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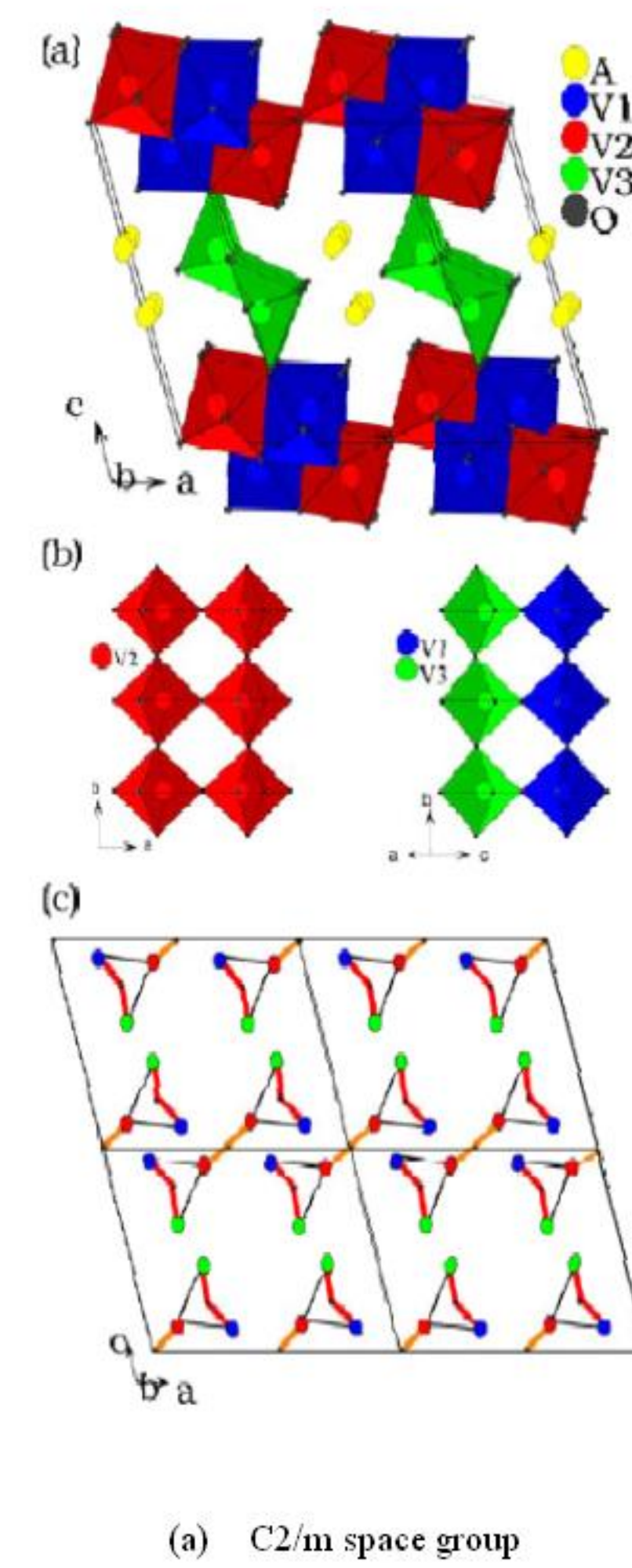
ABSTRACT :

The vanadium bronzes $\beta\text{-A}_x\text{V}_6\text{O}_{15}$ (A=Sr, Ca, Na, Li, Ag ...), first discovered in the 80's, are mixed-valence compounds V^{4+} ($3d^1$)/ V^{3+} ($3d^0$) presenting very exciting physics such as low dimensional electronic properties, metal-insulator transition and superconductivity. The stoichiometric samples $(\text{Na,Sr})_{x=1}\text{V}_6\text{O}_{15}$ present a complex phase transition associating a metal-insulator transition (MIT), a structural transition and a charge ordering occurring at $T_{\text{MI}}=170$ K for $\text{Sr}_1\text{V}_6\text{O}_{15}$ and $T_{\text{MI}}=135$ K for $\text{Na}_1\text{V}_6\text{O}_{15}$. Superconducting properties have been also evidenced recently at high pressure below $T=9$ K in $\text{NaV}_6\text{O}_{15}$.

