

# Quantum Hall Breakdown a microscopic model and a hydrodynamic analogy

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# Talk Outline

Breakdown Of the Integer Quantum Hall Effect (IQHE).

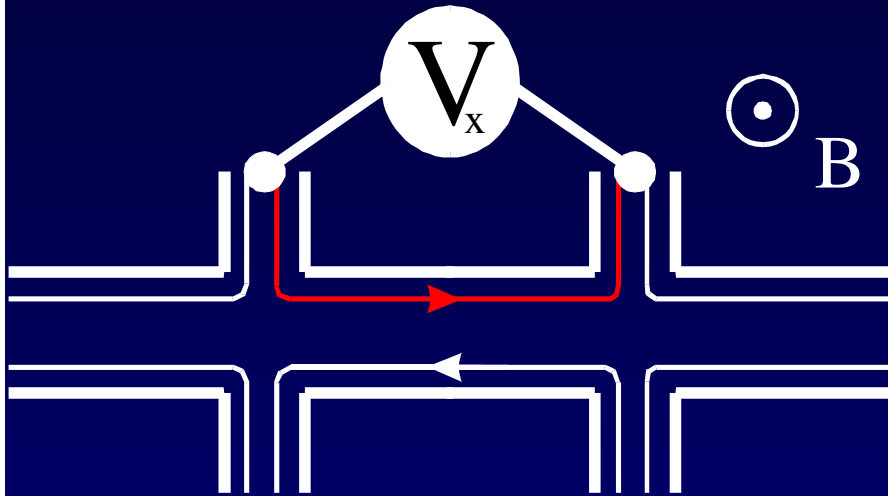
(1) Experimental observations

(2) Microscopic model of the breakdown of the IQHE

(3) Fluid analogy

(4) Future work

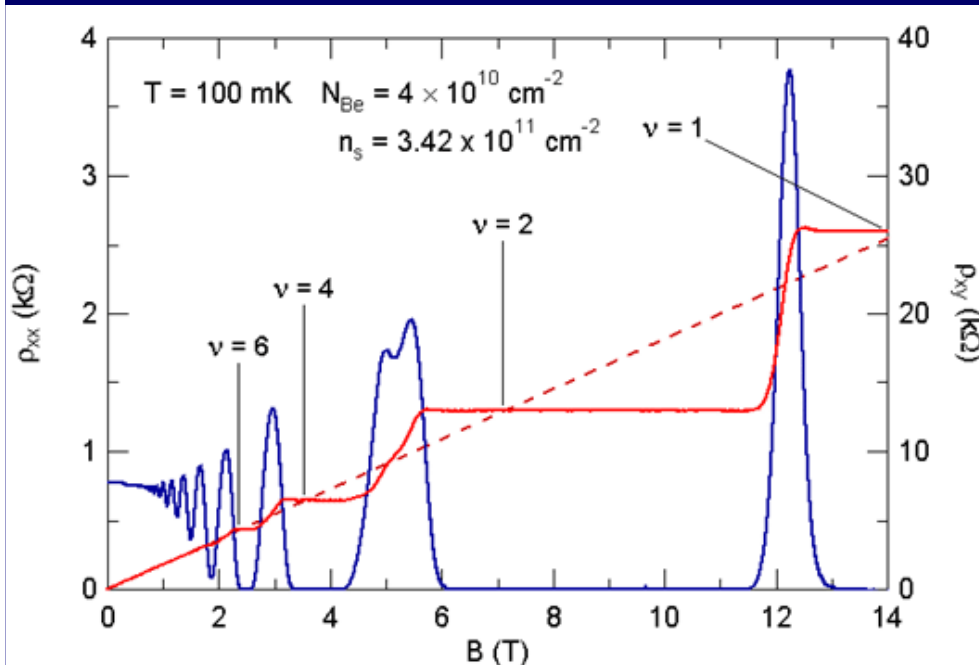
# Observations of Quantum Hall Breakdown



