

# Dynamical vertex approximation — a step beyond dynamical mean field theory

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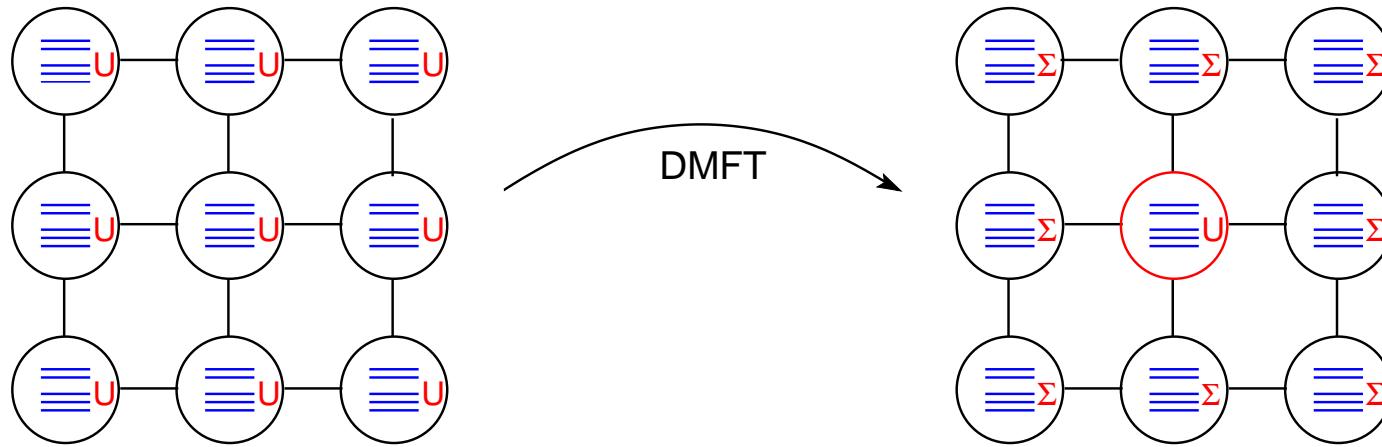
*CORPES07, April 26, 2007*

- Motivation
- Method
- Results for 3D, 2D, and 1D Hubbard model
- Conclusion and outlook

\* together with A. Toschi, A. Katanin (MPI-FKF), PRB 75, 045118 (2007)

