



Instant superconductor: Just Add Water!

Electronic Structure, Magnetism and Superconductivity in Na_xCoO_2

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Acknowledgements:

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M.J. Mehl (NRL)

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Nesting, spin fluctuations, and odd-gap superconductivity in $\text{Na}_x\text{CoO}_2 \cdot y\text{H}_2\text{O}$

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(Naval Research Laboratory)

Collaborators:

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Acknowledgment:

Dan Agterberg



Outline

Introduction

-- What's interesting about Na_xCoO_2 ?

PART I

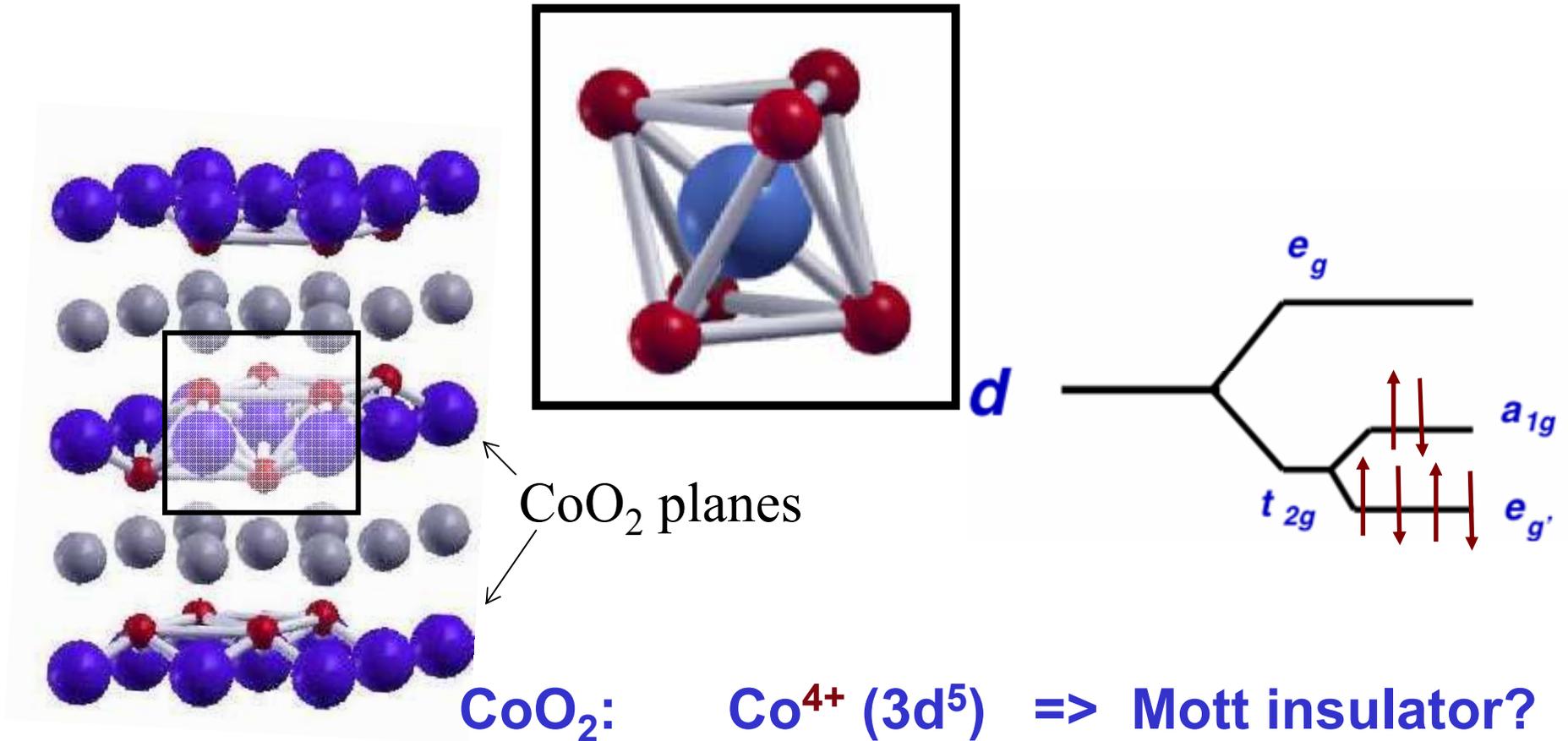
Electronic properties

PART II

Spin fluctuations and superconductivity



The Distorted Octahedral Environment of Co Ions

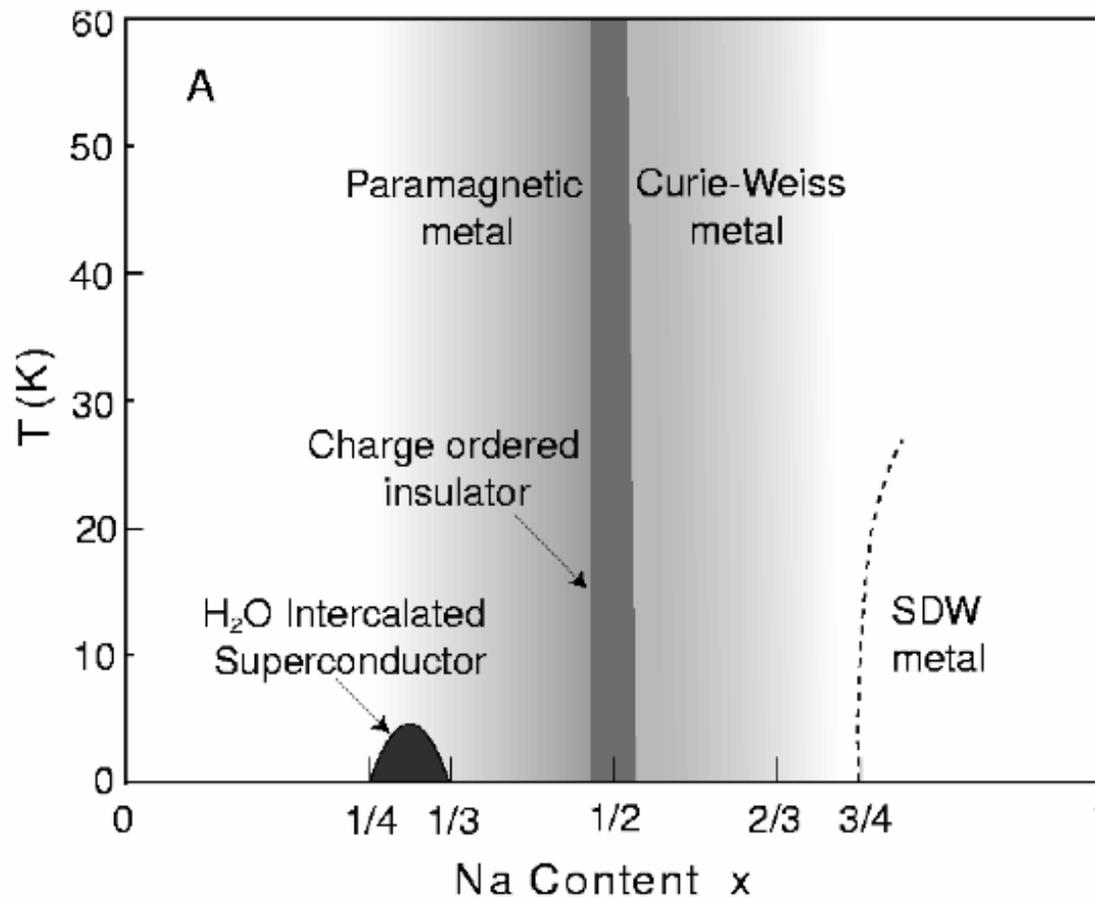


NaCoO₂: Co³⁺ (3d⁶) => band insulator

But Na_xCoO₂ behaves almost *oppositely*...



Na content phase diagram



- At $x=0$ system is a nonmagnetic metal
- At $x=1$ compound does not form
- At $x=0.5$ system is insulating

EXPECTED

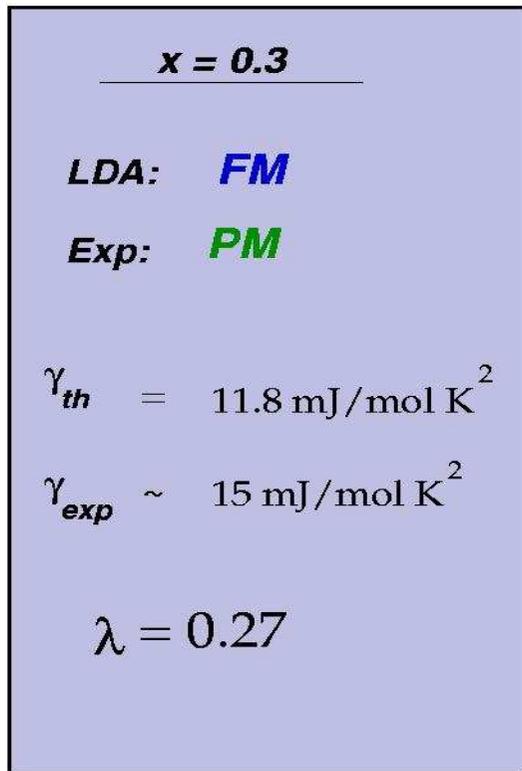
- At $x = 0$, system is a magnetic insulator
- At $x=1$, system is a band insulator
- For $x < 0.5$, system is a magnetic metal
- For $x > 0.5$, system is a simple metal