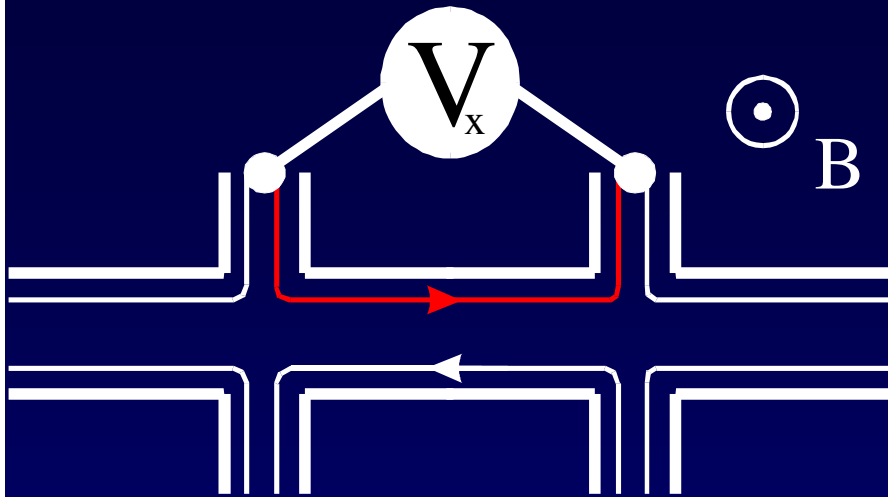
IQHE - almost dissipationless flow of current along sample

-  $V_x \cong 0$

- Hall resistance quantised,  $R_H = h/\nu e^2$

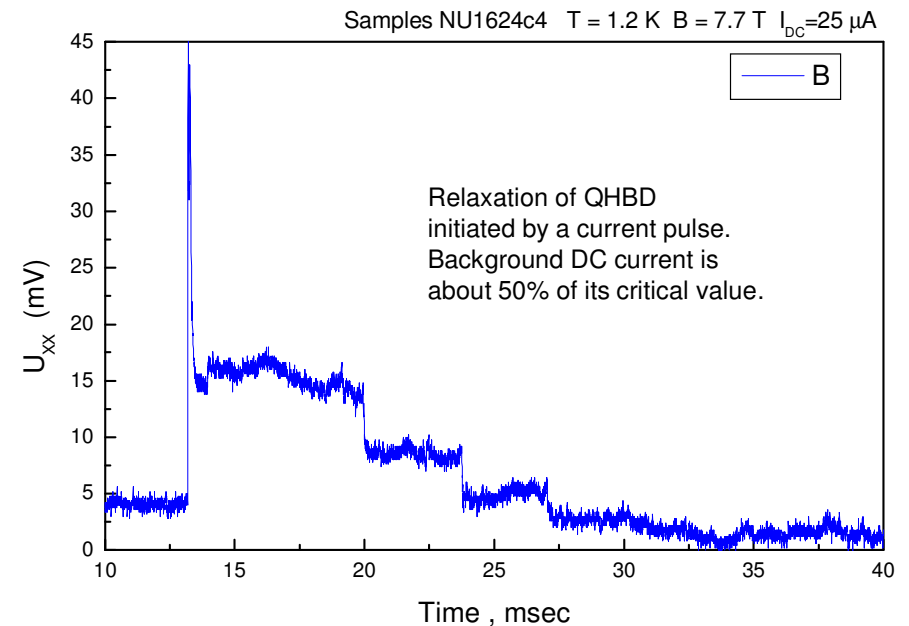
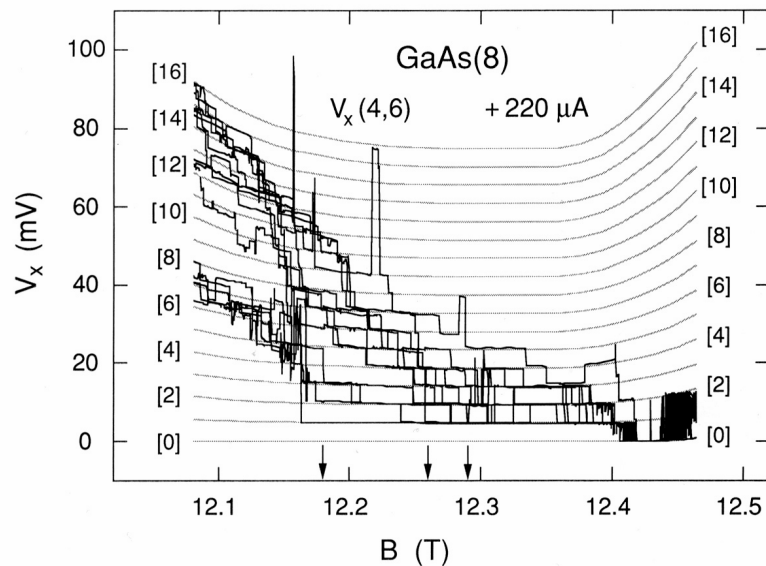


