



Underdoped cuprates phenomenology in the 2D Hubbard model within COM(SCBA)

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CORPES '07 – April 18, 2007



Outline

- 1 Composite Operator Method (COM)
 - Philosophy of COM
 - Dyson equation in COM
 - Ingredients of COM



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The simplest interacting model

$$H = -\mu \sum_{\sigma} c_{\sigma}^{\dagger} c_{\sigma} + U c_{\uparrow}^{\dagger} c_{\downarrow}^{\dagger} c_{\downarrow} c_{\uparrow}$$



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can be *formally* diagonalized as

$$H = -\mu \sum_{\sigma} \xi_{\sigma}^{\dagger} \xi_{\sigma} + \left(\frac{U}{2} - \mu \right) \sum_{\sigma} \eta_{\sigma}^{\dagger} \eta_{\sigma}$$

where the **composite operators**

$$\xi_{\sigma} = \left(1 - c_{\bar{\sigma}}^{\dagger} c_{\bar{\sigma}} \right) c_{\sigma} \quad \text{and} \quad \eta_{\sigma} = c_{\bar{\sigma}}^{\dagger} c_{\bar{\sigma}} c_{\sigma}$$



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are **eigenoperators** of the Hamiltonian

$$i \frac{\partial}{\partial t} \xi_{\sigma} = [\xi_{\sigma}, H] = -\mu \xi_{\sigma}$$

$$i \frac{\partial}{\partial t} \eta_{\sigma} = [\eta_{\sigma}, H] = -(\mu - U) \eta_{\sigma}$$



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satisfy **non-canonical** anti-commutation relations

$\{\xi_{\sigma}, \xi_{\sigma'}\} = \{\eta_{\sigma}, \eta_{\sigma'}\} = \{\xi_{\sigma}^{\dagger}, \eta_{\sigma'}\} = 0$	$\{\xi_{\sigma}, \eta_{\sigma'}\} = \delta_{\sigma\bar{\sigma}'} c_{\sigma} c_{\bar{\sigma}}$
$\{\xi_{\sigma}^{\dagger}, \xi_{\sigma'}\} = \delta_{\sigma\sigma'} (1 - c_{\bar{\sigma}}^{\dagger} c_{\bar{\sigma}}) + \delta_{\sigma\bar{\sigma}'} c_{\sigma}^{\dagger} c_{\bar{\sigma}}$	$\{\xi_{\sigma}^{\dagger} \xi_{\sigma}, \xi_{\sigma'}\} = -\delta_{\sigma\sigma'} \xi_{\sigma}$
$\{\eta_{\sigma}^{\dagger}, \eta_{\sigma'}\} = \delta_{\sigma\sigma'} c_{\bar{\sigma}}^{\dagger} c_{\bar{\sigma}} - \delta_{\sigma\bar{\sigma}'} c_{\sigma}^{\dagger} c_{\bar{\sigma}}$	$\{\eta_{\sigma}^{\dagger} \eta_{\sigma}, \eta_{\sigma'}\} = -\eta_{\sigma'}$
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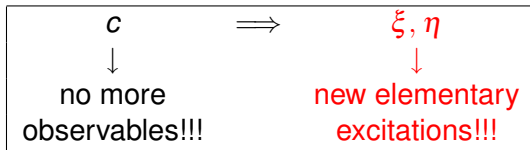
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are the new elementary excitations





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$$H = H[\varphi] \quad \varphi: \text{bare particles}$$



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$$m = \langle \{ J, \psi^\dagger \} \rangle$$

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where the **self-energy** Σ reads as

$$\Sigma = \langle \langle \delta J | \delta J^\dagger \rangle \rangle_{\text{irr}}$$



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Unknowns

- 1 *higher-order correlators* in
 - ε (*mean fields*)
 - l (*commutations*)
- 2 *residual* self-energy Σ



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- 1 Quest and use of a **composite operatorial basis**:
 - **higher-order fields** emerging from the equations of motion
 - **eigenoperators** of relevant **interacting terms**
 - **eigenoperators** of the problem reduced to a **small cluster**



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Mancini & Avella; Adv. Phys. **53**, 537 (2004)



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The basis: $\psi = \begin{pmatrix} \xi \\ \eta \end{pmatrix}$ where $\xi = (1 - n)c$ and $\eta = nc$



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1 $n = \langle c^\dagger c \rangle$



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2 $\Delta = \langle \xi^\alpha \xi^\dagger \rangle - \langle \eta^\alpha \eta^\dagger \rangle$

The constraints

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2 $\Delta = 0$ Hubbard I



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The unknowns

- 1 μ
- 2 $\Delta = \langle \xi^\alpha \xi^\dagger \rangle - \langle \eta^\alpha \eta^\dagger \rangle$
- 3 $p = \frac{1}{4} (\langle n^\alpha n \rangle + \langle \mathbf{s}^\alpha \mathbf{s} \rangle) - \langle \lambda^\alpha \lambda^\dagger \rangle$

where $\lambda = c_\uparrow(i)c_\downarrow(i)$

The constraints

- 1 $n = \langle c^\dagger c \rangle$
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- 3 $p \Leftarrow$ 1-loop Roth



