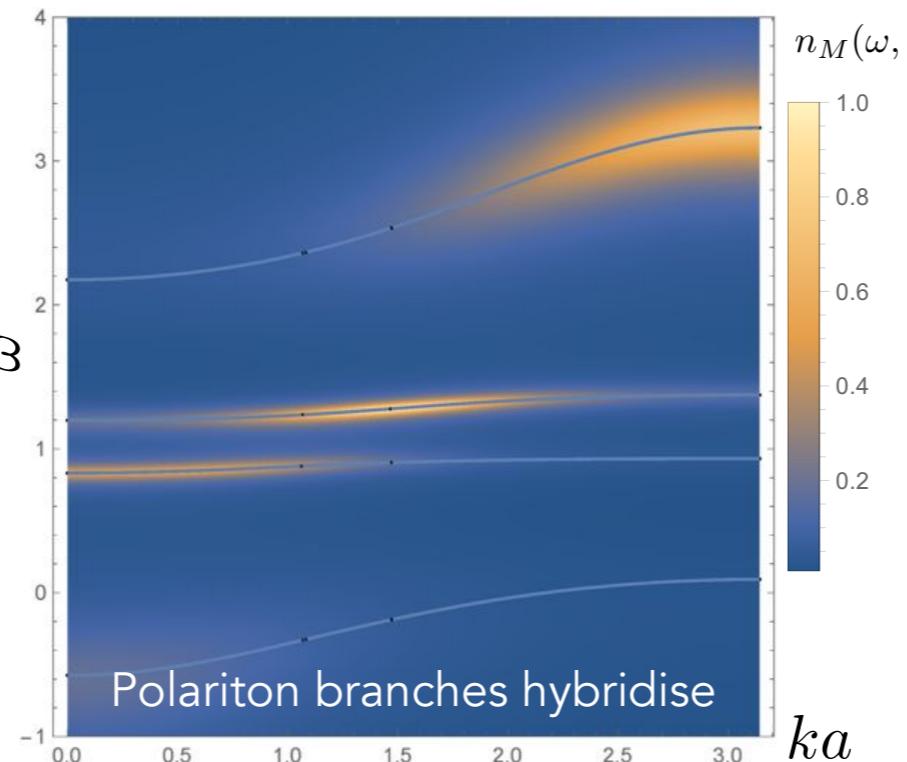
Interaction-induced transparency

Probe-photons occupation:

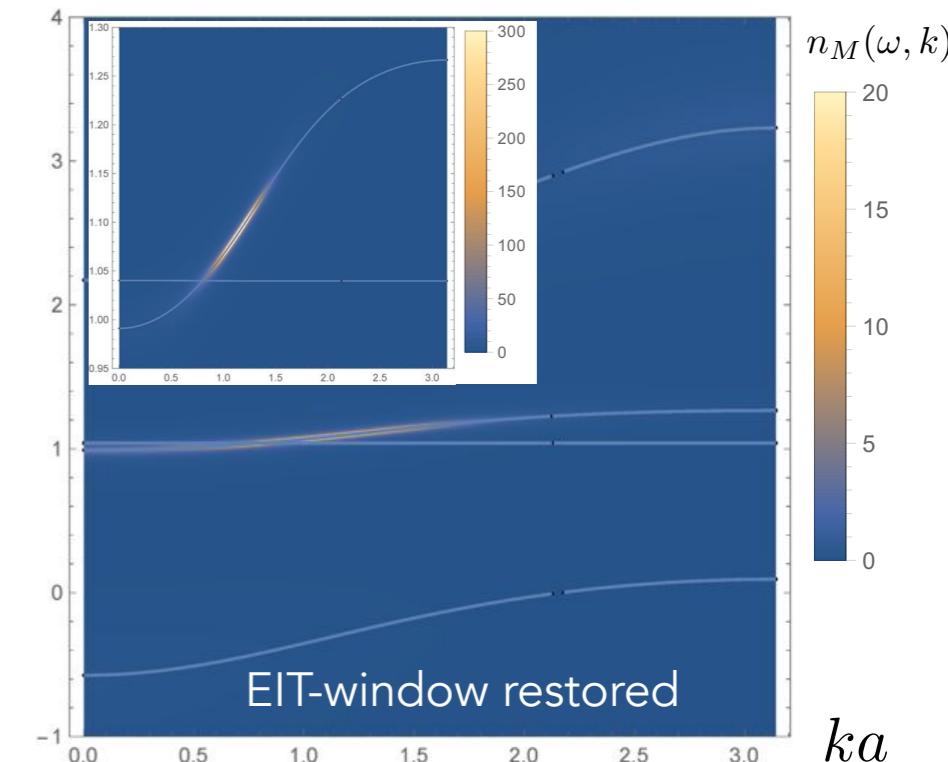
Atomic levels



Dark phase



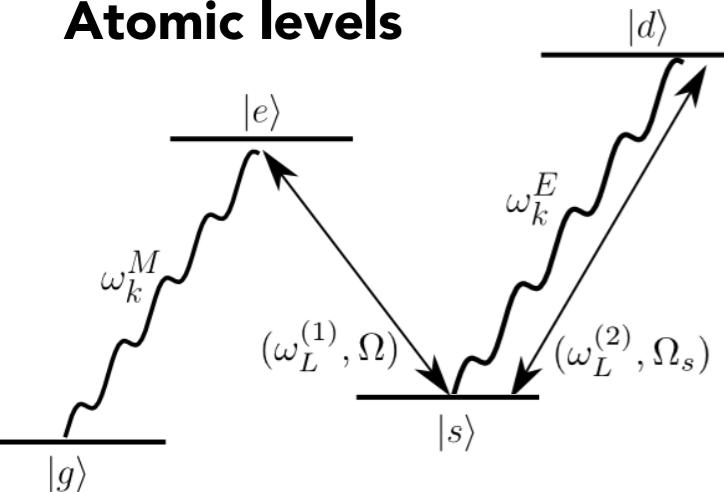
Bright phase



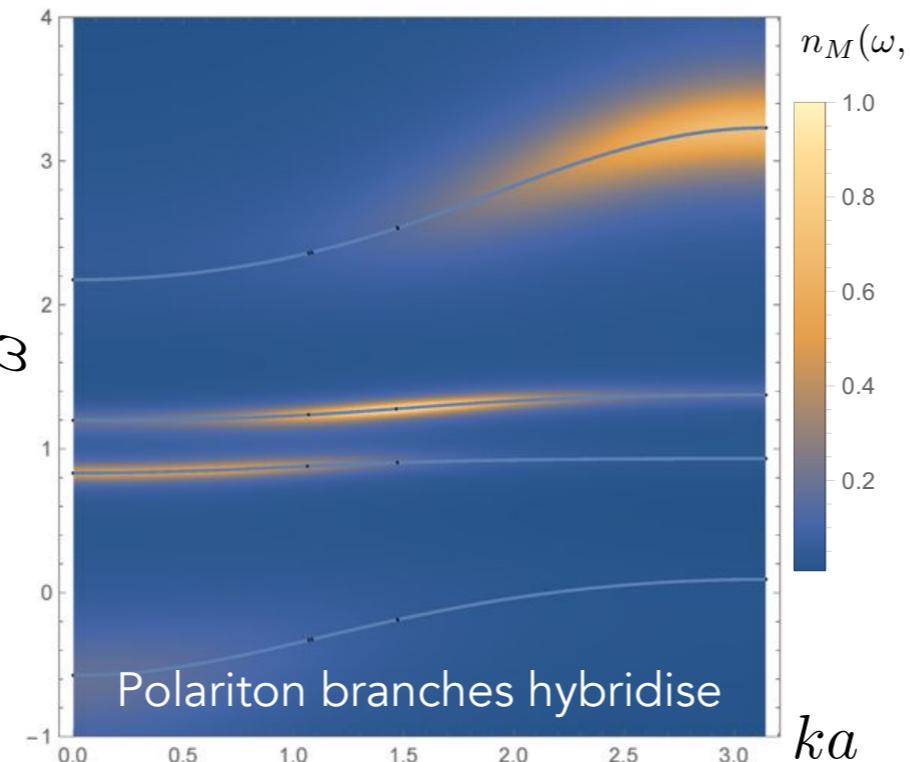
Interaction-induced transparency

Probe-photons occupation:

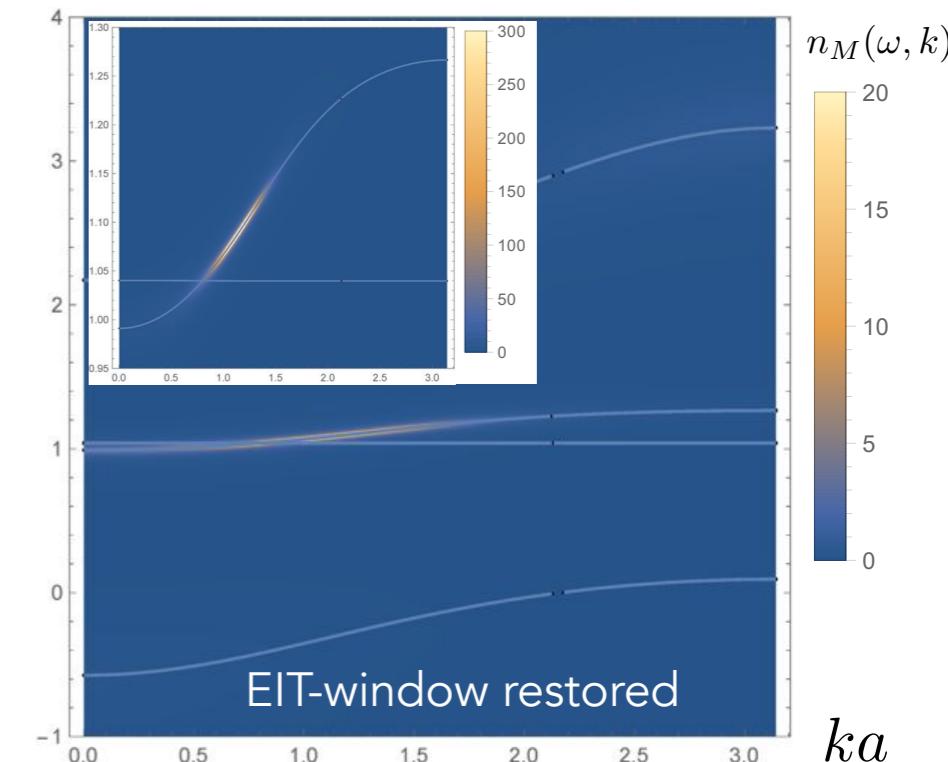
Atomic levels



Dark phase

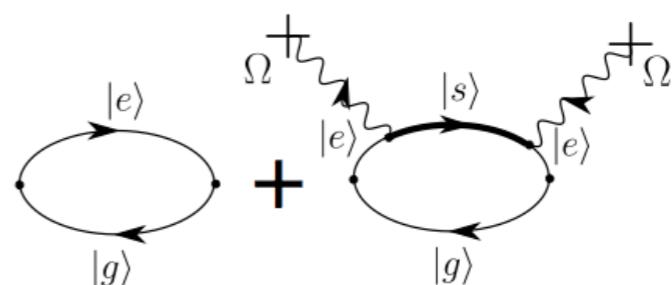


Bright phase



Electromagnetically Induced Transparency (EIT):