# Observations of Quantum Hall Breakdown

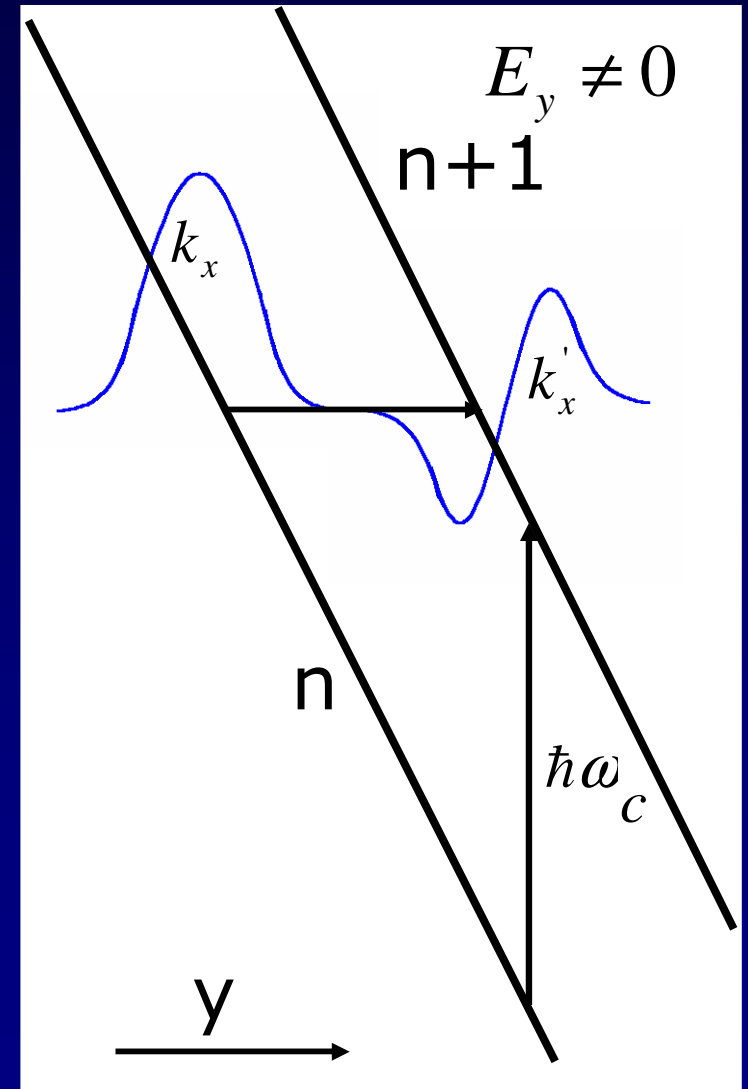
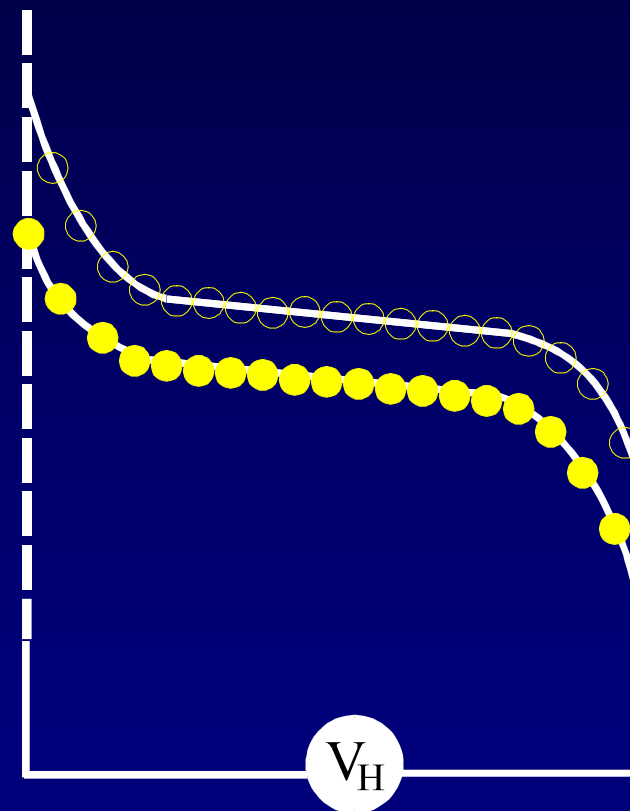
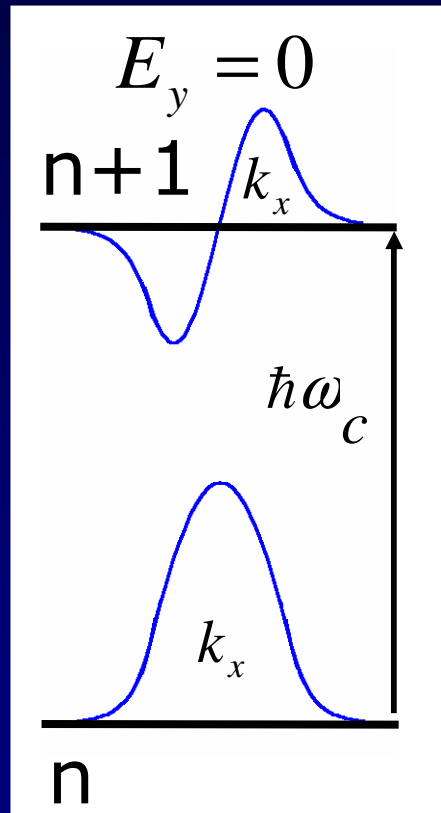