OBSERVED

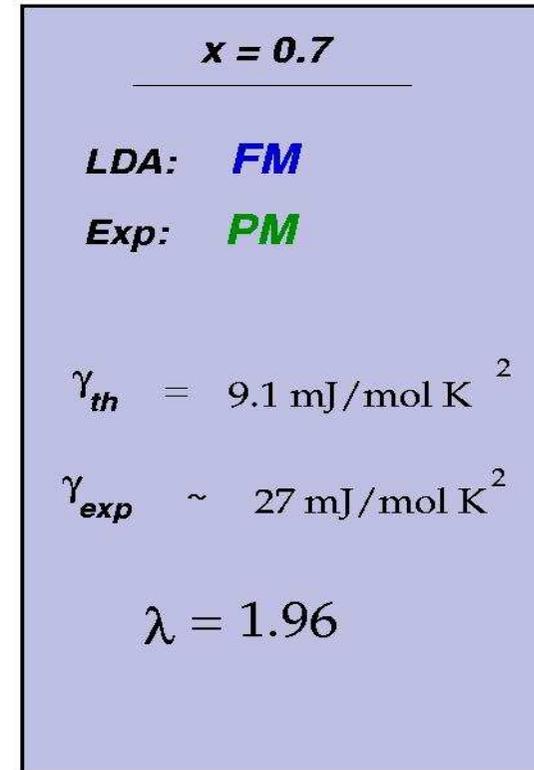
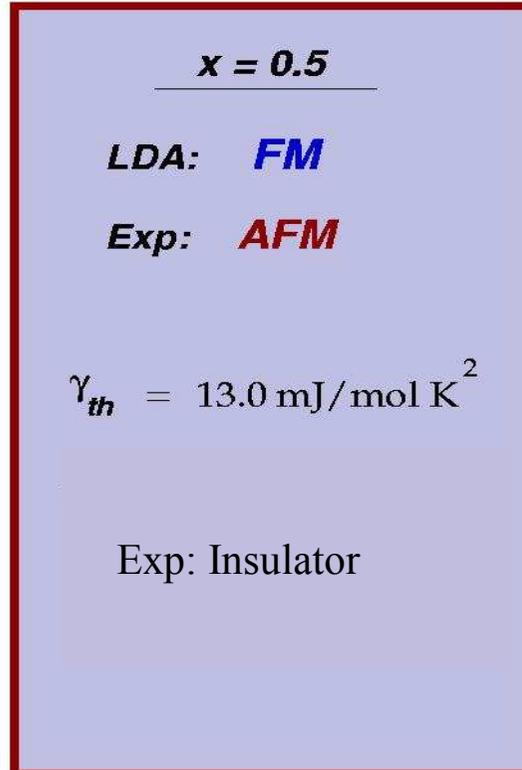
- For $x < 0.5$, system is a simple metal
- For $x > 0.5$, system go through a sequence of magnetic metallic phases



The calculated phase diagram of Na_xCoO_2



itinerant



localized

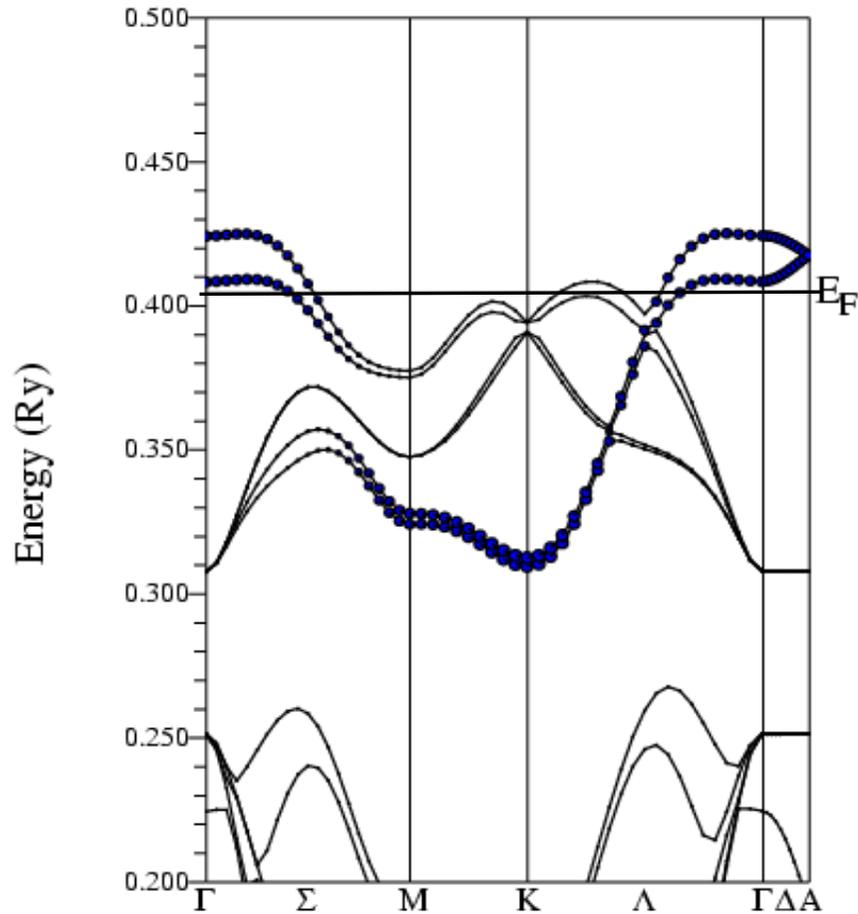
LDA typically finds *smaller* magnetic moments than experiment

