

DECOHERENCE BY CONTROLLED SPIN BATHS

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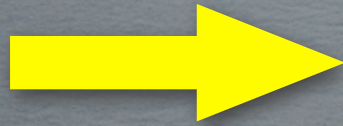


BEC - Trento

cond-mat/0605051

Motivation

Decoherence



Paradigm models

Engineered Quantum Baths



- Harmonic oscillators
- Spin baths



Spin baths

N.V. Prokof'ev and P.C.E. Stamp,
Rep. Prog. Phys. **63**, 669 (2000)

W.H. Zurek (1982)

F.M. Cucchietti, J.P. Paz, and W.H. Zurek, (2005)

L. Tessieri and J. Wilkie, (2003)

independent spins



Highly symmetric interactions and baths

The diagram shows a ring of 12 yellow spheres, each with a black arrow pointing outwards, representing independent spins. A central blue sphere with a red arrow pointing upwards represents an interacting spin. A teal horizontal bar is overlaid on the diagram, containing the text 'Highly symmetric interactions and baths'.

D.V. Khveshchenko, (2003)

C.M. Dawson *et al.*, (2005)

S. Paganelli, F. de Pasquale, and S.M. Giampaolo, (2002)

H.T. Quan *et al.* (2006).

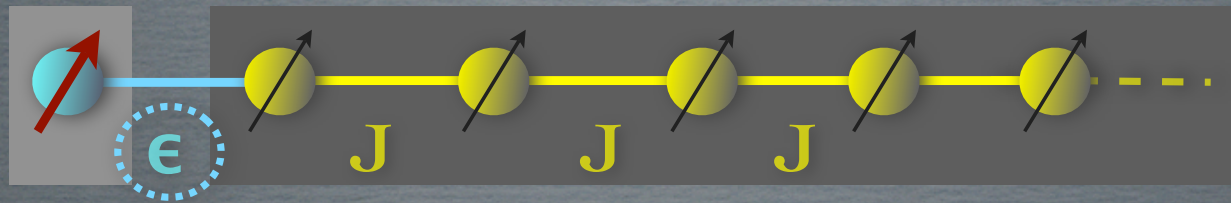
F.M. Cucchietti, S. Fernandez-Vidal, and J.P. Paz, (2006)

interacting spins

The setup

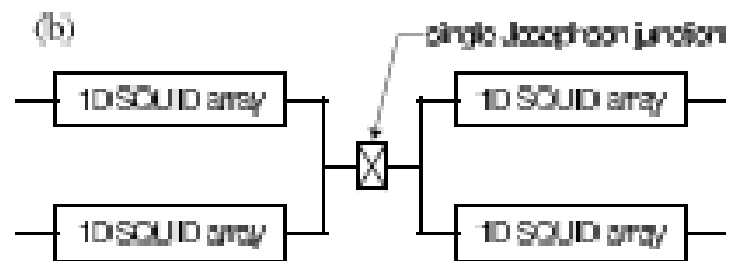
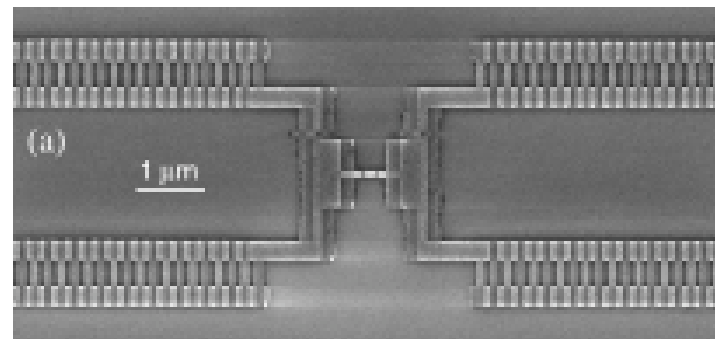
System

Bath



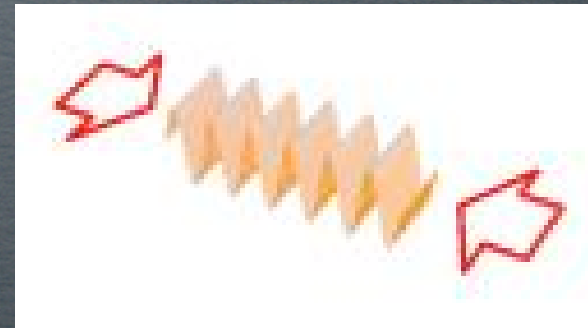
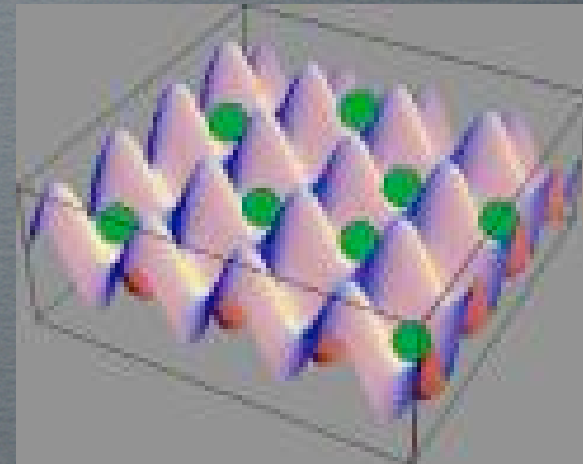
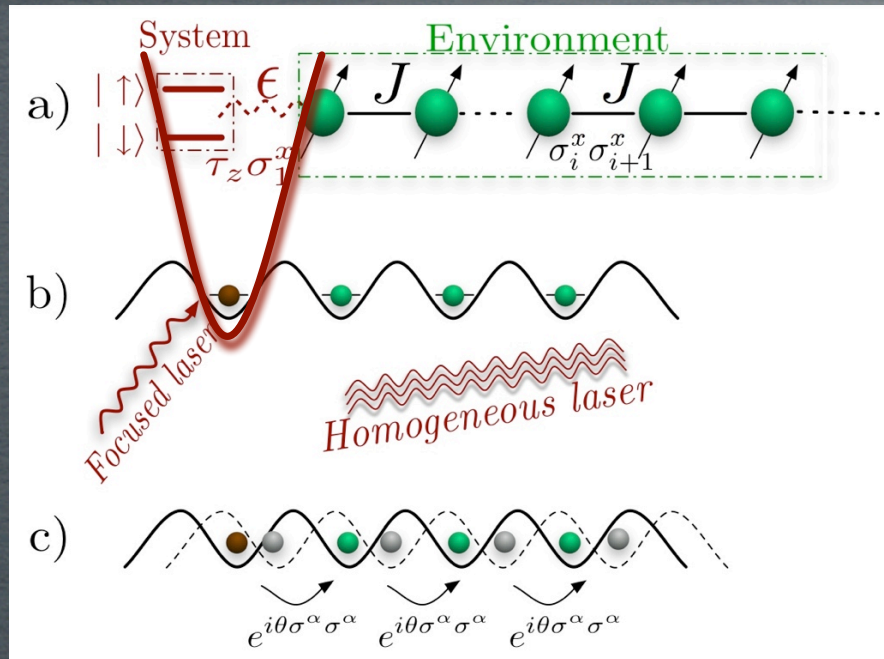
A two-level system (the quantum system) is coupled to a single spin of a one-dimensional spin-1/2 chain (the environment).