Above a critical current the flow is no longer dissipationless.

NIST Experiment;  
US Resistance Standard



# Quantum Picture of Generation of Electron-Hole Pairs



# Quantum Picture of Generation of Electron-Hole Pairs

Magneto-exciton - electron in  $(n+1)$ th Landau level bound to hole in  $n$ th Landau level;

Momentum,  $\mathbf{q}$ , is good quantum number;

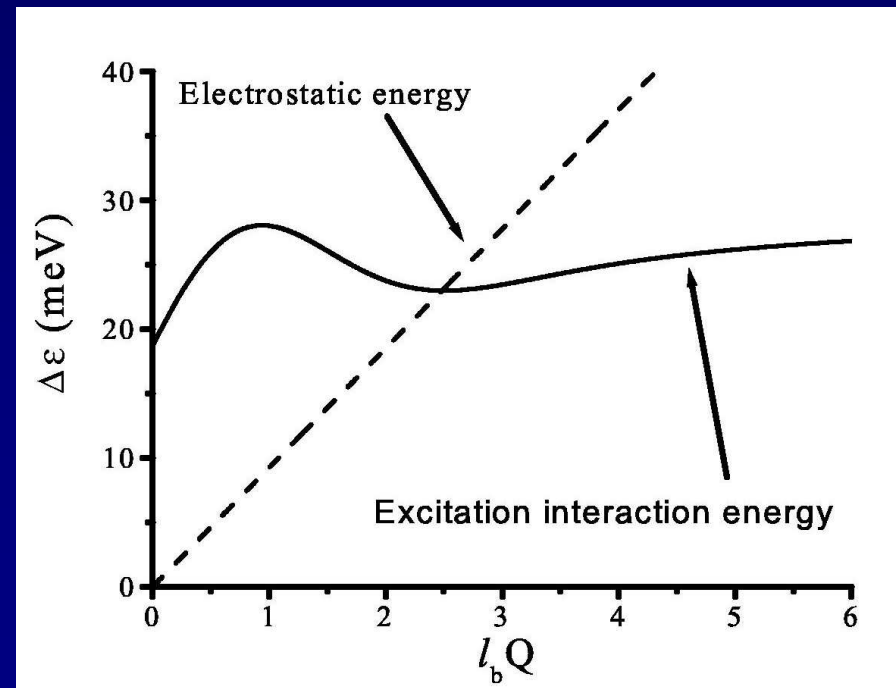
Electric Dipole  $e l_c^2 \mathbf{q}$ ;

