

Photoemission final states from the low-energy to high-energy limit: Resolving electronic structure in 3-dim k -space

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- Problem of 3-dim wavevector \mathbf{k} in ARPES:

Final-state k_{\perp} dispersions and lifetimes

- Low energies

Final states by VLEED

Properties of the final states (non-FE and self-energy effects)

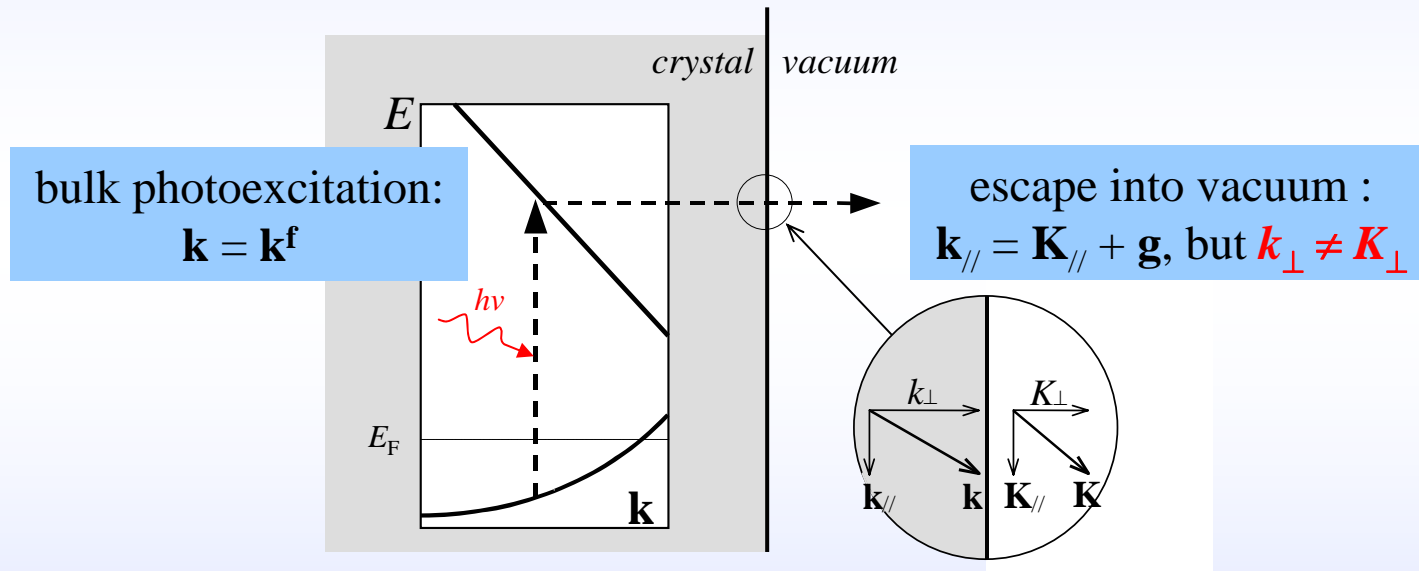
3-dim band mapping by VLEED+ARPES

- High energies

Where do the final states become FE-like?

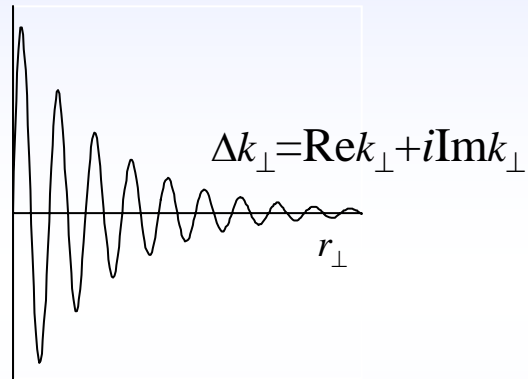
Problem of 3-dim \mathbf{k} in ARPES

- Final-state dispersions

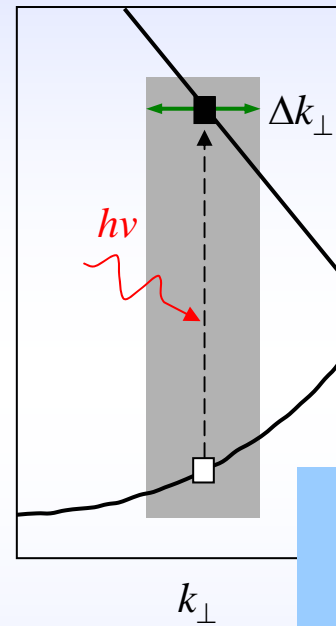


- resolving 3-dim \mathbf{k} requires final-state $E(k_{\perp})$
- low energies: non-free-electron and excited-state self-energy effects in the final states

- Final-state damping



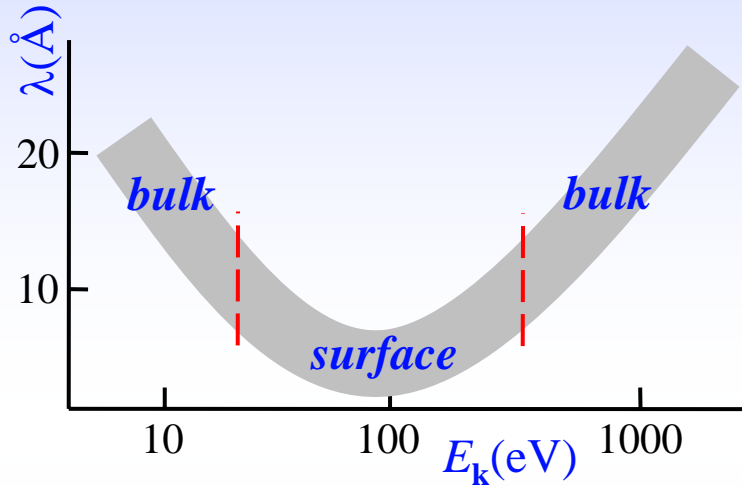
lifetime damping in $r_{\perp} \Rightarrow$
broadening $\Delta k_{\perp} = 2 \text{Im} k_{\perp}$



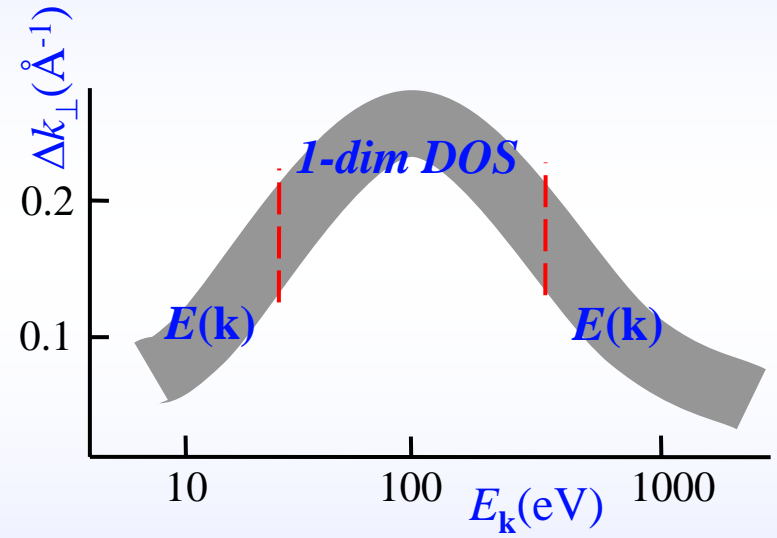
PE signal by *averaging*
of $E(k_{\perp})$ within Δk_{\perp}

- Δk_{\perp} broadening = *intrinsic* k_{\perp} resolution
- spectral peaks \neq true quasiparticle $E(\mathbf{k})$: *intrinsic accuracy* of 3-dim band mapping

- Two sides "universal" curve



$$\Delta k_{\perp} = \lambda^{-1}$$



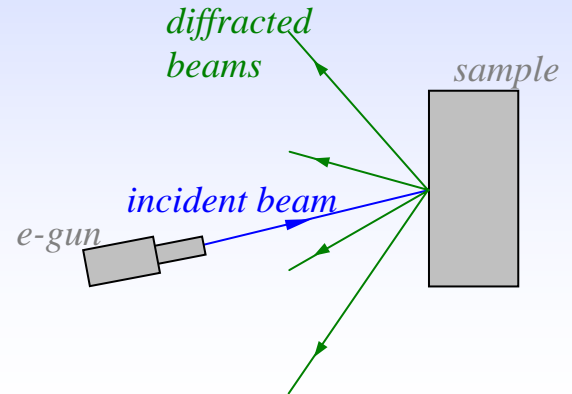
- ARPES regimes: $E(\mathbf{k})$ ($\Delta k_{\perp} \ll k_{\perp}^{\text{BZ}}$) \rightarrow 1-dim DOS ($\Delta k_{\perp} \gg k_{\perp}^{\text{BZ}}$) \rightarrow $E(\mathbf{k})$ (Feibelman & Eastman, 1974)
- low energies: non-free-electron and self-energy effects in the final states by **Very-Low-Energy Electron Diffraction (VLEED)**
- high energies: free-electron final states

Low energies: Final states by VLEED

- What is VLEED?

