

# **Entanglement**

**and**

# **Rydberg interaction**

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# Outline

1. Basics of entanglement
2. Entanglement and 2-atom gate using Rybderg blockade

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1. Basics of entanglement
2. Entanglement and 2-atom gate using Rydberg blockade

**Goal:** control Rydberg interaction between few (10 – 100) atoms  
(No ensemble average)

**Quantum state engineering**

# **Outline**

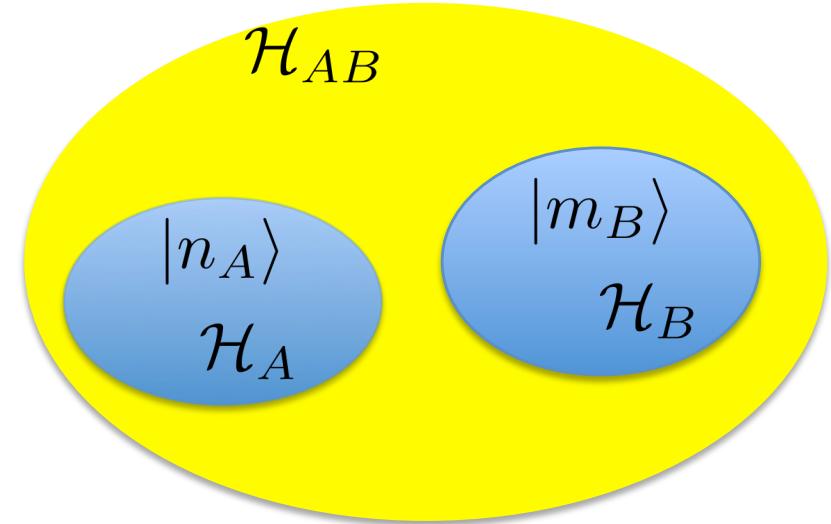
1. Basics of entanglement
2. Entanglement and 2-atom gate using Rybderg blockade

## What is entanglement?

2 systems A, B

2 d° of freedom

2 modes



Superposition principle  $\Rightarrow$   $|\phi_{AB}\rangle = \sum_{nm} c_{nm} |n_A\rangle \otimes |m_B\rangle$

Separable:  $|\phi_{AB}\rangle = |\phi_A\rangle \otimes |\phi_B\rangle$

Non-separable = **entangled**:  $|\phi_{AB}\rangle \neq |\phi_A\rangle \otimes |\phi_B\rangle$

**MOST states are entangled!!**

## Two-level systems

$|\uparrow\rangle$

2-level systems: atom, quantum circuits,  
spin, polarization of photon

$|\downarrow\rangle$

**Rotation** (e.g. polarizer, Raman, microwaves...)

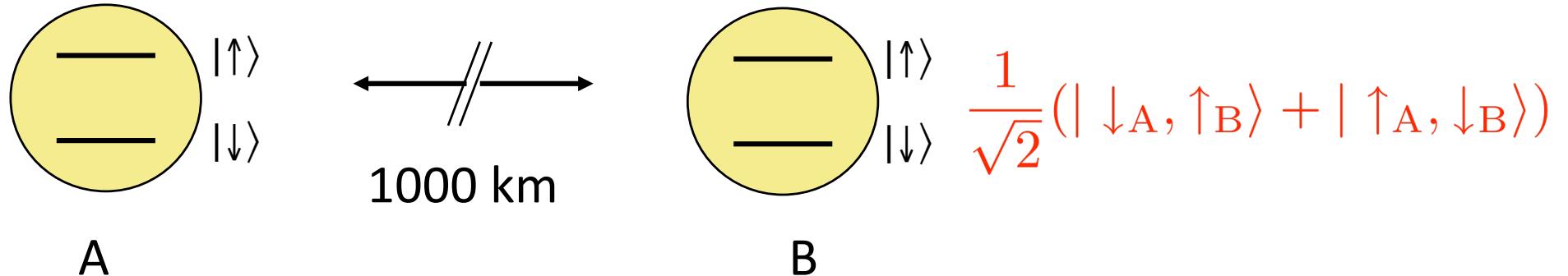
$$R(\theta, \varphi) = \begin{pmatrix} \cos \frac{\theta}{2} & ie^{i\varphi} \sin \frac{\theta}{2} \\ ie^{-i\varphi} \sin \frac{\theta}{2} & \cos \frac{\theta}{2} \end{pmatrix}_{|\uparrow\rangle, |\downarrow\rangle} \quad \begin{matrix} \Omega \text{ Rabi frequency} \\ \theta = \Omega t \end{matrix}$$

**Bell states**

$$|\psi_{\pm}\rangle = \frac{1}{\sqrt{2}} (|\uparrow, \downarrow\rangle \pm |\downarrow, \uparrow\rangle)$$

$$|\phi_{\pm}\rangle = \frac{1}{\sqrt{2}} (|\uparrow, \uparrow\rangle \pm |\downarrow, \downarrow\rangle)$$

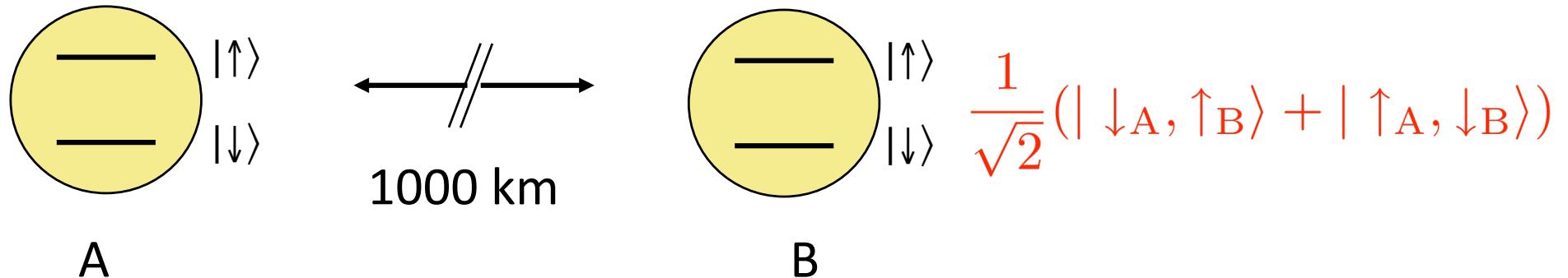
## What is so special about entanglement?



1. Perfect **correlations** (even at very large distance...!)

$$P(\uparrow_A \mid \downarrow_B) = \underbrace{\frac{P_{\uparrow_A, \downarrow_B}}{P_{\downarrow_B}}}_{\text{Cond. proba.}} = \frac{|\langle \uparrow, \downarrow | \psi_+ \rangle|^2}{P_{\downarrow_B}} = \frac{1/2}{1/2} = 1$$

## What is so special about entanglement?



1. Perfect **correlations** (even at very large distance...!)

$$P(\downarrow_A \mid \uparrow_B) = 1$$

$$P(\downarrow_A \mid \downarrow_B) = 0$$

$$P(\uparrow_A \mid \downarrow_B) = 1$$

$$P(\uparrow_A \mid \uparrow_B) = 0$$

2. You CAN NOT assigne a state to A or B = **system as a whole**

$$\rho_A = \text{Tr}_B \rho_{AB} = \frac{1}{2} |\uparrow\rangle\langle\uparrow| + \frac{1}{2} |\downarrow\rangle\langle\downarrow| \quad \text{Statistical mixture}$$

3. You CAN NOT “dis-entangled” by a **LOCAL** rotation

**Non Local Quantum correlations**

## Classification of entangled states

**Bipartite**  $\mathcal{H}_{AB} = \mathcal{H}_A \otimes \mathcal{H}_B$

Pure state  $|\phi_{AB}\rangle \neq |\phi_A\rangle \otimes |\phi_B\rangle$

Mixed state  $\rho \neq \sum p_i |a_i\rangle\langle a_i| \otimes |b_i\rangle\langle b_i| \quad (p_i \geq 0)$

**Multipartite**  $\mathcal{H}_{AB} = \mathcal{H}_A \otimes \mathcal{H}_B \otimes \mathcal{H}_C \otimes \dots$

Pure state  $|\phi_{ABC\dots}\rangle \neq |\phi_A\rangle \otimes |\phi_B\rangle \otimes |\phi_C\rangle \otimes \dots$

Bi-separable  $|\phi_{ABC\dots}\rangle = |\phi_1\rangle \otimes |\phi_2\rangle$   
 $\underbrace{\phantom{\phi_1}}_n \quad \underbrace{\phantom{\phi_2}}_{N-n}$

**Fully entangled:** not bi-separable w/r to any bi-partition of the system

## Examples of inequivalent entangled states

GHZ-state

$$|\text{GHZ}\rangle = \frac{1}{\sqrt{2}} (|\uparrow, \uparrow, \uparrow\rangle + |\downarrow, \downarrow, \downarrow\rangle)$$

W-state

$$|W\rangle = \frac{1}{\sqrt{3}} (|\uparrow, \downarrow, \downarrow\rangle + |\downarrow, \uparrow, \downarrow\rangle + |\downarrow, \downarrow, \uparrow\rangle)$$

**No LOCAL** rotation can transform GHZ state into W state

⇒ defines 2 **inequivalent classes** of entanglement

2 systems

1 classe

3 systems

2 classes    PRA **62**, 062314 (2000)

4 systems

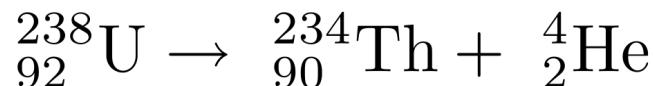
9 classes    PRA **65**, 052112 (2002)

Beyond ...?