# Motivation

## Dynamical mean field theory

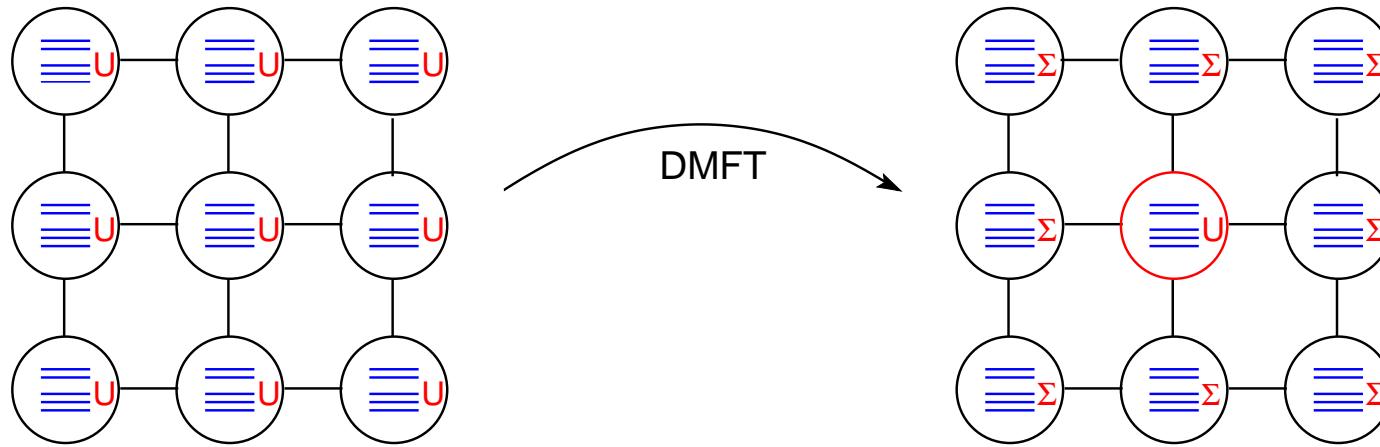


$\Sigma$  all topologically distinct, but **local** diagrams

**Success story:** quasiparticle renormalizations, magnetism, kinks ...

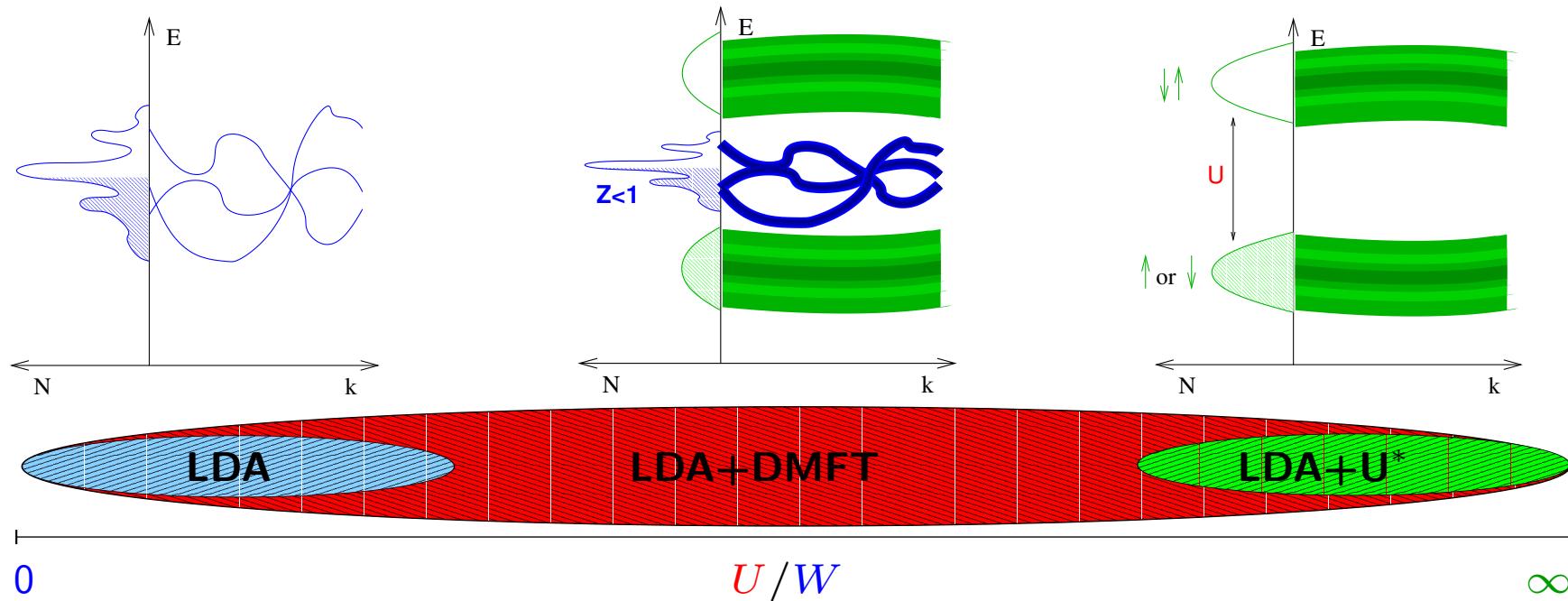
# Motivation

## Dynamical mean field theory



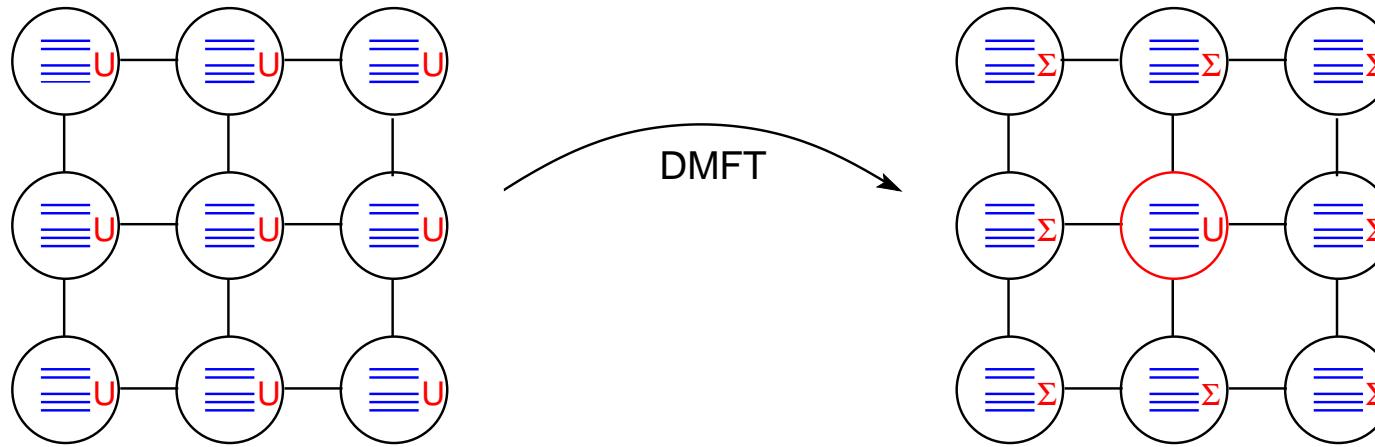
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$\Sigma$  all topologically distinct, but **local** diagrams

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**Not included:**

non-local correlations

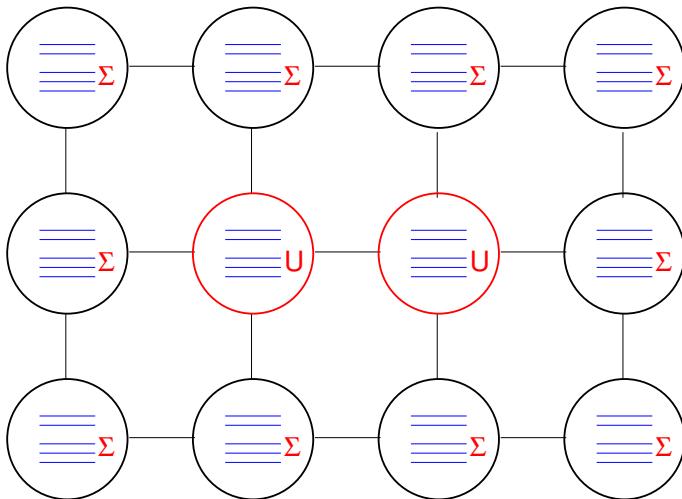
$p$ -,  $d$ -wave superconductivity, spin Peierls

magnons, (quantum) critical behavior ...

ARPES with  $k$ -dependent  $\Sigma$

# beyond DMFT

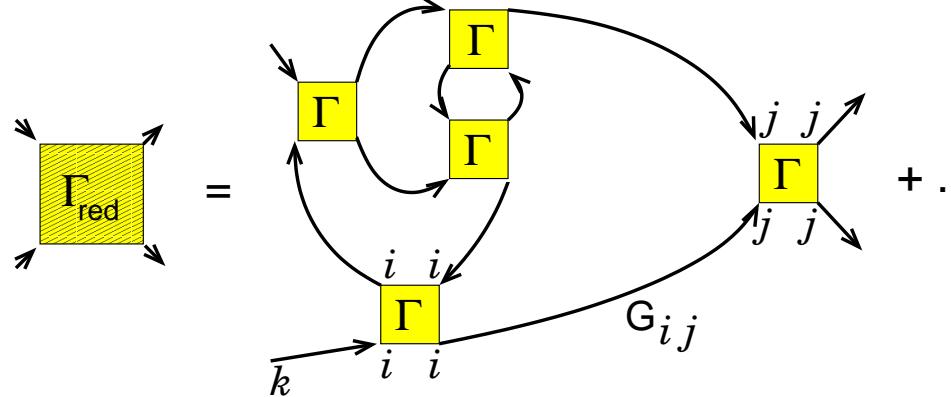
## cluster extensions of DMFT



- non-local **short-range** correlations
- $d/p$ -wave superconductivity

Hettler *et al.*'98, Lichtenstein Katsnelson'00,  
Kotliar *et al.*'01, Potthoff'03

## diagrammatic extensions of DMFT



## dynamical vertex approximation

- non-local **long-range** correlations
- (para-)magnons, phase transitions ...

Toschi, Katanin, KH cond-mat/0603100  
cf. Kusunose cond-mat/0602451  
Slezak *et al.* cond-mat/0603421

# Dynamical vertex approximation (DΓA)

**DMFT:** all (topological distinct) **local** diagram for  $\Sigma$

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$n = 2 \rightarrow$  DΓA: from **2-particle** irreducible vertex  $\Gamma$   
construct  $\Sigma$  (**local** and **non-local** diagrams)