VLEED = measurement of $R(E)$ elastic reflectivity in the energy range below ~ 40 eV:

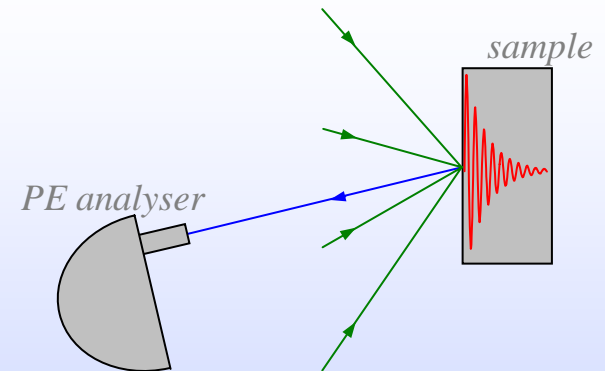
- weak $V_i \Rightarrow$ sensitivity to $E(\mathbf{k})$



- Why VLEED?

One-step ARPES theory:
$$I^{\text{ph}} = \left| \left\langle \Phi^{\text{LEED}*} \left| \hat{\mathbf{A}} \cdot \hat{\mathbf{p}} \right| \Phi^{\text{i}} \right\rangle \right|^2$$

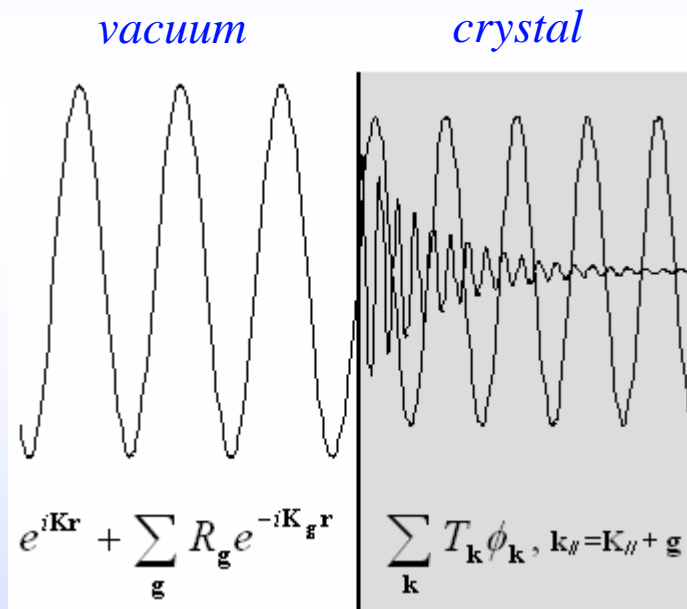
- direct connection to ARPES:
ARPES final state = time-inversed LEED state
- final-state energies in ARPES



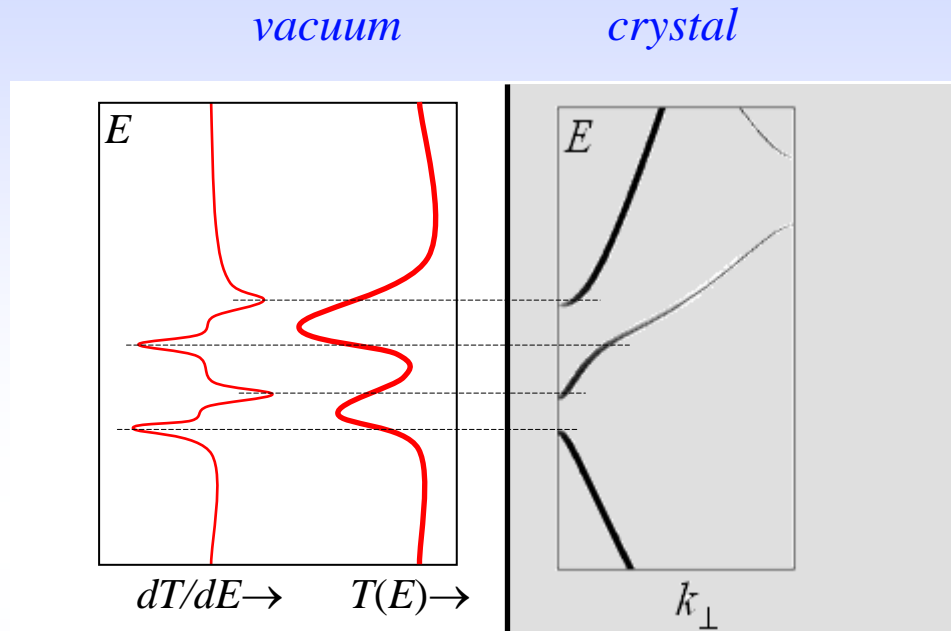
Connection of VLEED to $E(\mathbf{k})$

- elastic case (inelastic scattering = 0)

$R(E) = 1 - T(E)$ by matching vacuum and crystal wavefunctions:



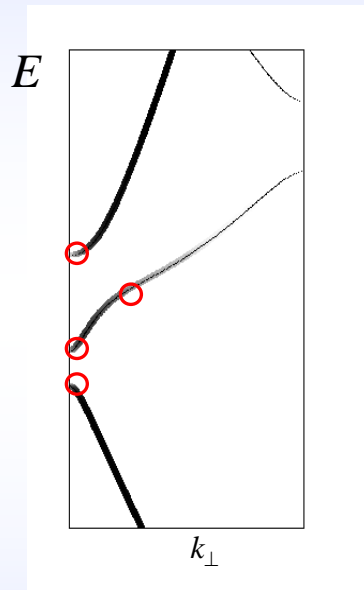
VLEED spectrum $T(E) \rightarrow \leftarrow E(k_{\perp})$ along $\mathbf{k}_{\parallel} = \mathbf{K}_{\parallel} + \mathbf{g}$



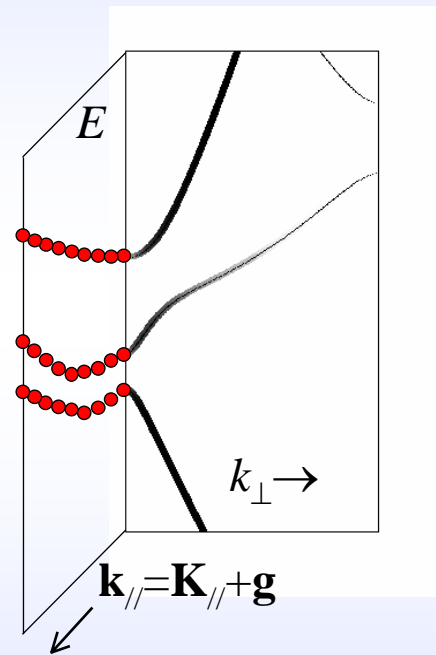
$T(E)$ minima/maxima \rightarrow \leftarrow band gaps/dispersions ranges in $E(k_{\perp})$
 dT/dE extremes \rightarrow \leftarrow critical points in $E(k_{\perp})$

VLEED structures = critical points in $E(k_{\perp})$

Band mapping techniques



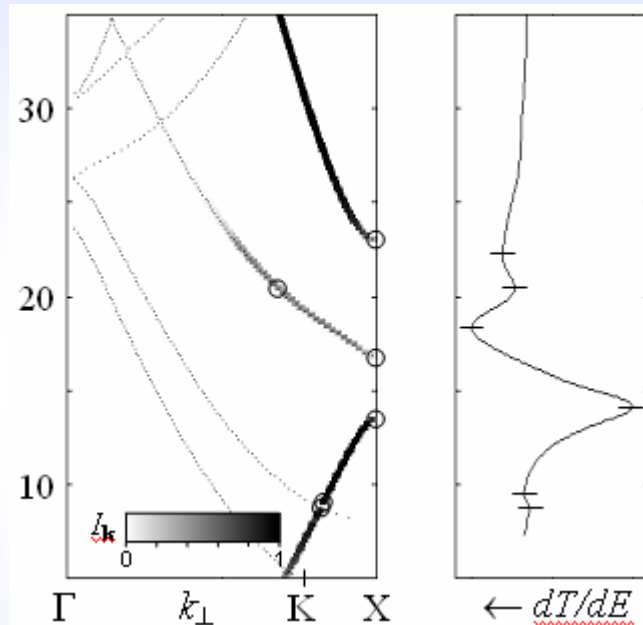
$E(k_{\perp})$: fitting experimental critical points