## Preparation of entangled states (1)

### 1. Conservation laws

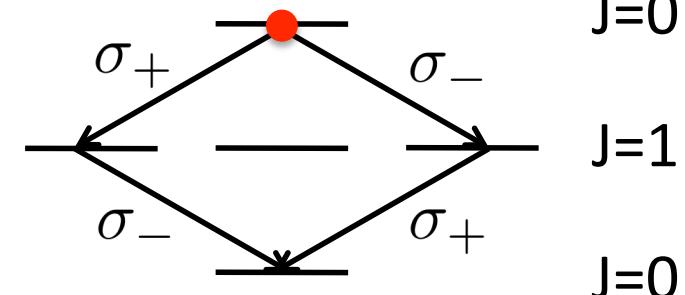
Radio-active decay



$\alpha$    

$$\frac{1}{\sqrt{2}} (|p, -p\rangle + | - p, p\rangle)$$

Atomic cascade (e.g. Ca)



$$\frac{1}{\sqrt{2}} (|\sigma_+, \sigma_-\rangle + |\sigma_-, \sigma_+\rangle)$$

### 2. Some “symmetry”: e.g. identical particles

Bosons 1 and 2 in two states  $\phi_a$  and  $\phi_b$

$$\phi_{\text{boson}}(1, 2) = \frac{1}{\sqrt{2}} (\phi_a(1)\phi_b(2) + \phi_a(2)\phi_b(1))$$

Fermions: Pauli principle

$$\phi_{\text{fermion}}(1, 2) = \frac{1}{\sqrt{2}} (\phi_a(1)\phi_b(2) - \phi_a(2)\phi_b(1))$$

## Preparation of entangled states (2)

### 3. Interaction between the sub-parts

$$H = H_A + H_B + H_{\text{int}}$$

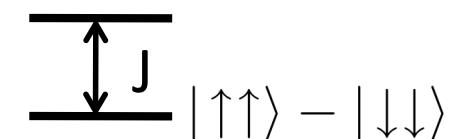
$$|\psi_A\rangle \otimes |\psi_B\rangle \xrightarrow[\text{evolution}]{e^{-i\frac{\hat{H}t}{\hbar}}} (|\psi_A\rangle \otimes |\psi_B\rangle) \neq |\psi'_A\rangle \otimes |\psi'_B\rangle$$

Examples of  $H_{\text{int}}$ :

- Spin orbit coupling  $C \vec{L} \cdot \vec{S}$
- Hyperfine structure  $A \vec{J} \cdot \vec{I}$  (ground state of H)

Magnetic interaction  $\frac{\vec{\mu}_A \cdot \vec{\mu}_B}{R^3}$

Generic Heisenberg  $H_{\text{int}}$ :  $J \vec{S}_a \cdot \vec{S}_b$  e.g. spin 1/2



Generation of “cluster” states:  $H_{\text{int}} = J(1 + \sigma_z^{(a)})(1 + \sigma_z^{(b)})$  (Ising)

$$(|\uparrow\rangle + |\downarrow\rangle)_A (|\uparrow\rangle + |\downarrow\rangle)_B \xrightarrow{} |\downarrow\downarrow\rangle + |\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle + e^{i\varphi} |\uparrow\uparrow\rangle \quad \varphi = \frac{4J}{\hbar} t$$

## Preparation of entangled states (3)

### 4. Projective measurement

Prepare

$$|\uparrow\uparrow\uparrow\uparrow\dots\rangle$$

Apply “small” rotation     $|\uparrow\rangle \rightarrow |\uparrow\rangle + \epsilon|\downarrow\rangle$      $\epsilon \ll 1$

$$|\uparrow\uparrow\uparrow\uparrow\dots\rangle \longrightarrow (|\uparrow\rangle + \epsilon|\downarrow\rangle)_A (|\uparrow\rangle + \epsilon|\downarrow\rangle)_B (|\uparrow\rangle + \epsilon|\downarrow\rangle)_C \dots$$

$$\propto |\uparrow\uparrow\uparrow\uparrow\dots\rangle + \epsilon (|\downarrow\uparrow\uparrow\uparrow\dots\rangle + |\uparrow\downarrow\uparrow\uparrow\dots\rangle + |\uparrow\uparrow\downarrow\uparrow\dots\rangle) + O(\epsilon^2)$$

Non-destructive measurement of  $|\downarrow\rangle$  (e.g. with cavity)

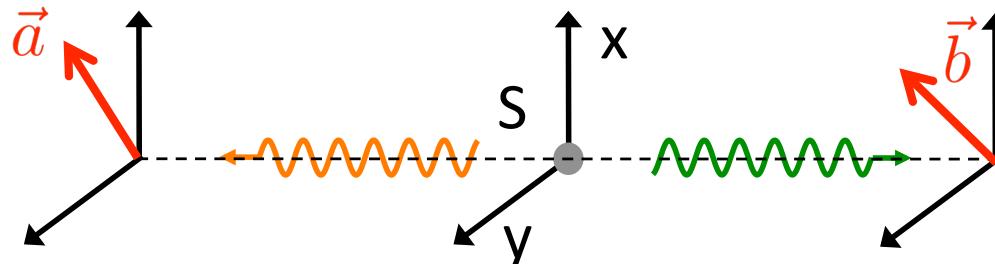
Project onto:  $|W\rangle = \frac{1}{\sqrt{N}} (|\downarrow\uparrow\uparrow\uparrow\dots\rangle + |\uparrow\downarrow\uparrow\uparrow\dots\rangle + |\uparrow\uparrow\downarrow\uparrow\dots\rangle)$

Heralded entanglement (probability  $\epsilon^2$ )

## Bi-partite entanglement and non-locality

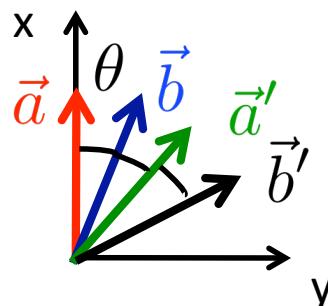
### Bell inequality tests (1964)

$$|\psi\rangle = \frac{1}{\sqrt{2}} (|xx\rangle + |yy\rangle)$$



$$\begin{aligned} E(\vec{a}, \vec{b}) &= \langle \psi | (\vec{\sigma}_A \cdot \vec{a}) \otimes (\vec{\sigma}_B \cdot \vec{b}) | \psi \rangle \\ &= -\cos(\vec{a}, \vec{b}) \end{aligned}$$

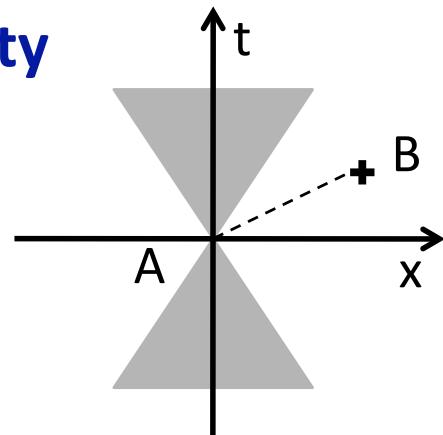
$$S(\theta) = E(\vec{a}, \vec{b}) + E(\vec{a}', \vec{b}) + E(\vec{a}', \vec{b}') - E(\vec{a}, \vec{b}') = \cos 3\theta - 3 \cos \theta$$



Entangled state  $\Rightarrow S(\pi/4) = -2\sqrt{2}$

Classical local correlations  $\Rightarrow |S(\theta)| \leq 2$

### Non-locality



$$x_{AB} > ct \Rightarrow A \text{ and } B \text{ not causally related}$$

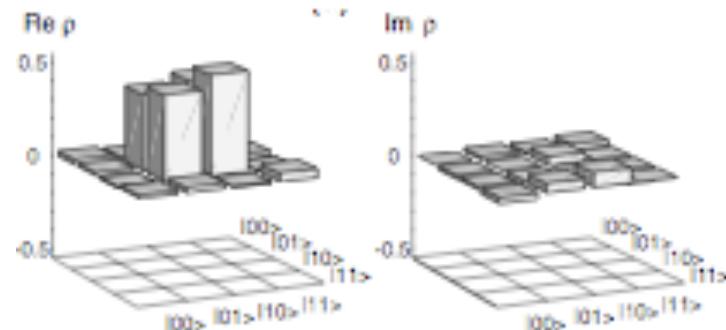
Entanglement between A and B possible

## Quantification of entanglement: quantum state tomography

For N atoms  $\hat{\rho} = \sum_{i_A, i_B, i_C \dots} \lambda_{i_A, i_B, i_C \dots} \sigma_{i_A} \otimes \sigma_{i_B} \otimes \sigma_{i_C} \dots$

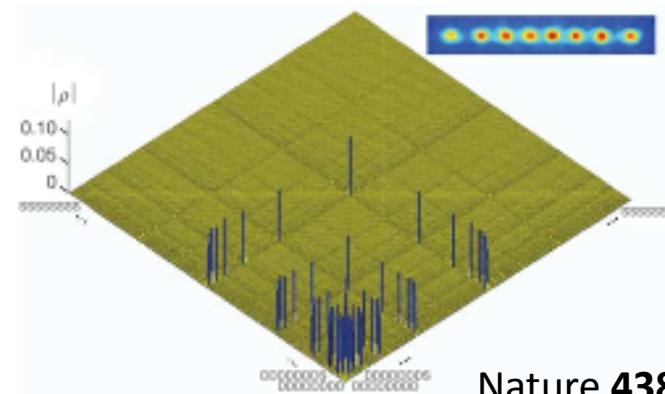
where  $\lambda_{i_A, i_B, i_C \dots} = \langle \sigma_{i_A} \otimes \sigma_{i_B} \otimes \sigma_{i_C} \dots \rangle$  ( $4^N - 1$  coefficients)

Measure  $\langle \sigma_{i_A} \otimes \sigma_{i_B} \otimes \sigma_{i_C} \dots \rangle \Rightarrow$  reconstruct the density matrix



e.g. Bell states of 2 ions  
PRL 92, 220402 (2004)

Le **largest state** characterized:  
8 ions W states, 65531 coefficients  
10 hours of measurements!  
Too hard for “large” systems



Nature 438, 643 (2005)

## Quantification of entanglement (1)

**Bipartite** ~OK: Schmidt decomposition, entropy, concurrence...

### Multipartite

Pure states: generalized Bell inequalities (PRL **104**, 240502 (2010))

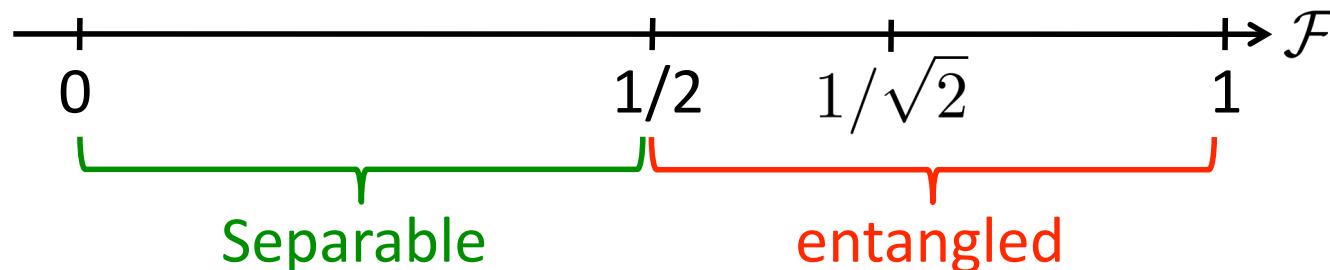
Mixed states: no general criteria (even if you know  $\rho$ )  
⇒ use weaker criteria

Fidelity: want  $|\psi_t\rangle$ , measure  $\rho_{\text{exp}}$

$$\mathcal{F} = \langle \psi_t | \rho_{\text{exp}} | \psi_t \rangle$$

Separable state ⇒  $\mathcal{F} \leq \frac{1}{2}$  C.A. Sackett *et al.*, Nature **404**, 256 (2000)

Quantum non-locality (Bell for 2 at.)