Physical realizations - JJAs



D. Haviland's group

Physical realizations - optical lattices



Optical lattices as simulators
of interacting spin systems

- E. Janè *et al* (2003)
- D. Jaksch *et al* (1999)
- O. Mandel *et al* (2003)

The Model

$$\mathcal{H} = \mathcal{H}_{TL} + \mathcal{H}_E + \mathcal{H}_{IN}$$

$$\mathcal{H}_{TL} = \omega_1 |1\rangle\langle 1|$$

$$\mathcal{H}_{TL} = -\epsilon |1\rangle\langle 1| \sigma_1^z$$

Pure dephasing

Unruh (1995)

Palma, Suominen, and Ekert, (1996)

$$\rho(t) = \begin{pmatrix} \rho_{00}(0) & \rho_{01}(t) \\ \rho_{10}(t) & \rho_{11}(0) \end{pmatrix}$$

$$\rho_{10}(t) = \rho_{10}(0)D(t)$$

$$D(t) = \langle e^{i\mathcal{H}t} e^{-i(\mathcal{H}_{TL} + \mathcal{H}_E)t} \rangle$$

Loschmidt echo

$$\mathcal{L}(t) = |D(t)|^2$$

MODELS
FOR THE
ENVIRONMENT

HE

Ising chain in a transverse field

$$H = -\frac{J}{2} \sum_{i=1}^N (1 - \gamma) \sigma_i^x \sigma_{i+1}^x + (1 + \gamma) \sigma_i^y \sigma_{i+1}^y - h \sum_{i=1}^N \sigma_i^z$$

$$\langle \sigma_x \rangle \neq 0$$

$\gamma \neq 0$ Ising universality class

$\gamma = 0$ XY universality class

1

$$\lambda = \frac{h}{J}$$

Exact solution - free fermions

Pfeuty '70

HEISENBERG MODEL

$$H = \sum_i J \left(\sigma_i^x \sigma_{i+1}^x + \sigma_i^y \sigma_{i+1}^y + \Delta \sigma_i^z \sigma_{i+1}^z \right)$$

HEISENBERG MODEL

Constant couplings



Ising chain environment

$$\mathcal{L}(t) = \det \left(1 - \mathbf{r} + \mathbf{r} e^{i\mathbf{C}t} \right)$$