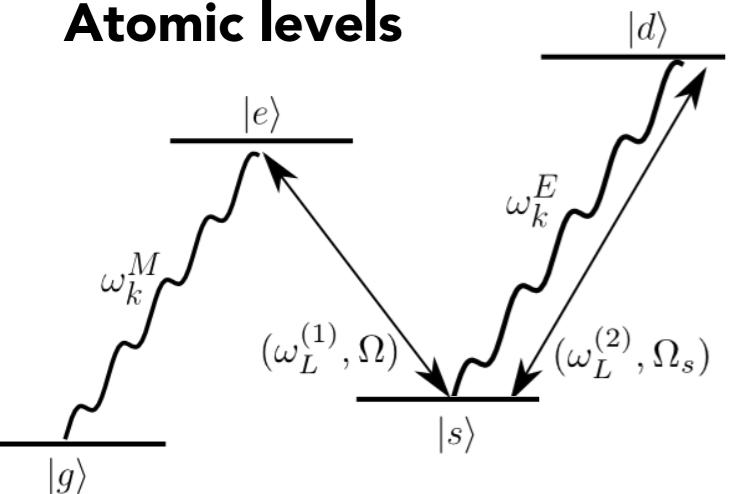
$$\chi_{\text{EIT}} =$$



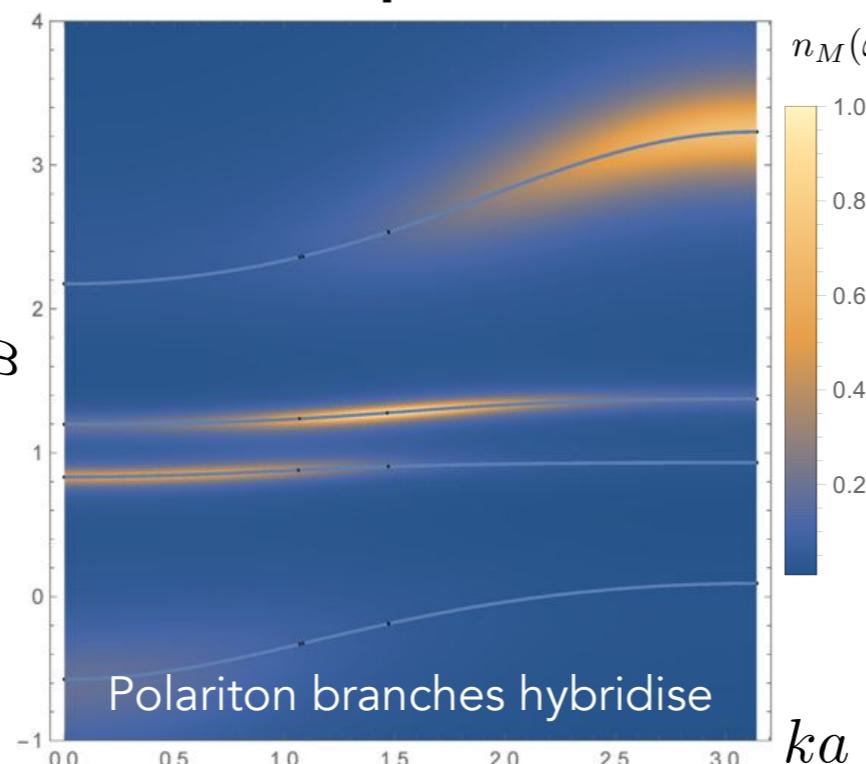
Interaction-induced transparency

Probe-photons occupation:

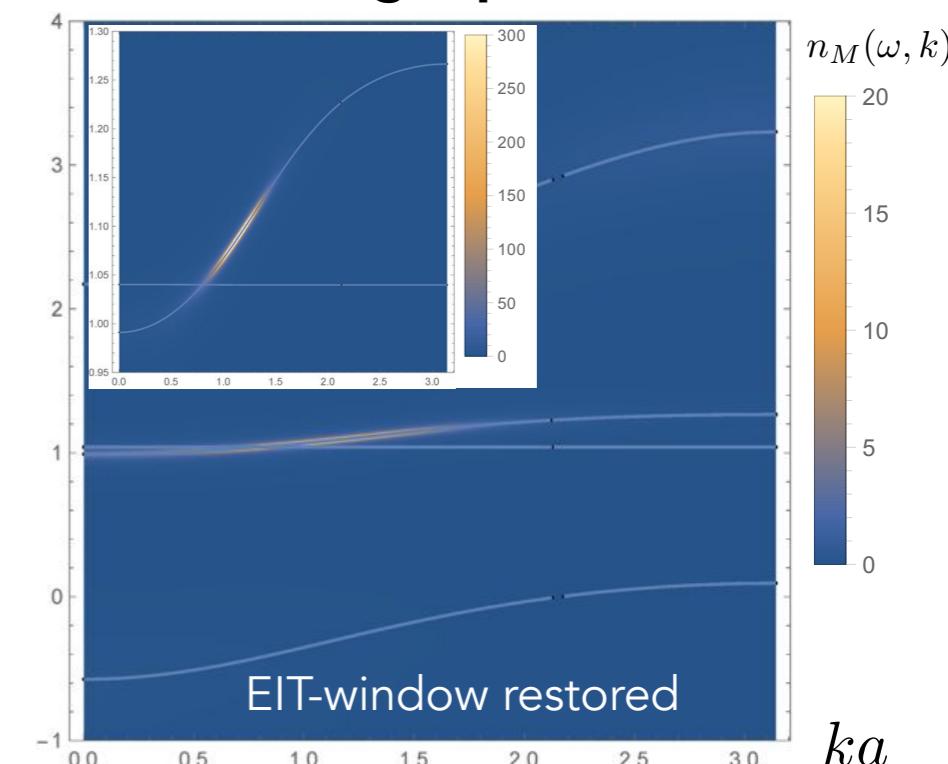
Atomic levels



Dark phase

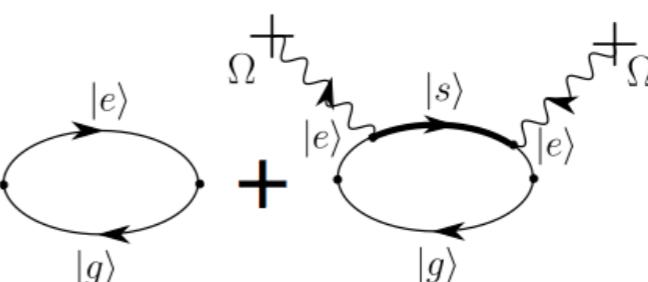


Bright phase



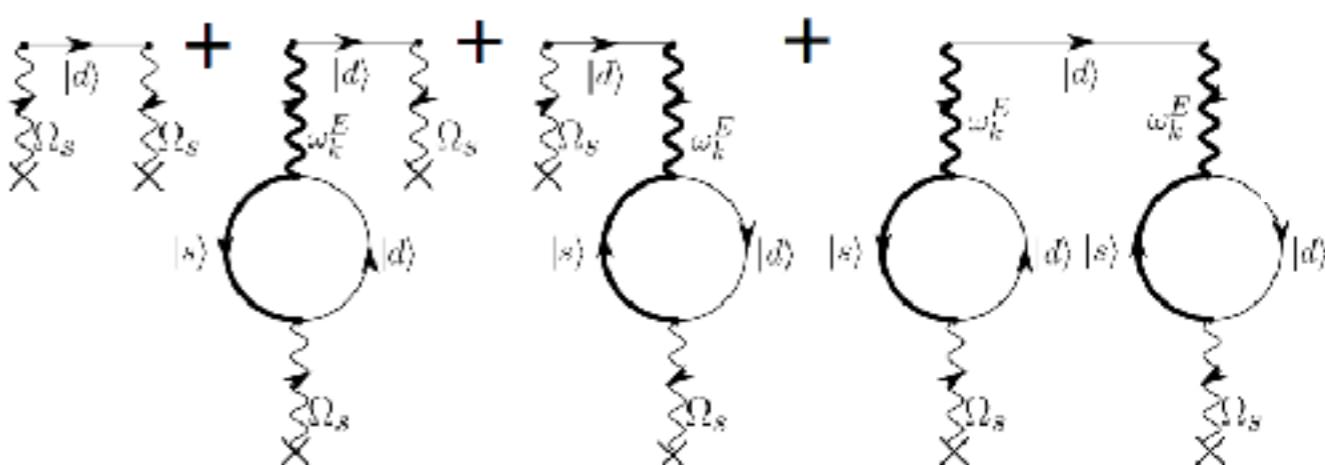
Electromagnetically Induced Transparency (EIT):

$$\chi_{\text{EIT}} =$$



Interaction-Induced Transparency (IIT):