## Quantification of entanglement (2)

More generally: entanglement witness operator  $\hat{\mathcal{W}}$

separable state  $\Rightarrow \langle \hat{\mathcal{W}} \rangle = \text{Tr}(\rho_{\text{exp}} \hat{\mathcal{W}}) \geq 0$

$\langle \hat{\mathcal{W}} \rangle < 0 \Rightarrow$  entangled state

**Examples:**

1. fidelity  $\hat{\mathcal{W}} = 1 - 2|\psi_t\rangle\langle\psi_t| \Rightarrow \langle \hat{\mathcal{W}} \rangle = 1 - 2\mathcal{F}$

2. Energy of spin 1/2:  $H = -J\vec{S}_a \cdot \vec{S}_b$  ( $J > 0$ )

Separable:  $\langle H \rangle = -J\langle \vec{S}_a \rangle \cdot \langle \vec{S}_b \rangle \Rightarrow -\frac{J}{4} \leq \langle H \rangle \leq \frac{J}{4}$

Singlet state:  $\frac{1}{\sqrt{2}}(|\uparrow, \downarrow\rangle - |\downarrow, \uparrow\rangle) \Rightarrow \langle H \rangle = \frac{3J}{4}$

Choose  $\hat{\mathcal{W}} = 1 - \frac{4\hat{H}}{J}$       Useful for macroscopic measurements

## Entanglement is fragile...

**Atom loss**     $|\text{GHZ}_N\rangle = \frac{1}{\sqrt{2}} (|\uparrow\uparrow\uparrow\dots\rangle + |\downarrow\downarrow\downarrow\dots\rangle)$

→  $\rho_{N-1} = \frac{1}{2} \underbrace{|\uparrow\uparrow\uparrow\dots\rangle\langle\uparrow\uparrow\uparrow\dots|}_{N-1} + \frac{1}{2} |\downarrow\downarrow\downarrow\dots\rangle\langle\downarrow\downarrow\downarrow\dots|$

GHZ strongly correlated

While:  $|W_N\rangle \longrightarrow |W_{N-1}\rangle$      $W$  weakly correlated

## Phase fluctuations (fluctuating B, E field)

$$|\text{GHZ}(t)\rangle = \frac{1}{\sqrt{2}} (|\uparrow\uparrow\uparrow\dots\rangle + e^{-iN\omega t} |\downarrow\downarrow\downarrow\dots\rangle)$$

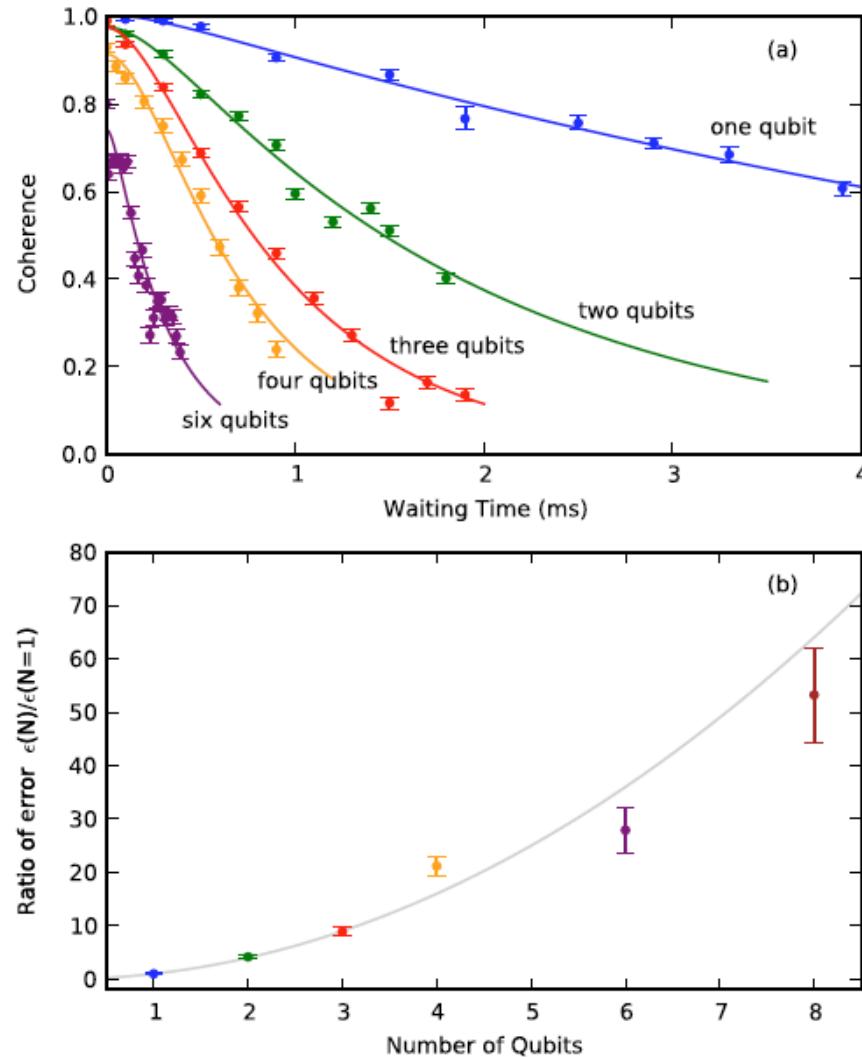
$$\rho_{\uparrow\uparrow\dots,\downarrow\downarrow\dots} = e^{-iN\omega t} \quad \text{and} \quad \langle e^{-iN\omega t} \rangle \propto e^{-N^2 \frac{\Delta\omega^2 t^2}{2}} \rightarrow 0$$

# Decoherence faster when **N** large

# Decoherence of a GHZ state

PRL **106**, 130506 (2011)

N-ion string



## But you can protect it...!

Start from  $|\text{GHZ}\rangle = \frac{1}{\sqrt{2}} (|\uparrow\uparrow\uparrow\uparrow\rangle + |\downarrow\downarrow\downarrow\downarrow\rangle)$

Energy 0      Energy  $4\hbar\omega_0$

Apply local rotation on “atom” B and D  
(does not change class of entanglement)

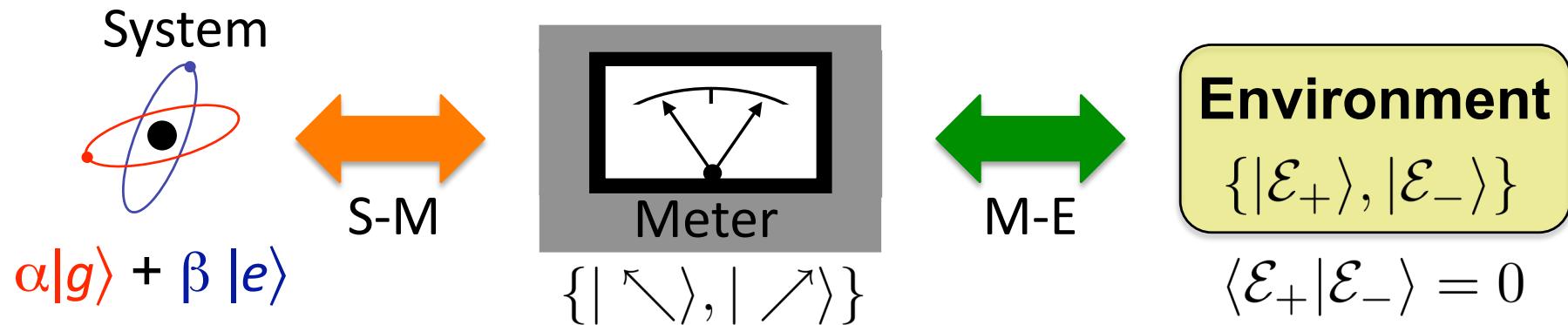
$\longrightarrow \frac{1}{\sqrt{2}} (|\uparrow\downarrow\uparrow\downarrow\rangle + |\downarrow\uparrow\downarrow\uparrow\rangle)$

$2\hbar\omega_0$        $2\hbar\omega_0$

evolution  $\longrightarrow \frac{1}{\sqrt{2}} (|\uparrow\downarrow\uparrow\downarrow\rangle + |\downarrow\uparrow\downarrow\uparrow\rangle) e^{-i2\omega_0 t} \Rightarrow$  insensitive to fluctuation of  $\omega_0$

Decoherence free subspace

## Entanglement, measurement and classicality



$$(\alpha|g\rangle + \beta|e\rangle)_S \otimes |0\rangle_M \otimes |\mathcal{E}_0\rangle_E \xrightarrow{\text{orange arrow}} (\alpha|g, \swarrow\rangle + \beta|e, \nearrow\rangle)_{SM} \otimes |\mathcal{E}_0\rangle_E$$

$$\xrightarrow{\text{green arrow}} |\psi\rangle_{SME} = \alpha|g, \swarrow, \mathcal{E}_+\rangle + \beta|e, \nearrow, \mathcal{E}_-\rangle$$

Entanglement S – M – E

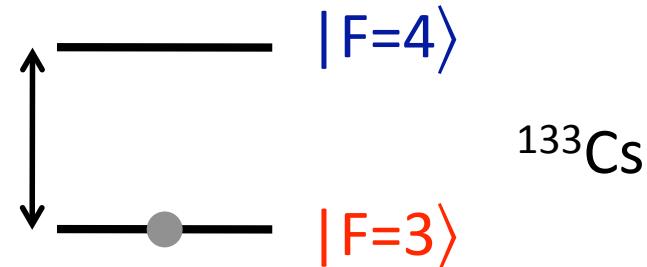
$$\rho_{SM} = \text{Tr}_{\mathcal{E}}(|\psi\rangle\langle\psi|) = |\alpha|^2|g, \swarrow\rangle\langle g, \swarrow| + |\beta|^2|e, \nearrow\rangle\langle e, \nearrow|$$

**Statistical mixture:** get  $|g\rangle$  (proba  $|\alpha|^2$ ) **OR**  $|e\rangle$  (proba  $|\beta|^2$ )

## Entanglement = resource for quantum metrology

### Frequency standard

$$\omega_0 = 9\ 192\ 631\ 770 \text{ Hz}$$



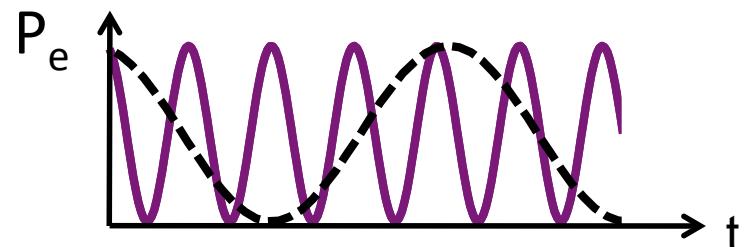
**Applications:** test of relativity, navigation using GPS...