$$\mathbf{C} = \begin{pmatrix} A & B \\ -B & -A \end{pmatrix}$$

$$r_{i,j} = \langle \Psi_i^\dagger \Psi_j \rangle$$

$$B_{j,k} = -J\gamma (\delta_{k,j+1} - \delta_{j,k+1})$$

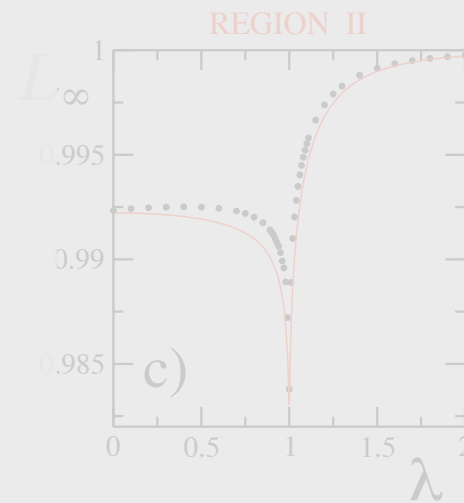
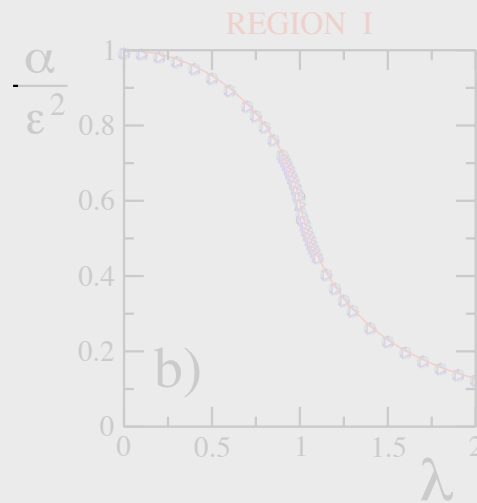
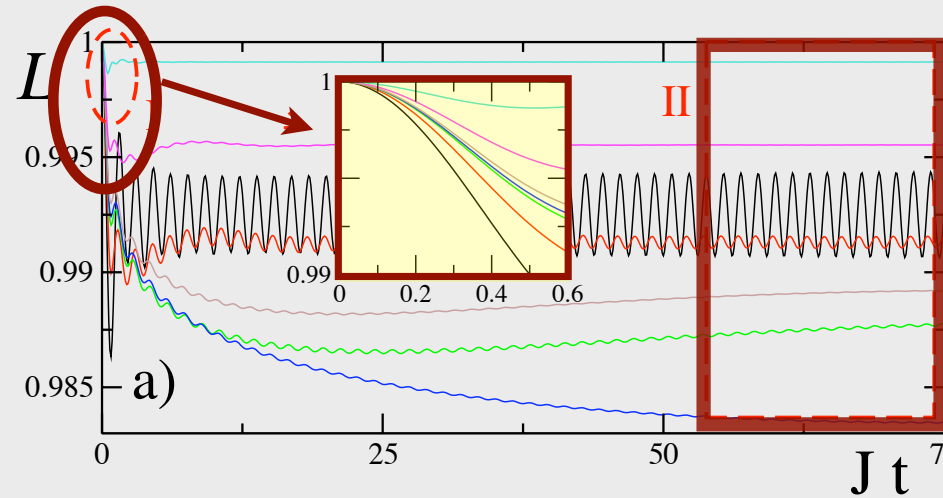
$$A_{j,k} = -J(\delta_{k,j+1} + \delta_{j,k+1}) - 2(\lambda + \epsilon_j)\delta_{j,k}$$