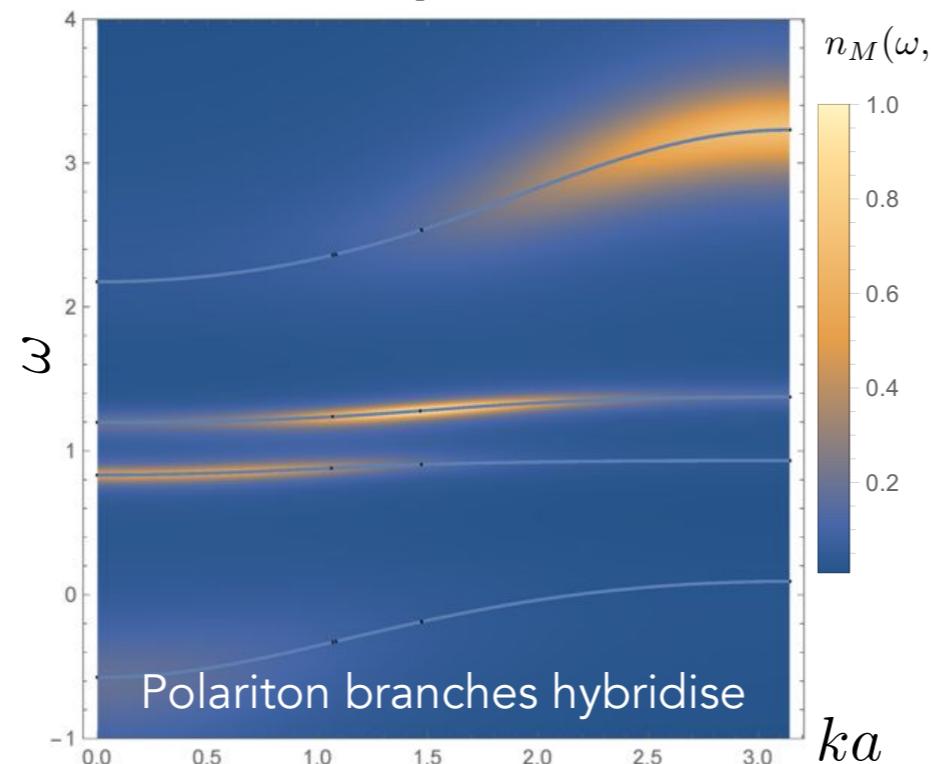
$$\chi_{\text{IIT}} =$$



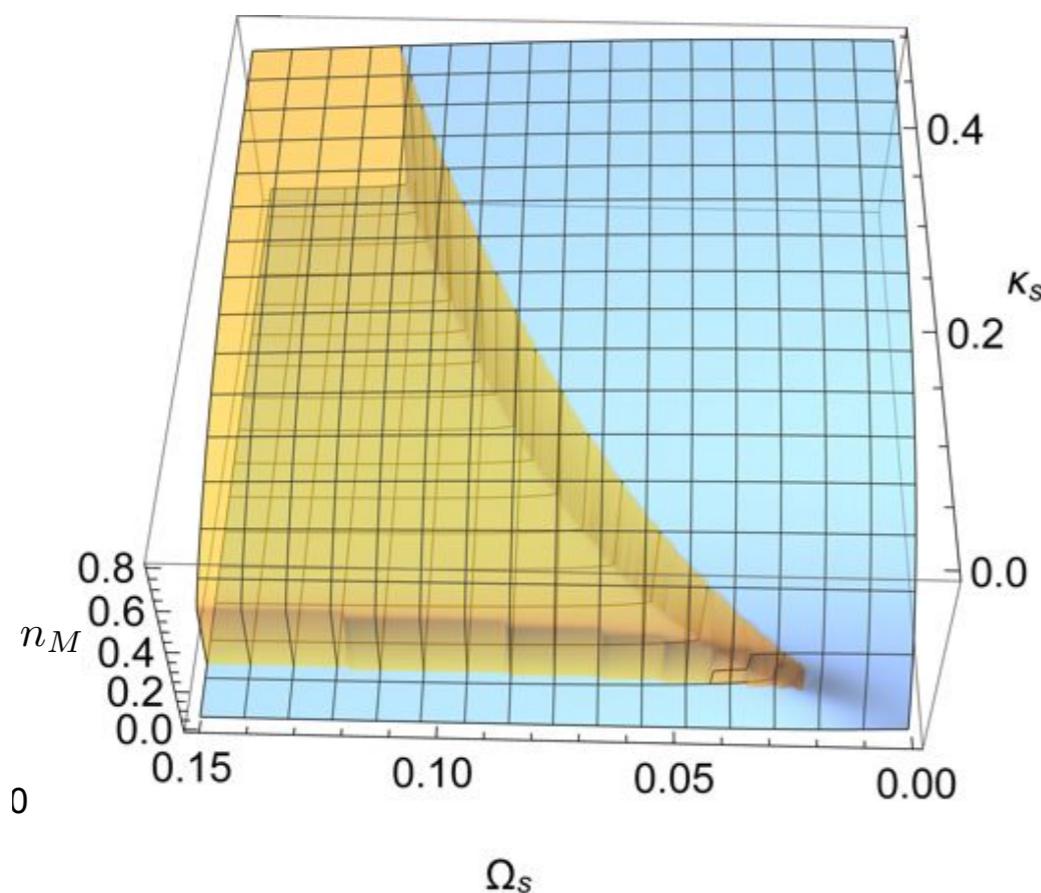
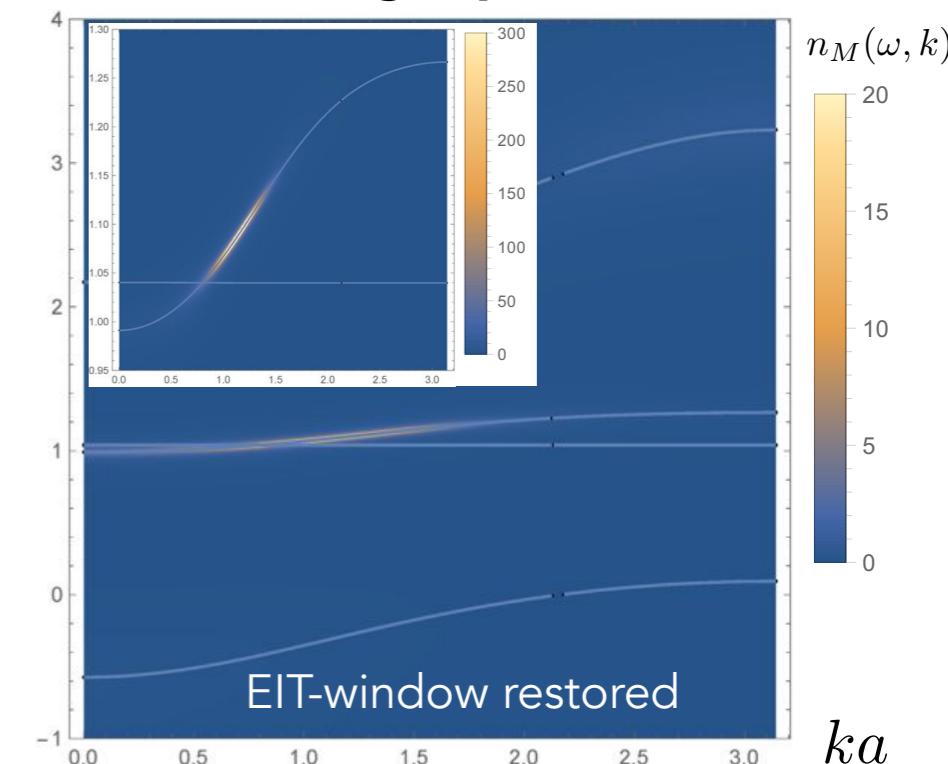
Bistability