Today, best atomic clock (ion based)  $\Delta\nu/\nu \sim 10^{-17}$

We can do better **using entanglement!**

Prepare N atom in  $|ggg\dots\rangle + |eee\dots\rangle$

$$\rightarrow |ggg\dots\rangle + e^{-iN\omega_0 t} |eee\dots\rangle$$

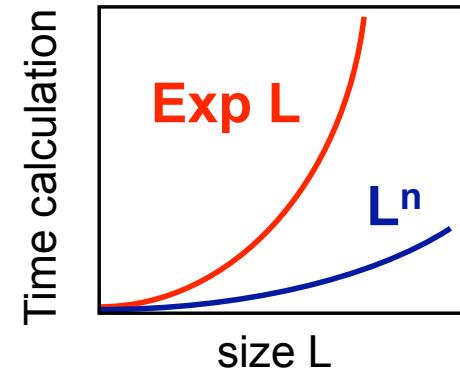


$\sqrt{N}$  more sensitive measurement

## Entanglement = resource for quantum computation

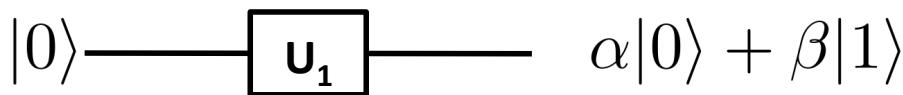
Help solving “hard” problems: factoring (Shor)  
searching (Grover)

Encode information on a **qubit**:  $|0\rangle$ ,  $|1\rangle$

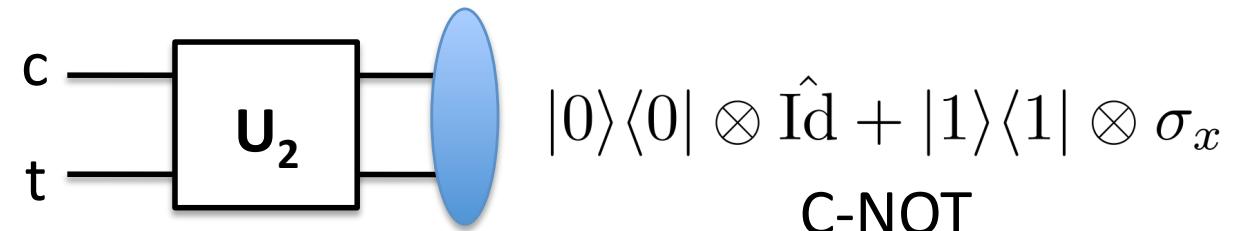


**Elementary bricks** (“circuit” approach):

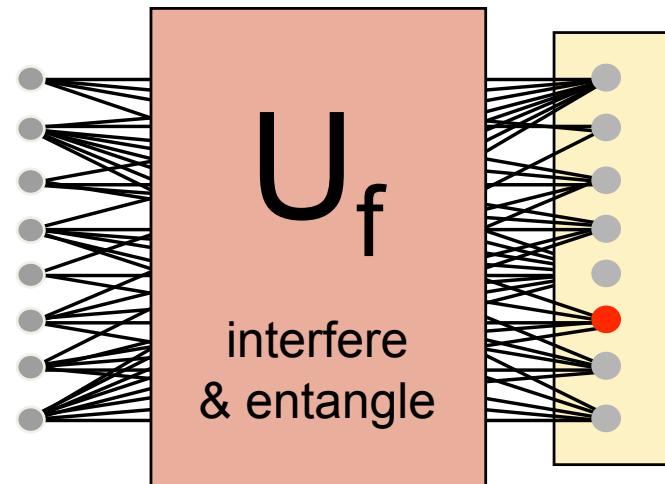
One-qubit gate



Two-qubit gate



$|0\rangle \otimes |0\rangle \otimes |0\rangle \dots$   
N qubits



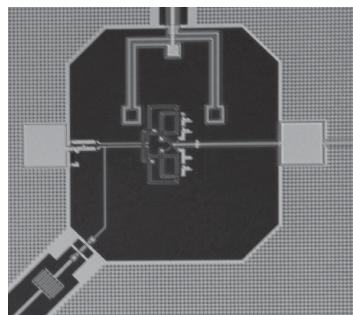
$|\Psi_f\rangle$  Highly  
entangled

## Engineered entangled states: state-of-the art (2013)



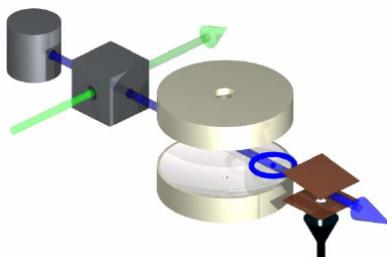
14 ions, GHZ ( $F=0.51$ )

PRL **106**, 130506 (2011)



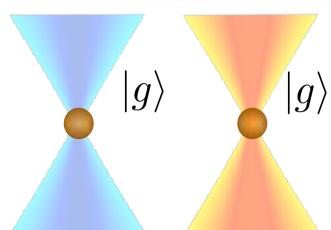
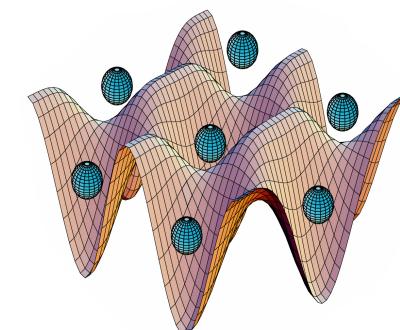
3 supracond. qubits W ( $F=0.78$ ), GHZ ( $F=0.62$ )

Nature **467**, 570, 574 (2010)



3 Rydberg atom GHZ ( $F=0.57$ )

Science **288**, 2024 (2000)



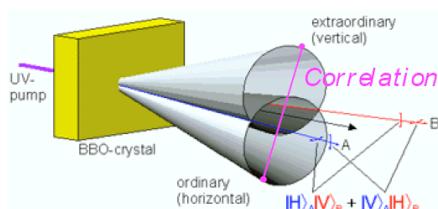
2 ground-state atoms

Bell ( $F\sim 0.75$ )

PRL **104**, 010502, 010503 (2010)

~ 50 atoms cluster state

Mandel *et al.*, **425**, 937 (2003)



8 photons (H,V) GHZ ( $F=0.71$ )

Nat. Phot. **6**, 227(2012)

## A short bibliography on entanglement

### Theory (from “easy” to hard...)

“Exploring the quantum”, S. Haroche and J-M. Raimond Cambridge

“Quantum computing and entanglement”, D. Bruss and C. Machiavello  
Les Houches summer school 2009 (Singapore)

“Experimental procedures for entanglement verification”,  
PRA **75**, 052318 (2007)

“Entanglement in many body system”, RMP **80**, 517 (2008)

### Experimental papers on quantification of entanglement

Q.A. Turchette *et al.*, PRL **81**, 3631 (1998)

C.A. Sackett *et al.*, Nature **404**, 256 (2000)

D. Leibfried *et al.*, Nature **438**, 639 (2005)

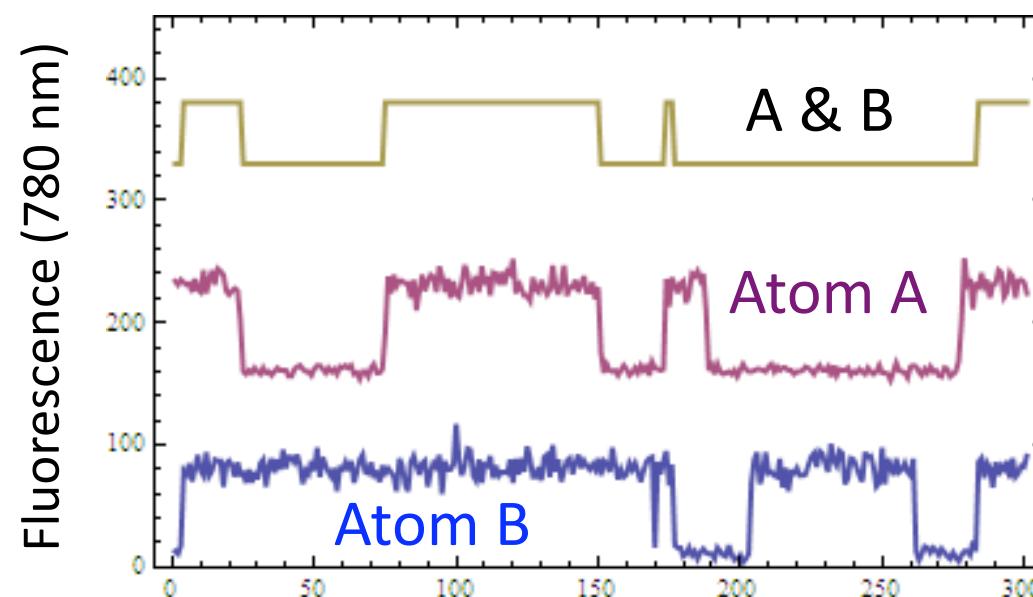
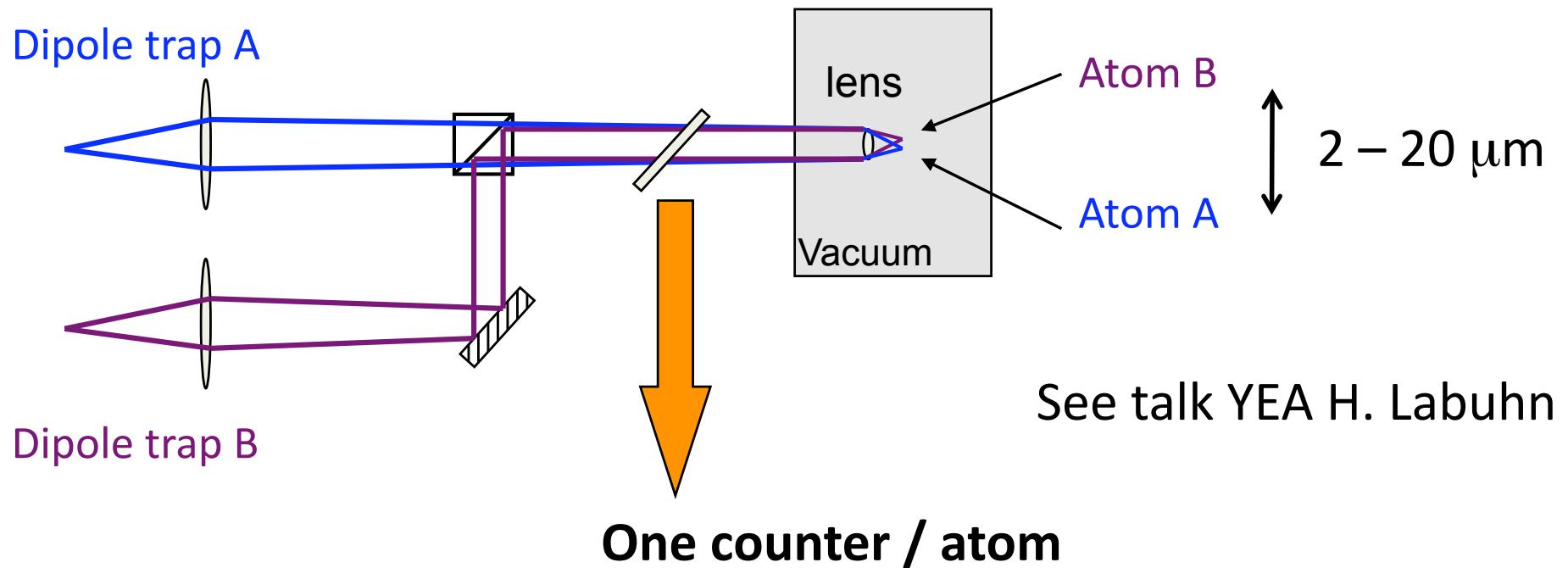
### Review on Rydberg and QIP