Ising chain environment

$$\mathcal{L}_I(t) \sim e^{-\alpha t^2}$$

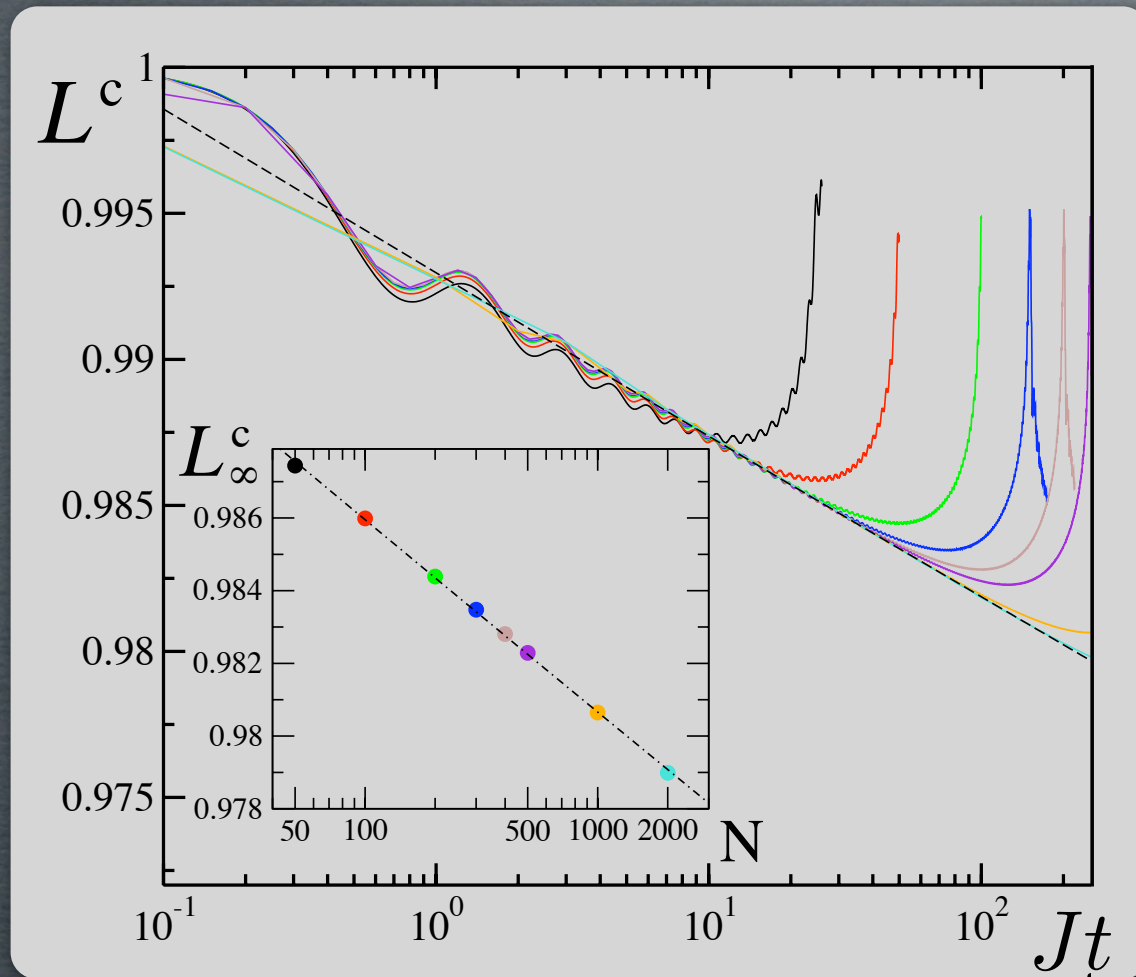
$\lambda < 1$ oscillations

$\lambda > 1$ constant



Ising chain environment

$$\lambda = \lambda_c$$



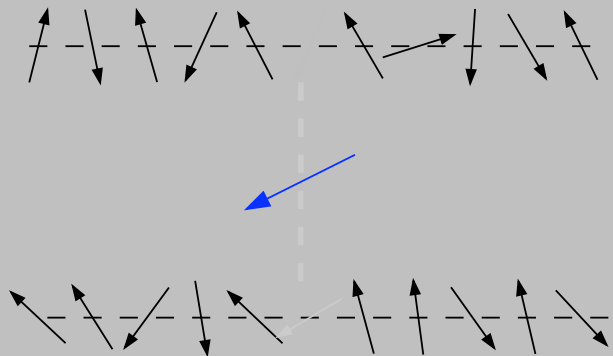
$$\mathcal{L}^c(t) = \frac{c_0}{(1 + c_1 \ln t)}$$

$$\mathcal{L}^c_\infty = \frac{l_\infty}{1 + \beta \ln N}$$

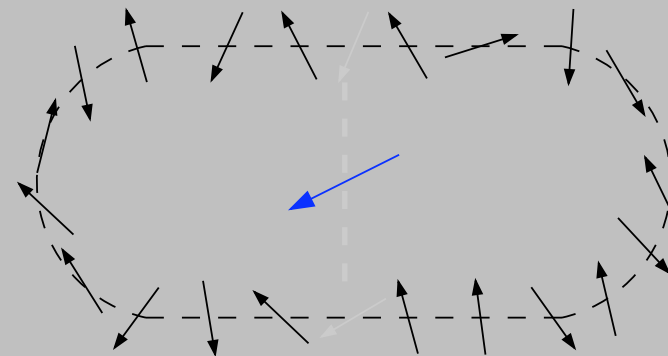
Ising chain - correlations in the bath

If the chain is not at the critical point, the decay of the coherences is independent of the number of spins in the bath

Configuration 1

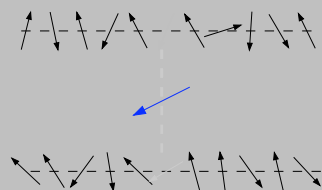


Configuration 2

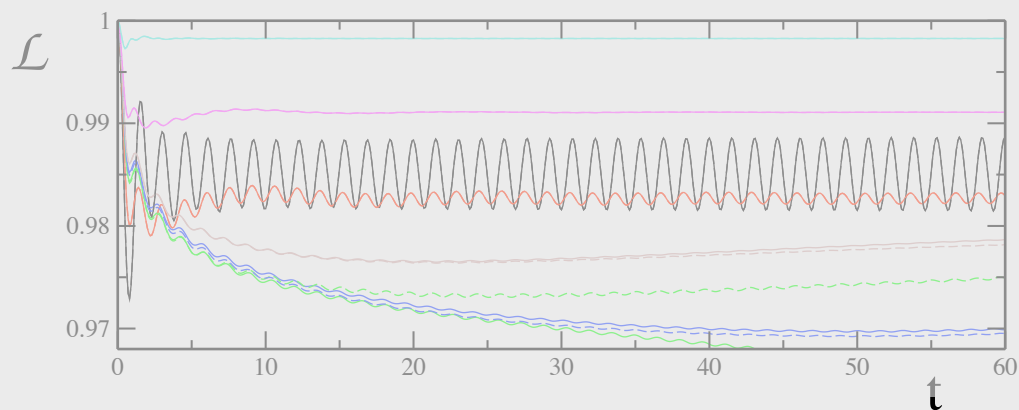
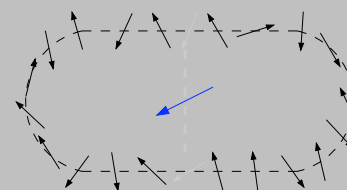