$E(\mathbf{k}_{\parallel})$: direct band mapping along symmetry lines

Effects of different bands

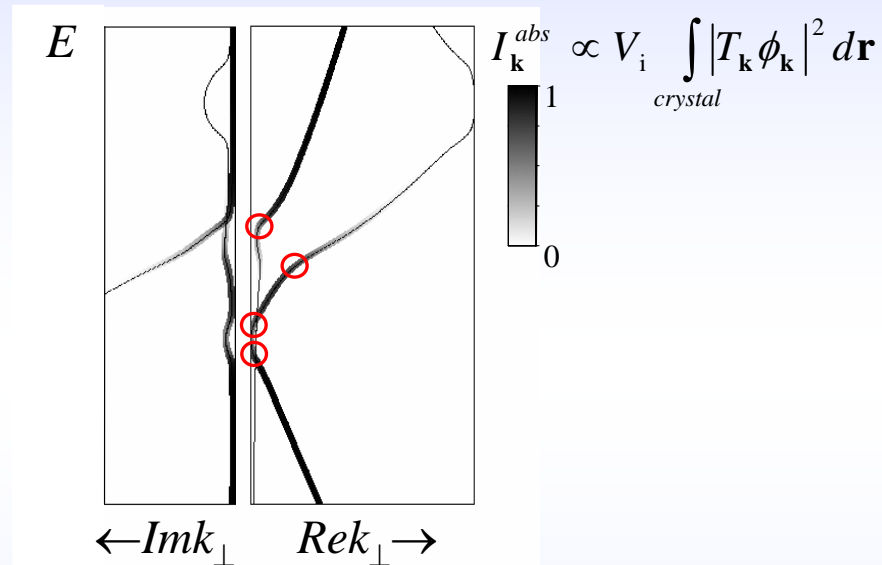
- normal-incidence
VLEED on Cu(110)



- VLEED spectrum \sim *conducting bands*: large *partial currents* $I_{\mathbf{k}} = |T_{\mathbf{k}}|^2 v_{\perp}$ (effective coupling to vacuum + transport into crystal)
- $I_{\mathbf{k}}^{\text{VLEED}} \sim I_{\mathbf{k}}^{\text{PE}}$: the same conducting bands effective in VLEED and ARPES

Effect of inelastic scattering ($V_i \neq 0$)

- damped Bloch waves with complex $k_{\perp} = \text{Re}k_{\perp} + i\text{Im}k_{\perp}$
- smooth $E(k_{\perp})$ dispersions

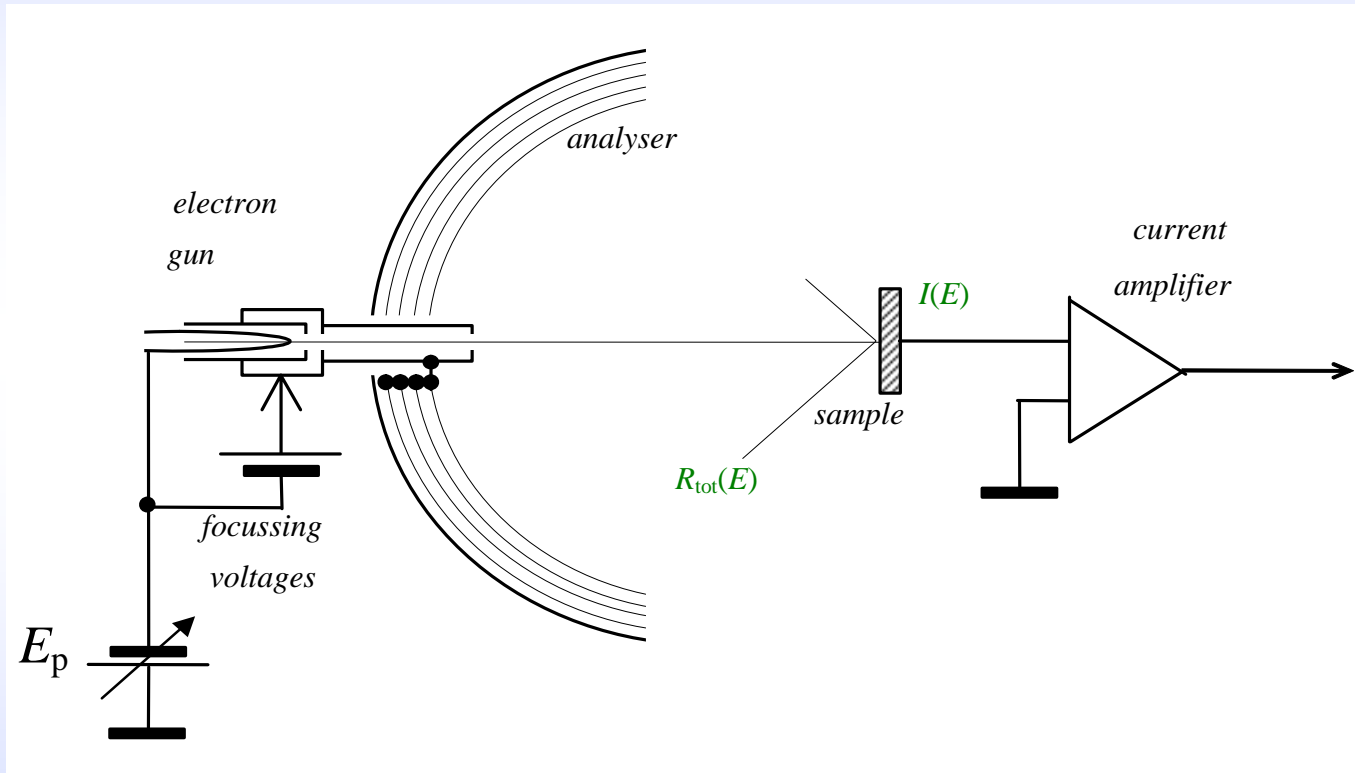


critical points = extremes
in $1/\text{curvature of } E(k_{\perp})$

- VLEED structures $\left\{ \begin{array}{l} \text{energies} \rightarrow \text{critical points in conducting bands } E(k_{\perp}) \\ \text{broadening} \rightarrow V_i = \hbar/\tau \end{array} \right.$

Experimental technique

- Experimental setup



- *retarding field* to maintain focusing down to $E_p \sim 0$
- $I_{gun} = const \Rightarrow$ measurements of $T(E) \propto$ - integral reflectivity $R(E)$ in *target current*
- data acquisition time < 1 min/spectrum

- **Retarding-field: Angle dependences**

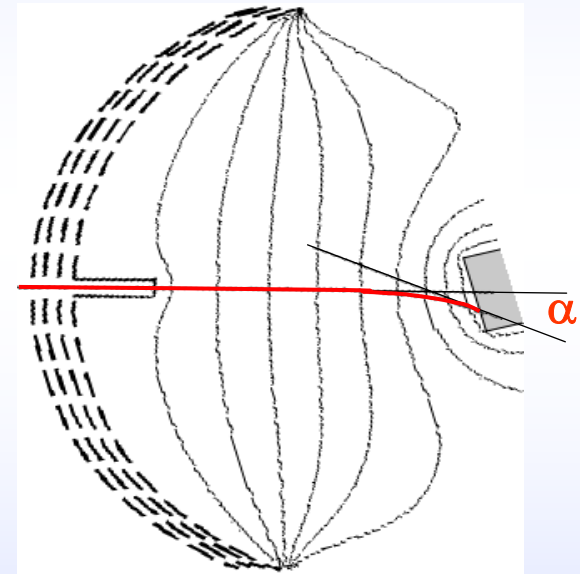
- α - and E -dependent deviation of off-normal trajectories and variation of incidence angle

- necessity to control $\mathbf{K}_{//}$:
- electrostatic ray-tracing

$$\sum_i \frac{\partial^2 U}{\partial x_i^2} = 0 \quad + \quad \frac{\partial^2 x_i}{\partial t^2} = -\frac{e}{m} \frac{\partial U}{\partial x_i}, \quad x_i = x, y, z$$