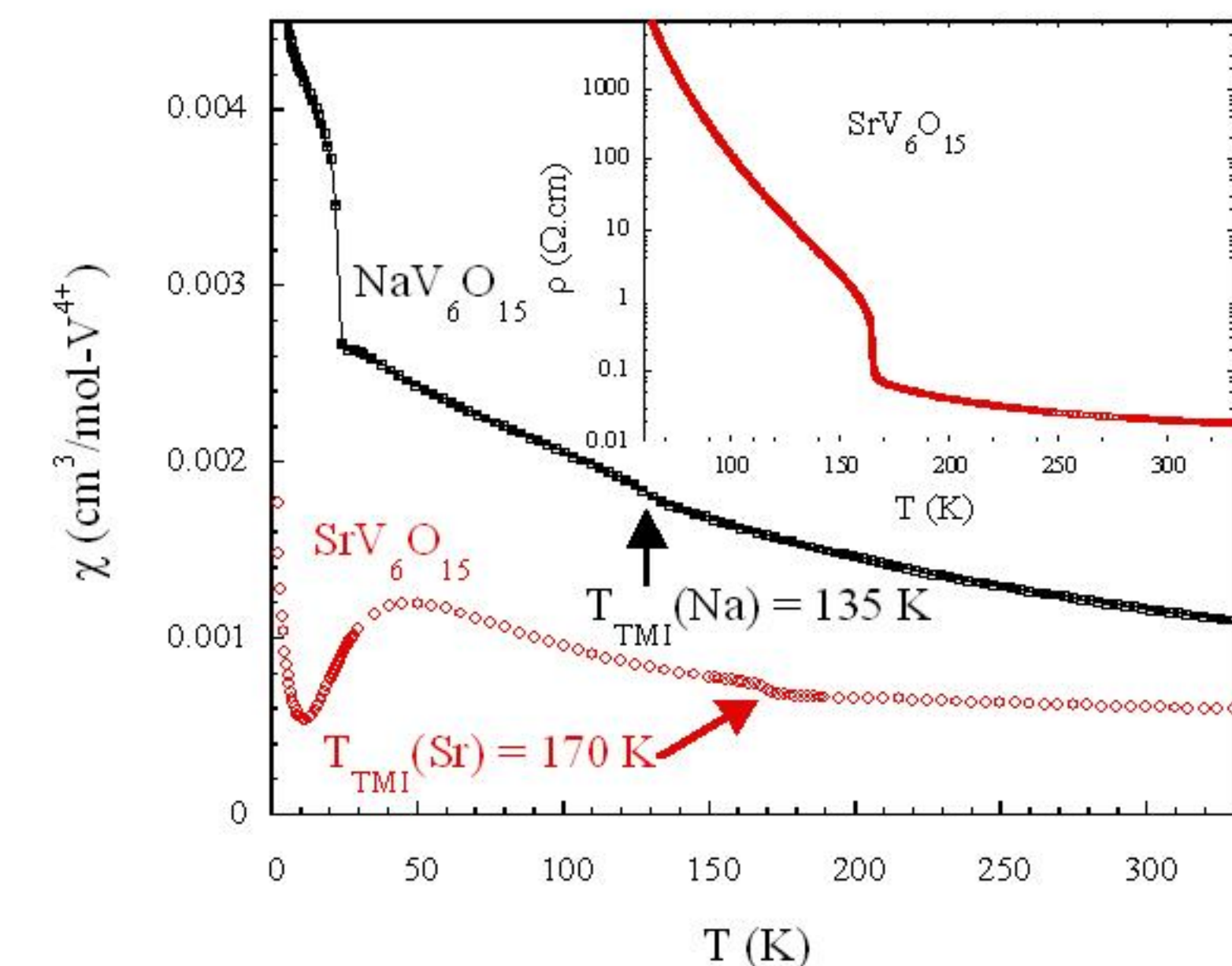
At room-temperature, they present « metallic » properties ($\rho_{\parallel} \approx 10^{-3} \Omega \cdot \text{cm}$ and $\rho_{\perp}/\rho_{\parallel} \approx 100$) with a quasi-1D Fermi surface (from band structure calculations). Surprisingly, $\text{NaV}_6\text{O}_{15}$ has a metallic resistivity associated with a Curie-Weiss magnetic susceptibility whereas $\text{SrV}_6\text{O}_{15}$ is semiconducting with a Pauli-like paramagnetic susceptibility suggested delocalized carriers ! The ordering of the Sr (for $x>0.9$) and Na (for $x<230$ K) atoms lead to a doubling of the lattice parameter along the b-axis. Only one type of $\text{V}^{4.83+}$ and $\text{V}^{4.66+}$ site are observed respectively for $\text{NaV}_6\text{O}_{15}$ and $\text{SrV}_6\text{O}_{15}$ above the MIT.