Ising chain - correlations in the bath

Configuration 1



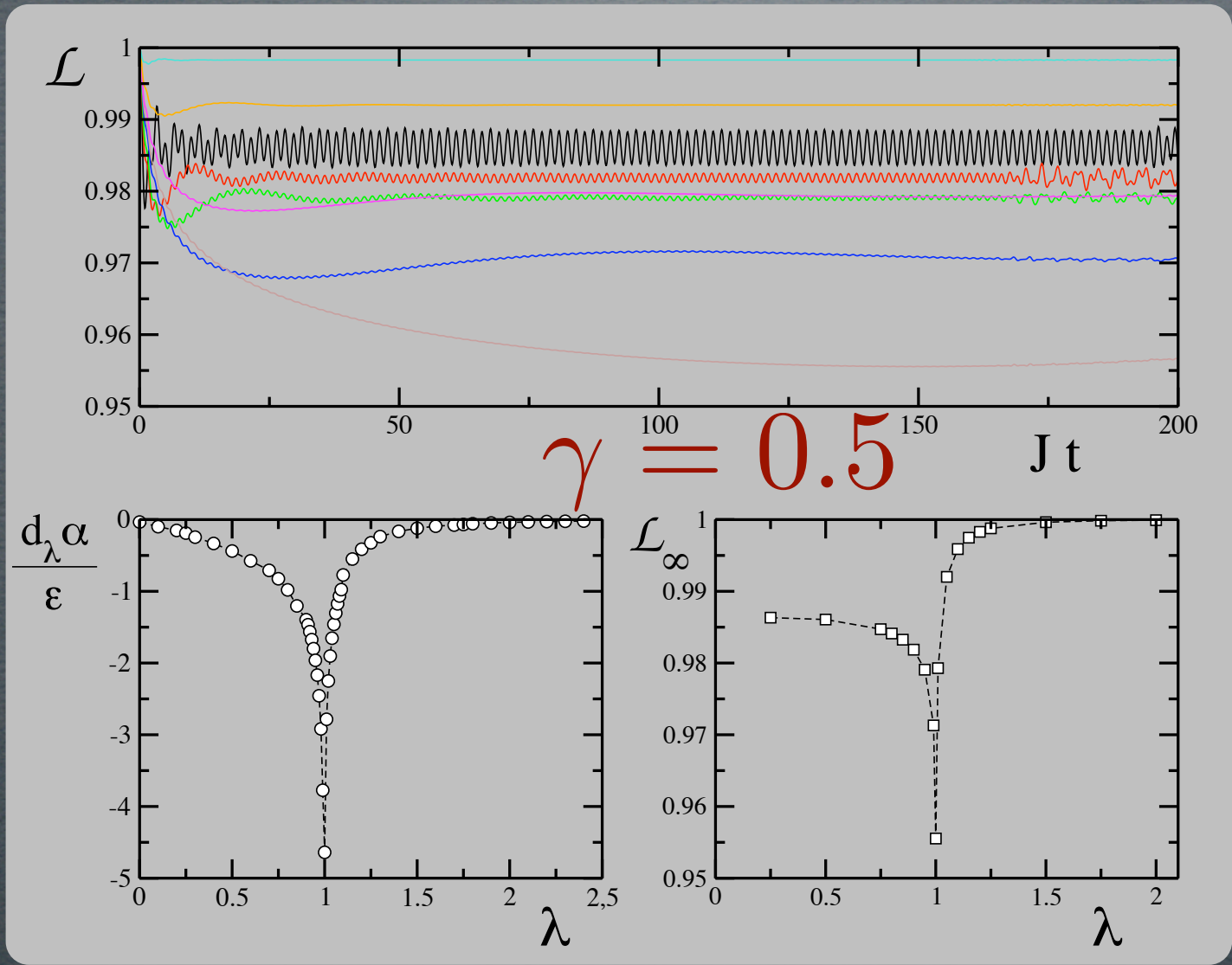
Configuration 2



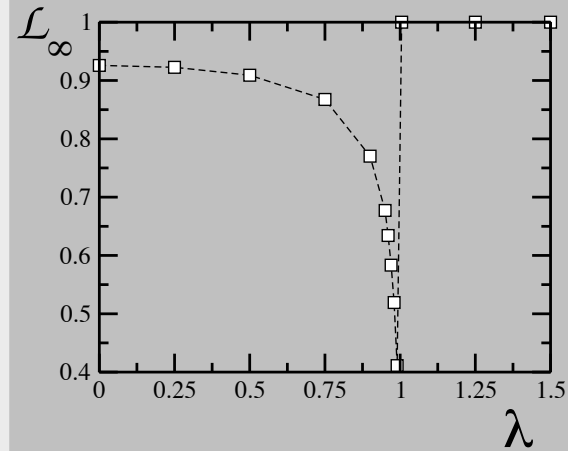
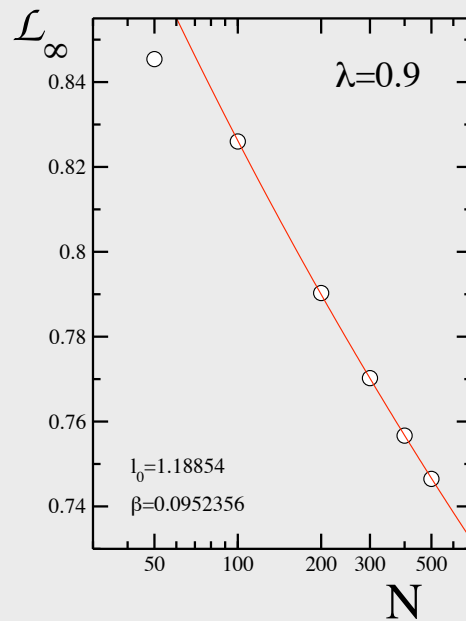
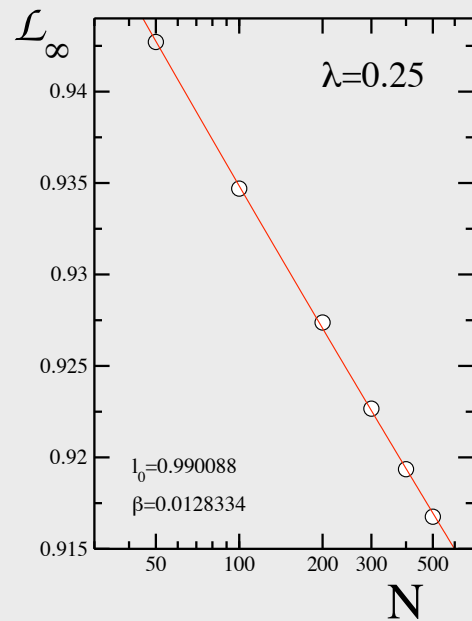
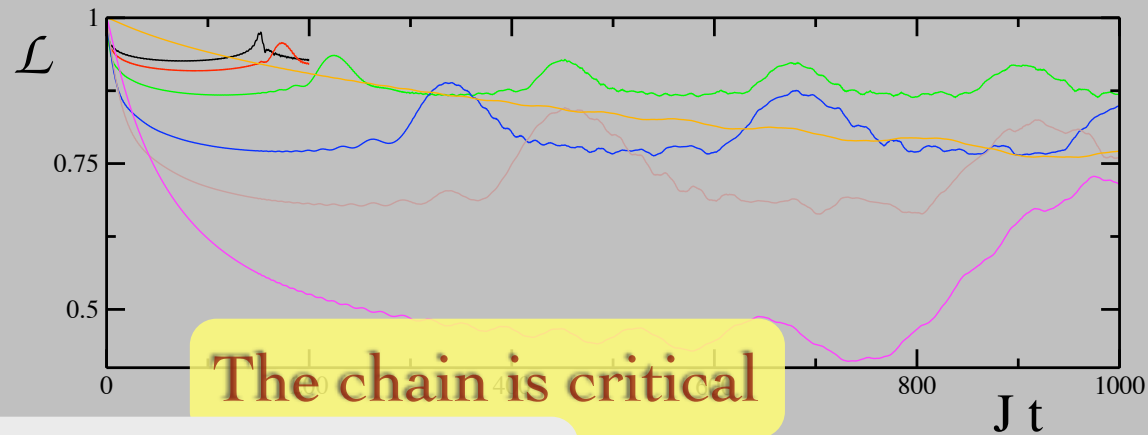
$$\lambda \neq \lambda_c$$