- $K_{//}(\alpha, E) = \sum_{m,n=0}^2 A_{mn} E^m \alpha^n$ fitted to

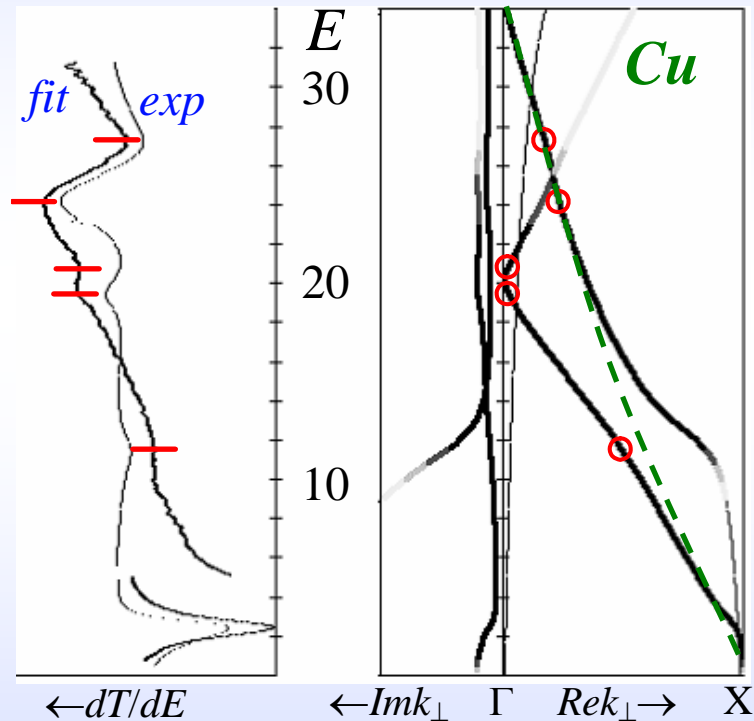
experimental points with well-defined $\mathbf{K}_{//}$



3. Experimental properties of the final states

Non-free-electron effects

- $E(k_{\perp})$ from normal-incidence VLEED on Cu(100)



- deviations from FE-like model:

- non-parabolic dispersions (local m^* , V_{000})

- *multiband* composition (different k_{\perp} available in the final state)

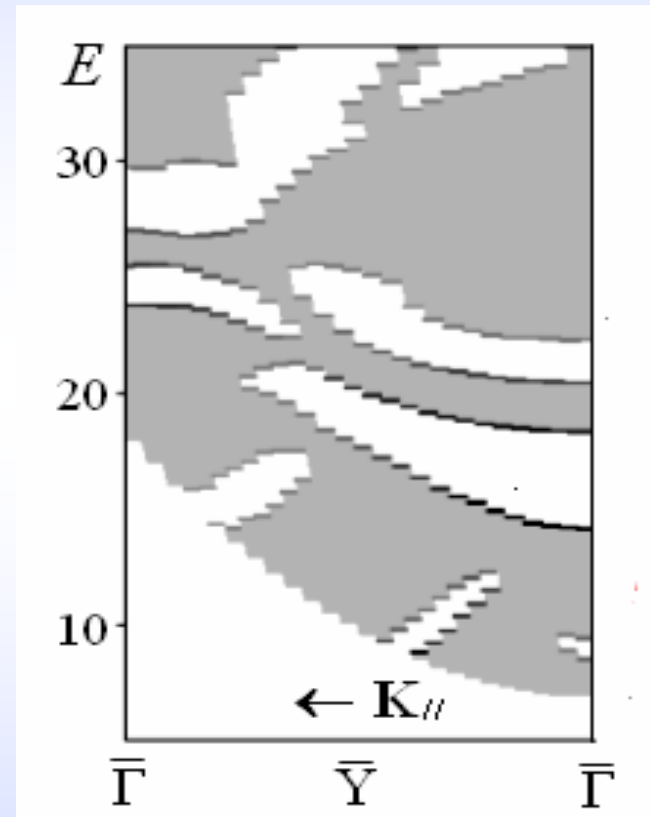
- $E(k_{//})$ from angle-dependent VLEED on Cu(110)

VLEED $\mathbf{K}_{//}$ dispersions:

$T(E)$ maxima (gray) = k_{\perp} dispersion ranges,

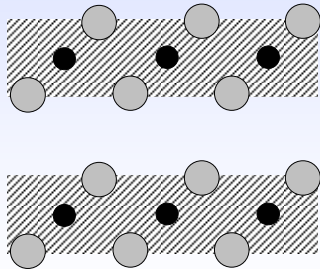
$T(E)$ minima (white) = band gaps

= *surface-projected* $E(\mathbf{k})$ of the final states



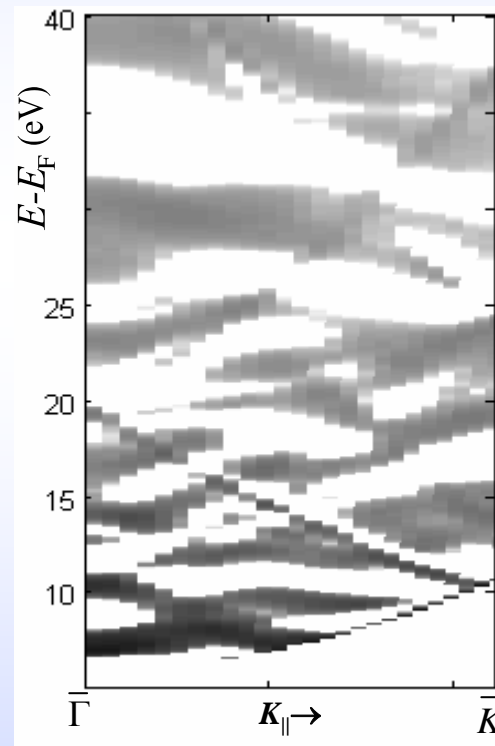
- deviations from FE-like dispersions closer to the BZ borders

- $E(\mathbf{k})$ of layered material TiTe_2



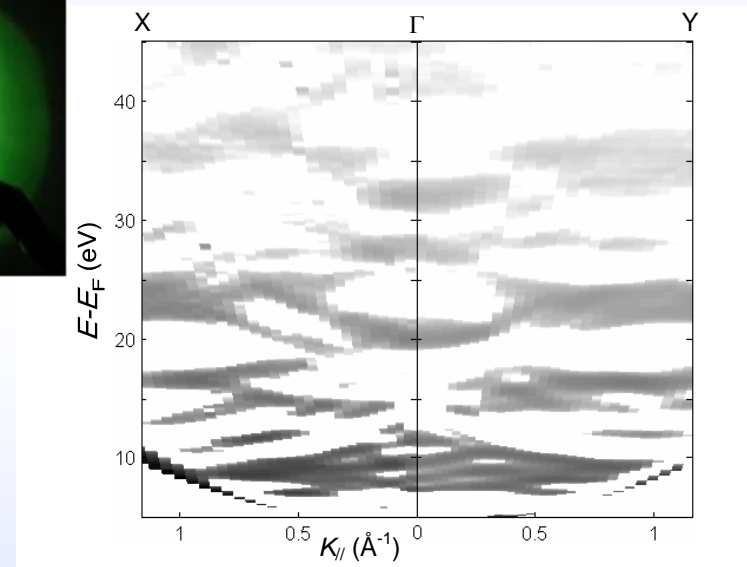
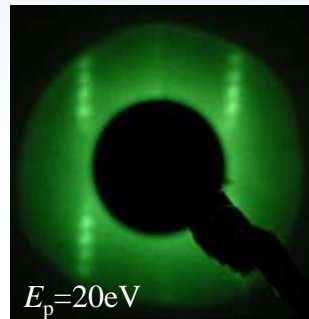
- anisotropic quasi-2-dim properties

$E(k_{\parallel}) = \text{VLEED } K_{\parallel} \text{ dispersion}$



- prominent non-FE effects
- similar strength of non-FE effects in other layered materials (graphite, VSe_2 , TiTe_2 , NbSe_2 ...) due to sharp modulations of the crystal potential