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The unknowns

- 1 μ
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- 3 $\rho = \frac{1}{4} (\langle n^\alpha n \rangle + \langle \mathbf{s}^\alpha \mathbf{s} \rangle) - \langle \lambda^\alpha \lambda^\dagger \rangle$

where $\lambda = c_\uparrow(i)c_\downarrow(i)$

The Algebra constraints

- 1 $n = \langle c^\dagger c \rangle$
- 2 $\Delta = \langle \xi^\alpha \xi^\dagger \rangle - \langle \eta^\alpha \eta^\dagger \rangle$
- 3 $\langle \xi \eta^\dagger \rangle = 0$



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$$① \langle \psi^{\alpha_{1,\sqrt{2},2,\sqrt{5},2\sqrt{2}}} \psi^\dagger \rangle$$

The constraints

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- 1 $\langle \psi^{\alpha_{1,\sqrt{2},2,\sqrt{5},2\sqrt{2}}} \psi^\dagger \rangle$
- 2 $\mathbf{a}_\mu = \mathbf{a}_\mu [\langle c^\dagger c^\alpha n \rangle, \langle c^\dagger c^\alpha \sigma_v c^\alpha n_v \rangle, \langle c^\dagger c^\alpha c^\dagger c^\dagger \rangle]$

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- 2 $\langle n_\mu n_\mu \rangle = n + 2(2\delta_{\mu 0} - 1)D$



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- 3 b_μ, c_μ, d_μ

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- 1 $\langle \psi^{\alpha_{1,\sqrt{2},2,\sqrt{5},2\sqrt{2}}} \psi^\dagger \rangle$
- 2 $\langle n_\mu n_\mu \rangle = n + 2(2\delta_{\mu 0} - 1)D$
- 3 $\lim_{|\mathbf{k}| \rightarrow 0} \omega(\mathbf{k}) \propto |\mathbf{k}|^\nu$



The basis: $\psi_\mu = \begin{pmatrix} n_\mu \\ \rho_\mu \end{pmatrix}$ where $n_\mu = (n, \mathbf{s})$ and $-t\rho_\mu = i\frac{\partial}{\partial t}n_\mu$

The approximation: $\delta J_\mu \approx 0 \Rightarrow \Sigma_\mu \approx 0$

The unknowns

- 1 $\langle \psi^{\alpha_{1,\sqrt{2},2,\sqrt{5},2\sqrt{2}}} \psi^\dagger \rangle$
- 2 $\mathbf{a}_\mu = \mathbf{a}_\mu [\langle c^{\dagger\alpha} c^\alpha n \rangle, \langle c^{\dagger\alpha} c_\nu c^\alpha n_\nu \rangle, \langle c_\dagger^\alpha c_\dagger^\alpha c_\dagger^\alpha c_\dagger^\alpha \rangle]$
- 3 b_μ, c_μ, d_μ

The constraints

- 1 $\langle \psi^{\alpha_{1,\sqrt{2},2,\sqrt{5},2\sqrt{2}}} \psi^\dagger \rangle$
- 2 $\langle n_\mu n_\mu \rangle = n + 2(2\delta_{\mu 0} - 1)D$
- 3 $\lim_{|\mathbf{k}| \rightarrow 0} \omega(\mathbf{k}) \propto |\mathbf{k}|^\nu$

Avella et al.; Phys. Rev. B **67**, 115123 (2003)



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The SCBA approximation:

$$\delta J = BF \Rightarrow \Sigma = \langle\langle BF|F^\dagger B^\dagger \rangle\rangle_{\text{irr}} \approx \langle\langle B|B^\dagger \rangle\rangle \langle\langle F|F^\dagger \rangle\rangle$$



The SCBA approximation:

$$\delta J = BF \Rightarrow \Sigma = \langle\langle BF|F^\dagger B^\dagger\rangle\rangle_{\text{irr}} \cong \langle\langle B|B^\dagger\rangle\rangle \langle\langle F|F^\dagger\rangle\rangle$$

The *bosons*: $B = n, \mathbf{s}, \lambda$



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Self-consistency Cycle

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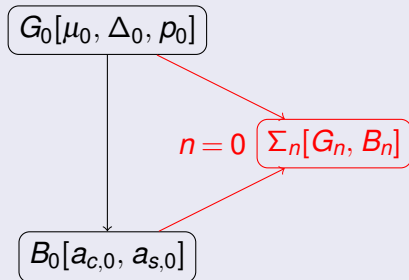
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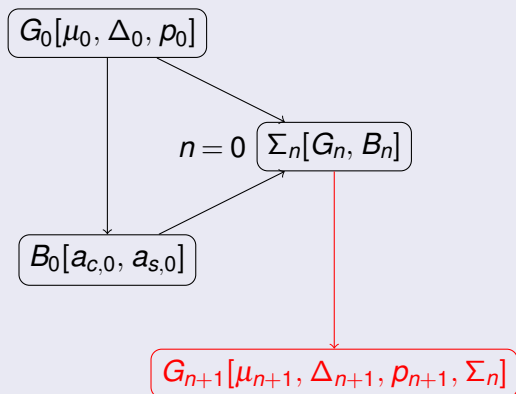