Dispersion relation:

$$\varepsilon(\mathbf{q}) = \hbar \omega_c + \frac{e^2}{4\pi\epsilon_0 \kappa l_c} \Delta_{n,(n+1)}(\mathbf{q})$$

$$\omega_c = \frac{eB}{m^*} \quad l_c = \sqrt{\frac{\hbar}{eB}} = \sqrt{\frac{\hbar}{m^* \omega_c}}$$

Lerner and Lozovik (1980), Kallin and Halperin (1984), MacDonald (1985).



# Quantum Picture of Generation of Electron-Hole Pairs

Spontaneous Magnetoexciton creation rate due to an ionised impurity- Golden rule

$$W = \frac{2\pi}{\hbar} \sum_{\mathbf{q}} \left| \langle \mathbf{q} | V_{imp}(r) | 0 \rangle \right|^2 \delta \left( \varepsilon(q) - eE_y l_c^2 q_x \right)$$

$|\mathbf{q}\rangle$  is a 1-magneto-exciton state

$V_{imp}(r)$  impurity potential, averaged over vertical sub-band wavefunction.

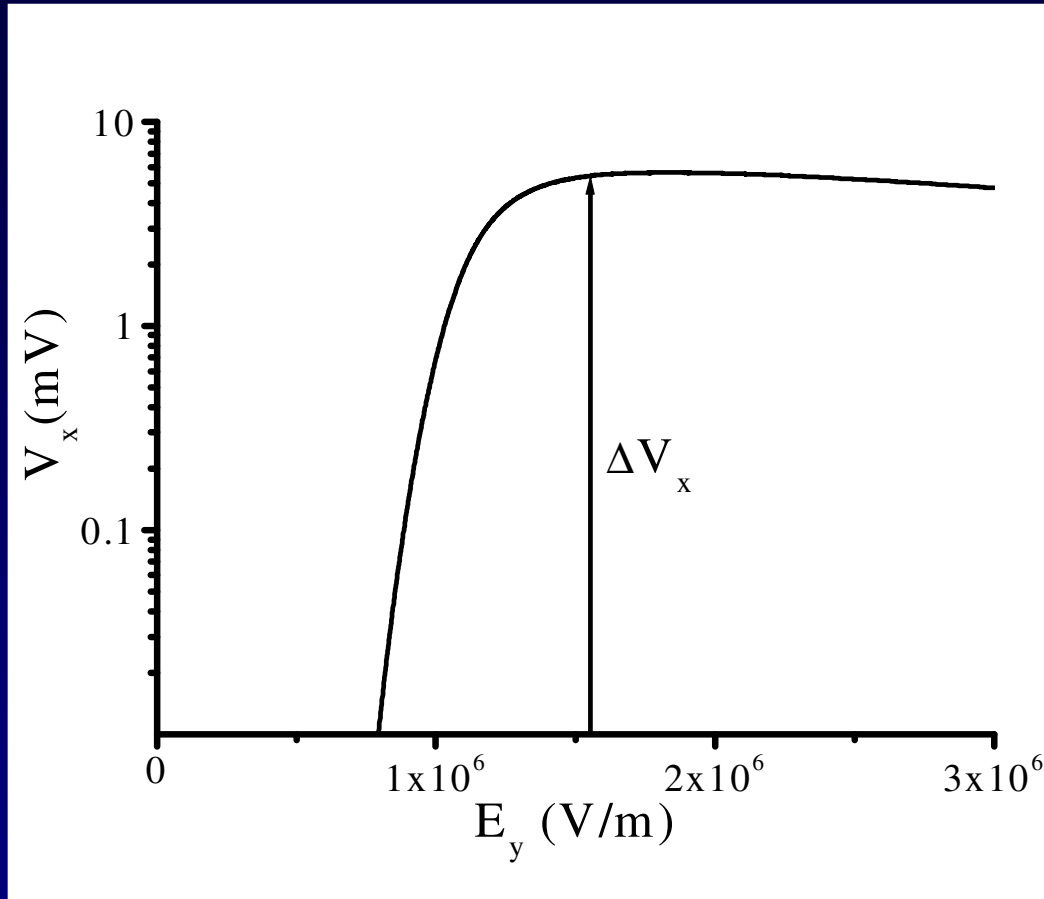
# Quantum Picture of Generation of Electron-Hole Pairs