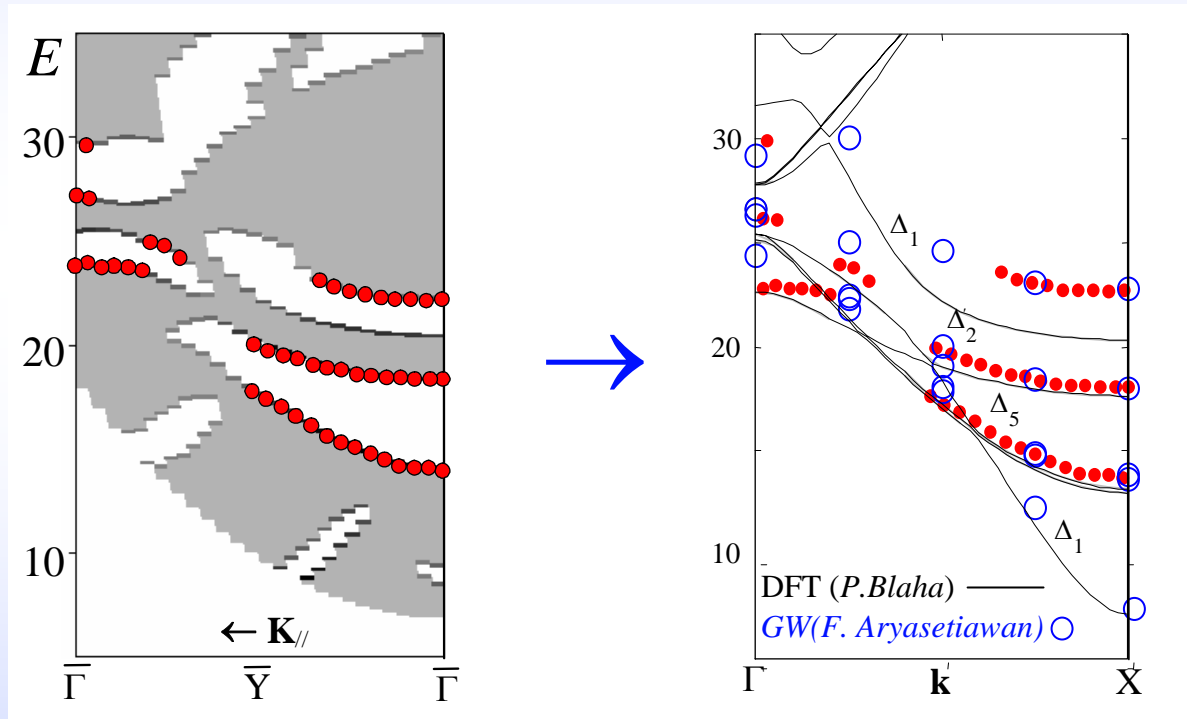
- $E(\mathbf{k})$ of high- T_c material $\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_8$



- 3D final-state dispersions
- highly structured final states \Rightarrow dramatic $h\nu$ and $\mathbf{K}_{||}$ dependences in ARPES

Self-energy effects

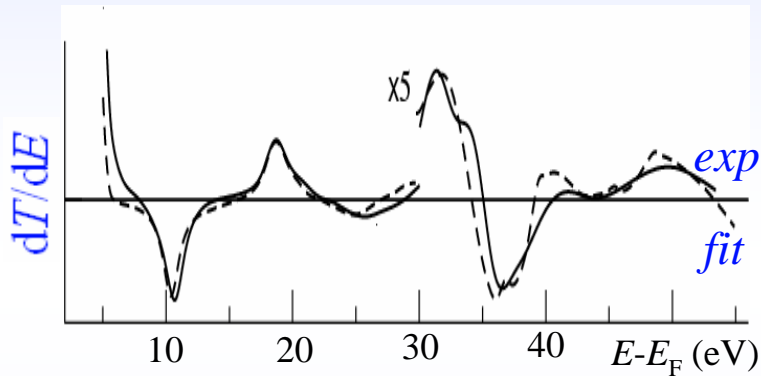
- $\Delta\Sigma(E,\mathbf{k})$ in Cu by angle-dependent VLEED on (110) surface



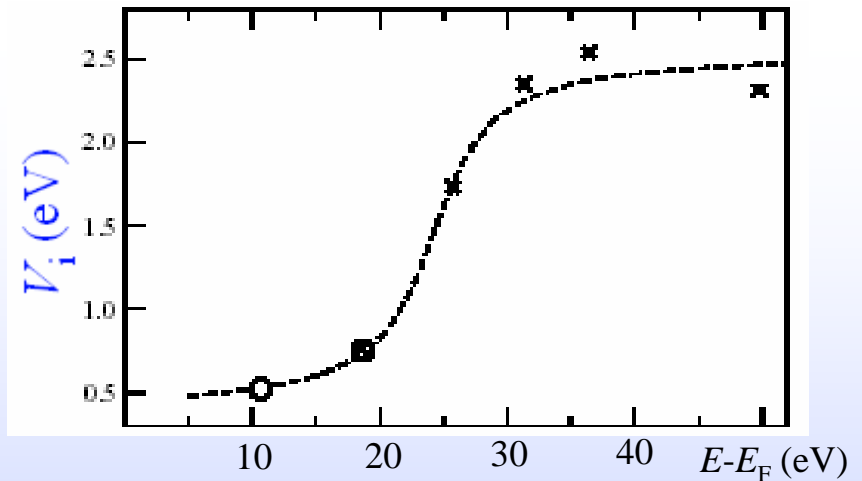
- Excited-state $\Sigma(E,\mathbf{k},\mathbf{k}') \neq \text{static } V_{xc} \Rightarrow \Delta\Sigma$
- Band- and \mathbf{k} -dependence of $\Delta\Sigma$

Lifetimes

- $V_i(E) = \hbar/\tau$ in graphite (*Barrett et al, 2005*)



• $V_i(E)$ by fitting the sharpness and relative amplitude of dT/dE structures



• Typically sharp increase of V_i at $\hbar\omega_p$

VLEED vs other spectroscopies for unoccupied states

- VLEED vs XAS and IPES

- resolution in 3-dim \mathbf{k}
- single electron state (XAS: core hole excitonic effects)
- direct connection with ARPES
- experimental simplicity

- VLEED vs SEE

- Recent study on Cu(110) (*Bovet et al, 2004*): **SEE and VLEED equivalent**
- Thermodynamic model of SEE (*Feder & Pendry, 1978*): $I_{\text{SEE}}(E) \propto T_{\text{VLEED}}(E) + bkg$

3-dim band mapping by VLEED+ARPES

- Connection between VLEED and ARPES

General connection:
$$I^{\text{ph}} = \left| \left\langle \Phi^{\text{LEED}*} \left| \hat{\mathbf{A}} \cdot \hat{\mathbf{p}} \right| \Phi^{\text{i}} \right\rangle \right|^2$$

Detailed connection between
the *partial absorbed currents*
and *partial photocurrents*:

$$I_{\mathbf{k}}^{\text{ph}} \propto I_{\mathbf{k}}^{\text{abs}} \cdot \left(\frac{1}{V_i} \frac{\partial k_{\perp}}{\partial E_i} \left| M_{\text{fi}} \right|^2 \right)$$

- VLEED coupling bands = PE dominant final bands

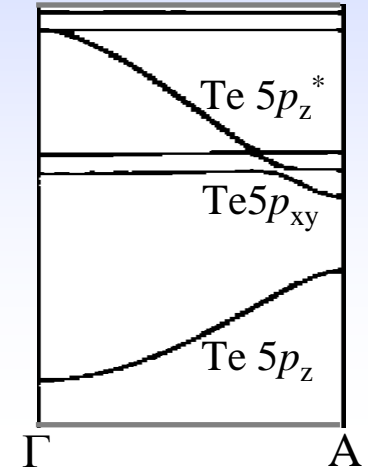
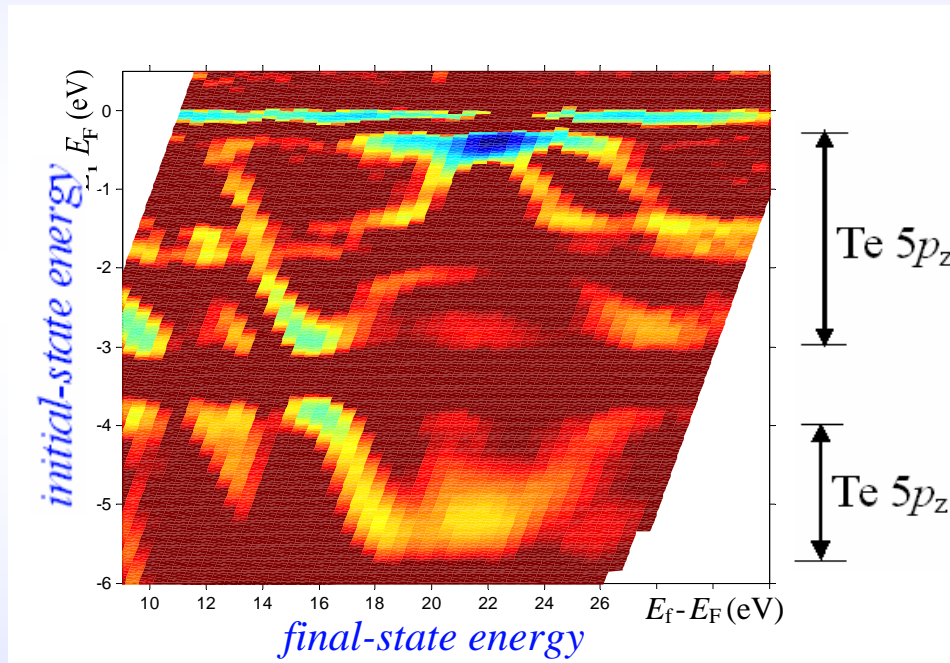