Probe-photons occupation:

Dark phase



Bright phase



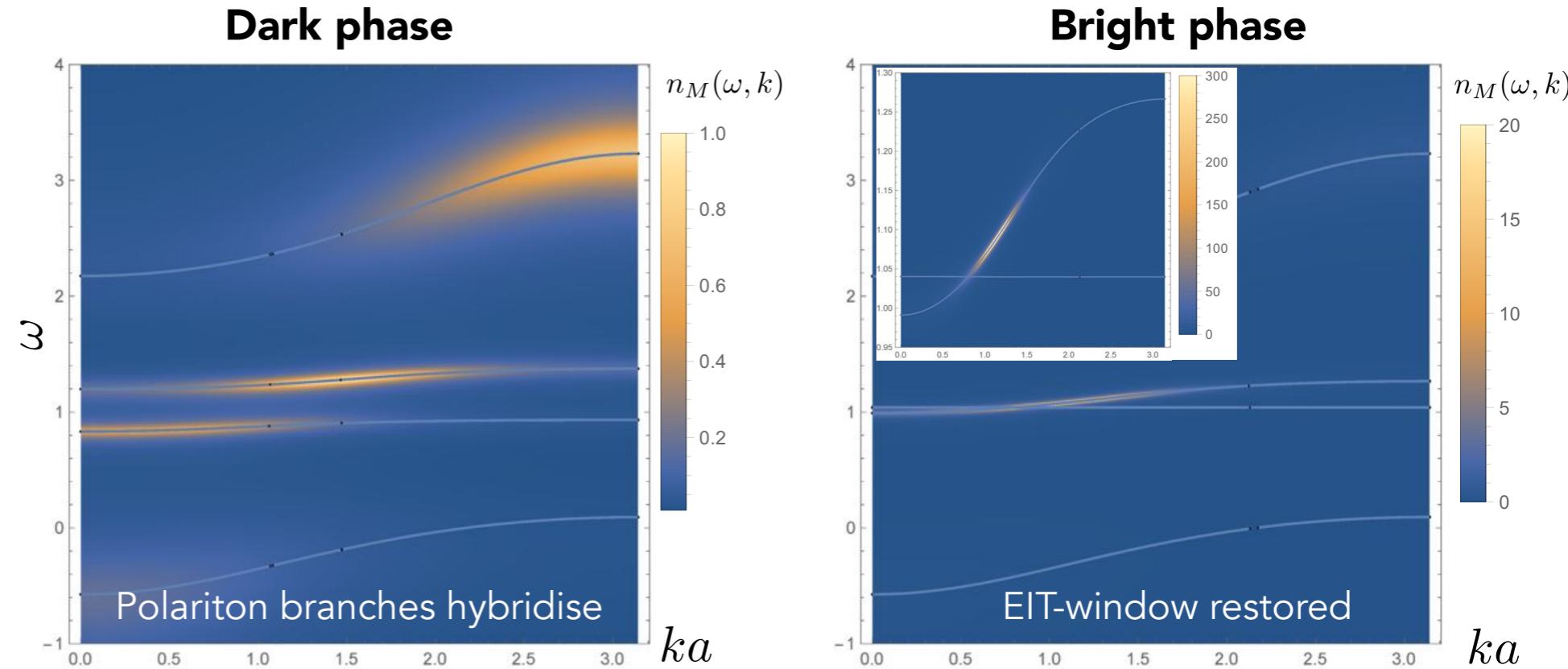
- **Transition between bright and dark phase**
First order, culminates in bi-critical point

- **Interactions can restore the EIT window**
Destructive interference between s-d excitation paths involving exchange photons