Saffman RMP **82**, 2313 (2010)

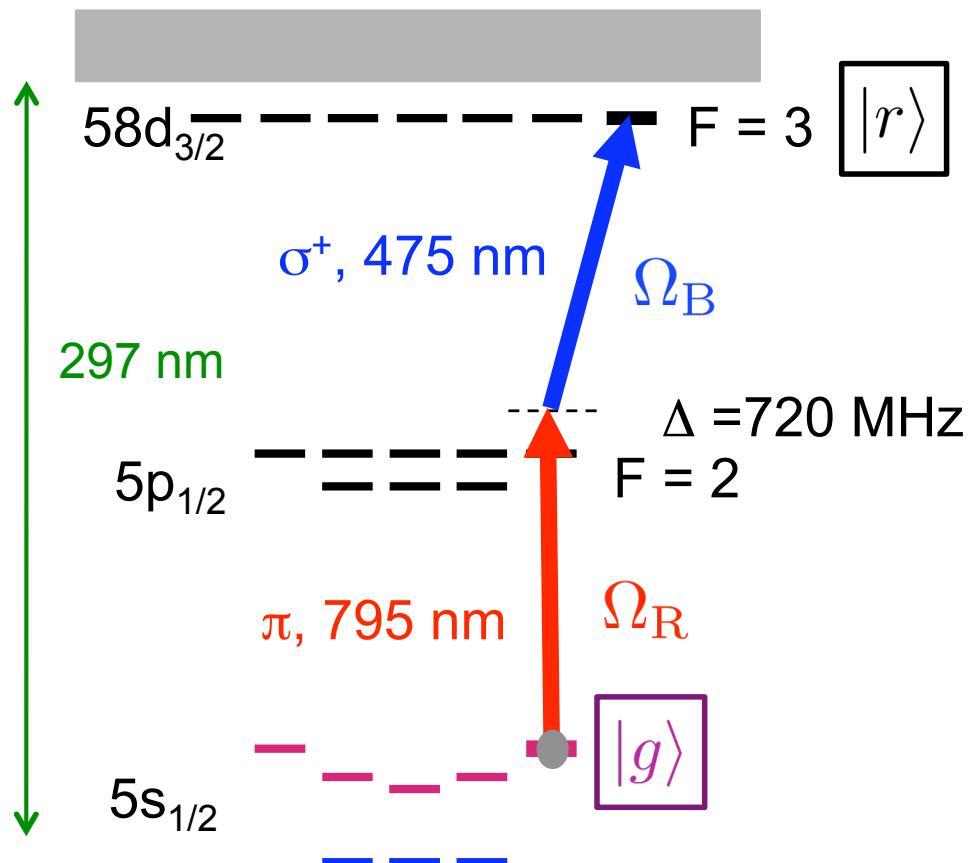
# **Outline**

1. Basics of entanglement
2. Entanglement and 2-atom gate using Rybderg blockade

## Trapping two atoms in tweezers: well-defined geometry



## Two-photon Rydberg excitation ( ${}^8\text{Rb}$ ): well-defined states



$$\Omega = \frac{\Omega_R \Omega_B}{2\Delta}$$

See talk YEA H. Labuhn

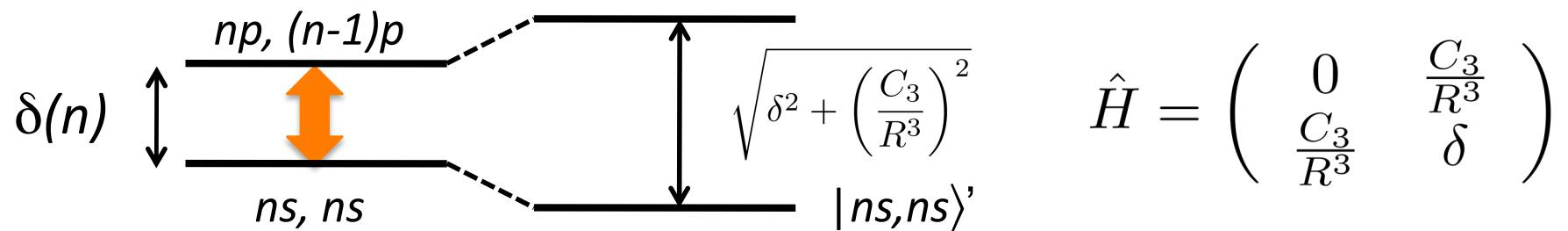
Detect atom loss = **Rydberg not trapped**

Miroshnychenko *et al.*, PRA **82**, 023623 (2010)

## vdW interaction between 2 atoms

2 atoms A and B with states  $ns$ ,  $(n-1)p$  and  $np$ ; distance  $R$

$$\hat{V} = -\frac{1}{4\pi\epsilon_0 R^3} \left( 2\hat{\mathbf{d}}_z^A \hat{\mathbf{d}}_z^B + \hat{\mathbf{d}}_+^A \hat{\mathbf{d}}_-^B + \hat{\mathbf{d}}_-^A \hat{\mathbf{d}}_+^B \right)$$



$$\langle nsjm; nsjm | \hat{V} | npj'm'; (n-1)pj''m'' \rangle = \frac{C_3}{R^3}$$

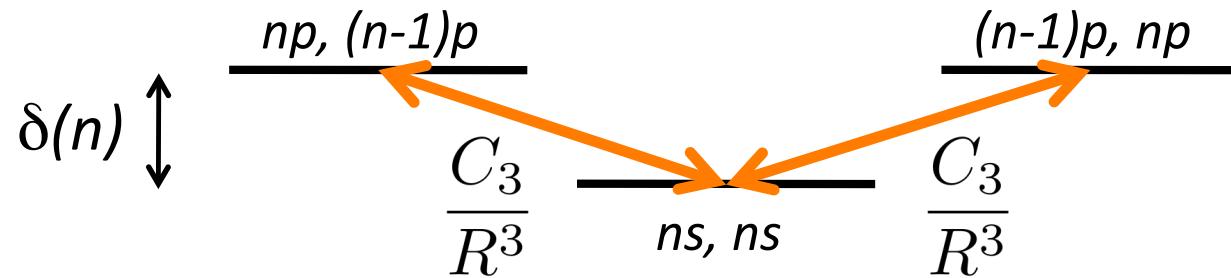
Ex.:  $n = 60$   
 $\delta = 5\text{-}10 \text{ GHz}, \quad \Rightarrow \quad \Delta E_{ns,ns} = \frac{1}{2\delta} \left( \frac{C_3}{R^3} \right)^2 = \frac{C_6}{R^6}$   
 $C_3/R^3 \approx 100 \text{ MHz}$

$$|ns, ns\rangle' = |ns, ns\rangle + \frac{1}{\delta} \frac{C_3}{R^3} |np, (n-1)p\rangle \approx |ns, ns\rangle$$

In the **vdW approximation**: shift of  $|r,r\rangle$  only...

## A note on vdW interaction between 2 atoms: role of symmetry

There are 3 two-atom states!



## A note on vdW interaction between 2 atoms: role of symmetry

There are 3 two-atom states!

$$\frac{1}{\sqrt{2}}(|np, (n-1)p\rangle + |(n-1)p, np\rangle) \quad \frac{1}{\sqrt{2}}(|np, (n-1)p\rangle - |(n-1)p, np\rangle)$$

The diagram illustrates the energy levels of a two-atom system. It features three horizontal black lines representing different states. On the left, there is a double-headed vertical arrow labeled  $\delta(n)$ , indicating a separation between two levels. In the center, there is a single horizontal line labeled  $\frac{C_3}{R^3} \sqrt{2}$  below it, with the text "ns, ns" written underneath. On the right, there is another horizontal line labeled "0" inside a small orange square. Orange arrows point from the central line to both the top and bottom lines, and from the bottom line to the "0" state.

## A note on vdW interaction between 2 atoms: role of symmetry

There are 3 two-atom states!

$$\frac{1}{\sqrt{2}}(|np, (n-1)p\rangle + |(n-1)p, np\rangle)$$
$$\frac{C_3}{R^3} \sqrt{2}$$
$$ns, ns$$
$$\delta(n)$$

Still two-state model, but:  $\frac{C_3}{R^3} \rightarrow \frac{C_3}{R^3} \sqrt{2}$

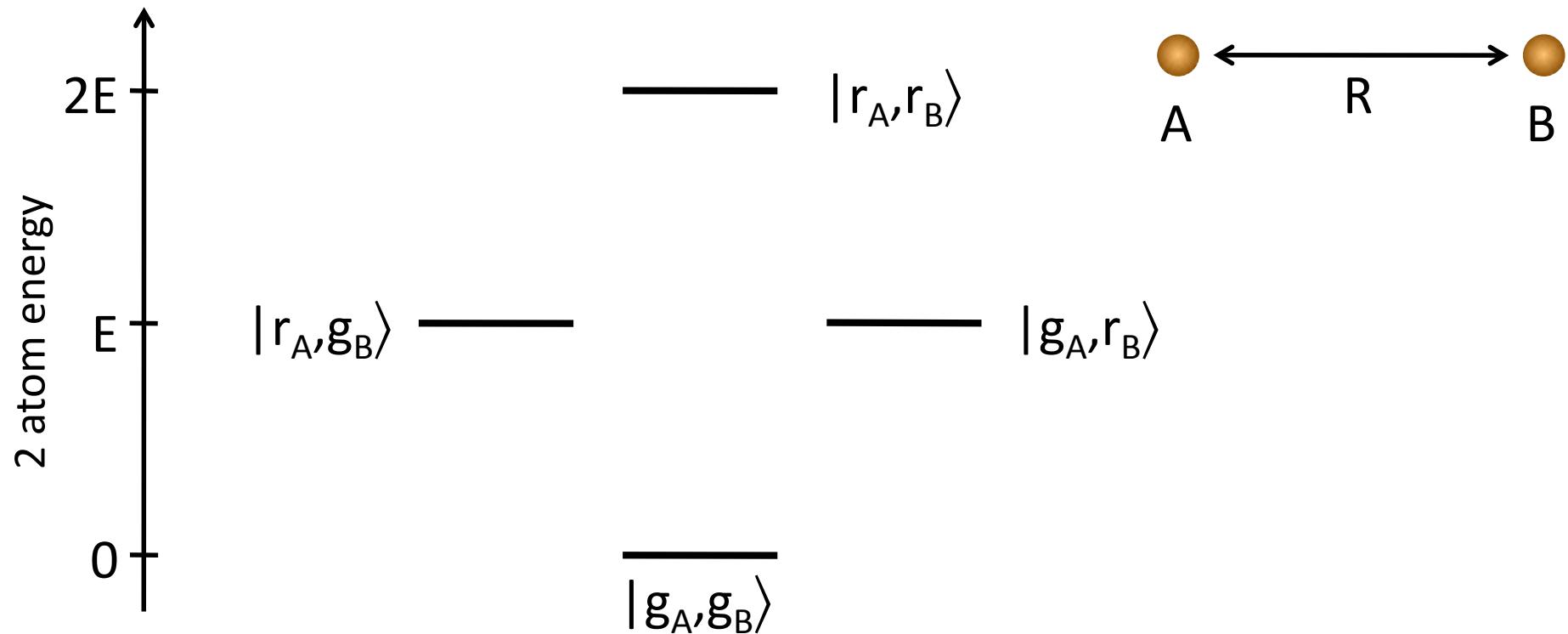
Typical interaction strength:  $C_6(43s) \approx -2.4 \text{ GHz} \cdot \mu\text{m}^6$