Final-state $\left\{ \begin{array}{l} E(\mathbf{k}) \\ \text{lifetimes} \end{array} \right. \rightarrow$ Valence-band $\left\{ \begin{array}{l} E(\mathbf{k}) \text{ resolved in 3-dim } \mathbf{k} \\ \text{controlled intrinsic accuracy} \end{array} \right.$

Mapping in k_{\perp} (photon energy variation): Quasi-2-dim TiTe_2

- ARPES data

$-d^2I/dE^2$ intensity map (log scale)

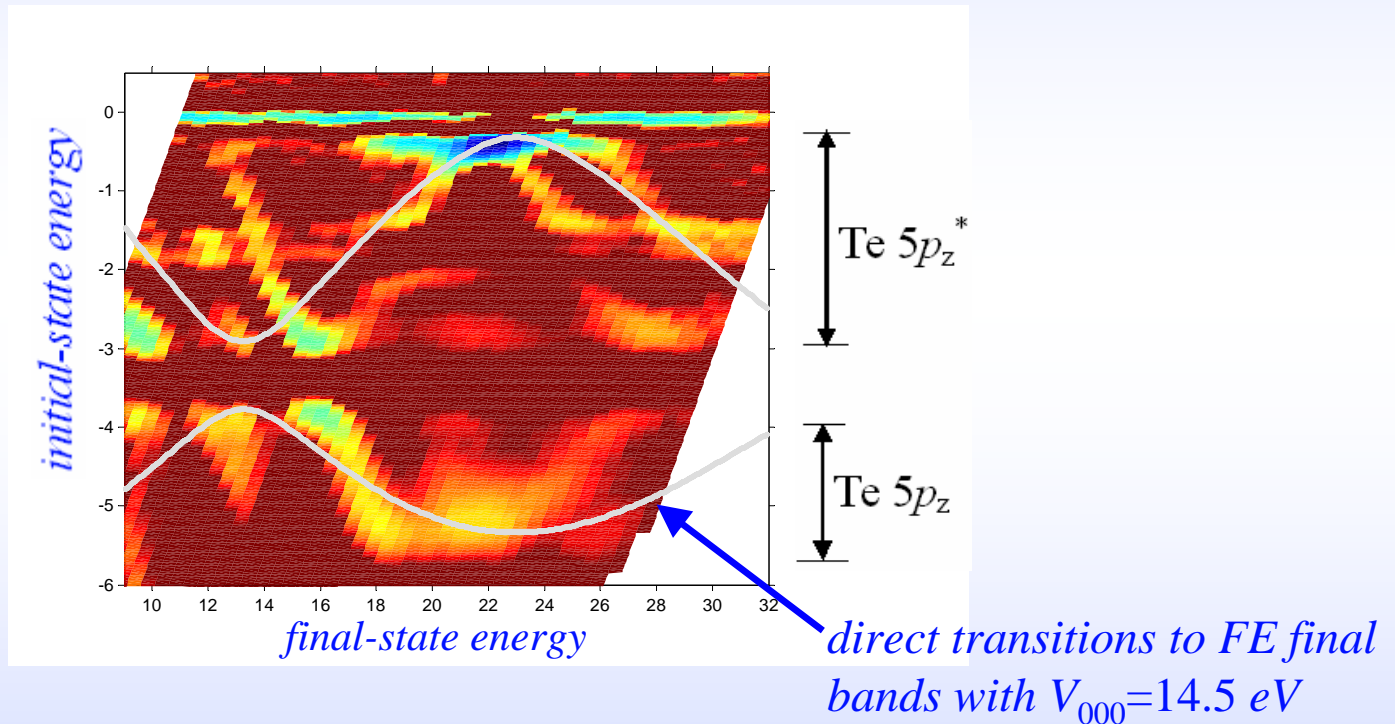


- non-parabolic and multiple Te $5p_z/p_z^*$ dispersions

Mapping in k_{\perp} (photon energy variation): Quasi-2-dim TiTe_2

- ARPES data

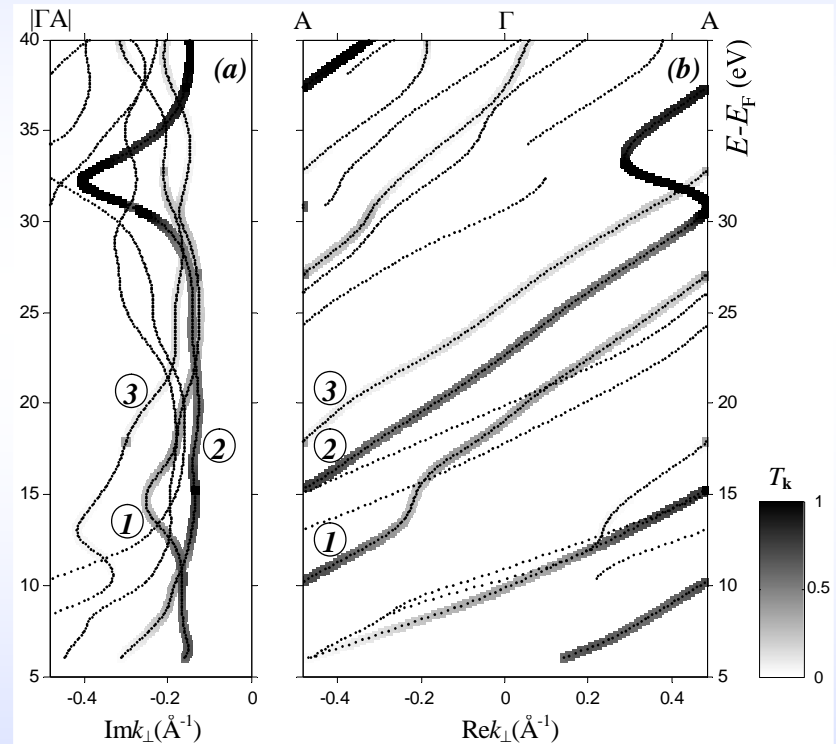
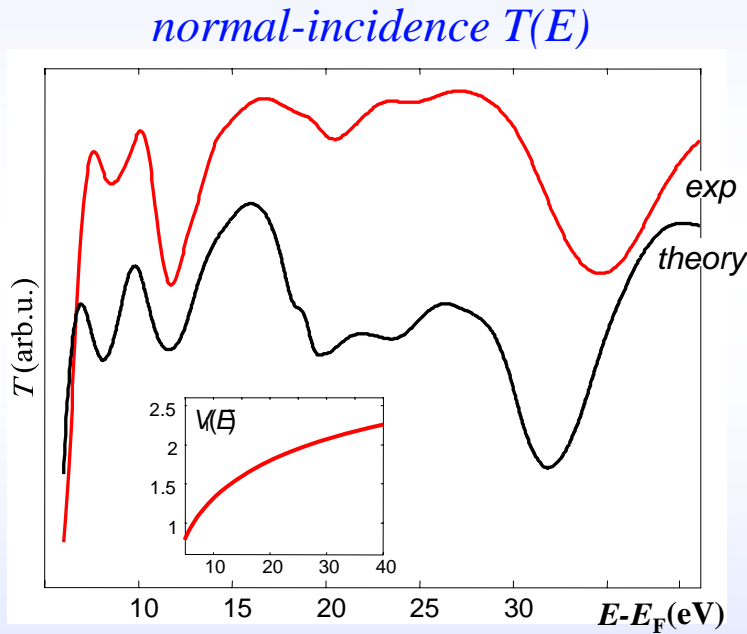
$-d^2I/dE^2$ intensity map (log scale)



- failure of the FE approximation to describe the experimental $\text{Te } 5p_z/p_z^*$ dispersions