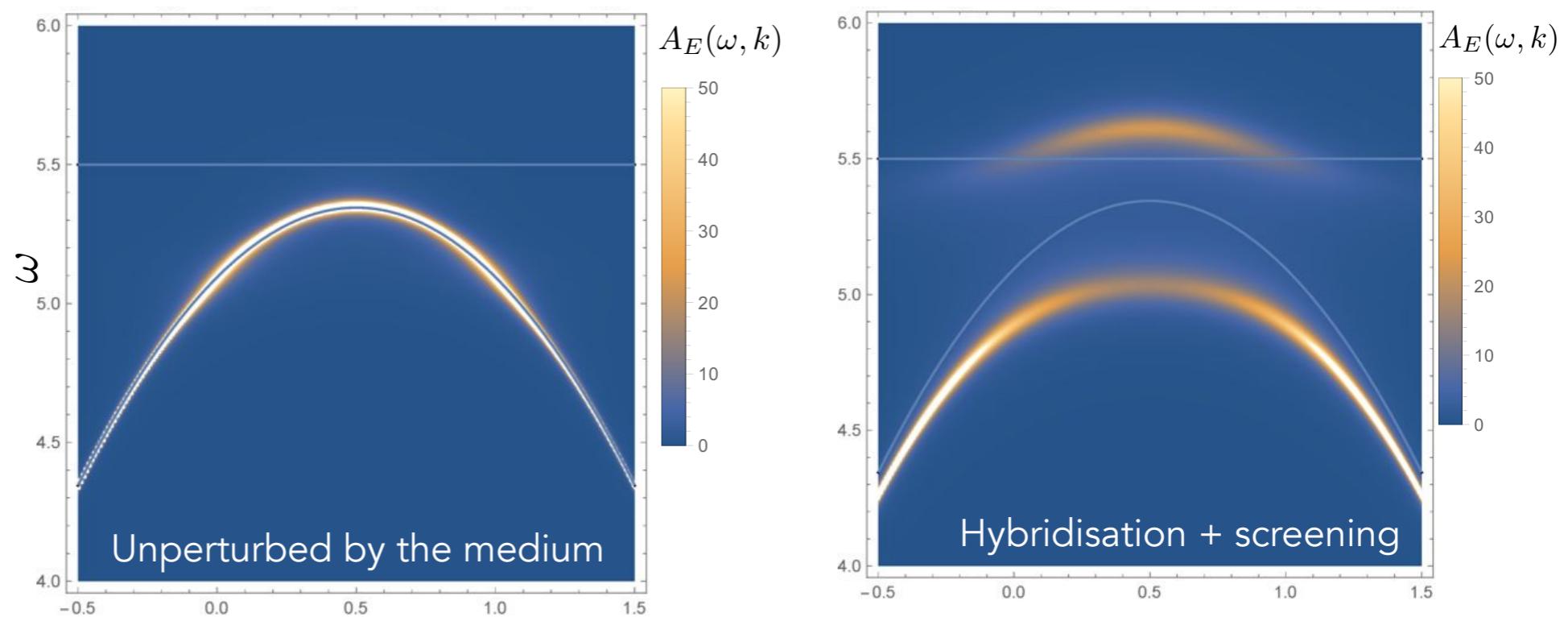
- **Bistability between dark and bright phase**
One interference path involves polariton interactions i.e. is nonlinear

Screening effects

Probe-photons occupation:

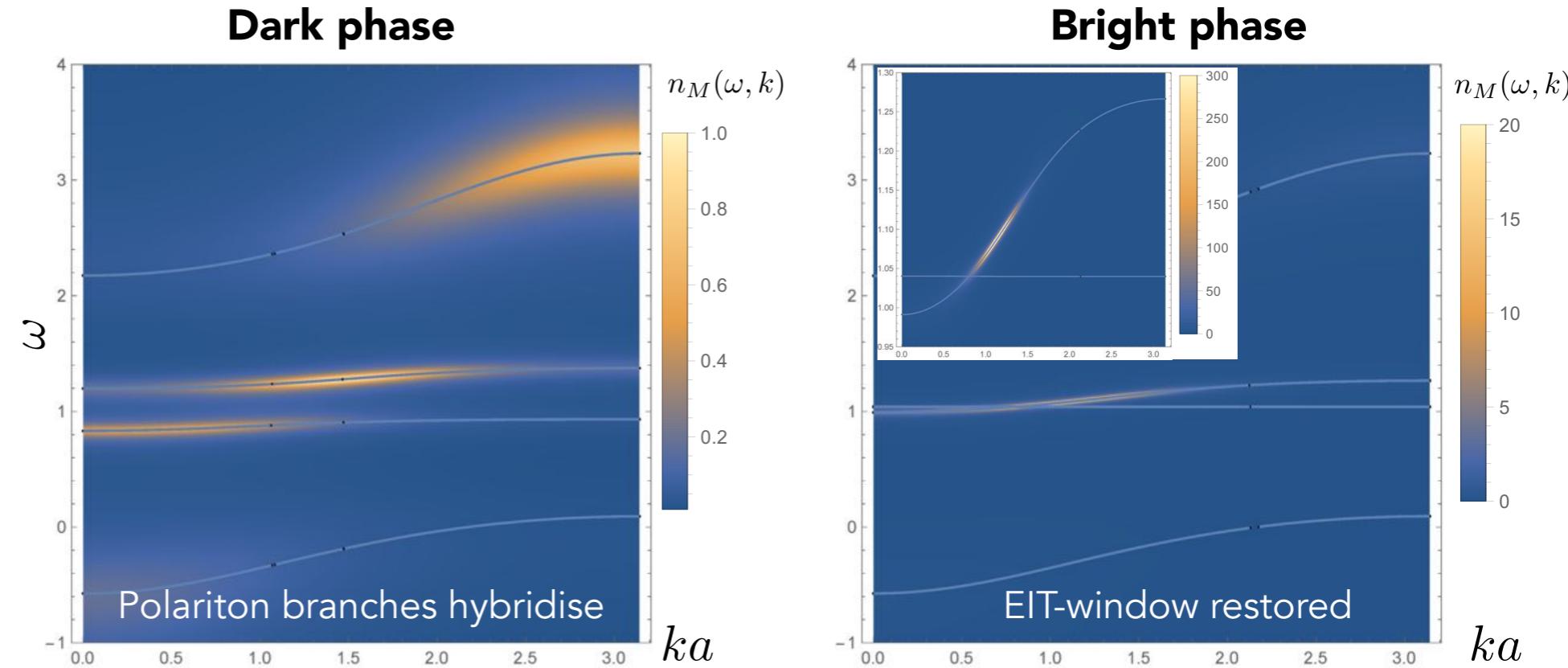


Exchange-photons spectrum:

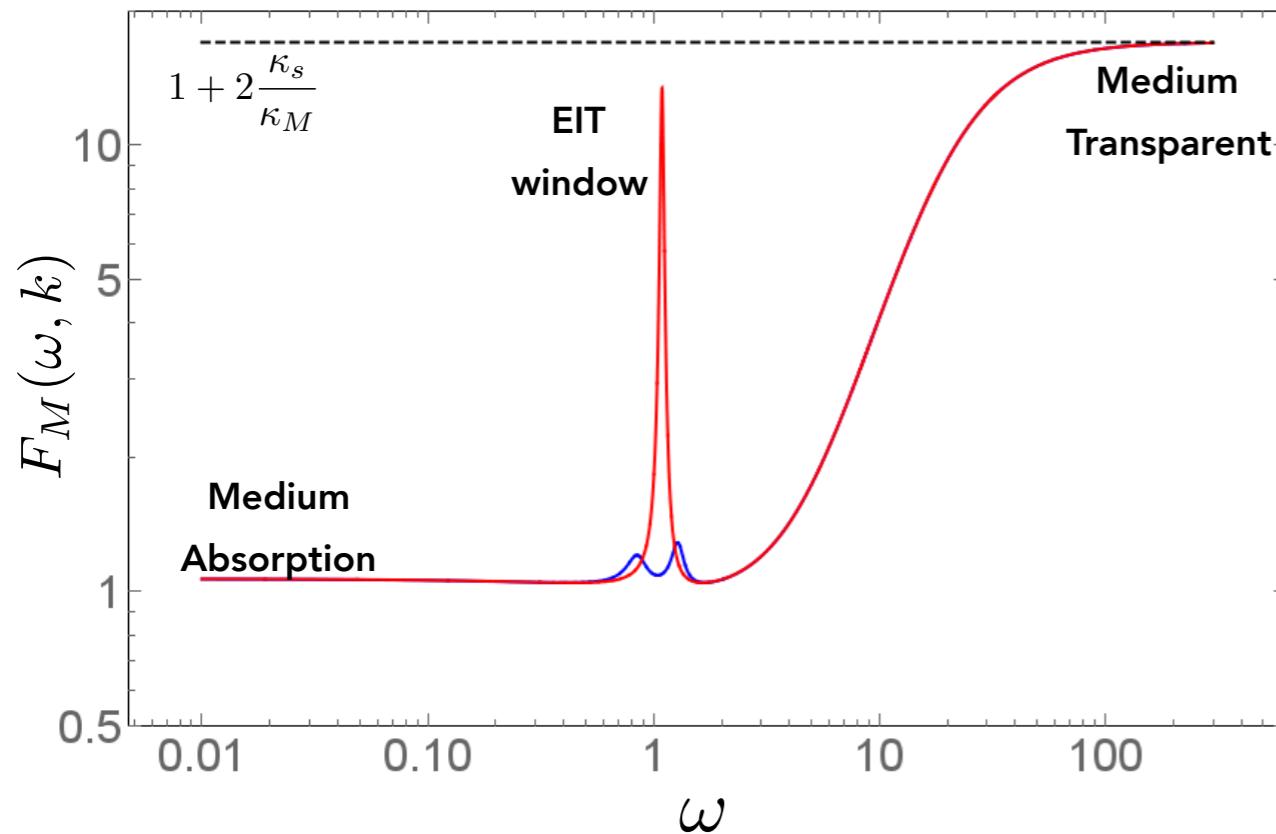


Intrinsic non-equilibrium nature

Probe-photons occupation:



Probe-photons distribution function:



No effective thermal equilibrium

- EIT window emerges between Markov regions of total transparency and total absorption
- EIT peak is a sharp Lorentzian: no effective thermalisation

Summary

1. Controlled diagrammatic approach to strongly interacting EIT polaritons

- A. Non-perturbative diagrammatic expansion in 1/range for non-equilibrium GFs
- B. Tuneability of photon dispersion allows to select interaction processes
- C. Identify parameter regimes for quantitative relevance

2. Non-equilibrium phase transition on the EIT-window

- D. Transition between a dark and bright phase where interactions restore the EIT-window
- E. Bistability culminating in bi-critical point
- F. The EIT effect makes the transition of non-equilibrium nature

Outlook

- i) Dynamics of inhomogeneously driven system: How to ideally enter the bright phase experimentally?
- ii) Universality class of the non-equilibrium transition
- iii) Include scattering diagrams and treat photon crystallization.