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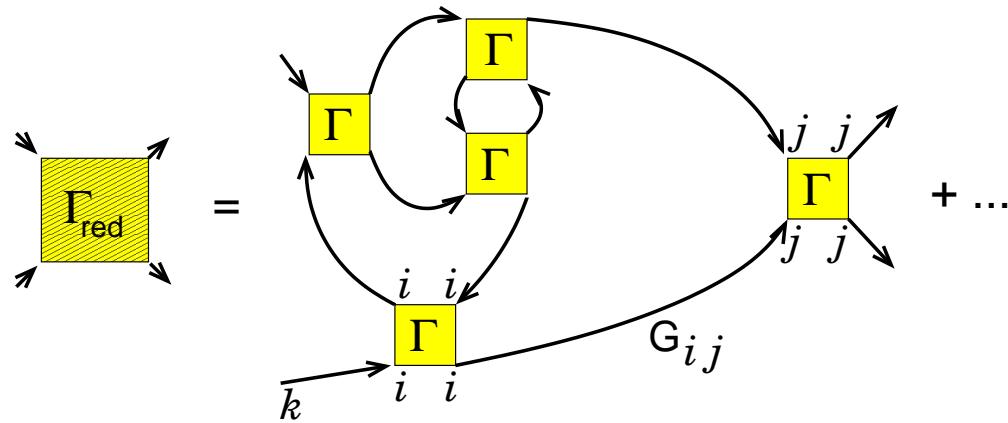
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non-local reducible vertex  $\Gamma_{\text{red}}$   
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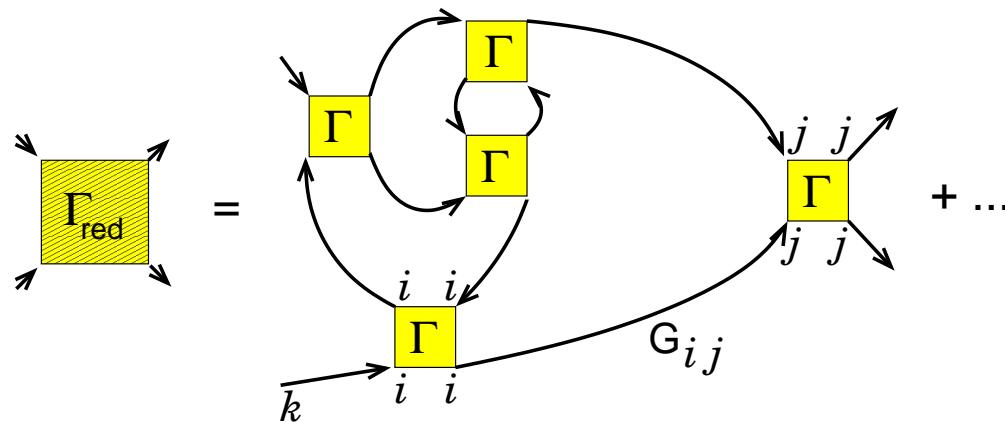
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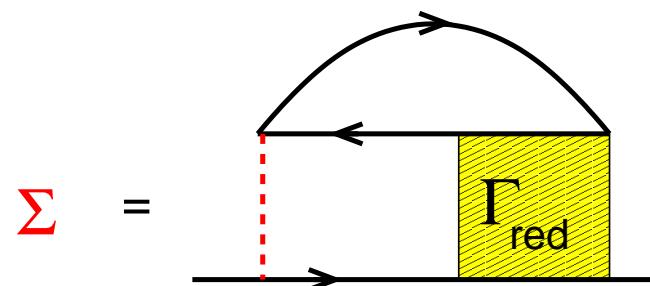
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$\Gamma_{\text{red}}$

→

non-local  $\Sigma$

exact relation (eq. of motion)

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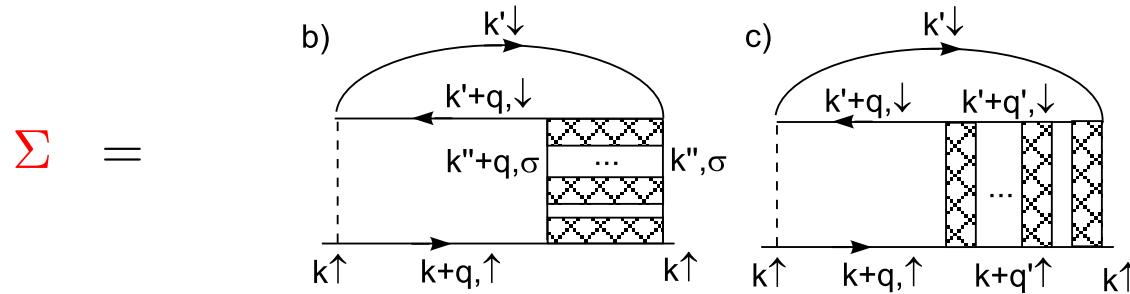
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## First step: restriction to ladder diagrams



lines: non-local  $G$

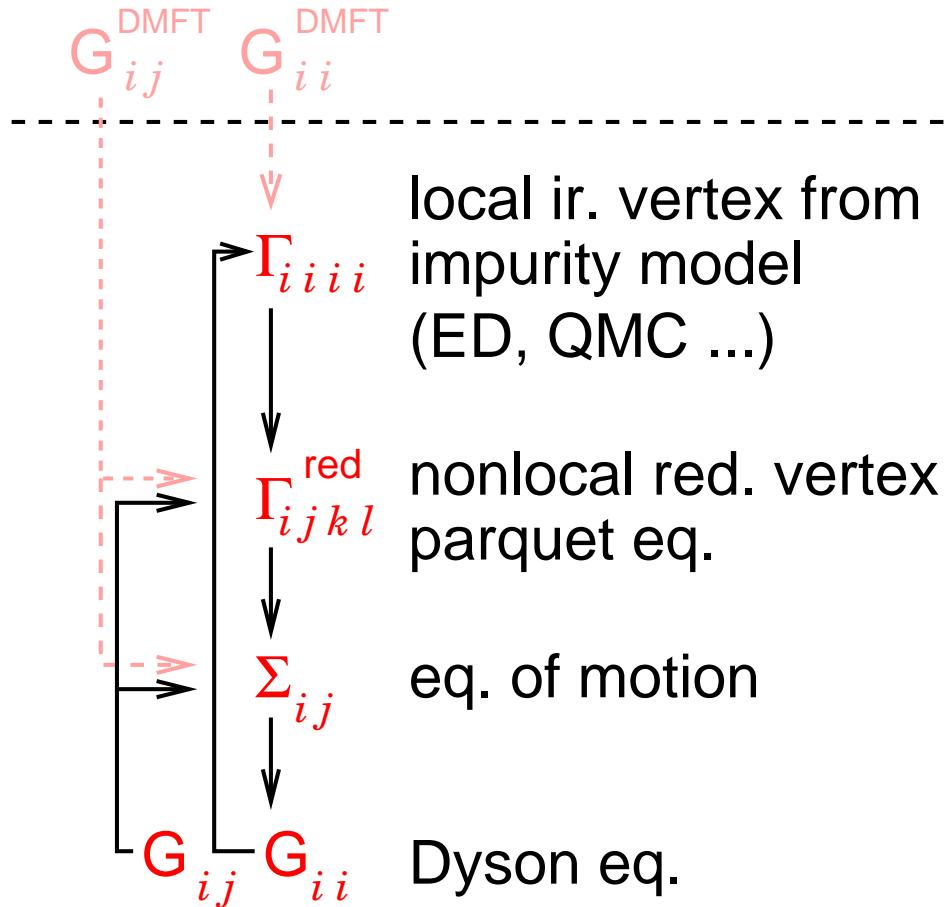
crosshatched: local irreducible vertex in spin/charge channels