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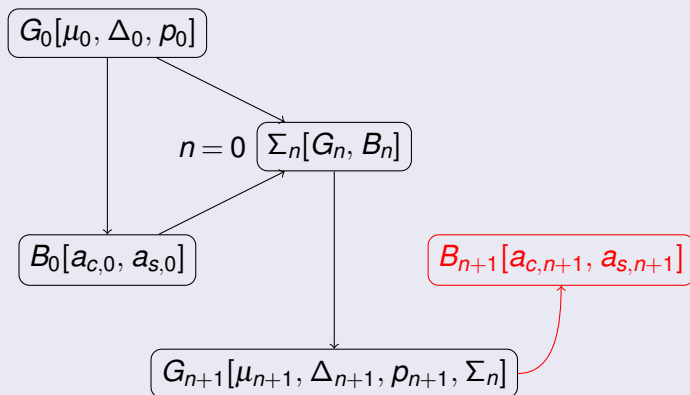


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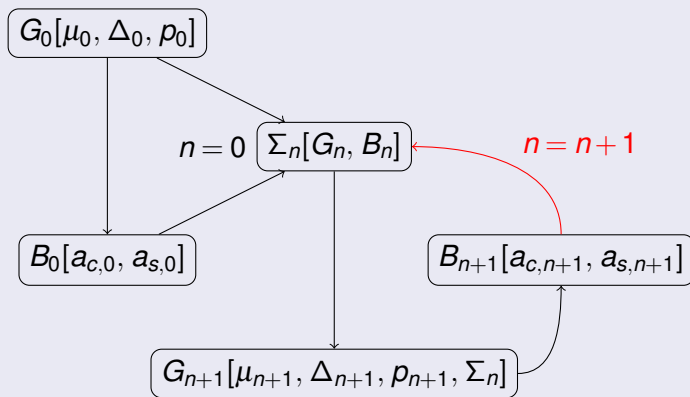


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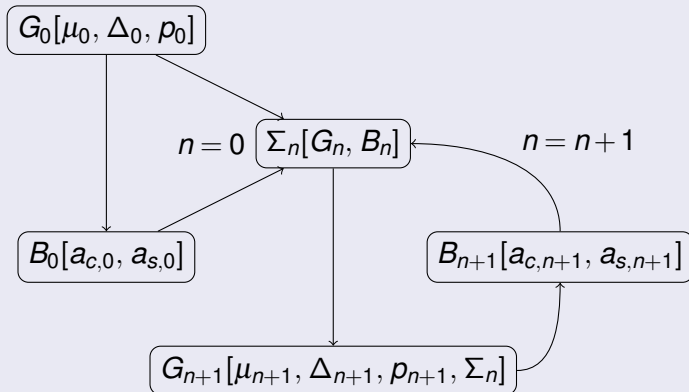


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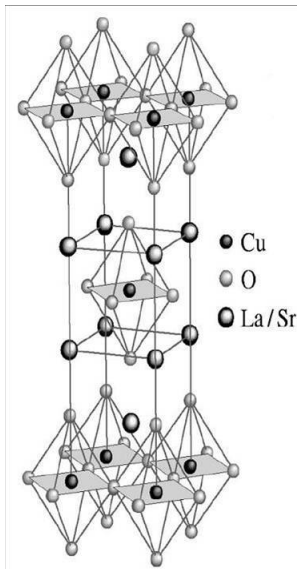
Self-consistency Cycle (Avella & Mancini; Phys. Rev. B **75**, ??? (2007))