$C_6(53d) \approx 15 \text{ GHz} \cdot \mu\text{m}^6$

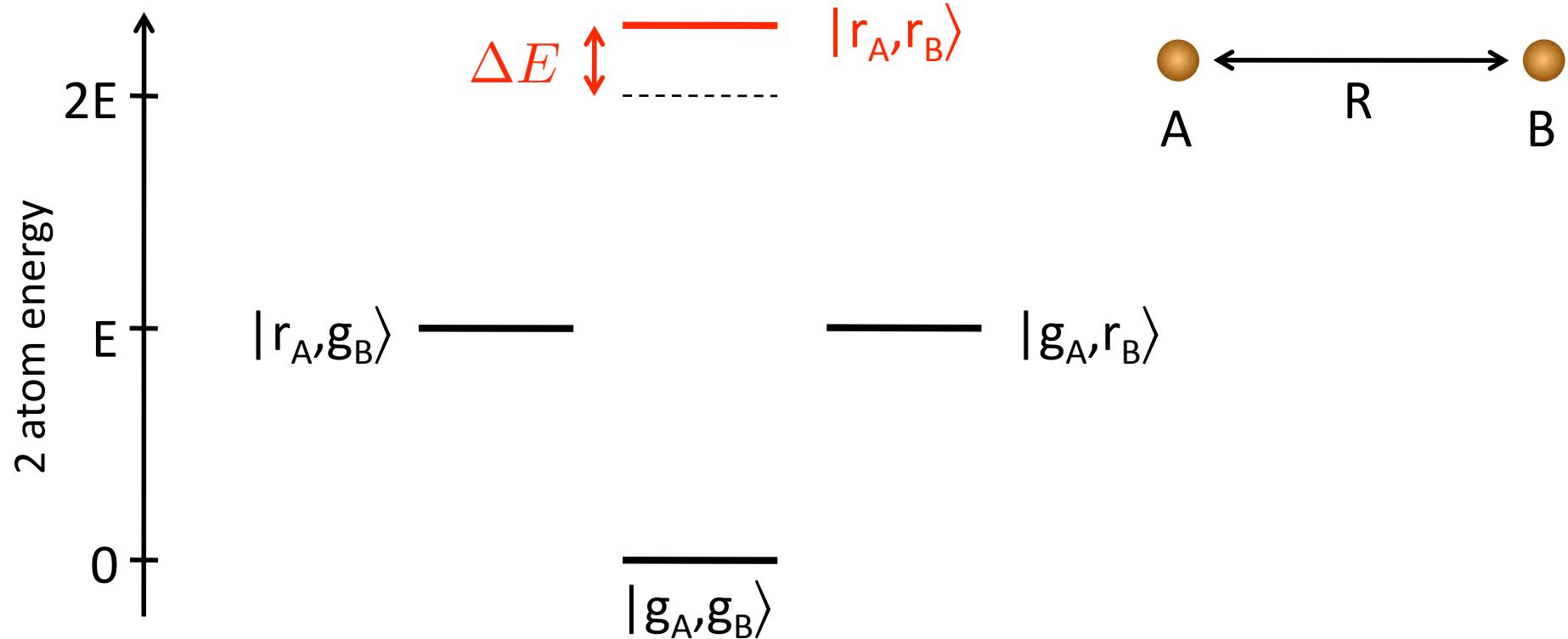
$C_6(82d) \approx 8500 \text{ GHz} \cdot \mu\text{m}^6$

$\Rightarrow V_{\text{vdw}} \sim 1 - 100 \text{ MHz for } R \sim \text{few } \mu\text{m}$

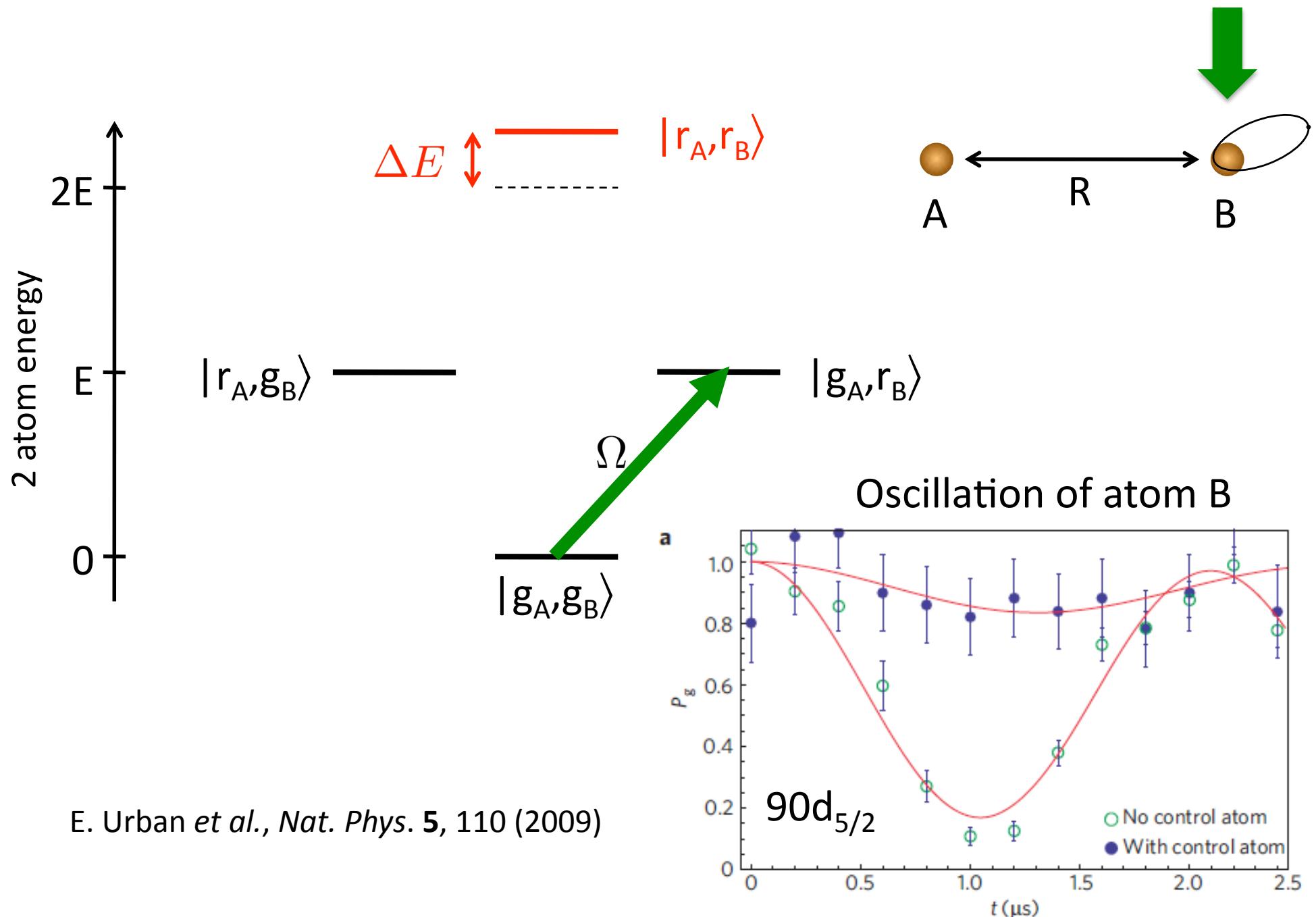
## Rydberg blockade: “addressable” version (U. Wisconsin)



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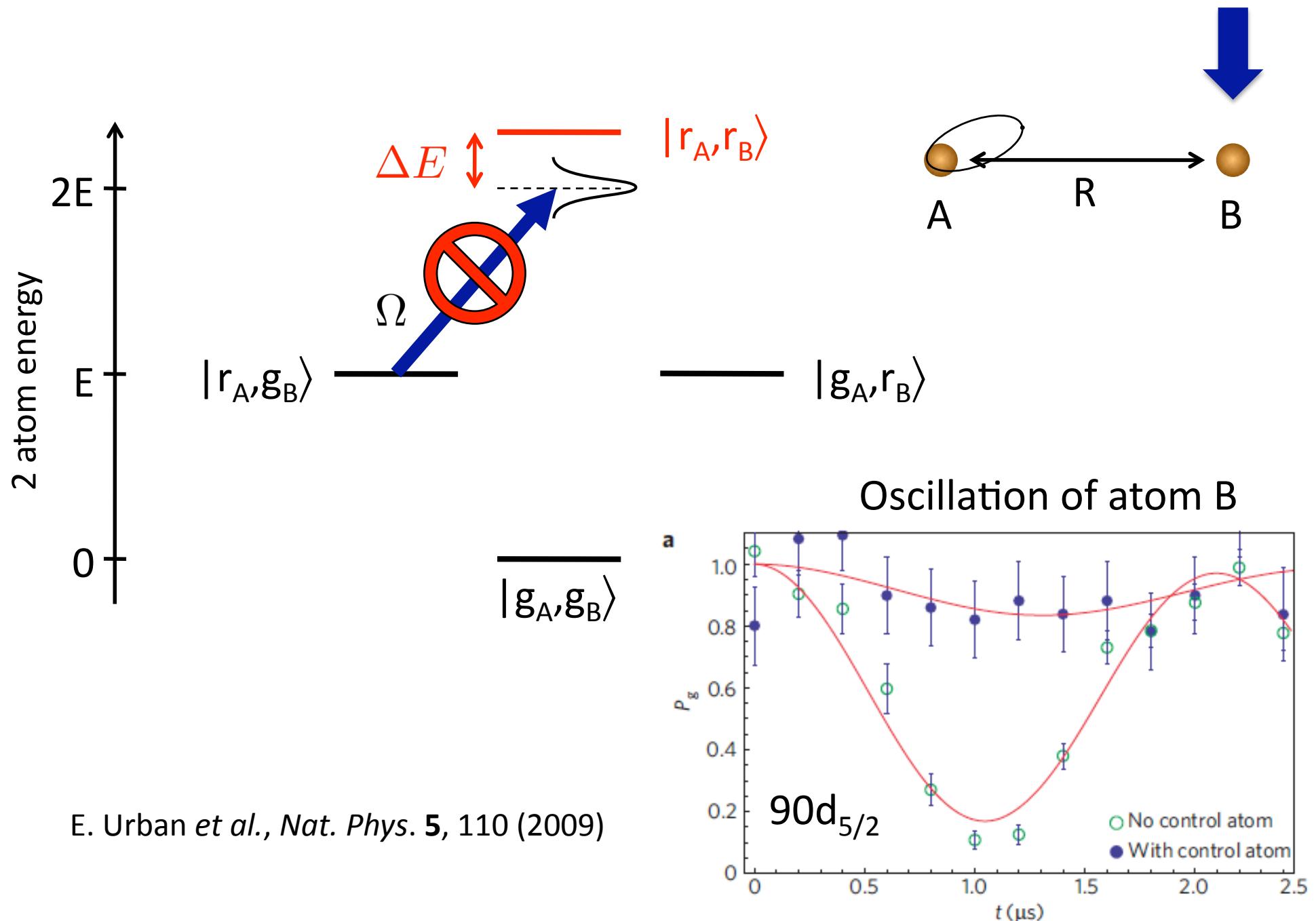