Ising chain environment - Universality

$$\gamma \neq 0$$



XY chain environment



Decoherence vs interaction & entanglement in the bath

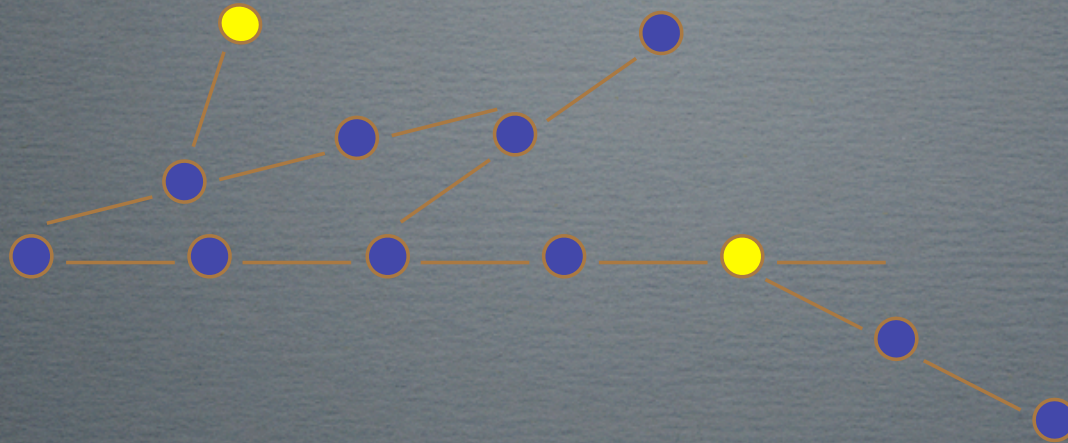
see the suggestion of

C. M. Dawson, A.P. Hines, R.H. McKenzie, and G.J. Milburn G.J., (2005)

How to measure entanglement?

- Entanglement between two spins in the network
- Multipartite entanglement
- Block entropy
- Localizable entanglement
- ...

Bipartite entanglement



The state of the two selected spins is mixed

Measure of mixed state entanglement for two spin-1/2 states

- Separable state

$$|\alpha\rangle = |00\rangle$$

- "NOT"

$$|\beta\rangle = |11\rangle$$

$$\langle\beta|\alpha\rangle = 0$$

- Entangled state state

$$|\alpha\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$

- "NOT"

$$|\beta\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$

$$\langle\beta|\alpha\rangle = 1$$

Concurrence between spins at sites i and j

- construct

$$R = \rho(i, j) \tilde{\rho}(i, j)$$

where

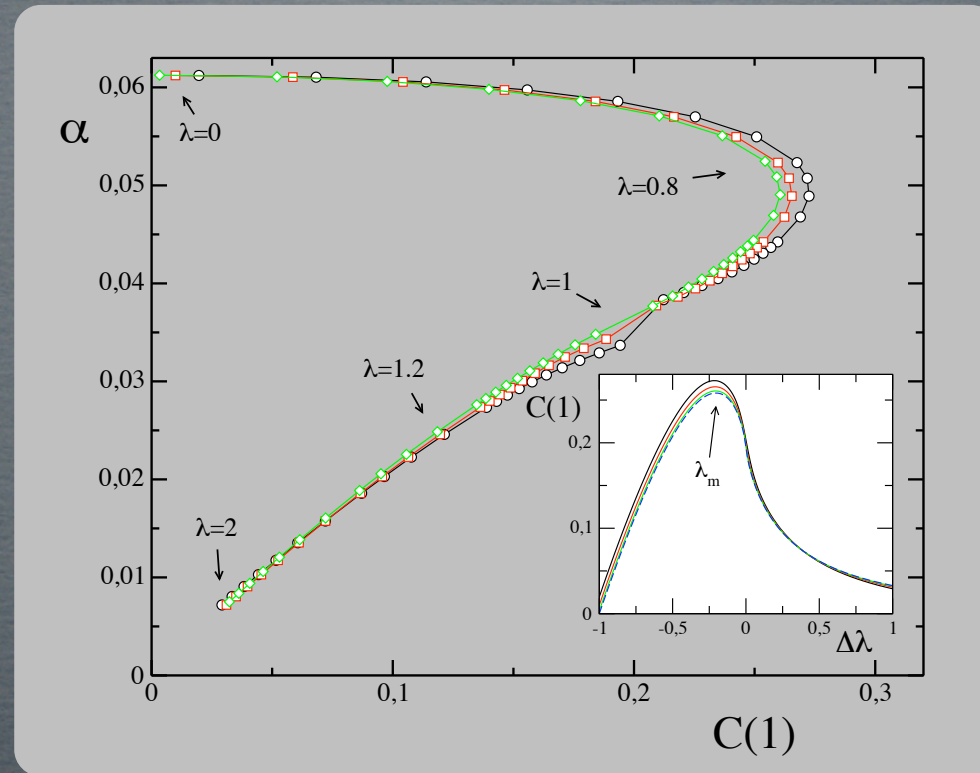
$$\tilde{\rho} \doteq \sigma^y \otimes \sigma^y \rho^* \otimes \sigma^y$$