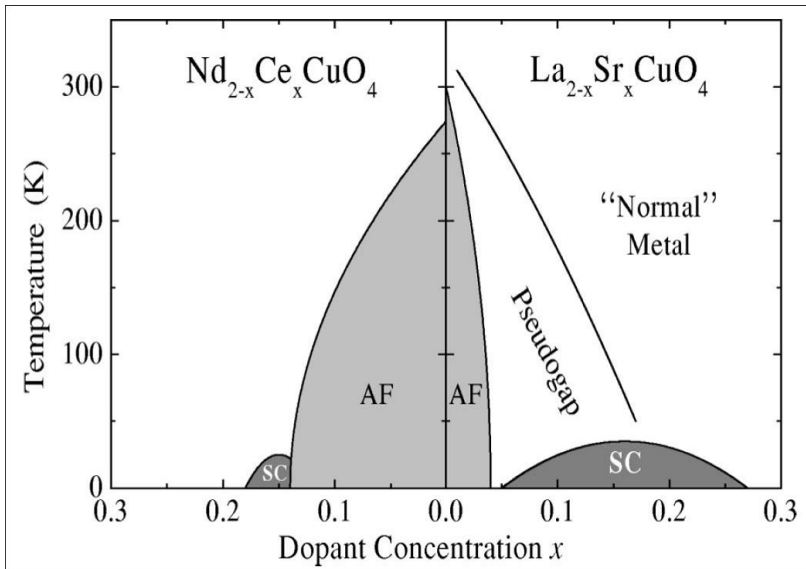


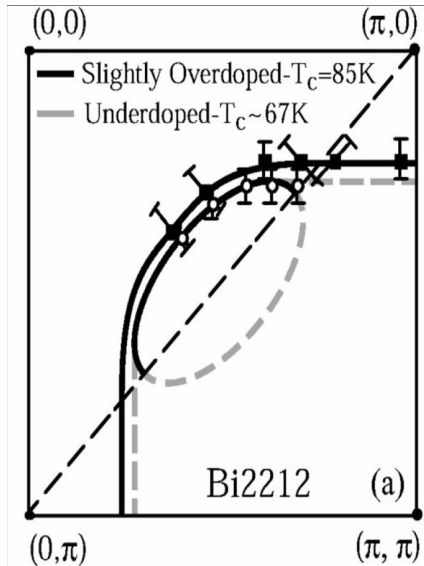


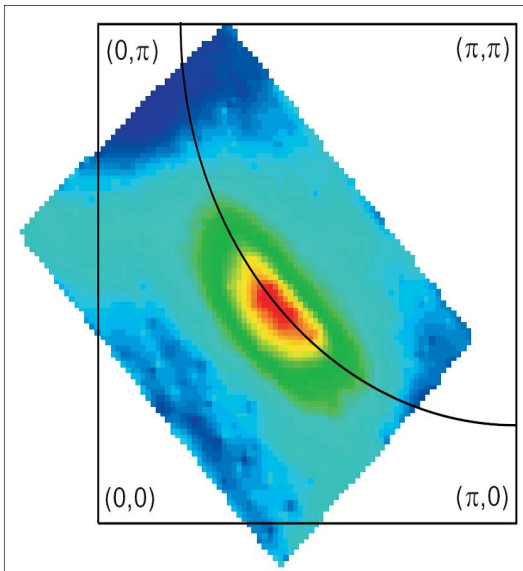
Outline

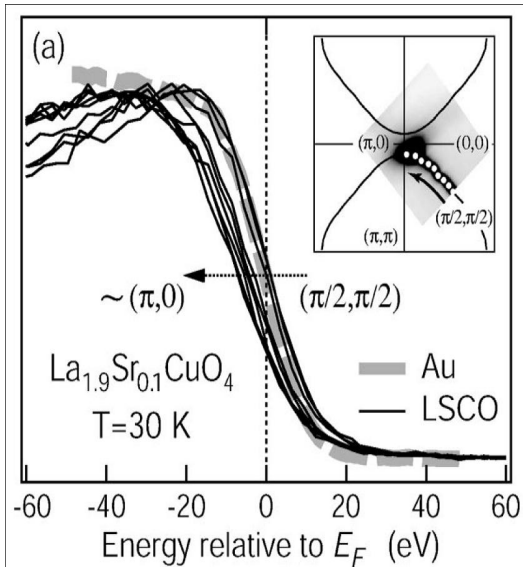
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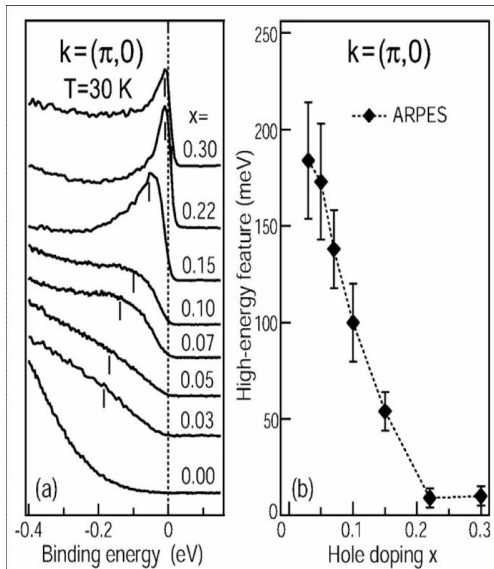


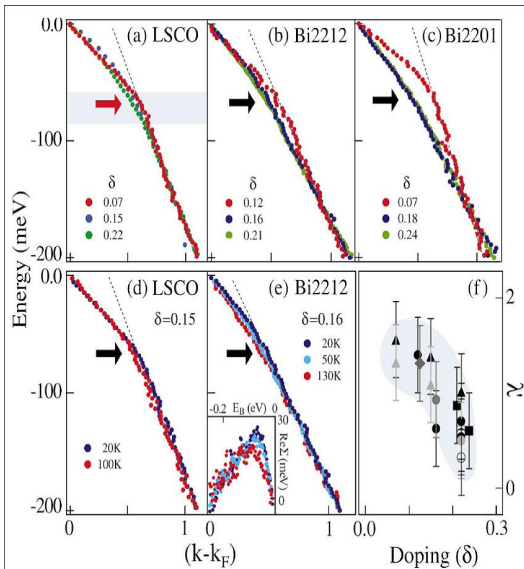


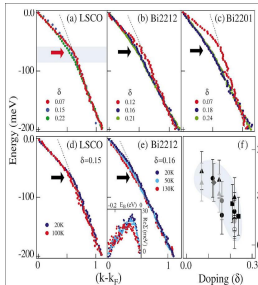
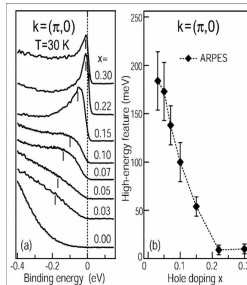
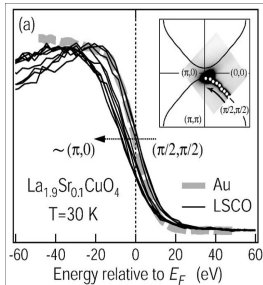
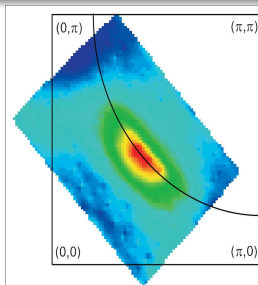
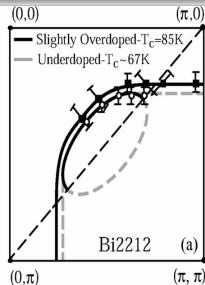
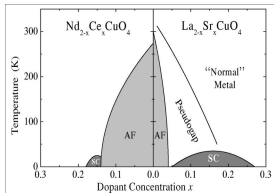