Exception: the vicinity of a quantum critical point

Consistent overestimation of magnetism suggests spin fluctuations

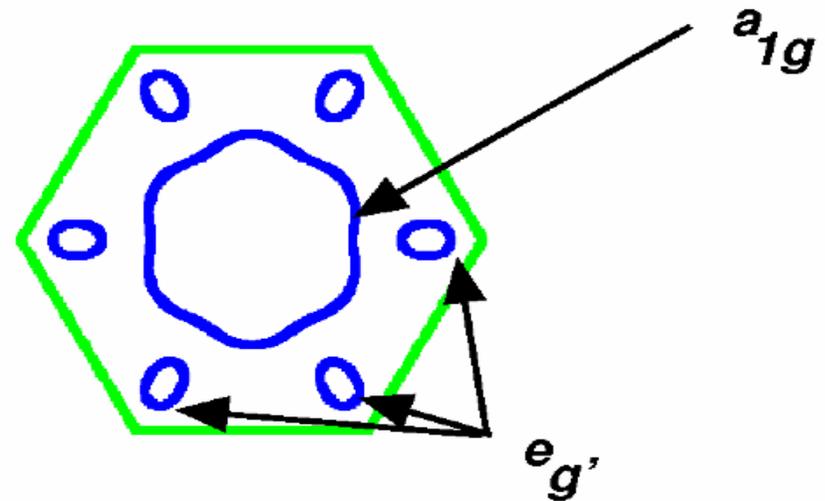


Multi-Orbital Nature of Fermi Surfaces



$$a_{1g} = (xy) + (yz) + (zx) = 3z^2 - r^2$$

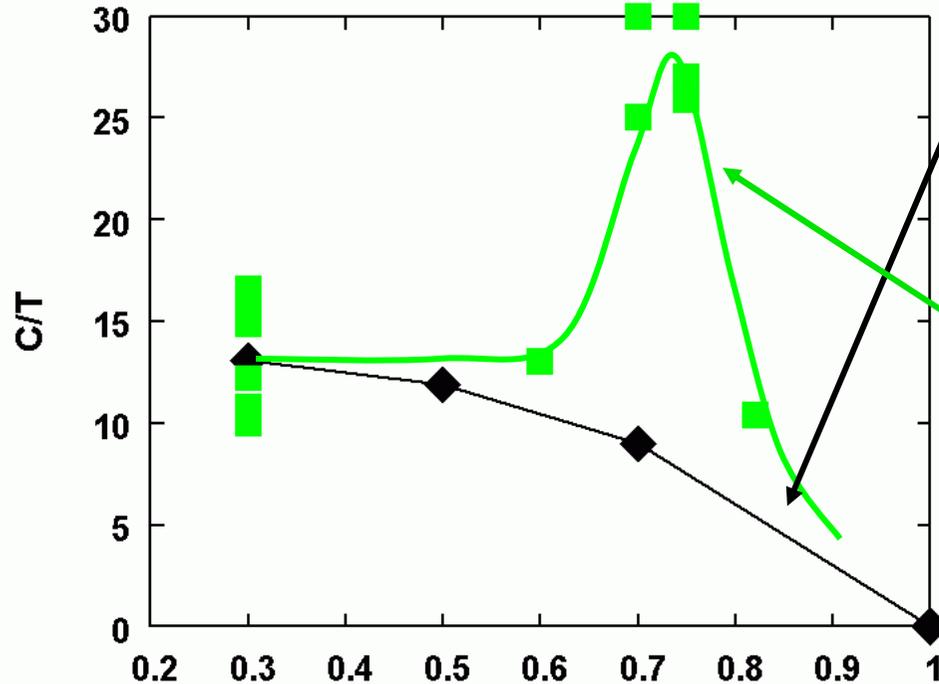
$$e_g' = (xy) + e^{\pm 2\pi i/3}(yz) + e^{\pm 4\pi i/3}(zx)$$



Two distinct Fermi surface types are predicted by calculation.



Specific heat



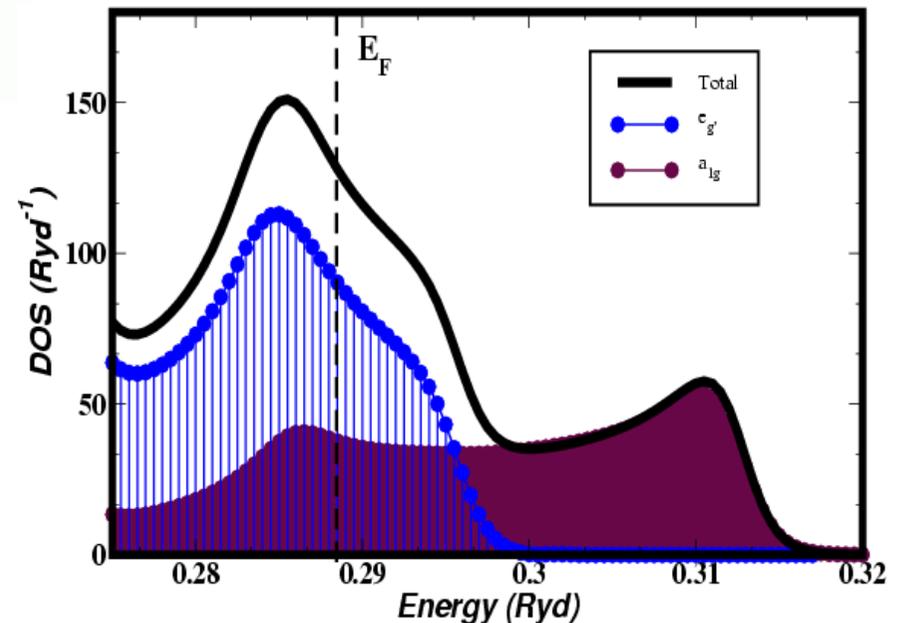
Calculations (not very accurate) show gradual decay towards (hypothetical) NaCoO_2

Exp. (not very reliable) show a strong enhancement near $x=0.7$

At $x=0.3$, 70% of DOS comes from the small e' pockets!