Close to a suitably located impurity, local electric field is enhanced;  
impurity potential enables spontaneous creation of magneto-excitons;

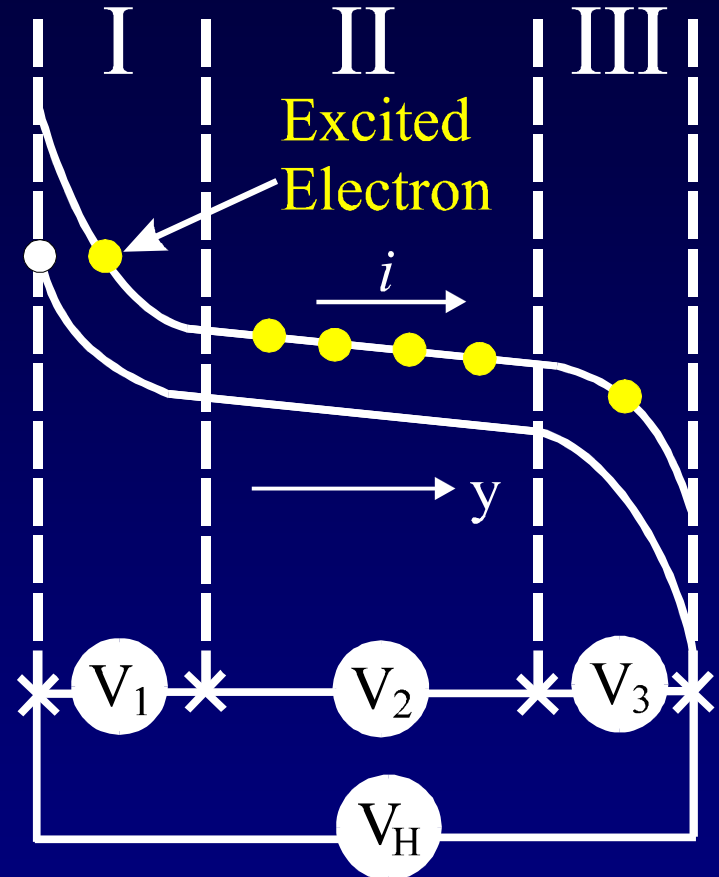
Rate of generation,  $W$ , leads to a net flow of charge across device,  $i=eW$ , and hence a dissipative voltage,  $V_x=hW/ve$ .