Below the Metal-insulator transition (MIT), a tripling of the lattice parameter is observed in both cases associated with a charge ordering. Nevertheless, magnetism is still present at low-T : a Curie-Weiss susceptibility is observed in the case of $\text{NaV}_6\text{O}_{15}$ whereas a spin-gap phase is revealed in the case of $\text{SrV}_6\text{O}_{15}$. A charge gap of 200 meV is evidenced below the TMI in $\text{Sr}_1\text{V}_6\text{O}_{15}$ by angle-integrated photoemission measurements associated with a stabilization of the correlated d band. One of the key point to understand this complex phase transition could be the Fermi wave vector k_{F} which can be extracted from ARPES measurements on single crystals.

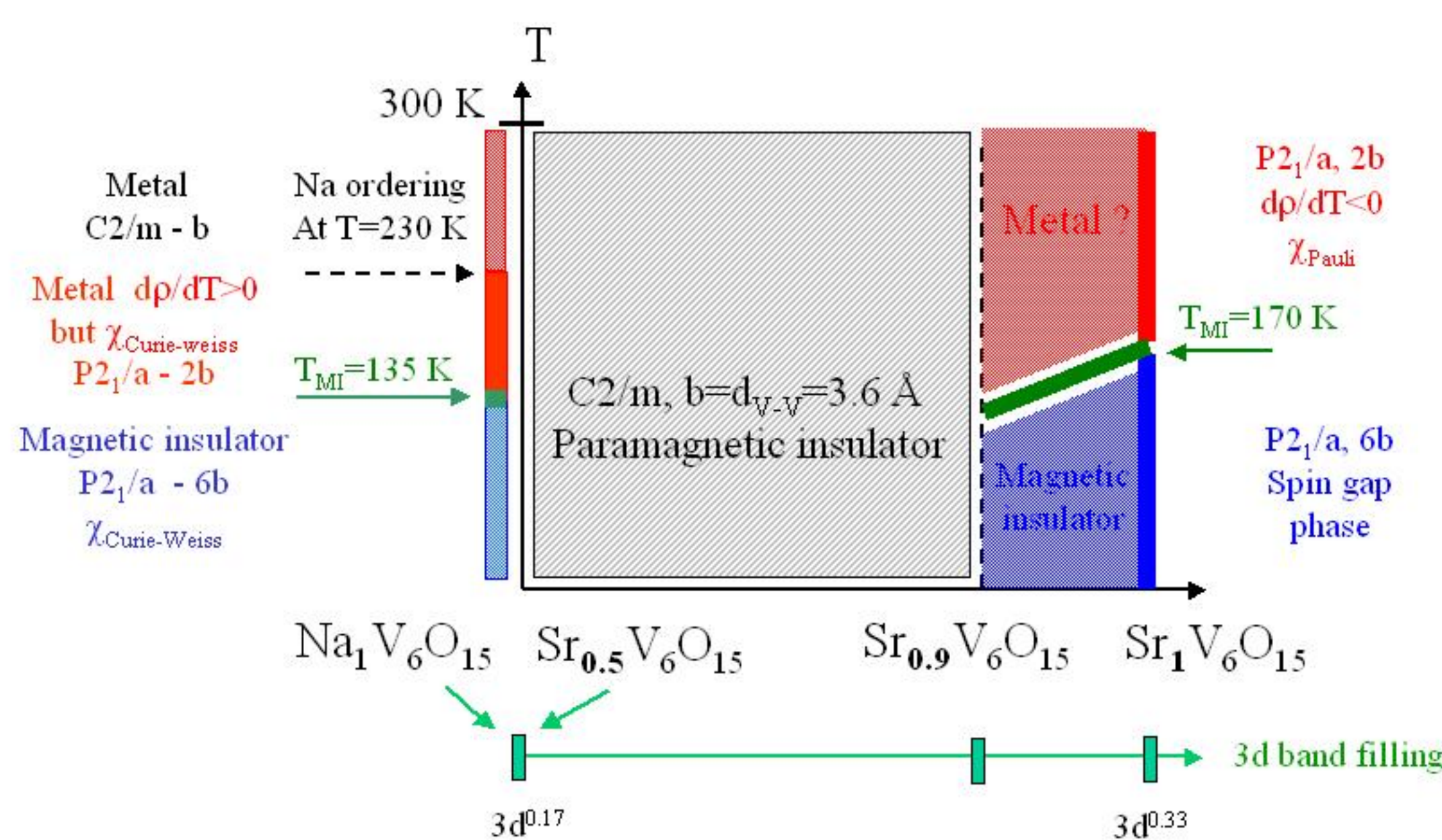
$\text{Sr}_x\text{V}_6\text{O}_{15}$ structure