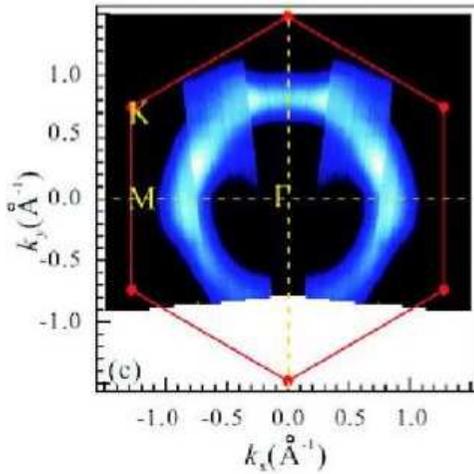
Practically all data at $x < 0.5$ are consistent with weak or no correlations

Practically all data at $x > 0.5$ are consistent with strong correlations

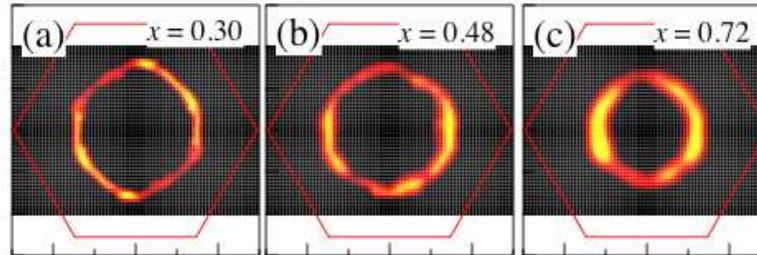




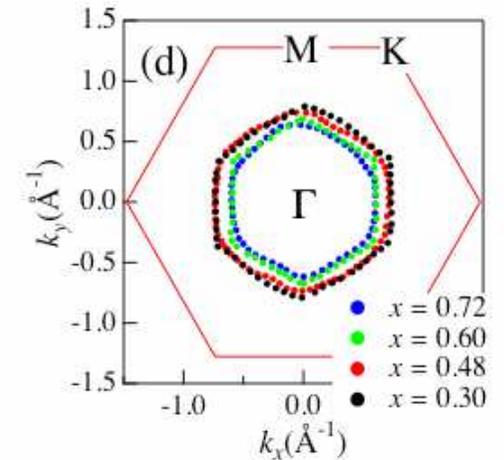
Comparison with Experiment



M.Z. Hasan *et al*



H. B. Yang *et al*



The large (a_{1g}) Fermi Surface is clearly seen by ARPES

The smaller (e_g') surfaces are absent

WHY?

- Correlations beyond LDA
- Surface effects (relaxation, surface bands, Na content)
- Matrix elements



How does correlation affect the electronic structure?

Strongly correlated systems are characterized by large U/t

What is U in Na_xCoO_2 ?

LMTO: 3.7 eV (for all 5 d-bands)

Narrow t_{2g} bands screened by
Empty e_g orbitals ... $U < 3.7\text{eV}$

(A.Liebsch)

LDA+U: Corrects on-site Coulomb repulsion

Gets good FS match for $U = 4\text{ eV}$ (P. Zhang, PRL 93 236402)