# Final Quantum Result



$$\nu = 2 \quad V_x = hW / 2e$$

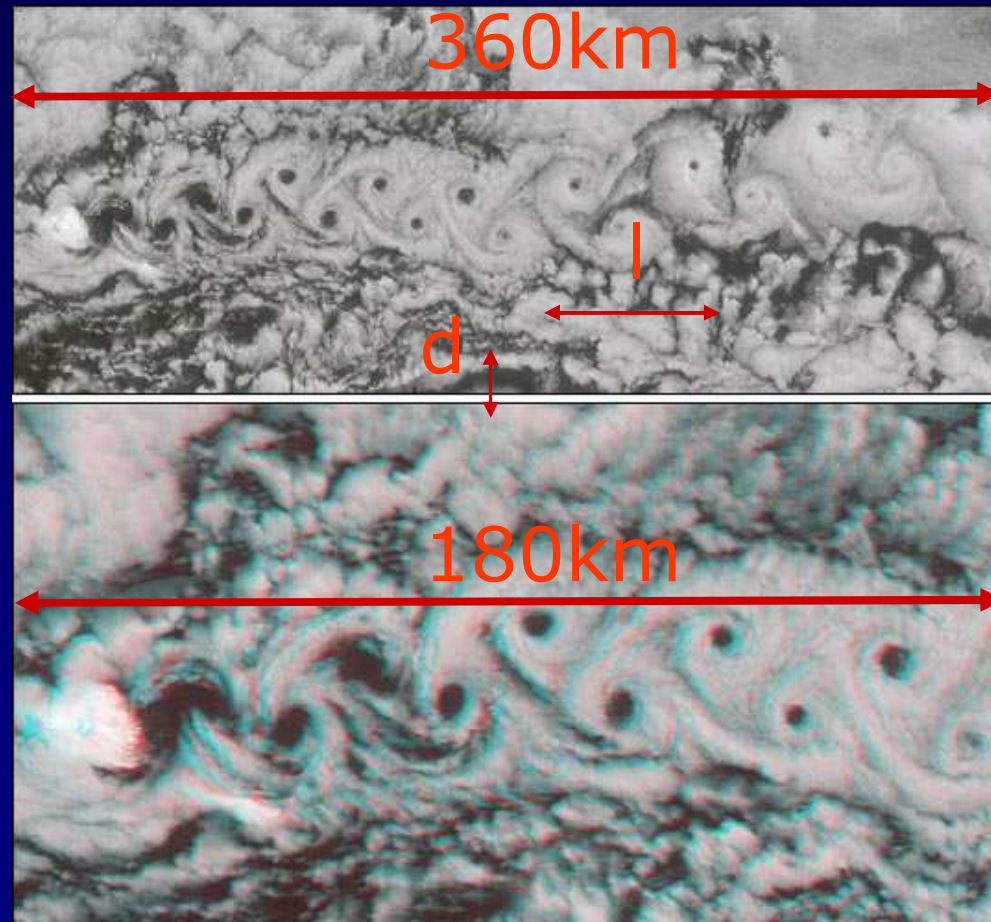


$$\Delta V_x = 5.6 \text{ mV}$$

Also see agreement between theory and experiment for step heights in hole gas devices with  $\Delta V_x \approx 1 \text{ mV}$ .

# An analogy with hydrodynamics: Von-Karman Vortex Street

$$d/l = 0.28$$



# Fluid Model of the Breakdown of the IQHE

$$m^* [\dot{\mathbf{v}} - [\mathbf{v} \times \boldsymbol{\Omega}]] = e[\mathbf{E} + \mathbf{v} \times \mathbf{B}] - \nabla \left( \frac{m^*}{2} |\mathbf{v}|^2 + \mu \right)$$

$$m^* \boldsymbol{\Omega} + e\mathbf{B} = 2\pi\rho \quad \text{M. Stone, Phys Rev. B } \mathbf{42}, 212 \text{ (1990)}$$

$$\dot{\rho} + \nabla \cdot (\rho \mathbf{v}) = 0 \quad \text{Continuity Equation}$$

$$v_x = \frac{E_y}{B} + \varepsilon_1 \cos(Qx - \omega t) \quad v_y = \varepsilon_2 \sin(Qx - \omega t) \quad \text{Excitation}$$