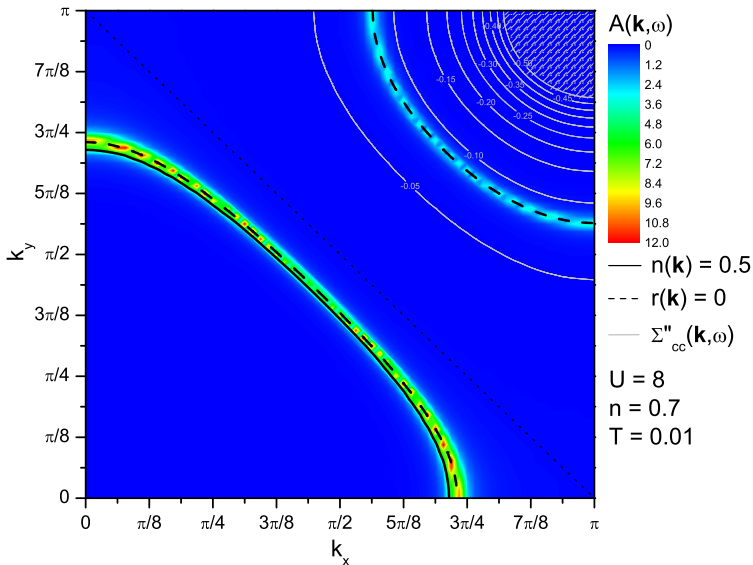


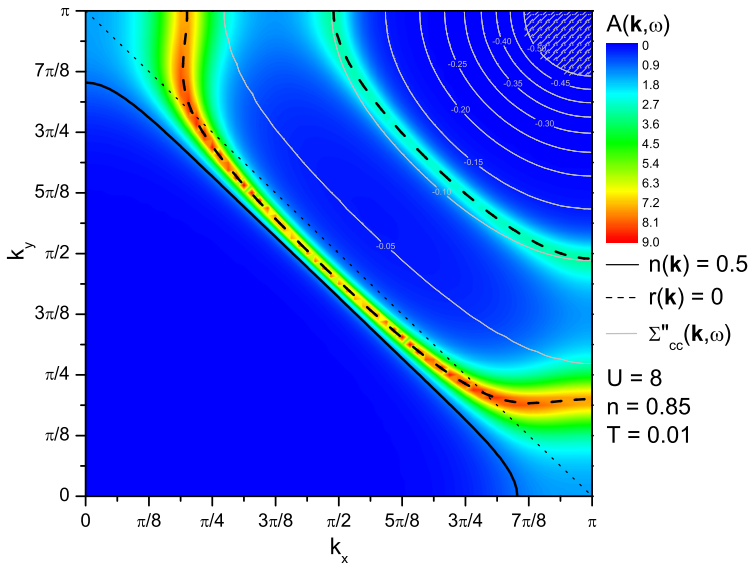


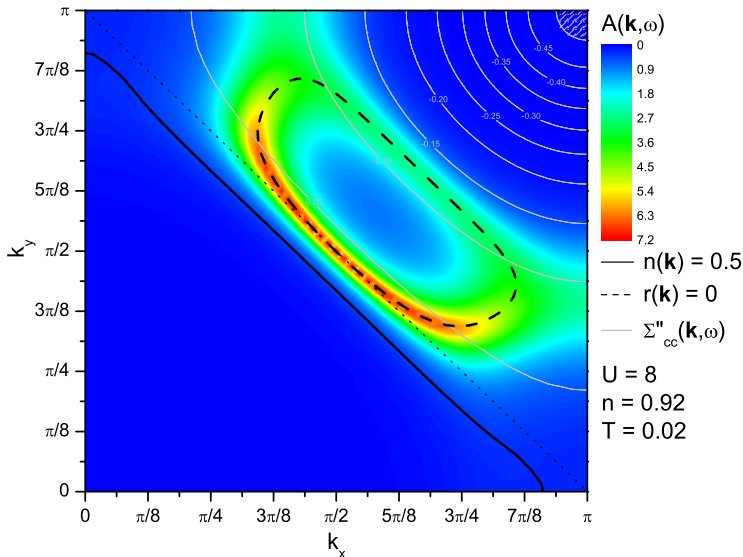


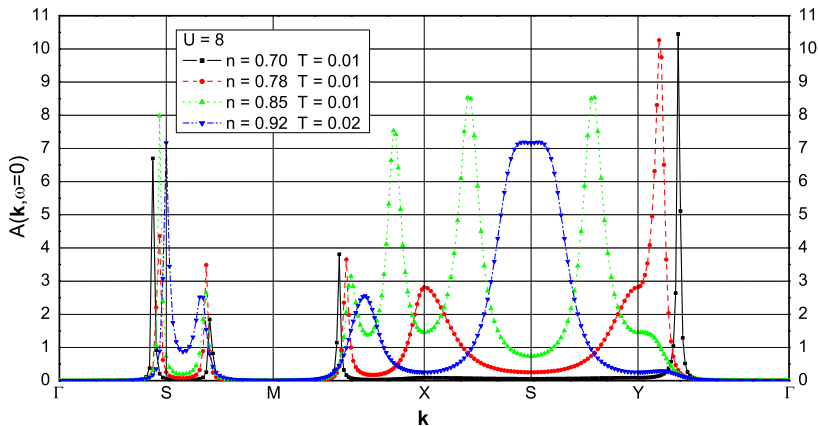
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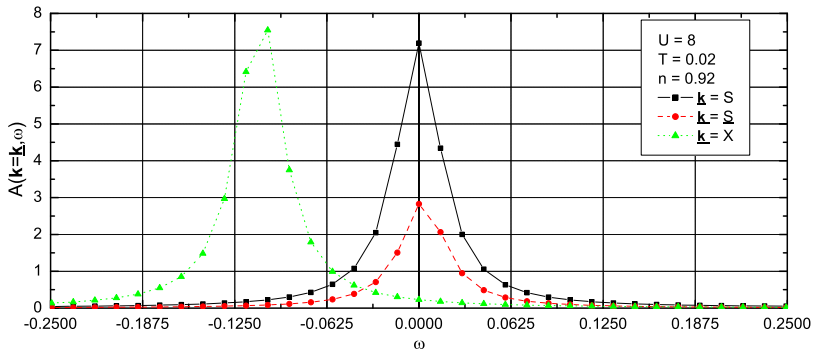