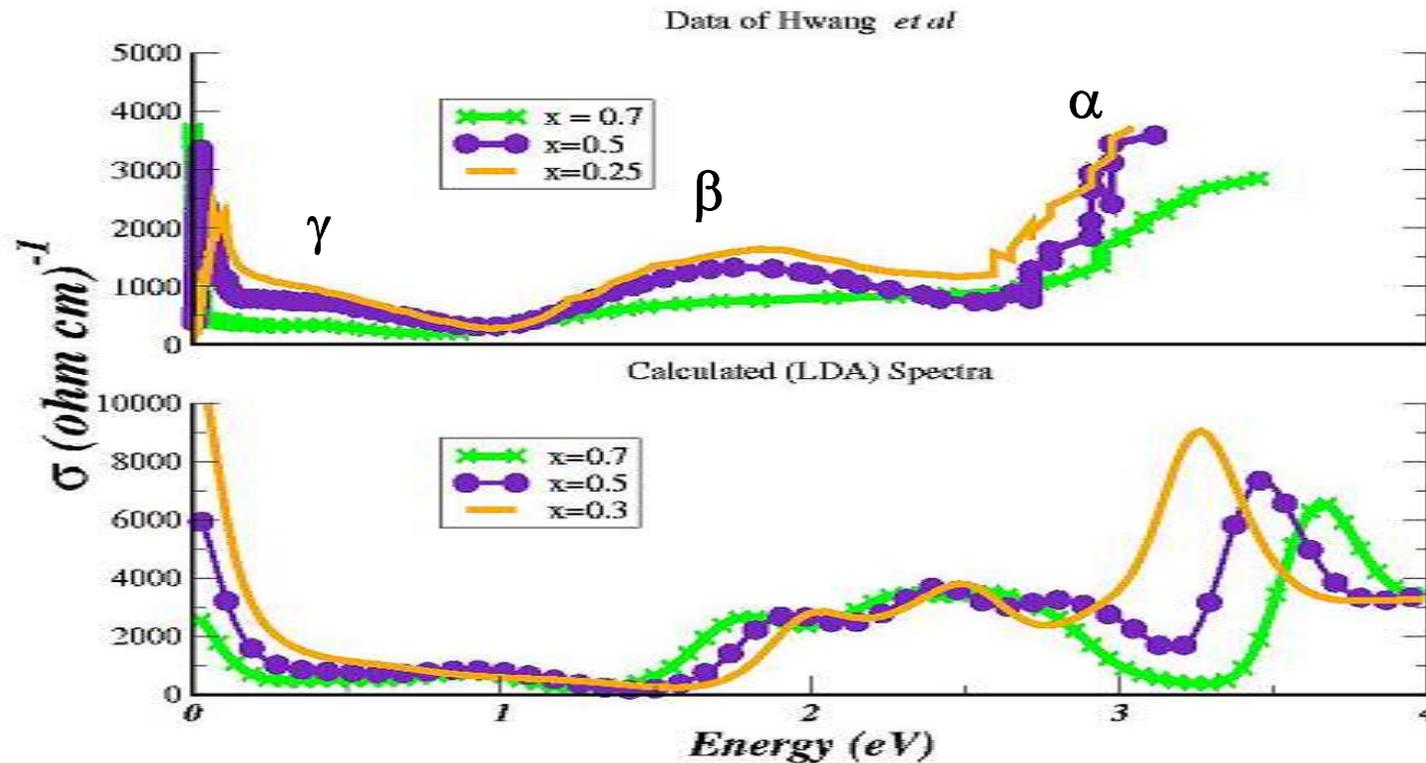
But $U = 4\text{ eV} > U_C = 3\text{ eV}$ for unobserved charge disproportionation
(K-W. Lee PRL 94 026403)

For $U < 2.5\text{ eV}$, small pockets remain

Spin fluctuations: Renormalize bands, similarly to phonons
Fermi surface is preserved, less weight



Optics: A Probe of Bulk Electronic Structure



There are three basic peaks: α , β , γ .

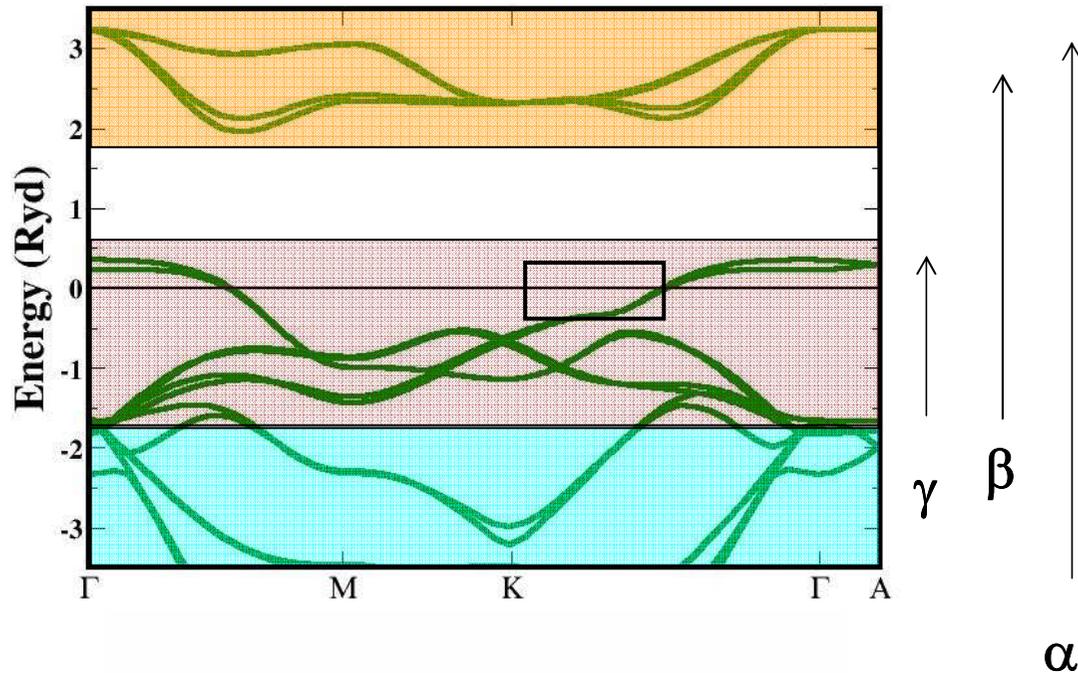
Peak shifts with changing Na content are reproduced.

Peak heights and energy positions are exaggerated.