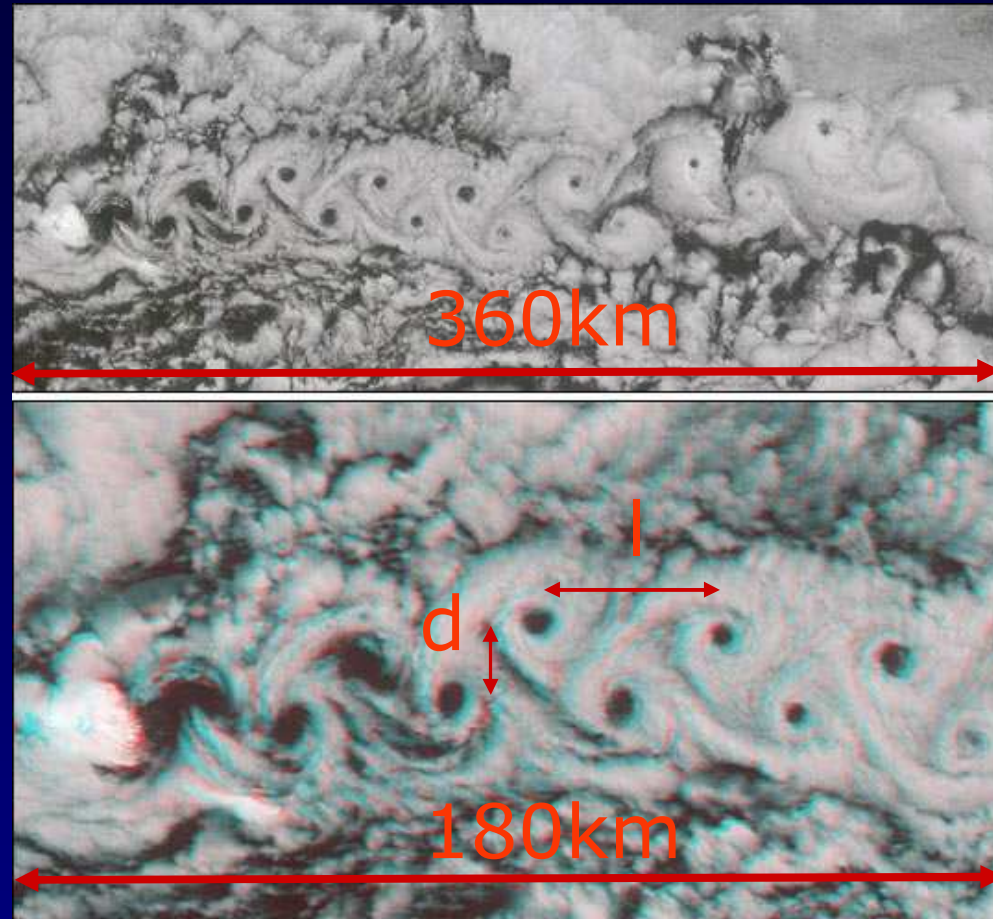
$$\hbar\omega = \hbar\omega_c + \lambda Q^2 + \dots - eE_y l_B (l_B Q)$$

Energy Dispersion for excitation

$$\Delta E = \hbar\omega_c + \frac{e^2}{4\pi\kappa l_B} \Delta_{n,(n+1)}(\mathbf{q}) - eE_y l_B^2(q_x) = 0$$

Quantum Result

# Von-Karman Vortex Street



$$d / l = 0.28$$

$$v_x = E_y / B$$

$$V_x = \frac{0.28\pi \hbar \omega_c}{e(l_B Q)^2}$$

$$l_B Q = 1.9$$

$$\hbar \omega = \Delta \varepsilon = \hbar \omega_c - e E_y l_B (l_B Q)$$

$$\Delta V_x = 4.9 \text{mV}$$

## Future work

1. More sophisticated treatment of the electrostatics within device;
2. More careful treatment of fluid analogy - Bogoliubov type analysis of Composite Bosons in a field.
3. Analysis of more complex experimental geometries
4. Low temperature scanning probe microscopy of quantum Hall devices under breakdown conditions (CJ Mellor).

# Future work

PHYSICAL REVIEW B

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## Tunneling between edge states in the quantum Hall regime limited by a mesoscopic island: A current-plateau phenomenon

Z. H. Liu, G. Nachtwei, J. Groß, R. R. Gerhardts, J. Weis, K. von Klitzing, and K. Eberl

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