Transport and magnetic measurements

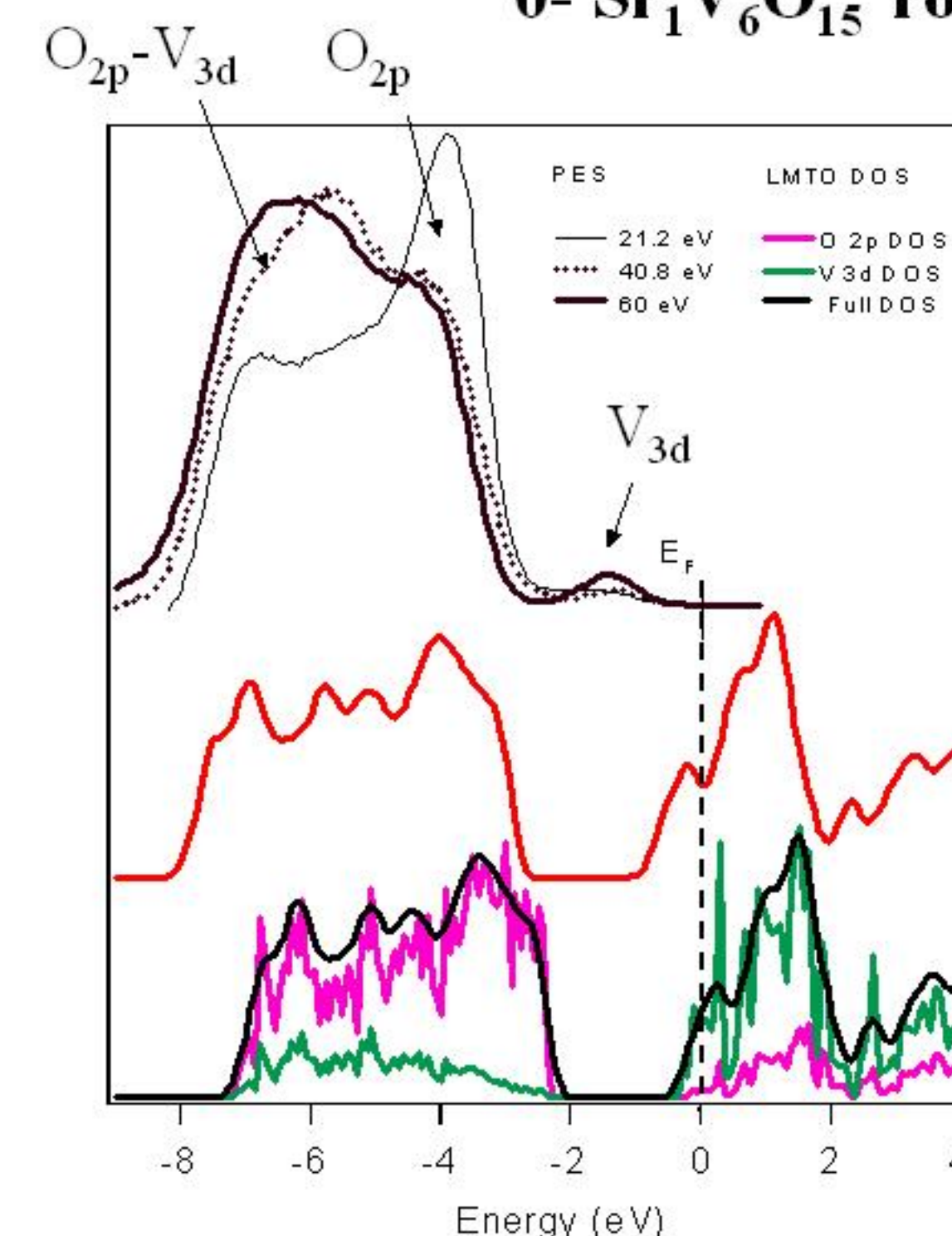


$\text{A}_x\text{V}_6\text{O}_{15}$ PHASE DIAGRAM



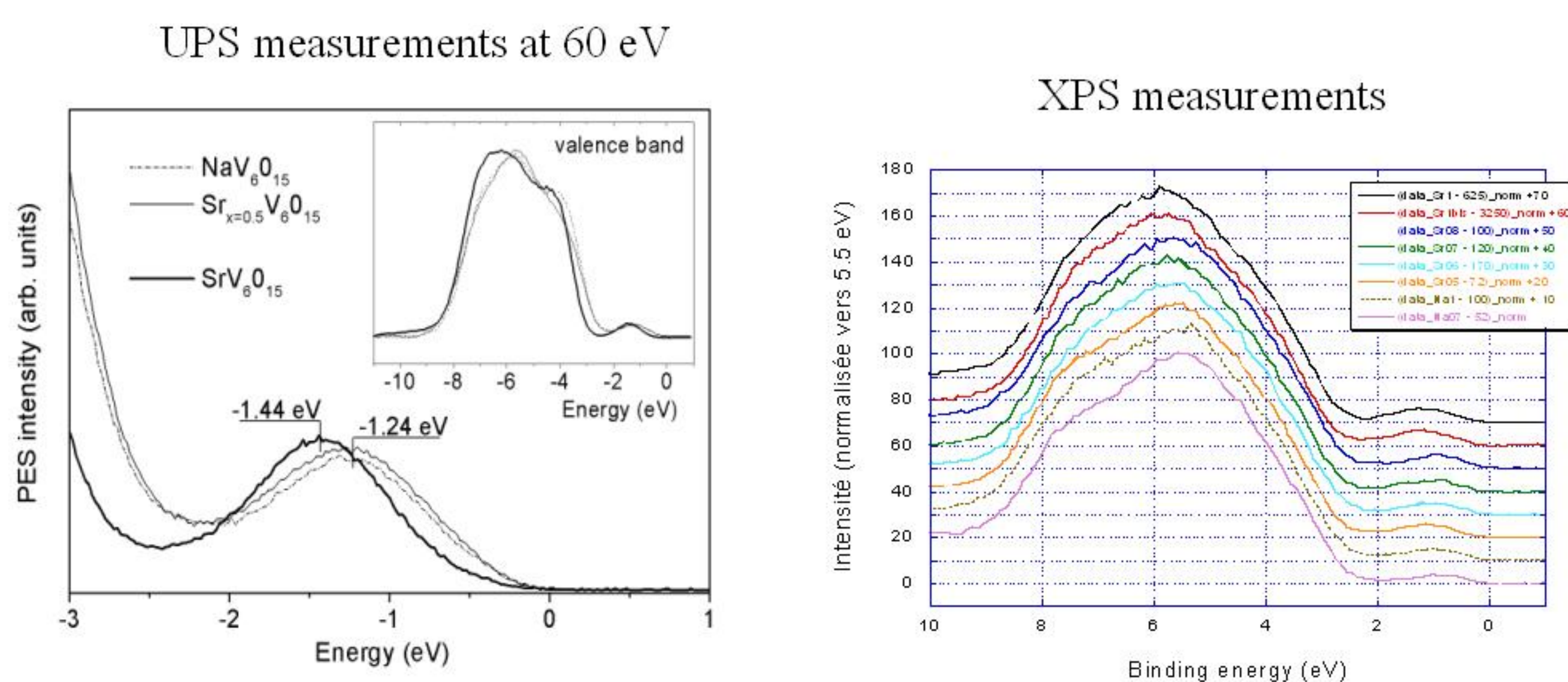
→ The MIT has been investigated by ARPES both as function of doping in $\text{Sr}_x\text{V}_6\text{O}_{15}$ and as function of temperature in $\text{Sr}_1\text{V}_6\text{O}_{15}$ and $\text{NaV}_6\text{O}_{15}$.