- the concurrence is defined as

$$C(i, j) = \max\{0, \lambda_1(i, j) - \lambda_2(i, j) - \lambda_3(i, j) - \lambda_4(i, j)\}$$

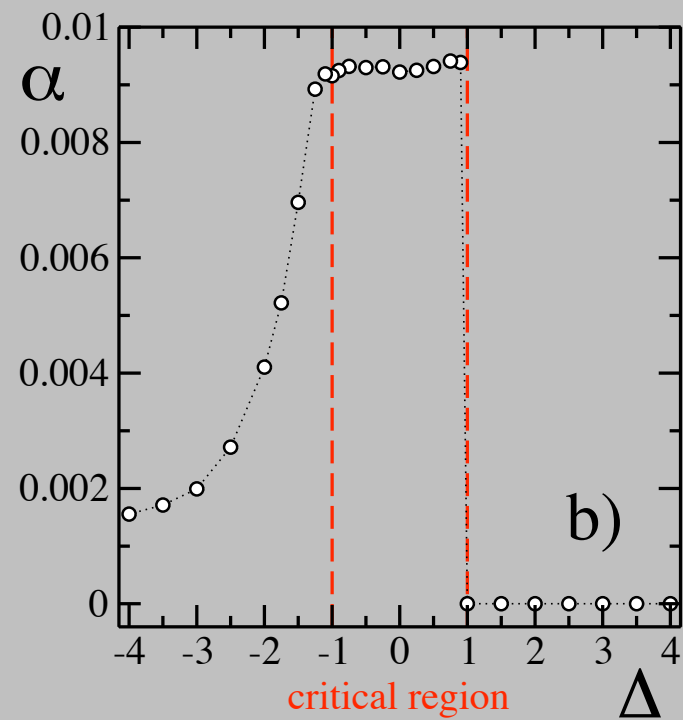
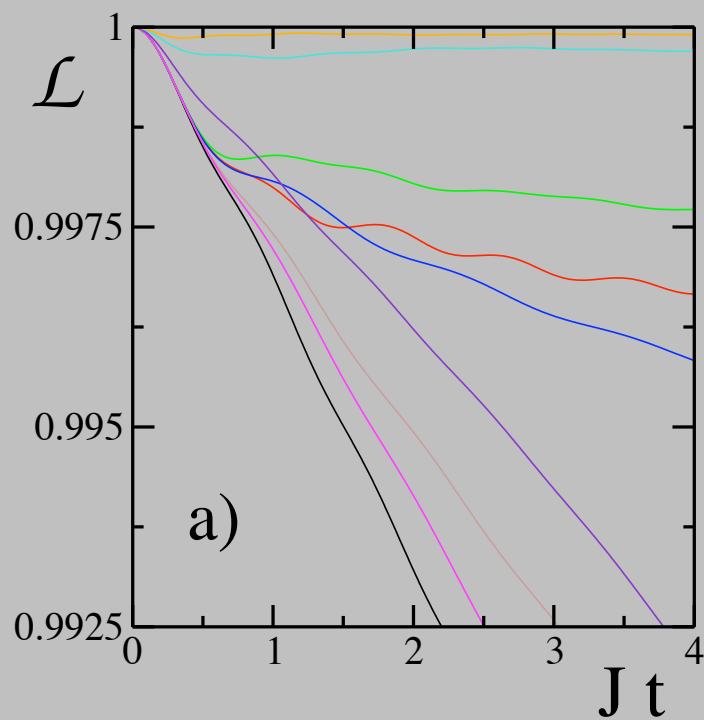
where λ_s are the eigenvalues of R in ascending order

Decoherence vs entanglement in the bath



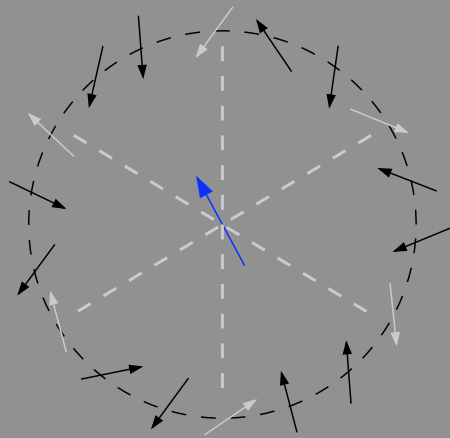
with the suggestion of C. M. Dawson *et al.*, (2005)