- Final states

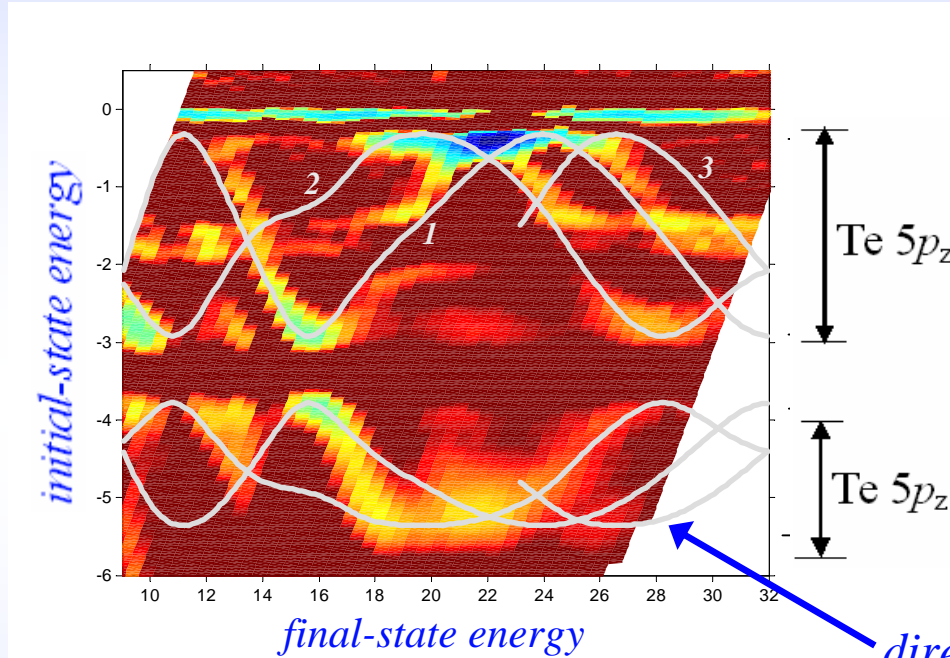
- final states by $\mathbf{k}\cdot\mathbf{p}$ + ELAPW calculations (*E.E. Krasovskii*) with $V_i(E)$ fitted to the experiment + $\Delta\Sigma$ corrections



- dramatic non-FE effects: non-parabolic dispersions + multiband composition
- $\Delta k_{\perp} = \text{Im}k_{\perp} \ll k_{\perp}^{\text{BZ}}$ – band structure regime

- ARPES data vs VLEED final states

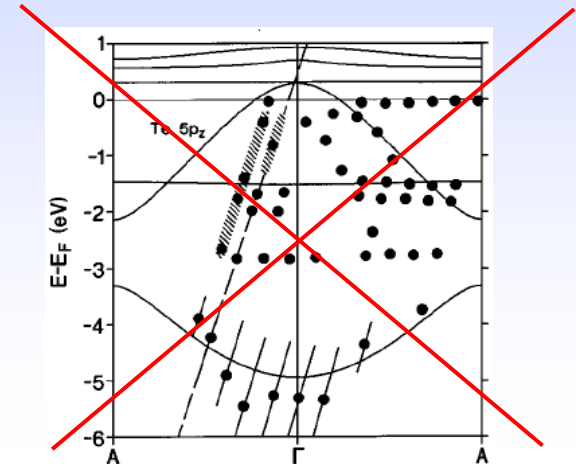
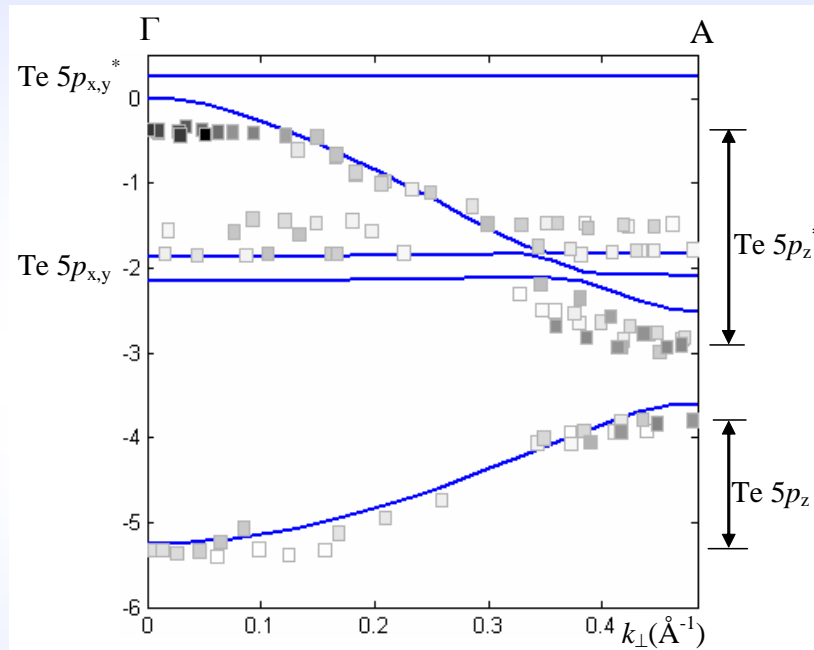
$-d^2I/dE^2$ intensity map (log scale)



direct transitions to the VLEED derived final bands

- the VLEED derived final states (including the non-FE and $\Delta\Sigma$ effects) reproduce the experimental non-parabolic and multiple $\text{Te } 5p_z/p_z^*$ dispersions
- remnant disagreement due to intrinsic shifts from the direct transitions
- multiple final bands = common phenomenon of 'umklapp bands' or 'secondary cone' emission

- Band mapping

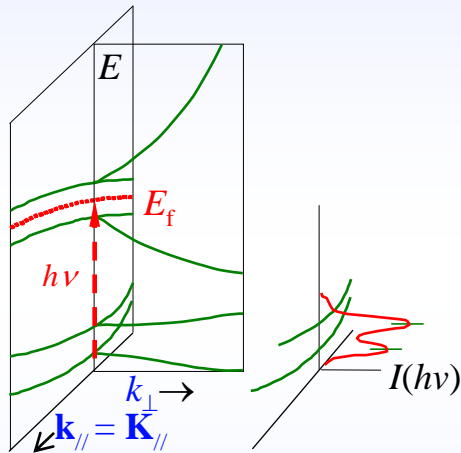


- consistent $E(k_{\perp})$ in contrast to the FE final states
- $\text{Te } 5p_z^*$ does not cross $E_F \rightarrow$ no FS electron pocket in Γ
- true final states incorporating the non-FE and $\Delta\Sigma$ effects are essential (VSe_2 , TiS_2 ...)

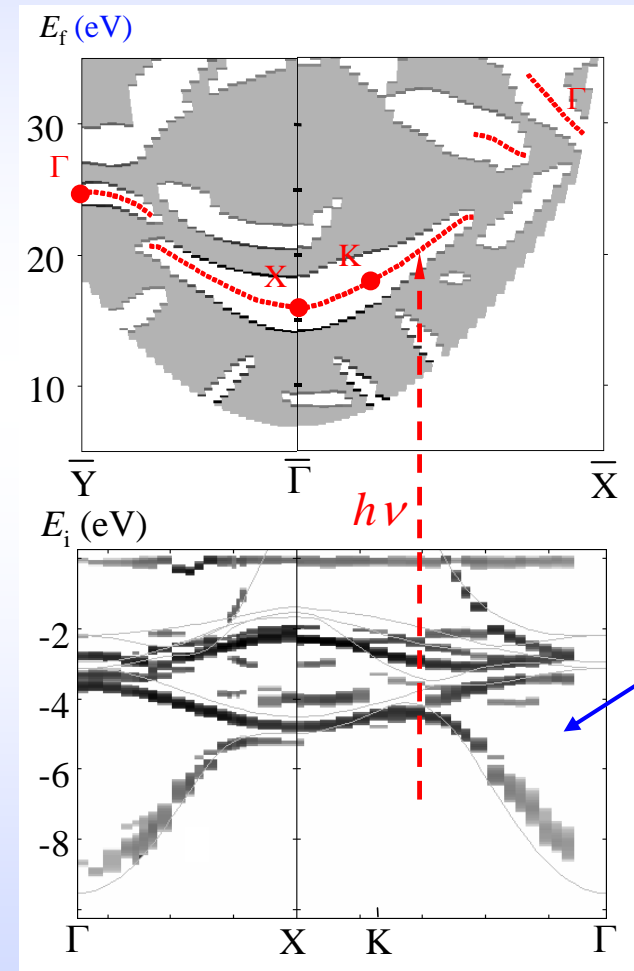
Mapping in $K_{//}$ (emission angle variation)

- Results on Cu

- Idea



- Angle-dependent VLEED + Constant-Final-State ARPES:
 - 1DOS maxima \Rightarrow intensity gain
 - many directions at one surface
 - direct image of valence band $E(\mathbf{k})$