0- $\text{Sr}_1\text{V}_6\text{O}_{15}$ room-T electronic structure



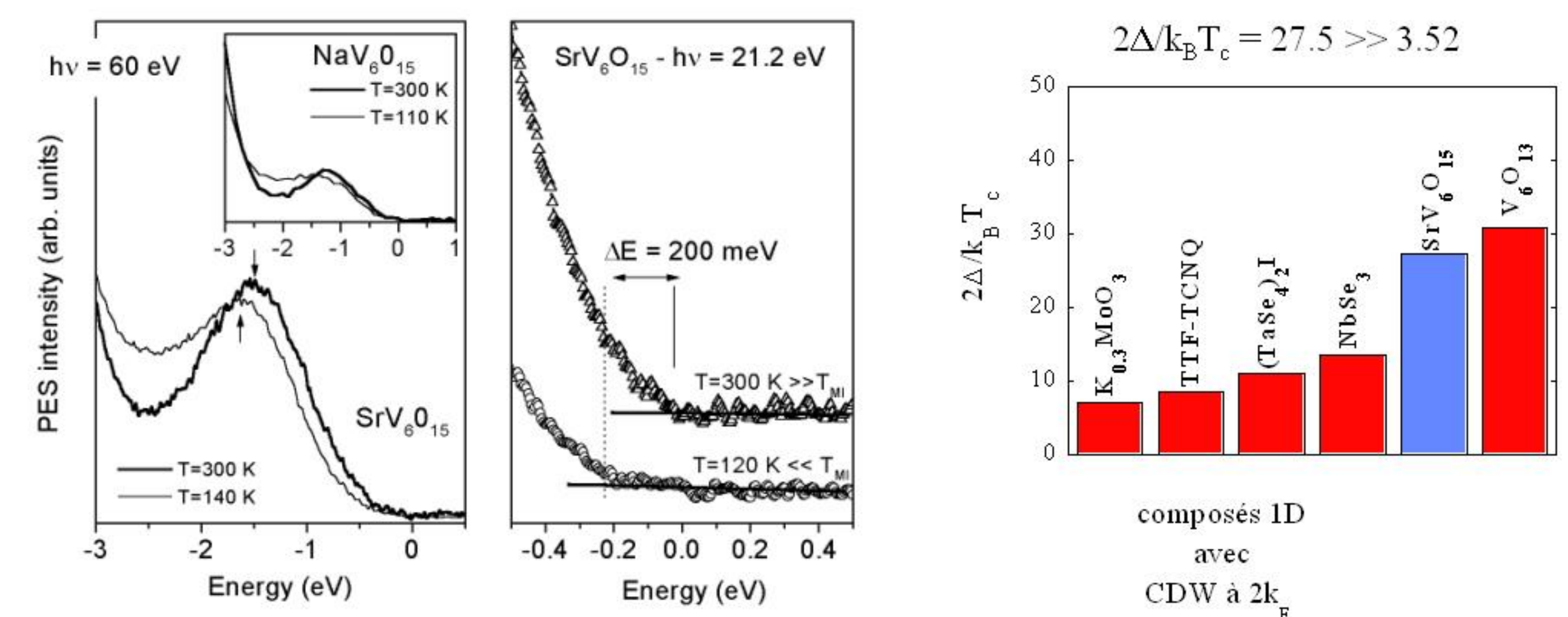
Angle-integrated photoemission spectra obtained on $\text{Sr}_1\text{V}_6\text{O}_{15}$ powder samples for different photon energies. The energy dependence allows to clearly identify O_{2p} bands for binding energies higher than 3 eV and V_{3d} spectral feature around 1.4 eV. A comparison with LMTD Band structure calculations is presented. As expected, the $\text{O}_{2p}\text{-V}_{3d}$ hybridization is corroborated by the strong increase of this feature with increasing the photon energy. On the other hand, the calculated V_{3d} density of occupied states is maximum at the Fermi energy whereas the experimental one shows a maximum around 1.4 eV. Even far from half filling ($3d^{0.33}$), the observation of a strong incoherent peak at a so high energy is probably due to strong correlations, this 3d band constituting the lower Hubbard band observed in numerous vanadium oxides. In addition, no spectral weight has been measured at E_{F} signifying a quasi-1D metallic behavior and/or polaronic effects (small polarons).