# Rydberg blockade: “addressable” version (U. Wisconsin)



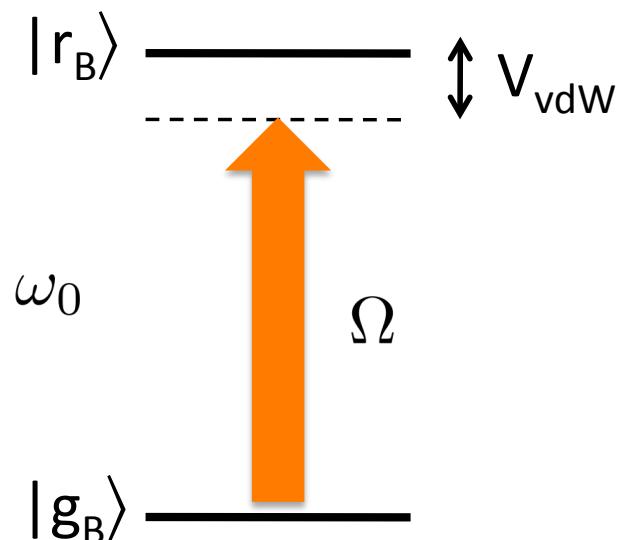
E. Urban *et al.*, Nat. Phys. 5, 110 (2009)

## Rydberg blockade: “addressable” version (U. Wisconsin)



## Rydberg blockade: “addressable” version

The “blockade radius” picture

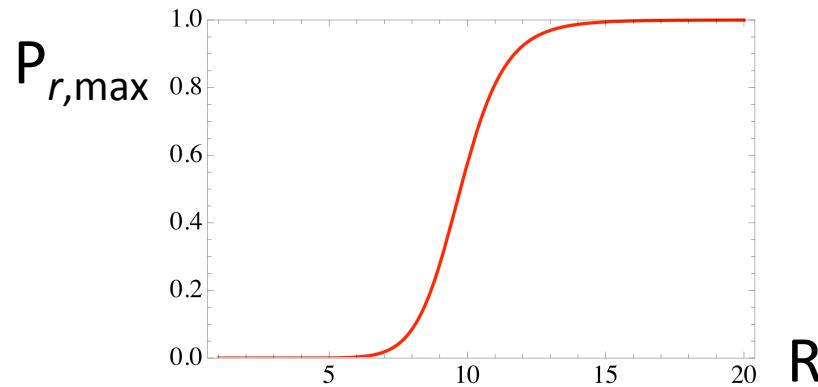


Atome B when  
A in  $|r_A\rangle$

Rabi formula:

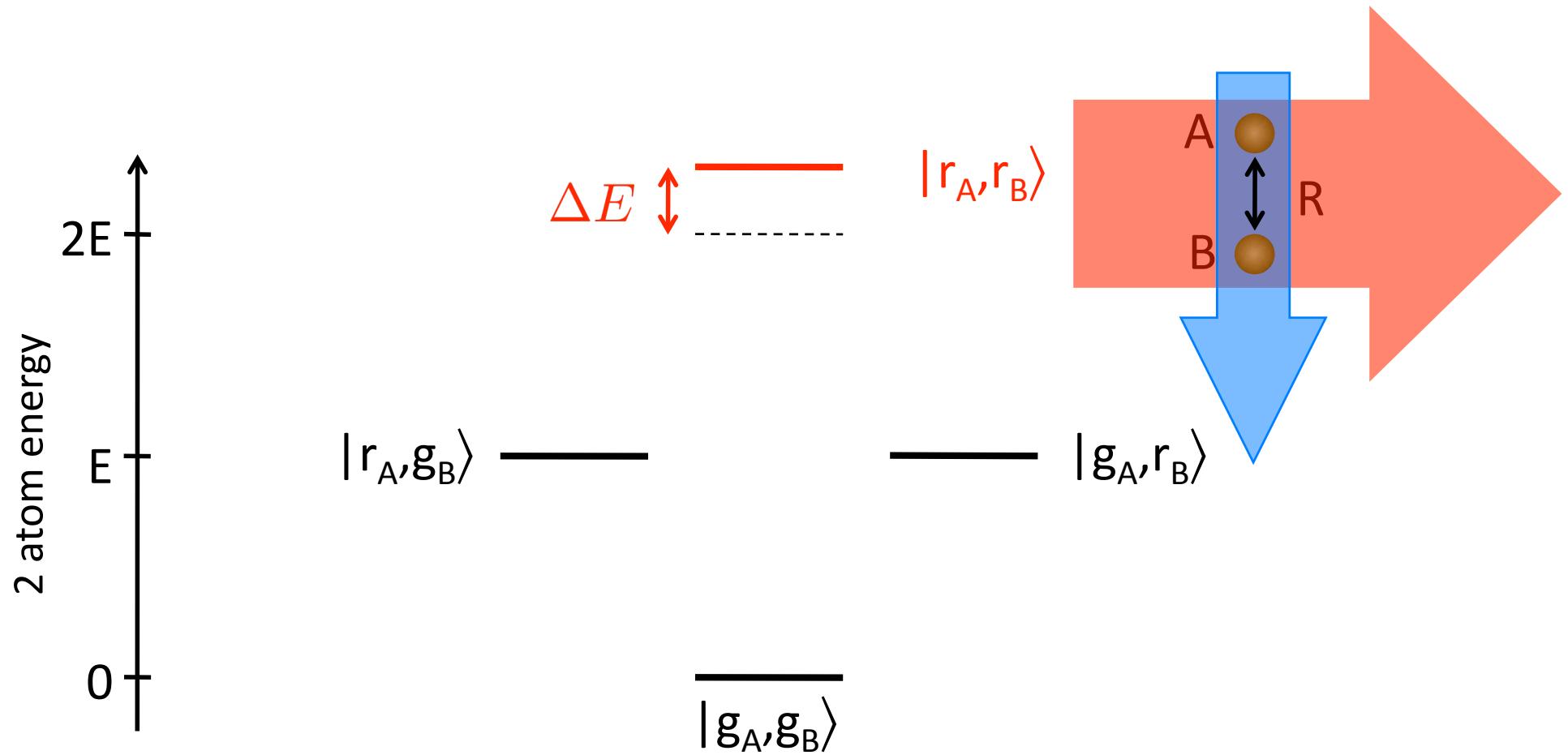
$$P_r(t) = \frac{\Omega^2}{\Omega^2 + V_{\text{vdW}}^2} \sin^2 \frac{\sqrt{\Omega^2 + V_{\text{vdW}}^2}}{2} t$$

Defines  $R_b$ :  $\hbar\Omega = V_{\text{vdW}}(R_b)$

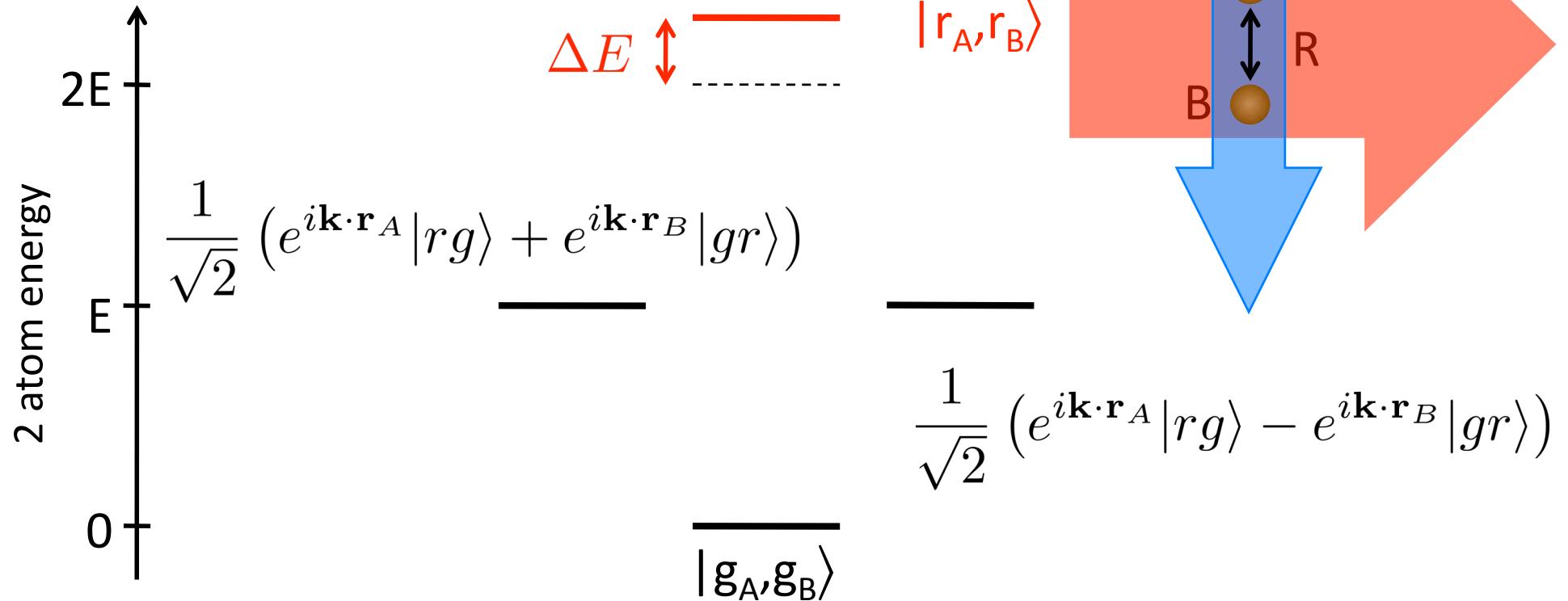


Blockade alone **does not** provide entanglement :  
need indistinguishability w/r excitation  
 $\Rightarrow$  **collective excitation**