$$\Gamma_{S,C}(\nu, \nu', \omega) = \chi_{0,\text{loc}}^{-1} - \chi_{S,C}^{-1}$$

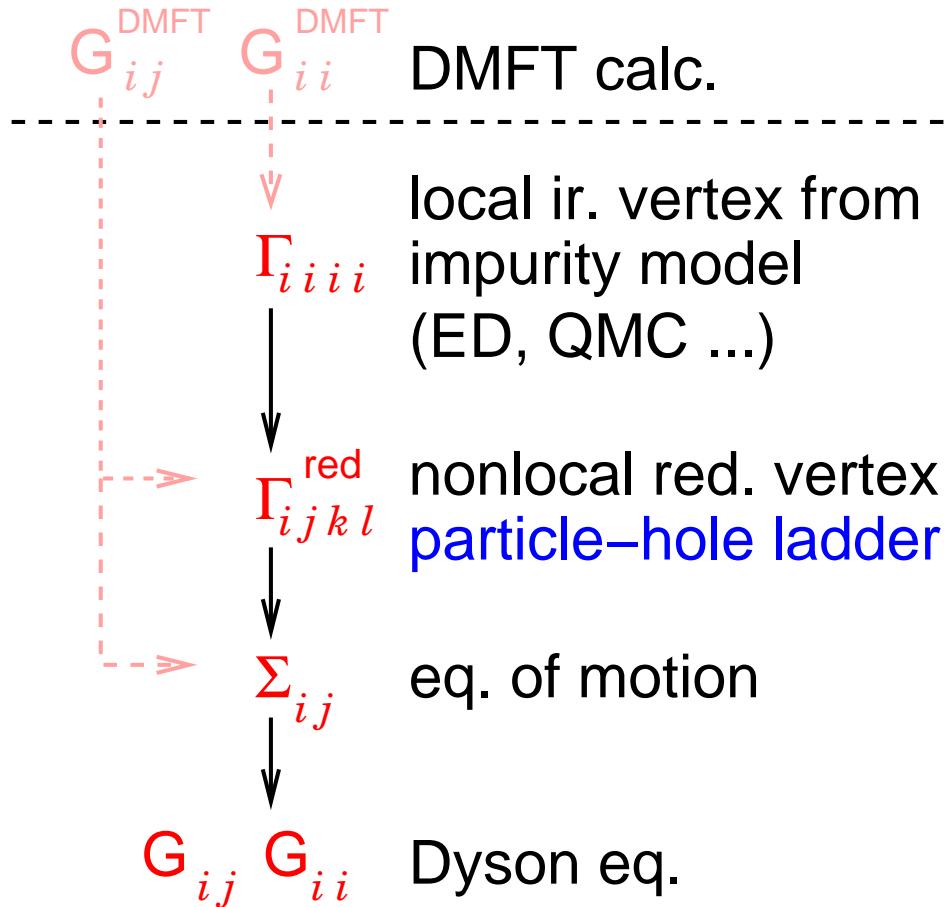
magnons, spin-fluctuations at (A)FM phase transition

$G_{ij}$  from DMFT

# D $\Gamma$ A algorithm



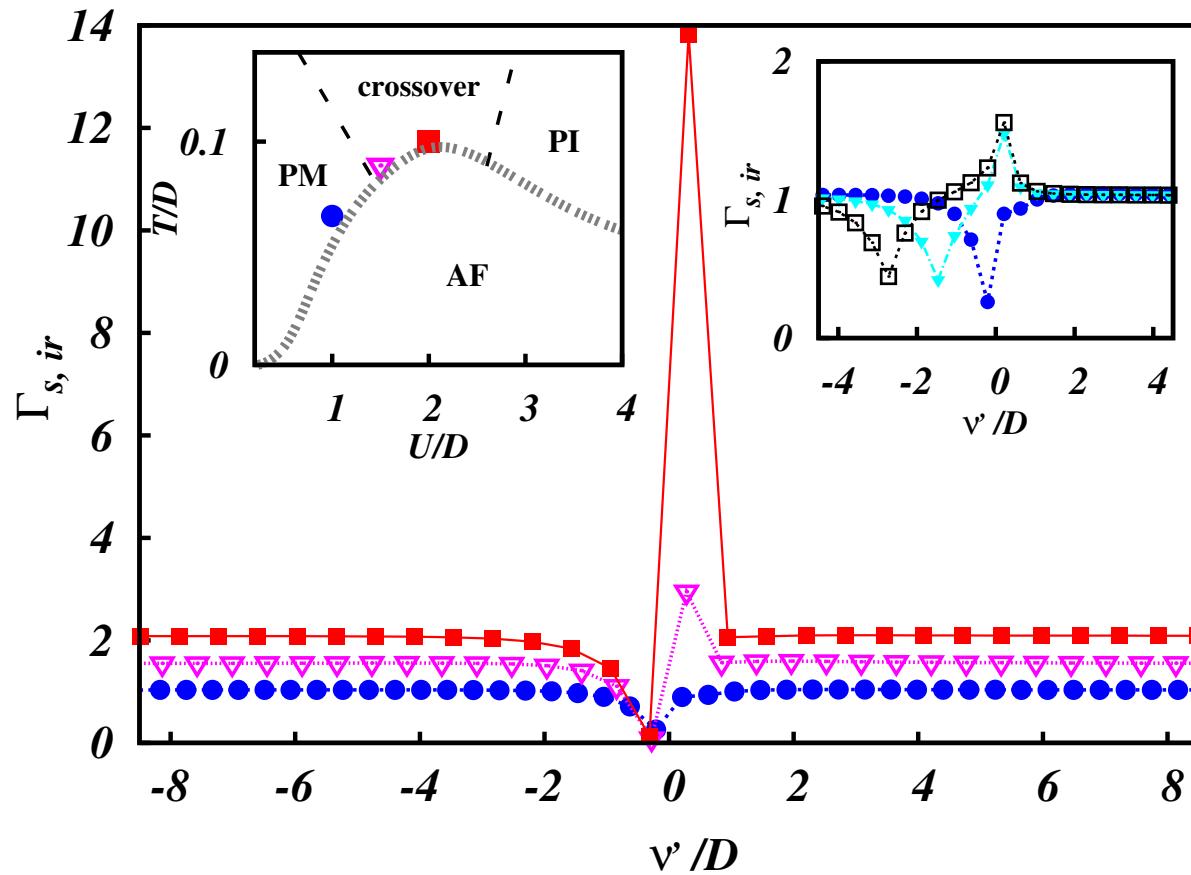
# D $\Gamma$ A algorithm (restriction to ph ladders)