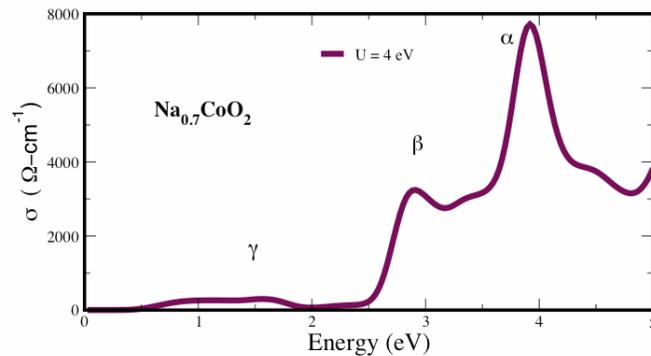
Optics: Effect of LDA+U

How does electronic correlation manifest itself?



Application of LDA+U
worsens agreement with
experiment.

Mott-Hubbard type
correlation is not exhibited
for any x!





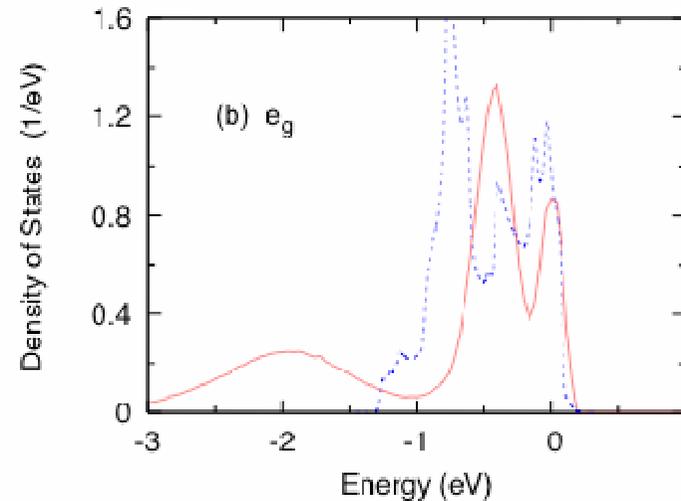
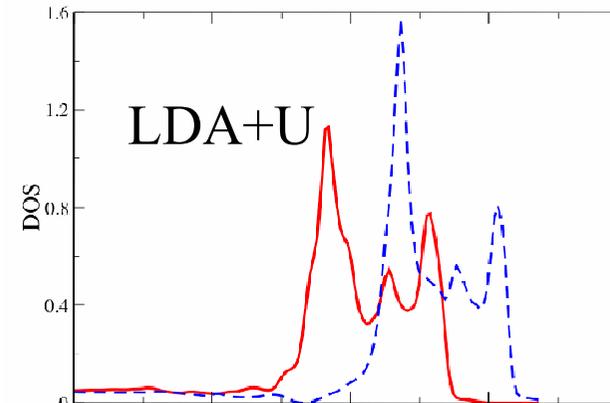
Dynamical Correlation: DMFT

Dynamical Mean Field Theory gives a very different picture of correlation effects:



Small e_g holes grow

A.Liebsch, '05

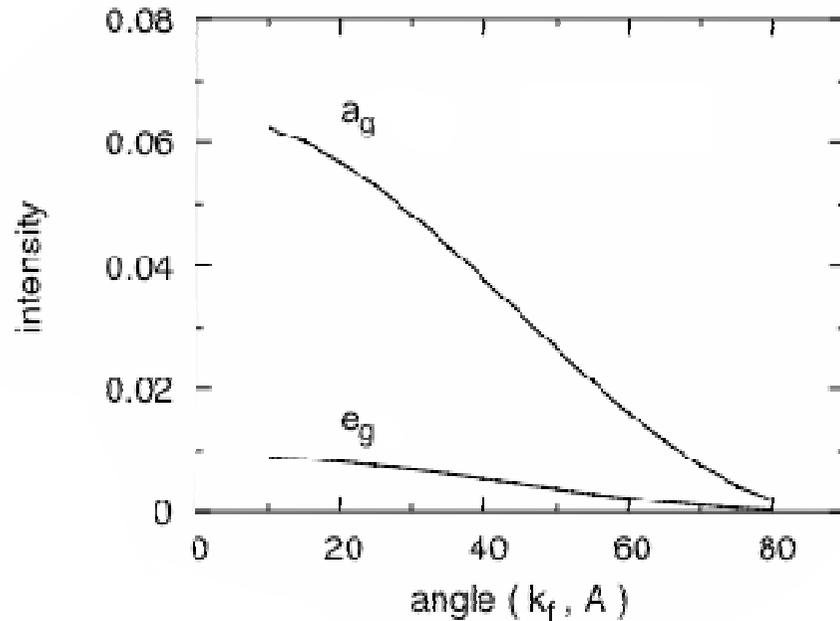


Some spectral weight shifts downward



Matrix elements and surface effects

ARPES measurements have either s or p polarized light



For p-polarized light, the dipole matrix element is substantially different for the two symmetries.