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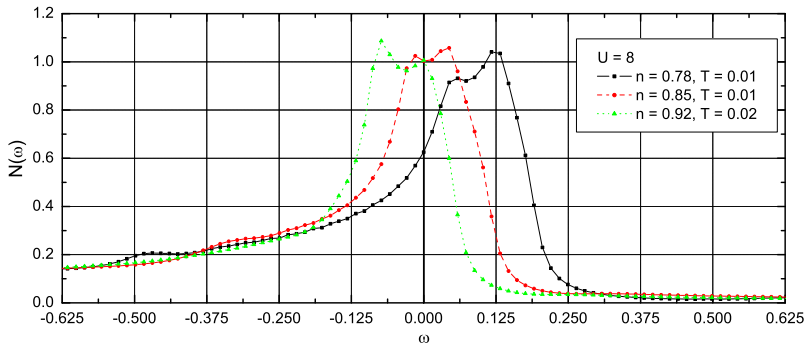






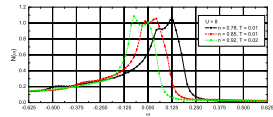
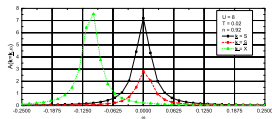
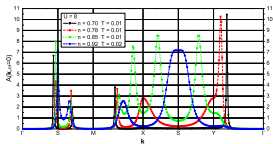
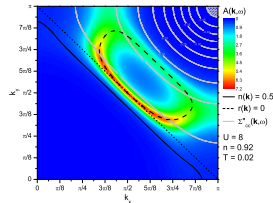
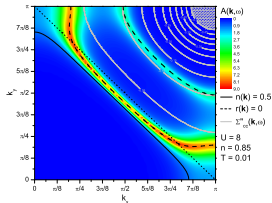
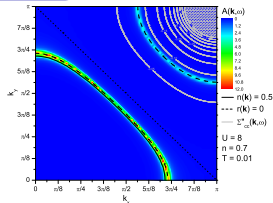