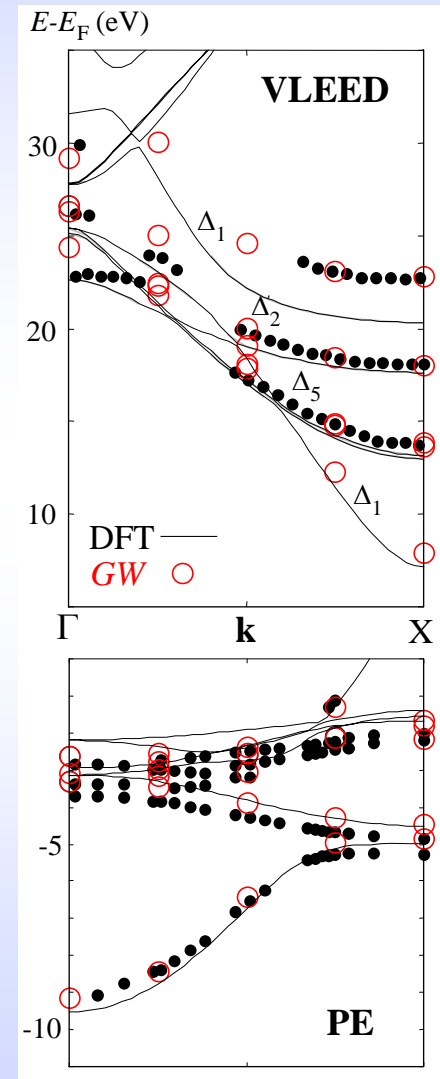


- \mathbf{k} - and band-dependent $\Delta\Sigma$ in the valence band

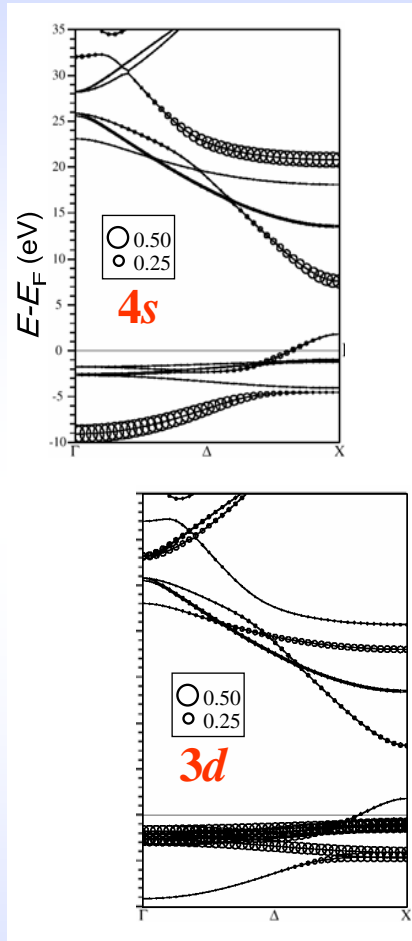
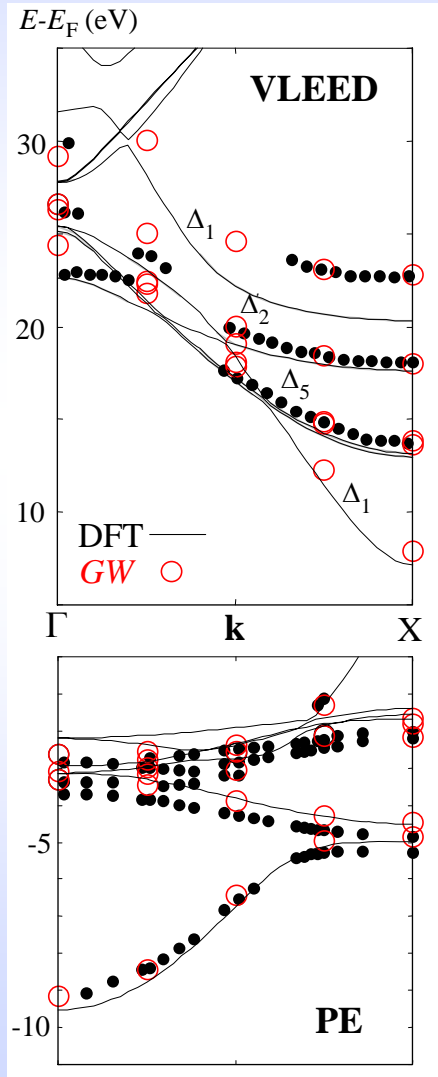
Self-energy effects in Cu

- VLEED experiment vs DFT and GW

- Significant band- and \mathbf{k} -dependent $\Delta\Sigma$ renormalization thru unoccupied and occupied $E(\mathbf{k})$ despite weakly correlated nature of Cu
- Agreement with GW quasiparticle $E(\mathbf{k})$



- Mechanism: Non-local exchange

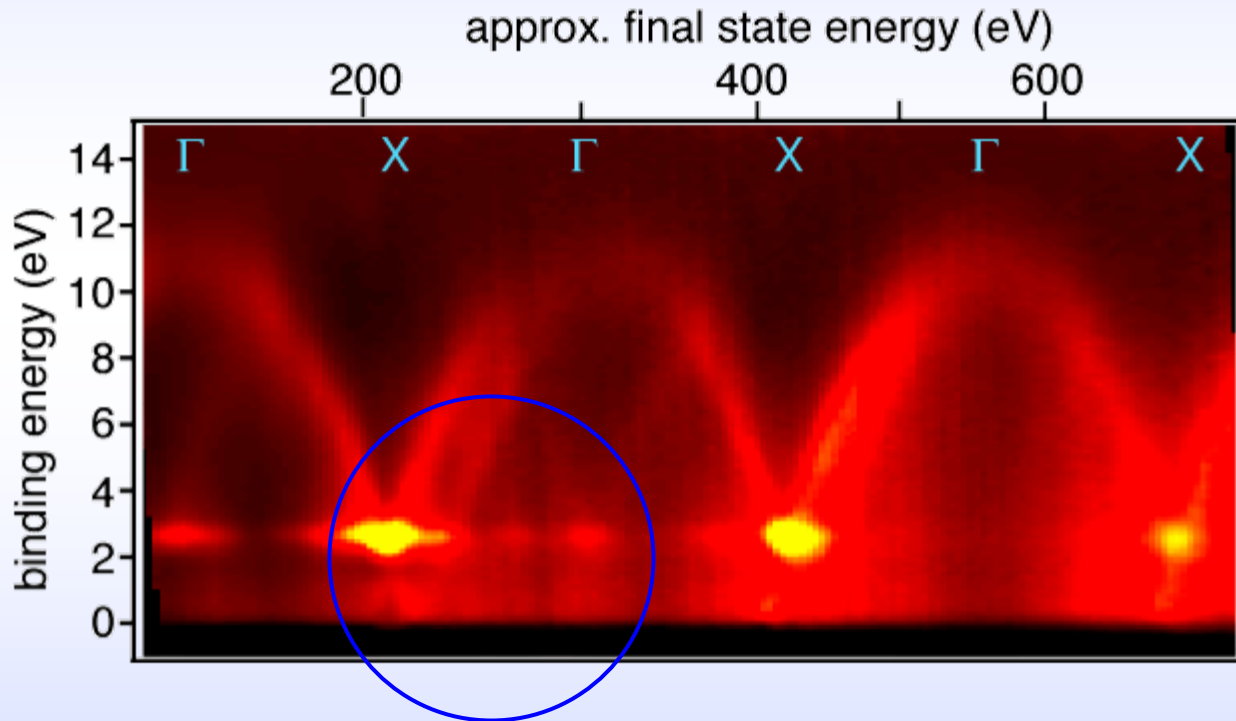


- 4s weight: anomalous $\Delta\Sigma > 0$
- 3d weight: anomalous $\Delta\Sigma < 0$

- Mechanism: Exchange integral $-\sum_q \int \phi_k^*(\mathbf{r}) \phi_q^*(\mathbf{r}) v(\mathbf{r}-\mathbf{r}') \phi_q(\mathbf{r}') \phi_k(\mathbf{r}') d\mathbf{r} d\mathbf{r}'$ couples to the valence d -electrons \Rightarrow more negative for d -states and less negative for s -states

High energies: Where do the final states become FE-like?

- Example: Normal-emission data on Al(100) (*Hoffman et al 2002*)



- non-FE multiband final states up to 400 eV

High energies: Soft-X-ray ARPES at Swiss Light Source

ADvanced RESonant Spectroscopies (ADDRESS) Beamline :

- energy range 400-1800 eV
- resolution 35 meV @ 1 keV
- 3×10^{11} to 1×10^{13} photons/s/0.01%BW
- RIXS endstation:
 - high-res spectrometer (70 meV @ 1 keV) by Politecnico di Milano
 - rotating platform to study \mathbf{q} -dependences
- ARPES endstation
 - operation in spring 2008
 - talk by L Patthey (Thursday, April 26)



Cooperations:

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E Krasovskii, Kiel

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