(A. Liebsch)

Bulk calculations suggest that surface relaxation of O ions could diminish or eliminate small FS pockets



Summary of Part I

- Na_xCoO_2 has an unusual magnetic phase diagram
- The system does not behave as a Mott-Hubbard insulator, despite a rather narrow t_{2g} bandwidth
- The LDA+U method *worsens* agreement with optical measurements
- Dynamical correlations show weight transfer from $a_{1g} \rightarrow e_g'$ *i.e.* holes grow!
- Calculations, in conjunction with experiment, suggest the presence of spin fluctuations



Part II: Superconductivity

What kind of superconductor is $\text{Na}_{0.35}\text{CoO}_2 \cdot y\text{H}_2\text{O}$?

Pairing state: Singlet? Triplet?

Order parameter: $s, p, d, f \dots$?



Experimental evidence for pairing state

...**singlet** order parameter with **s-wave symmetry** is realized in $\text{Na}_x\text{CoO}_2 \cdot y\text{H}_2\text{O}$ - *JPSJ* 72, 2453 (2003)

...an **unconventional superconducting symmetry** with line nodes - *cond-mat/0410517* (2004)

Unconventional superconductivity in $\text{Na}_x\text{CoO}_2 \cdot y\text{H}_2\text{O}$ - *cond-mat/0408426* (2004)

Possible singlet to triplet pairing transition in $\text{Na}_x\text{CoO}_2 \cdot y\text{H}_2\text{O}$ - *PR B70*, 144516 (2005)

Possible unconventional superconductivity in $\text{Na}_x\text{CoO}_2 \cdot y\text{H}_2\text{O}$ probed by muon spin rotation and relaxation - *PR B70*, 13458 (2005)

Evidence of **nodal superconductivity** in $\text{Na}_{0.35}\text{CoO}_2 \cdot 1.3\text{H}_2\text{O}$ - *PR B71*, 20504 (2005)

...**magnetic fluctuations** play an important role in the occurrence of superconductivity in $\text{Na}_x\text{CoO}_2 \cdot y\text{H}_2\text{O}$ - *JPSJ* 74 (2005)

... superconducting electron pairs are in the **singlet state** - *JPSJ* 74 (2005)

... superconducting Na_xCoO_2 a clear candidate for **d-wave pairing** - *cond-mat/0503010* (2005)



What pairing states can we exclude?

~~$\{1, k_x^2 + k_y^2, k_z^2\}$~~

~~$\begin{cases} k_x k_z \\ k_y k_z \\ k_z(k_x + ik_y) \end{cases}$~~

~~$\begin{cases} k_x^2 - k_y^2 \\ k_x k_y \\ (k_x + ik_y)^2 \end{cases}$~~

$\begin{cases} \hat{x}k_x + \hat{y}k_y, \hat{z}k_z \\ \hat{x}k_x - \hat{y}k_y \\ \hat{z}k_x(k_x^2 - 3k_y^2), k_z[(k_x^2 - k_y^2)\hat{x} - 2k_x k_y \hat{y}] \\ \hat{z}k_y(k_y^2 - 3k_x^2), k_z[(k_y^2 - k_x^2)\hat{y} - 2k_x k_y \hat{x}] \end{cases}$

~~$\begin{cases} \hat{x}k_z, \hat{z}k_x \\ \hat{y}k_z, \hat{z}k_y \\ \hat{z}(k_x + ik_y), k_z(\hat{x} + i\hat{y}) \end{cases}$~~

~~$\begin{cases} \hat{x}k_x - \hat{y}k_y \\ \hat{x}k_y + \hat{y}k_x \\ (\hat{x} + i\hat{y})(k_x + ik_y) \end{cases}$~~

μSR

- No static magnetic moments
- » **No states with L≠0**
- » **No non-unitary triplet states**

MP 63 240 (1991)

Two dimensionality

- c/a ratio ~ 3.5
- $\rho_{ab}/\rho_c \sim 10^3$
- » **k_z-dependent order parameter**
- unrealistic**

DOS Probes

- Non-exponential decay of C/T vs. T
- No coherence peak in 1/T₁
- Non-exponential decay of relaxation time
- » **Superconducting state not fully gapped**



How can pairing state be further resolved?

All remaining states
are triplet f

f states

$$\hat{z}k_x (k_x^2 - 3k_y^2)$$
$$\hat{z}k_y (k_y^2 - 3k_x^2)$$

Both f states are axial

Knight Shift can distinguish:

- Spin direction is \perp to vector order parameter
- KS constant across T_C for planar spins (axial order parameter)
- KS decreases across T_C for axial spins (planar order parameter)

Presently, results are contradictory



Evidence of Spin Fluctuations in $\text{Na}_{0.35}\text{CoO}_2 \cdot 1.4\text{H}_2\text{O}$

There is growing evidence that SF have a role in the superconductivity:

- Curie-Weiss like behavior of $1/T_1$ (above T_C), with negative θ
- Correlation of T_C with magnetic fluctuations as measured by NQR
- Direct neutron observation of spin fluctuations in related compounds
- LDA calculations indicate proximity to quantum critical point

Details of pairing/pair-breaking in a particular system depend on:

- i) Fermiology
- ii) spin fluctuation spectrum - $\text{Im}\chi(\mathbf{q},\omega)$