1 - Electronic structure as function of doping



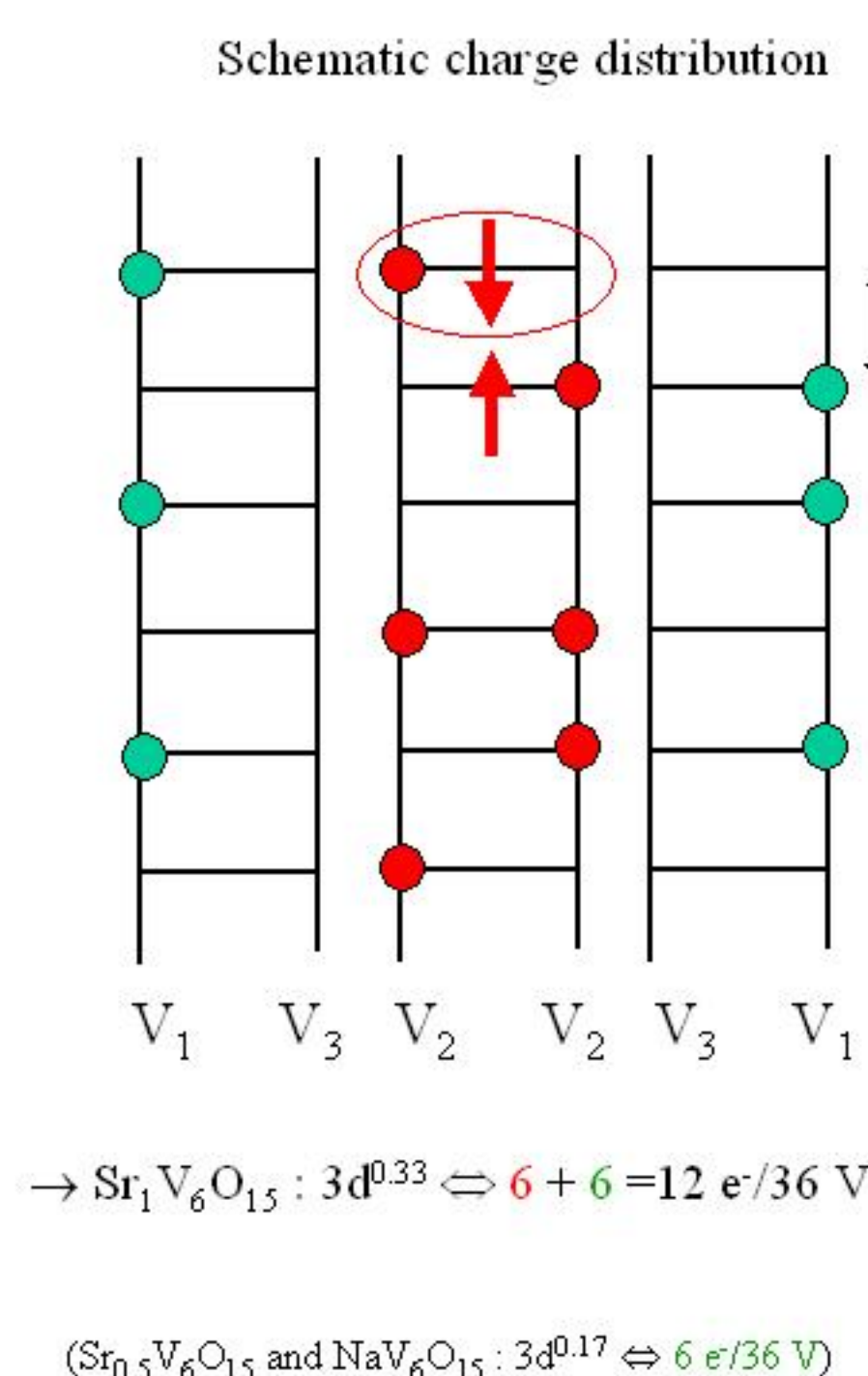
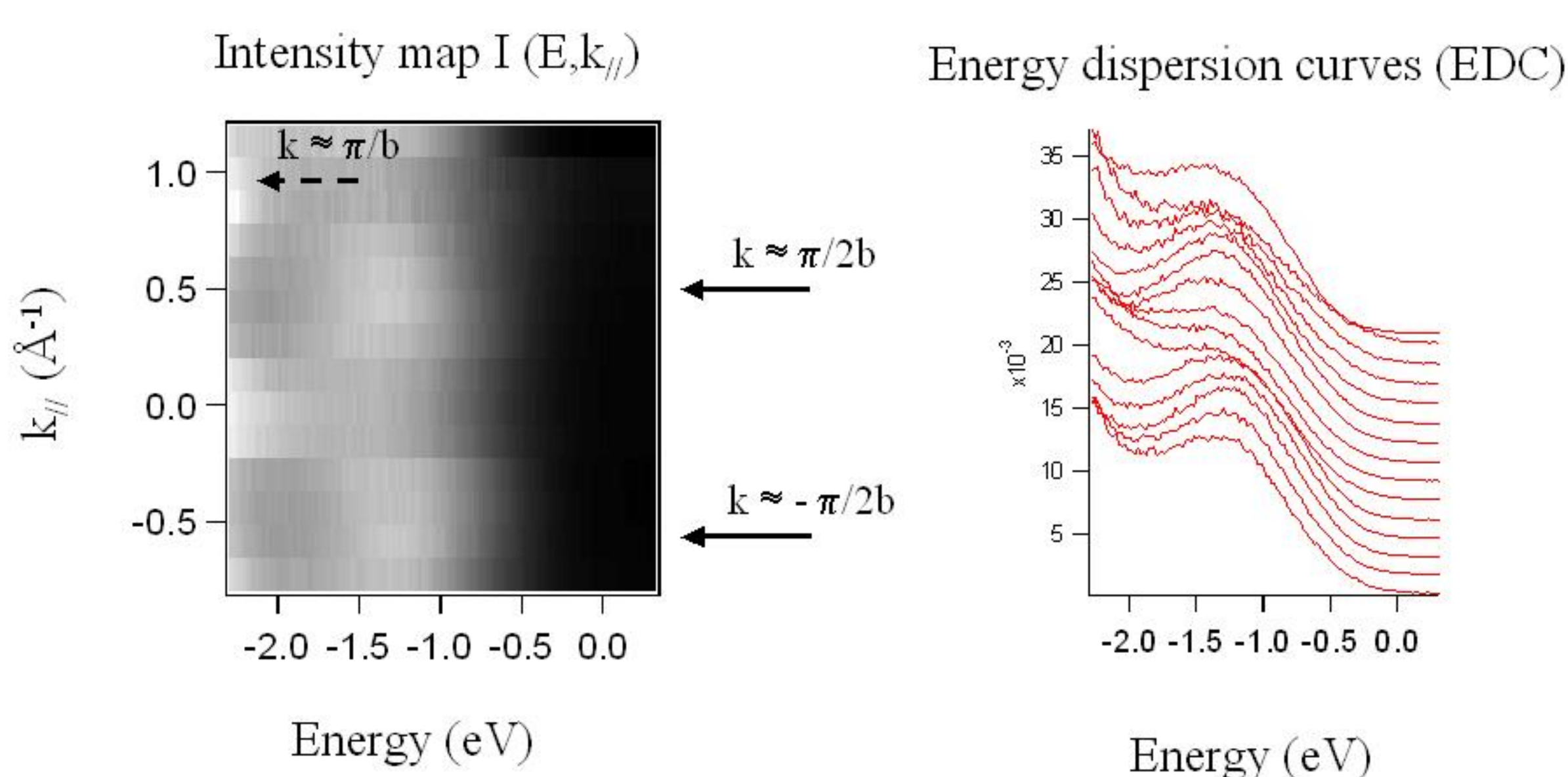
The position of the V_{3d} incoherent peak (partially) follows the 3d band filling : $E=1.24$ eV in the case of $\text{Sr}_{0.5}\text{V}_6\text{O}_{15}$ and $\text{NaV}_6\text{O}_{15}$ which have the same filling ($3d^{0.17}$) and $E=1.44$ eV for $\text{Sr}_1\text{V}_6\text{O}_{15}$ ($3d^{0.33}$). This is corroborated by XPS measurements suggesting that these measurements are bulk-like !