# Results: 3D Hubbard model

$$H = -\textcolor{blue}{t} \sum_{\langle i,j \rangle \sigma} c_{i\sigma}^\dagger c_{j\sigma} + \textcolor{blue}{U} \sum_i n_{i\uparrow} n_{i\downarrow}$$

cubic lattice, exact diagonalization as impurity solver

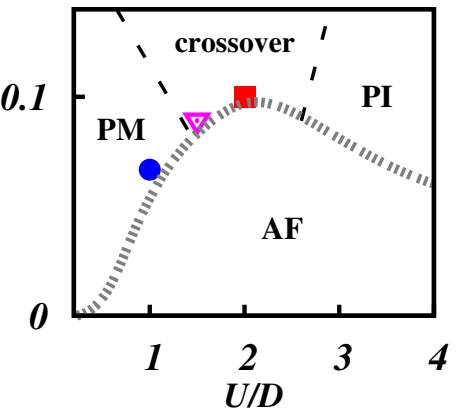


eff. bandwidth  $\equiv 2D$   
 $\omega = 0$   
 $\nu = \pi T$

$\Gamma_{s,\text{ir}}(\nu, \nu', \omega)$  strongly frequency dependent

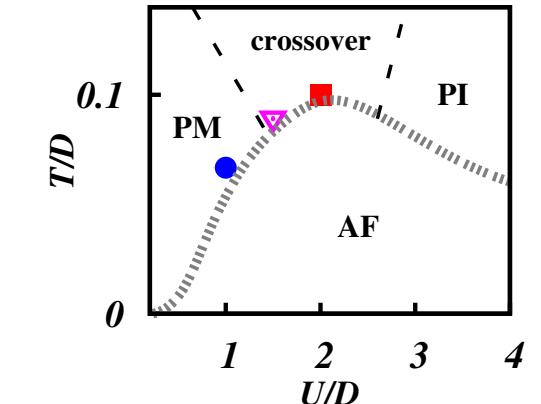
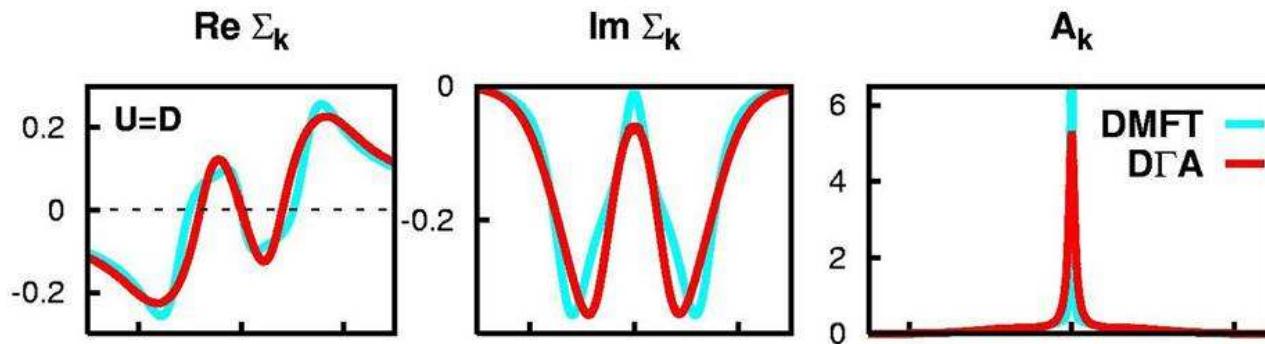
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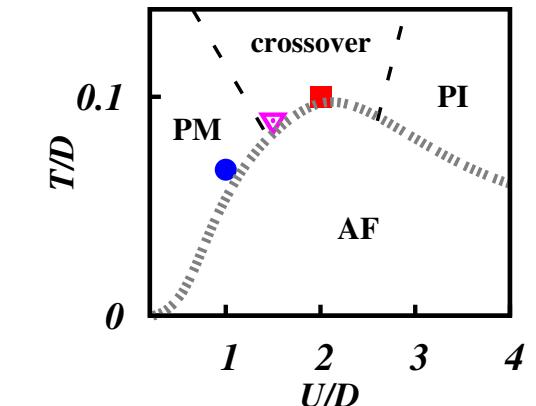
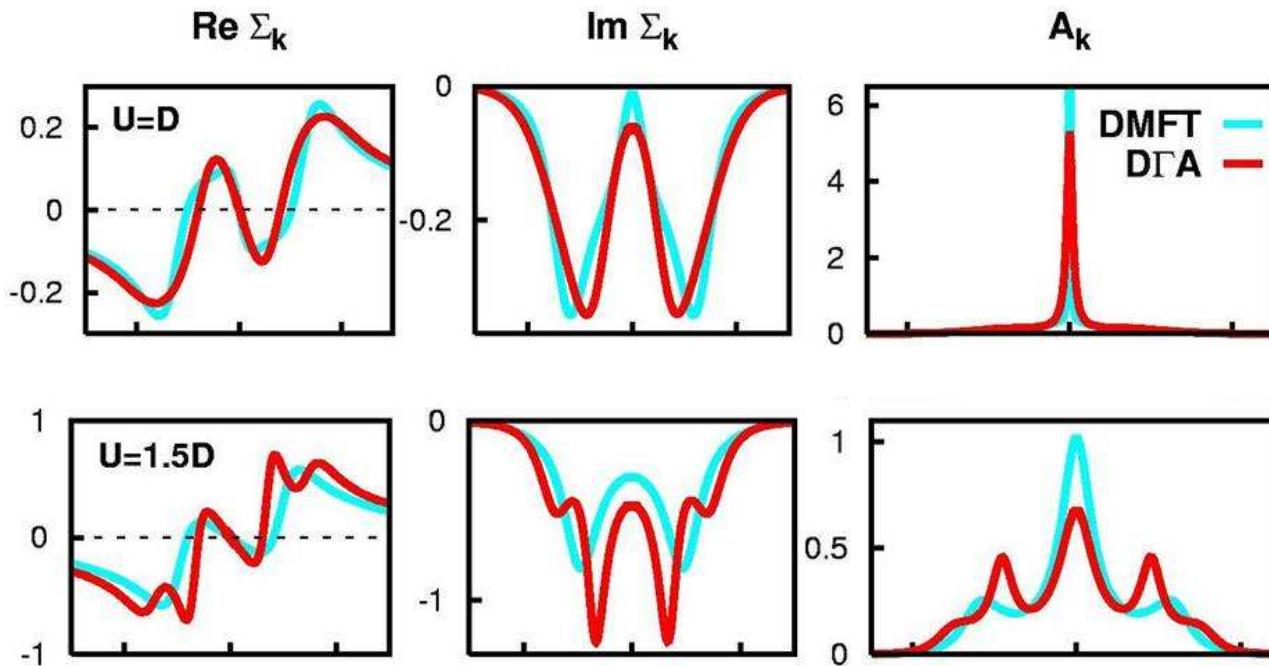
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weak damping  
of QP peak