► ZOOMS





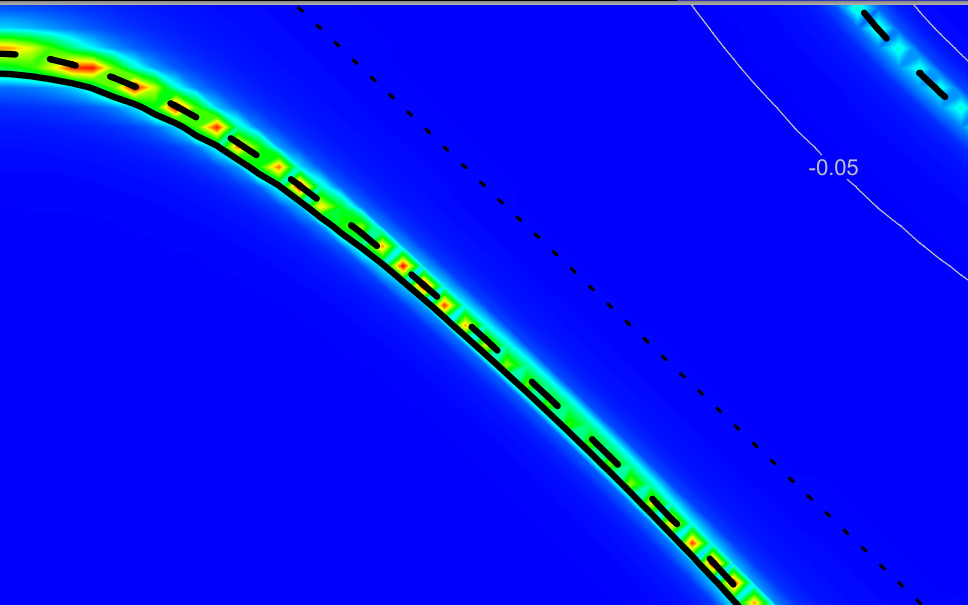
Composite Operator Method (COM)
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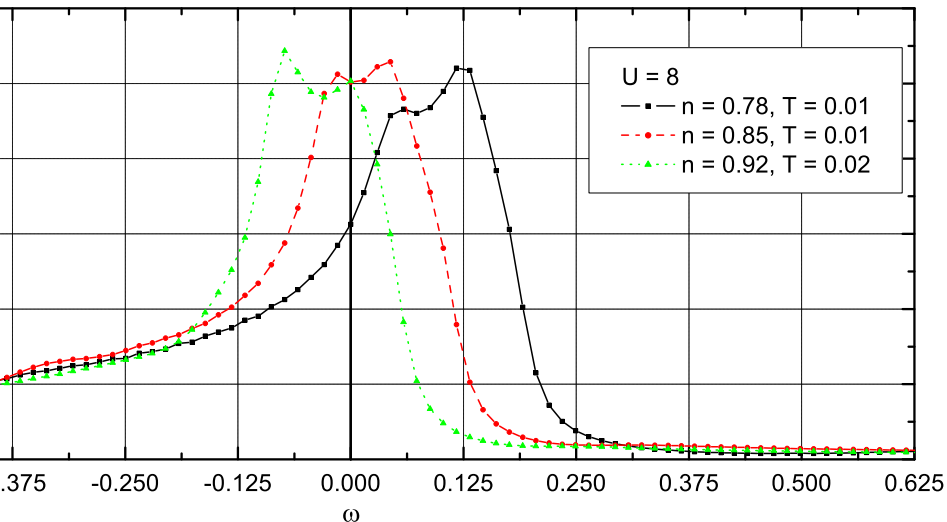
Hubbard Model
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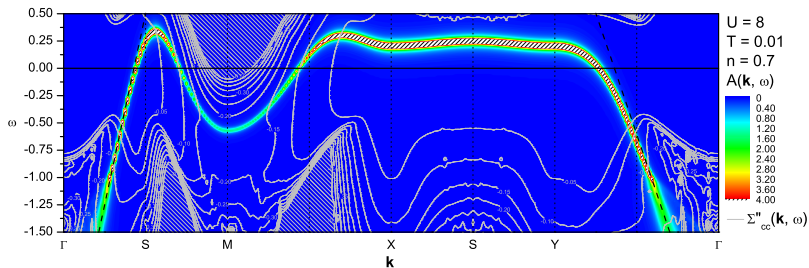
Cuprates
○○●○