2 - Charge gap across the MIT: opening of a gap $\Delta \approx 200$ meV



The V_{3d} spectral feature is stabilized below the metal-insulator transition (energy gain ≈ 150 meV). In addition, a precise investigation of the spectral weight close to the Fermi level evidences the opening of a gap in the occupied density of states in agreement with an increase with more than one order of magnitude of the resistivity and the insulating nature of the low-T phase. The strong BCS ratio measured in this compound is higher than those measured in standard $2k_{\text{F}}$ Peierls compounds. It is closer to polaronic compounds like Ti_4O_7 ...

3 - Angle-resolved results on $\text{Sr}_{0.5}\text{V}_6\text{O}_{15}$ single crystals



DISCUSSION

→ origin of the absence of spectral weight at E_{F} (non Fermi liquid, small polarons...)
→ mechanism of this (very) complex transition :
 $2k_{\text{F}}$ Charge density wave excluded by the presence of magnetism below T_{MI}
 $4k_{\text{F}}$ Charge density wave (in a correlated metal) ?
 $2b \rightarrow 6b$ ($\times 3$) a tripling of the lattice parameter is possible if $\rho=1/3=0.33 \text{ e}^-/\text{site}$ (CDW at $4k_{\text{F}}$) or $\rho=2/3=0.66 \text{ e}^-/\text{site}$ (CDW at $2k_{\text{F}}$)
Others...
⇒ to answer we need the k_{F} value measured by ARPES
In $\text{NaV}_6\text{O}_{15}$: $k_{\text{F}}=0.22 \text{ \AA}^{-1} \approx \pi/4b$ in the $\text{P}2_1/a$ phase (from Okazaki et al., cond-mat/0308368)
...Should be measured on $\text{Sr}_{0.5}\text{V}_6\text{O}_{15}$ and $\text{Sr}_1\text{V}_6\text{O}_{15}$ single crystals