Heisenberg chain environment - tDMRG

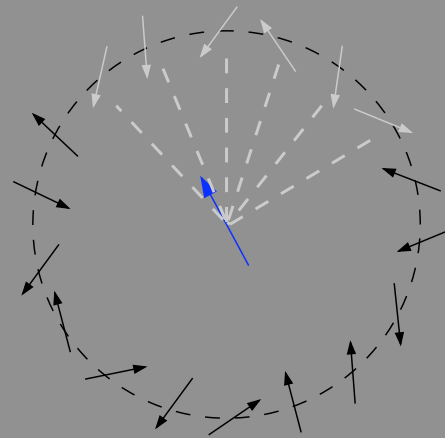


Coupling to m spins in the environment

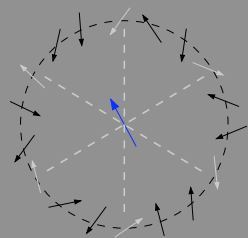
Configuration A



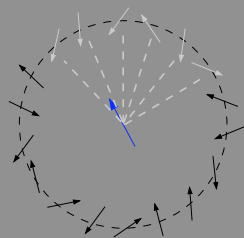
Configuration B



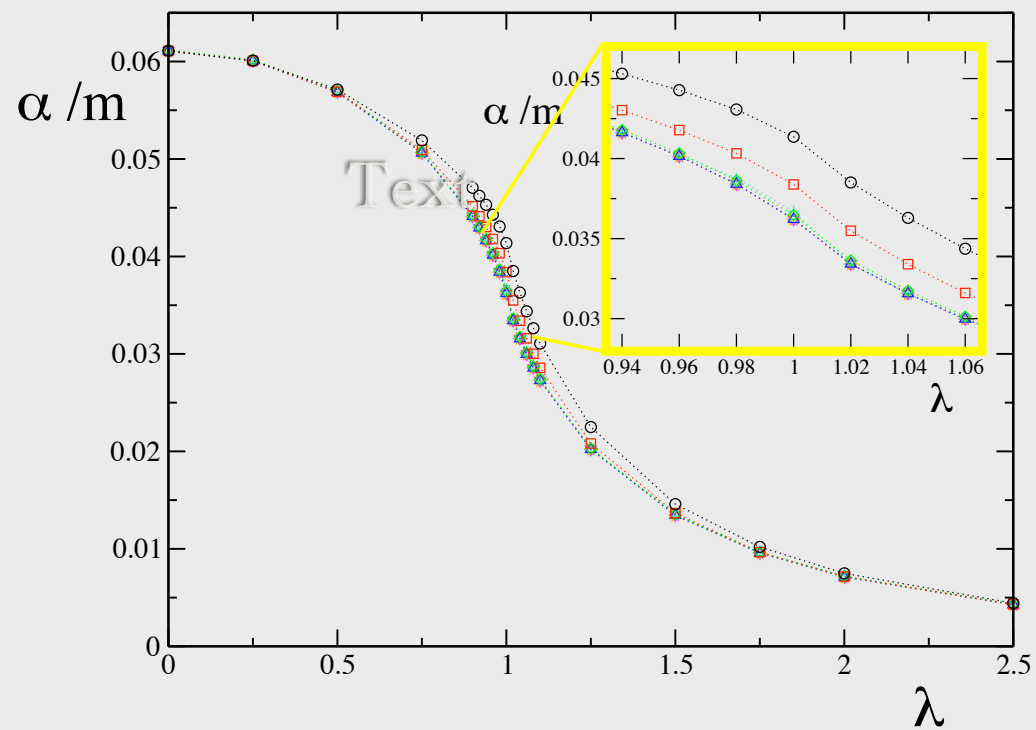
Configuration A



Configuration B

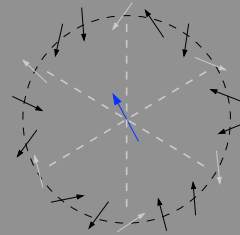


Weakly
dependent on m

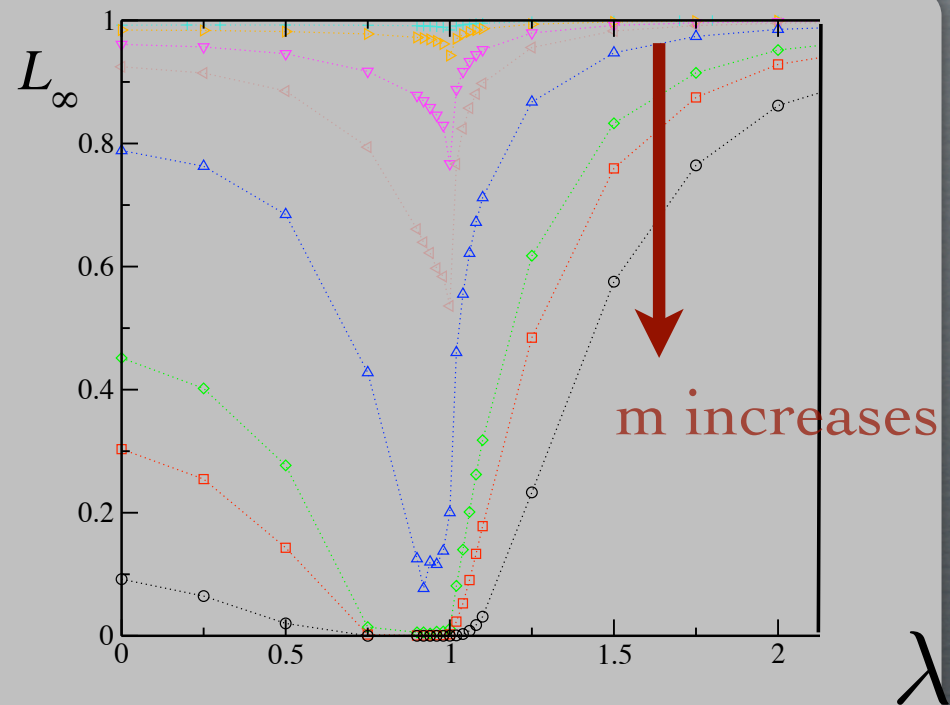
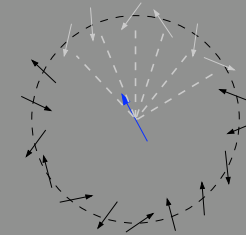


Strong dependence on the number of couplings

Configuration A



Configuration B



Conclusions

- ✓ Optical lattices to simulate quantum baths
- ✓ Exact solution for an interacting baths
- ✓ Numerical t-DMRG for Heisenberg bath