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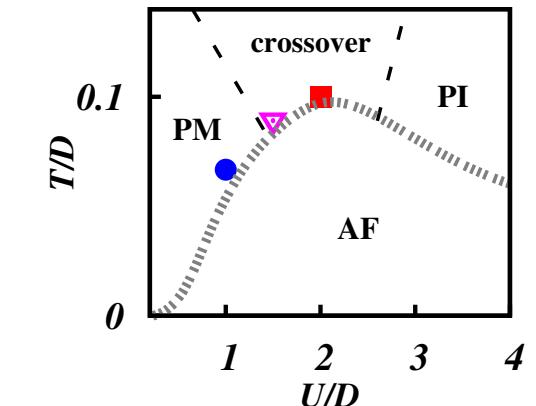
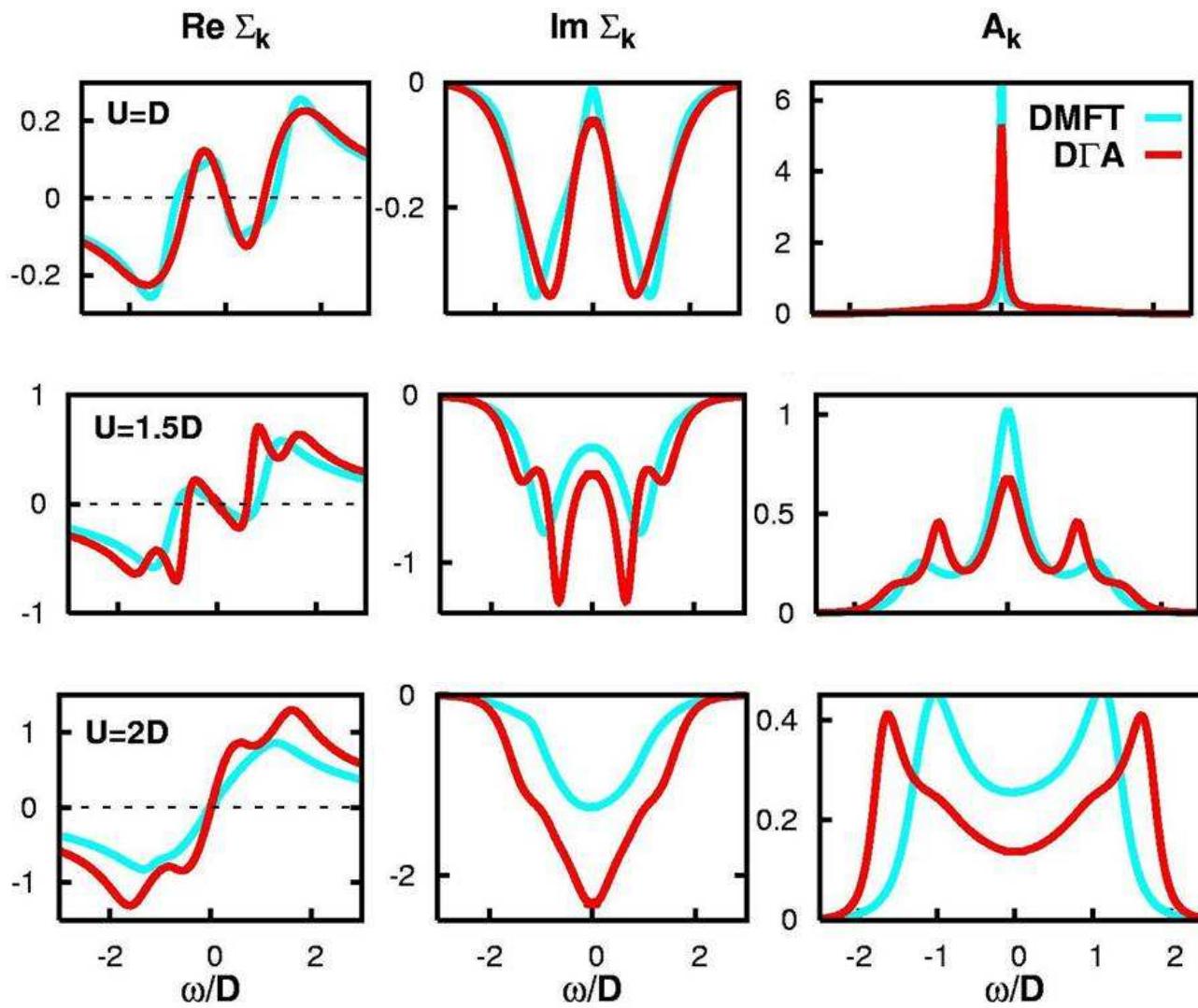


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QP-damping  
strongly enhanced

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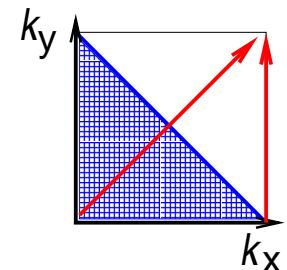
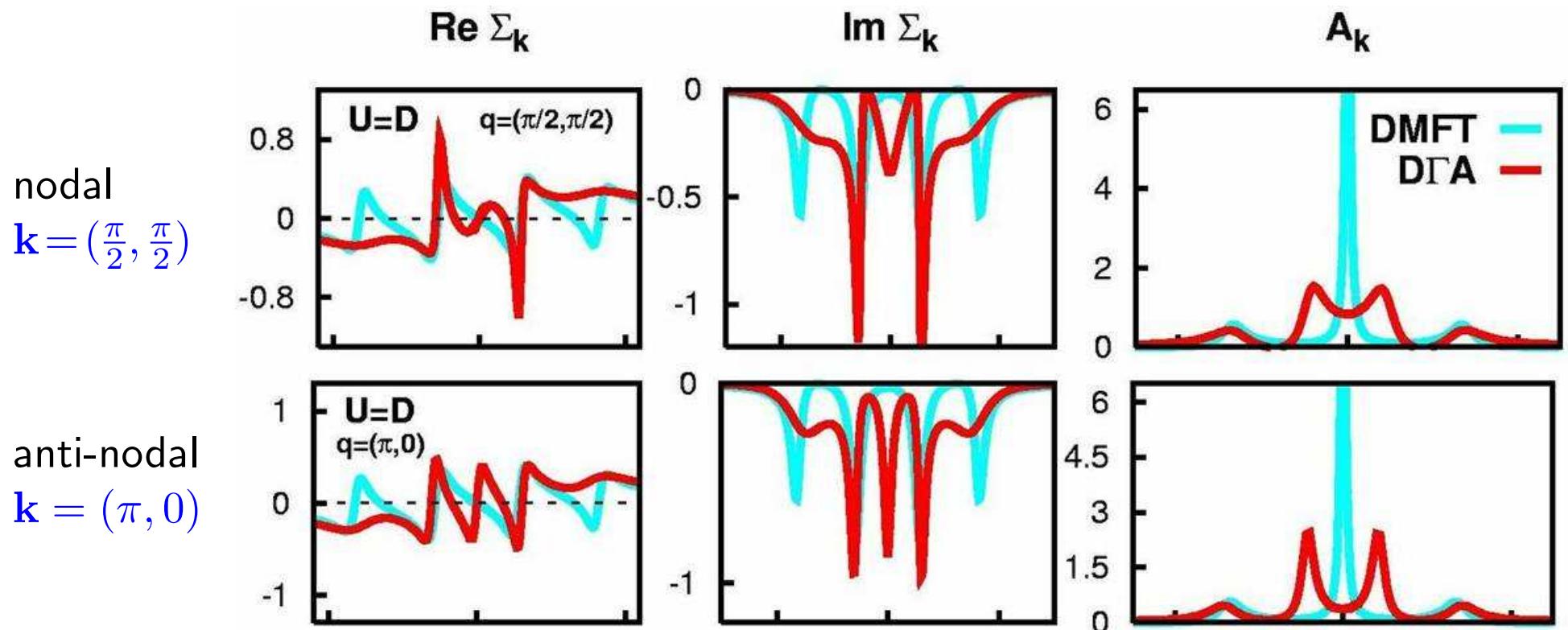