Summary

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COM results









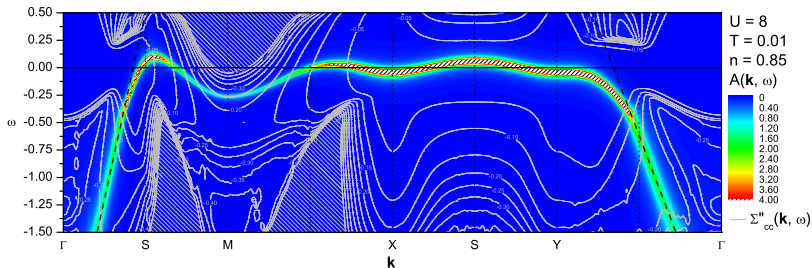
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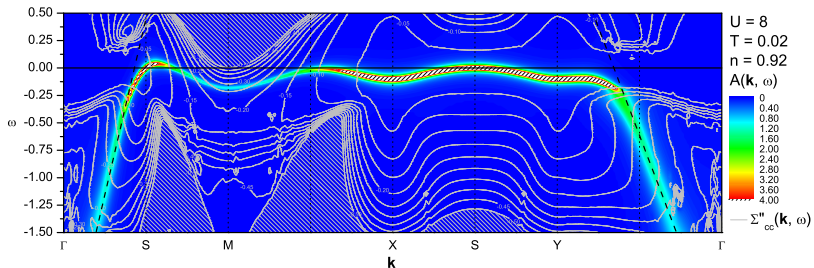
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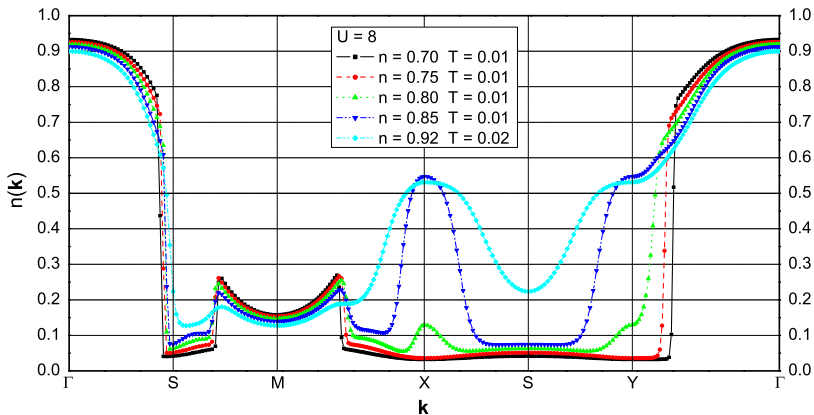
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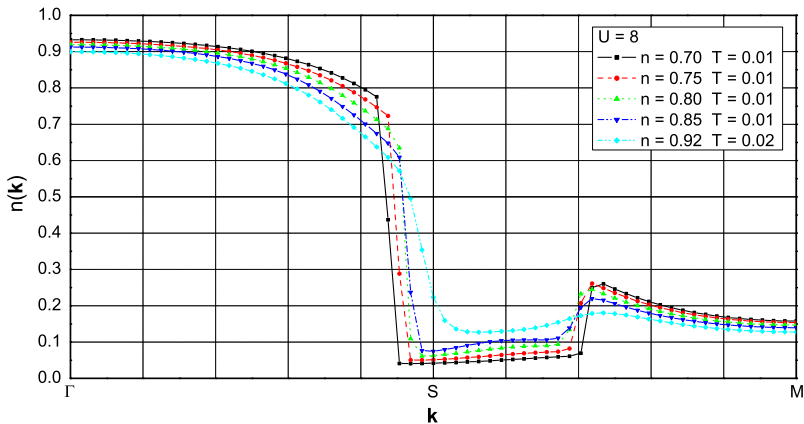
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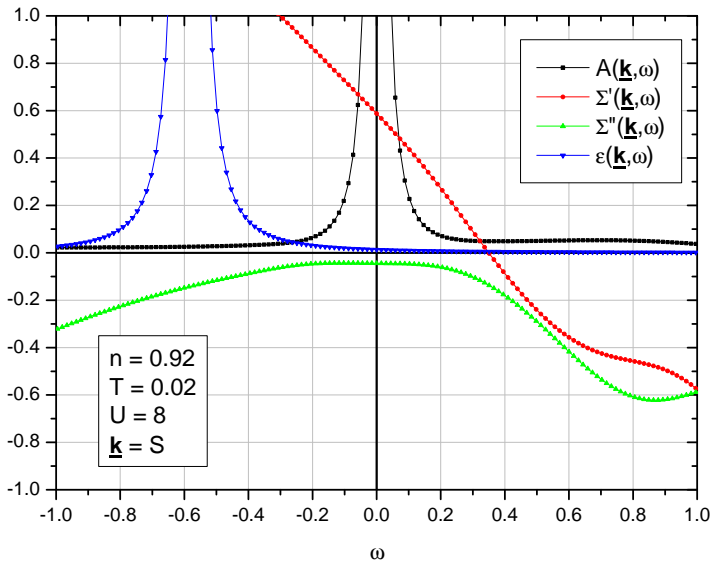
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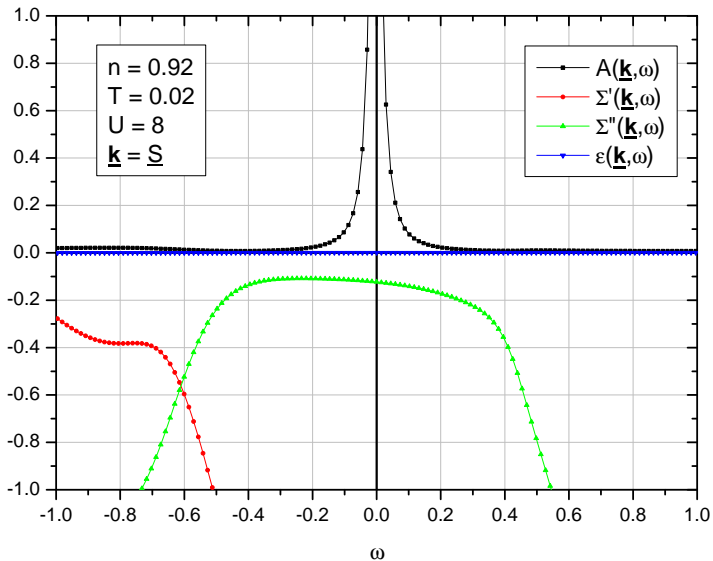


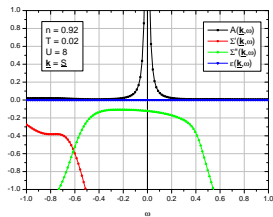
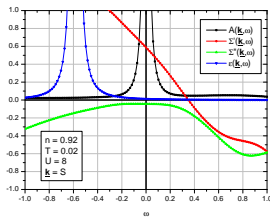
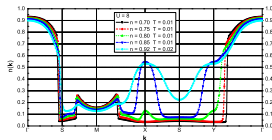
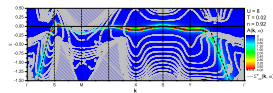
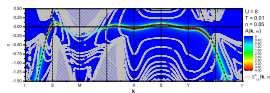
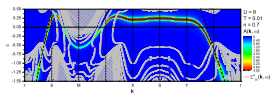














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- **A properly microscopically derived susceptibility can give results practically identical or, at least, very similar to those attainable by means of phenomenological susceptibilities specially tailored to describe experiments.**



Composite Operator Method (COM)
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Hubbard Model
○○○○○

Cuprates
○○○○

Summary



Outlook

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- Check for Luttinger sum rule



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Thank you for your attention!