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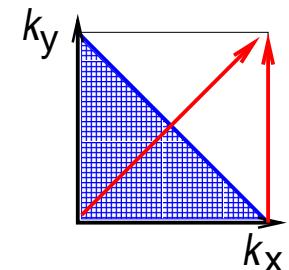
more insulating

# Results: 2D Hubbard model (half-filling)

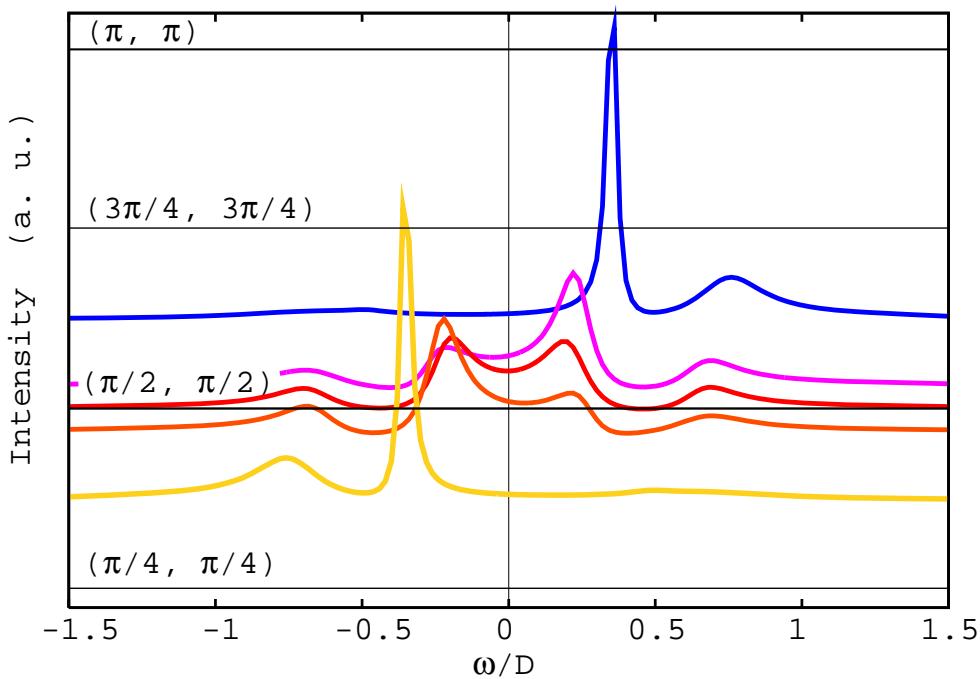


anisotropic pseudogap

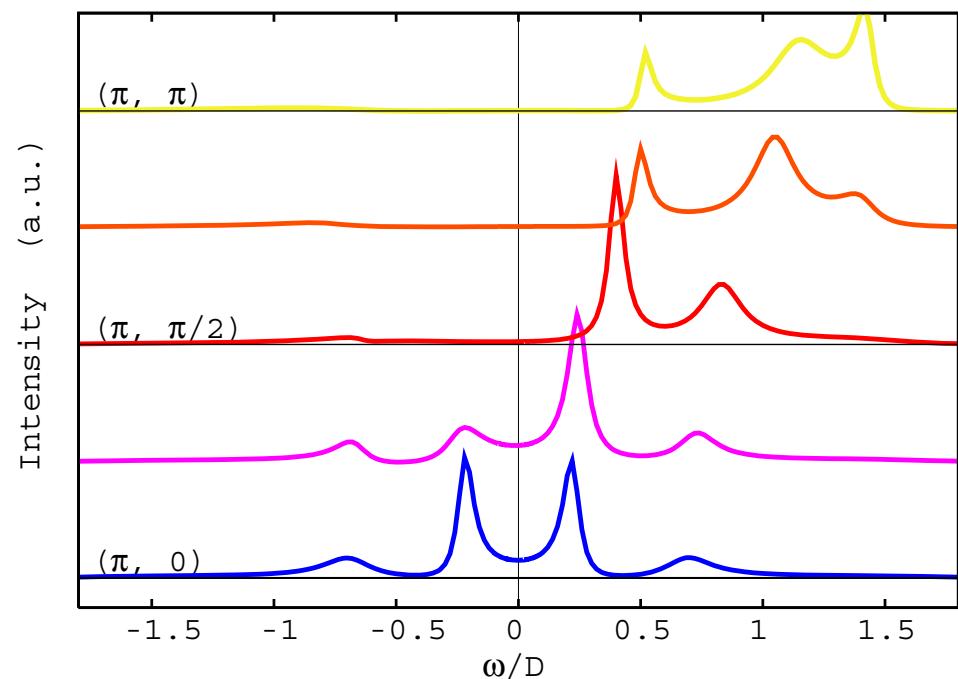
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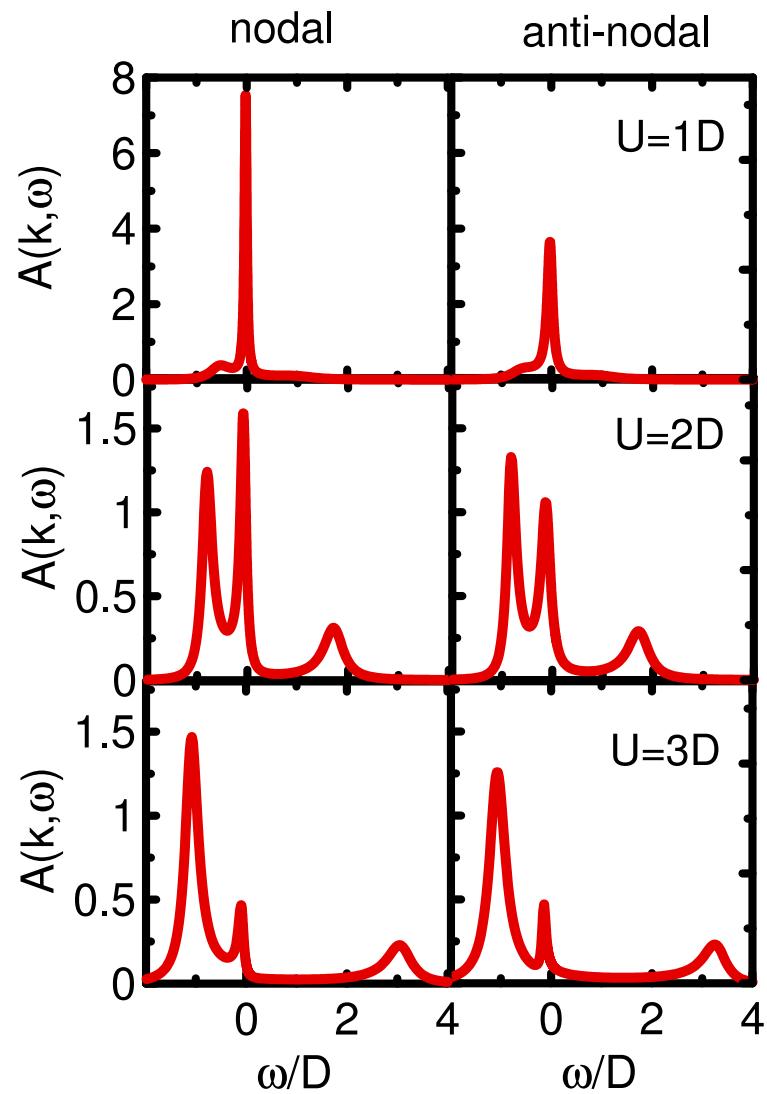
nodal



antinodal



# Results: 2D Hubbard model (off half-filling)



$t'/t = 0.3$   
 $n = 0.8$   
 $\beta = 100/D$

less anisotropic  
at strong coupling

# Results: 1D Hubbard model

Slezak, Jarrell, Maier, Deisz cond-mat/0603421

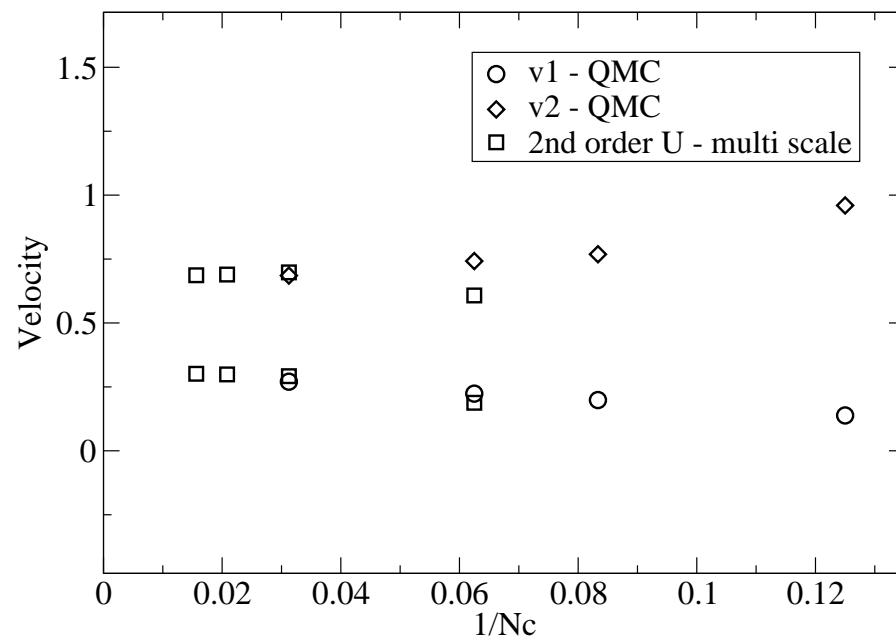
## Spin-charge separation

$$U = W = 1$$

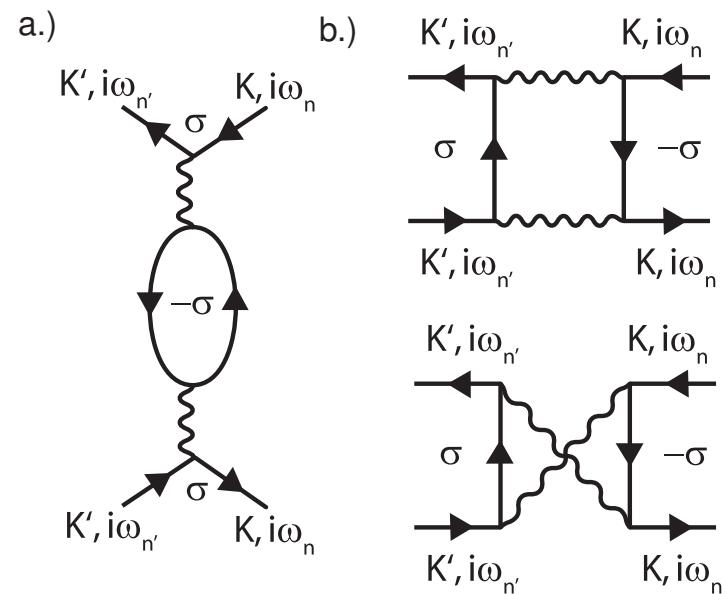
$$k = \pi/2$$

$$\beta = 31$$

$$n = 0.7$$

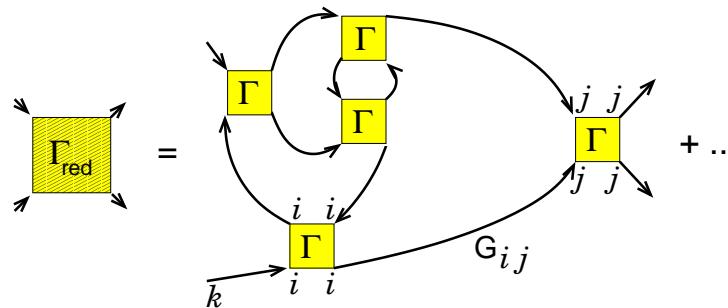


Here, only 2nd order diagrams for vertex  
 $(q = 0, \omega = 0)$   
but 8-site DCA for short-range  $\Sigma$



# Conclusion

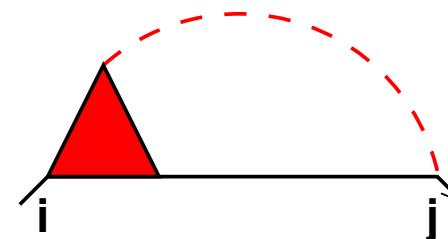
- DΓA assumption: local 2-particle irreducible  $\Gamma$



- DΓA can access short- and long-range correlations
- Results: pseudogap in 2D  
Mott-Hubbard transition modified by AF fluctuations

# Outlook

- Physics: magnons, interplay between AFM and superconductivity, QCP
- Realistic multi-orbital calculations possible
- *Ab initio* calculations with DΓA



includes DMFT, GW, and vertex corrections beyond