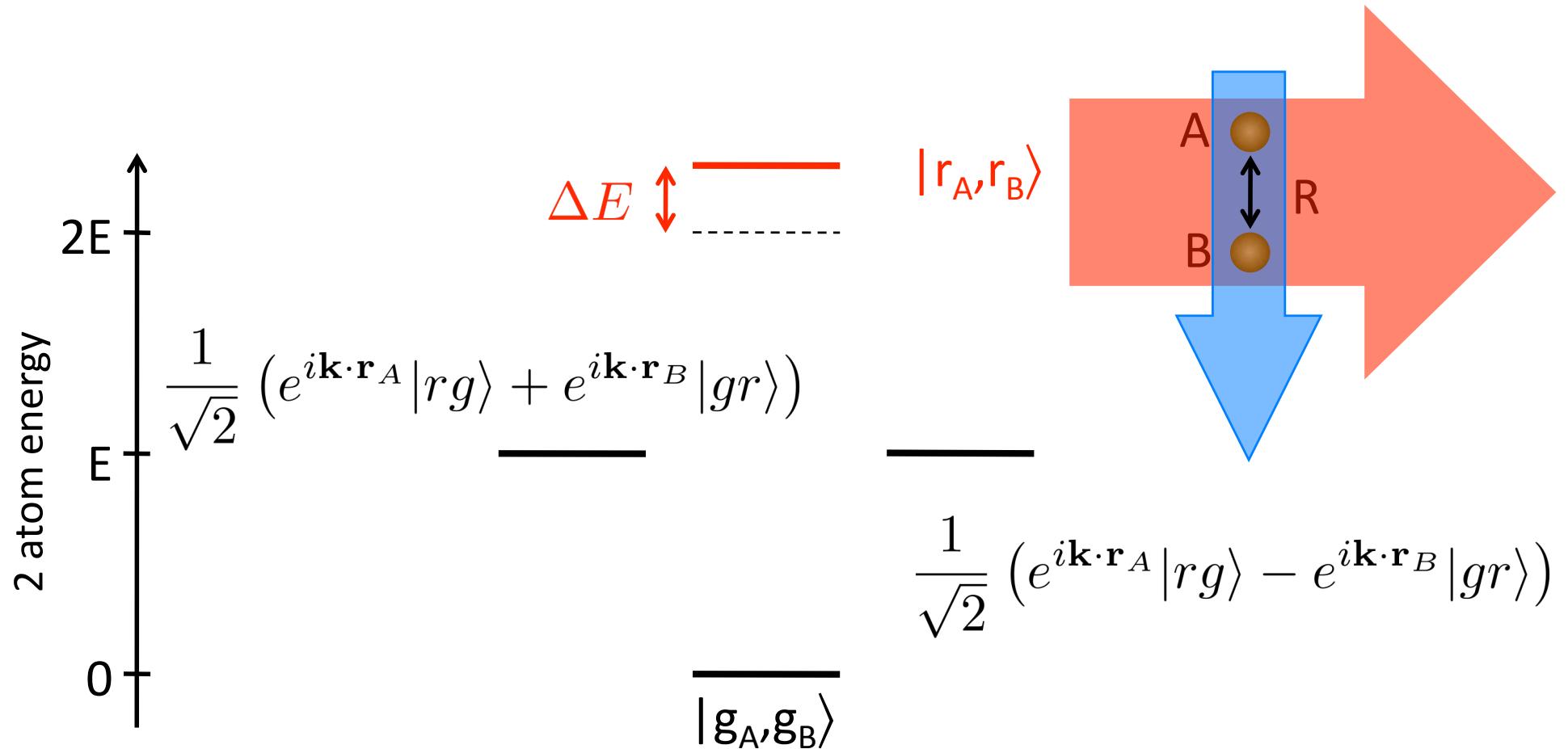
## Rydberg blockade: collective excitation (IO Palaiseau)



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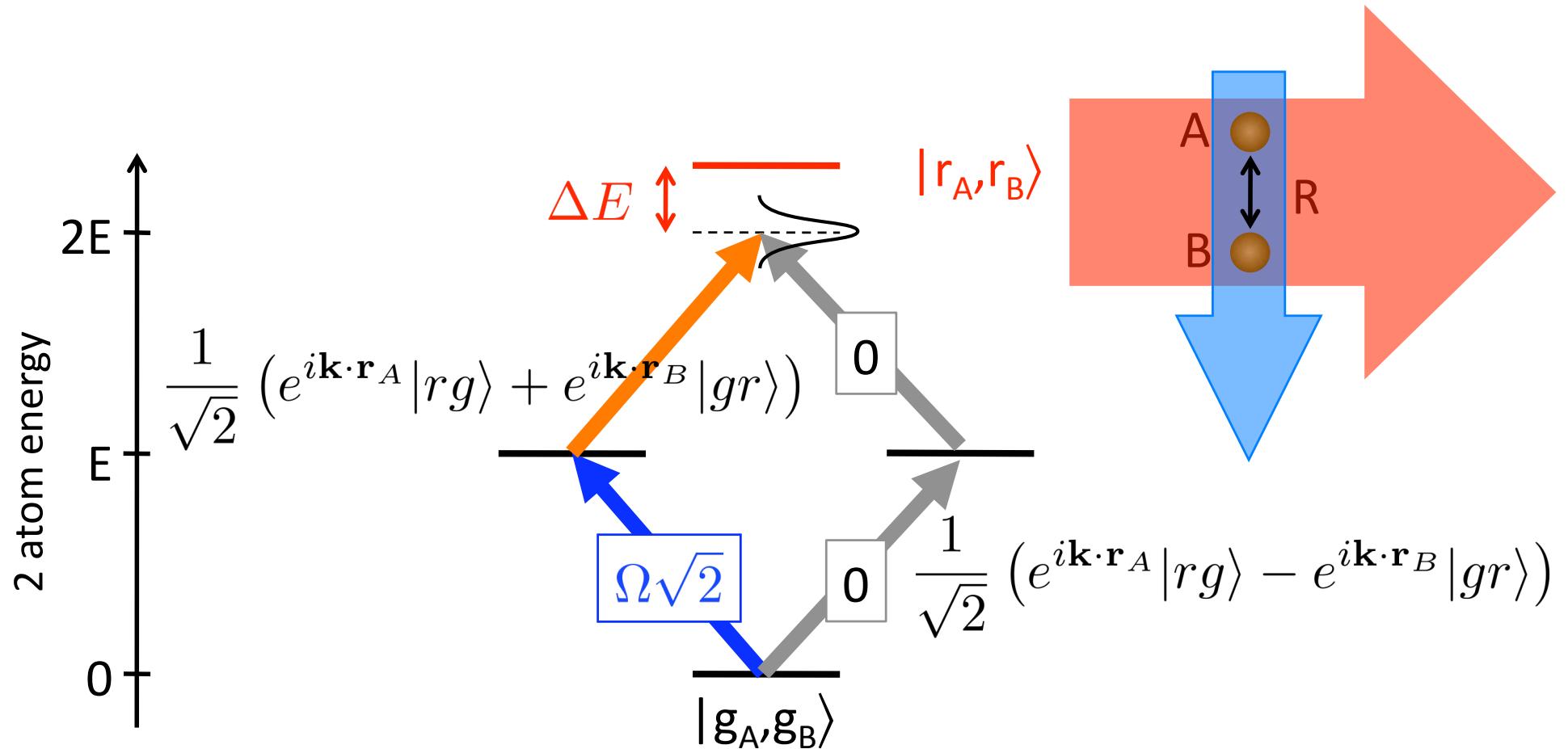
## Rydberg blockade: collective excitation (IO Palaiseau)



Recall (RW approx.):  $H_{\text{at-L}}^j = \frac{\hbar\Omega}{2} (|r\rangle\langle g| e^{i\mathbf{k} \cdot \mathbf{r}_j} + |g\rangle\langle r| e^{-i\mathbf{k} \cdot \mathbf{r}_j})$

$$H_{\text{at-L}} = H_{\text{at-L}}^A \otimes \hat{\text{Id}}_B + \hat{\text{Id}}_A \otimes H_{\text{at-L}}^B$$

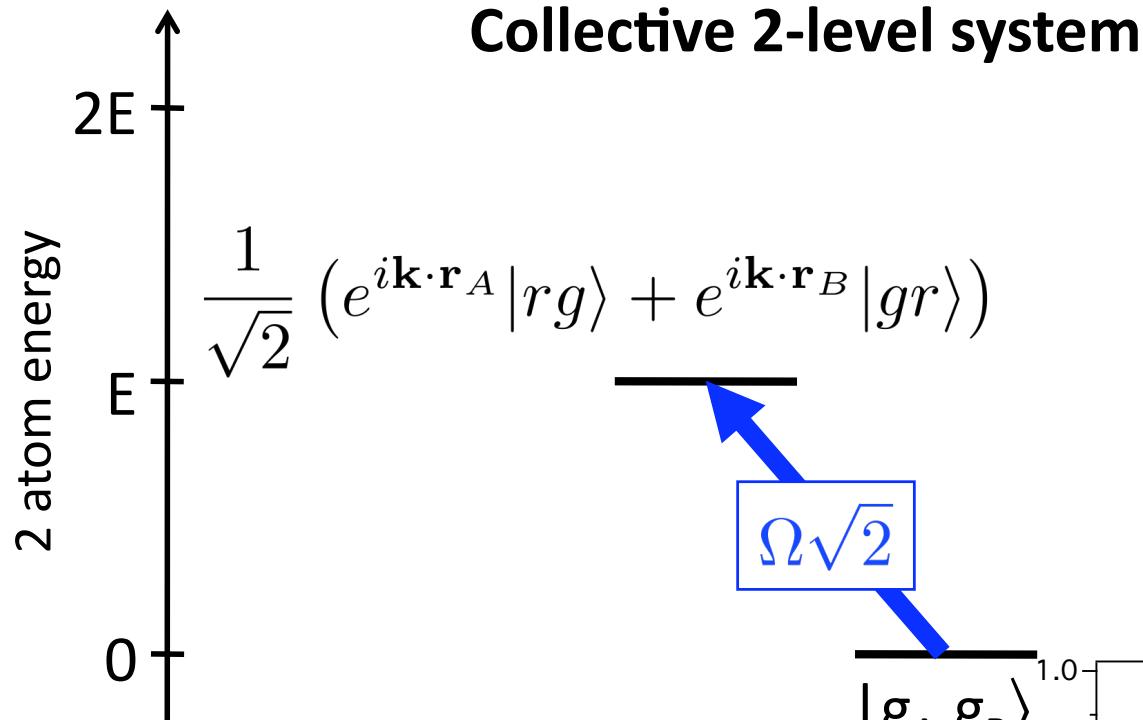
## Rydberg blockade: collective excitation (IO Palaiseau)



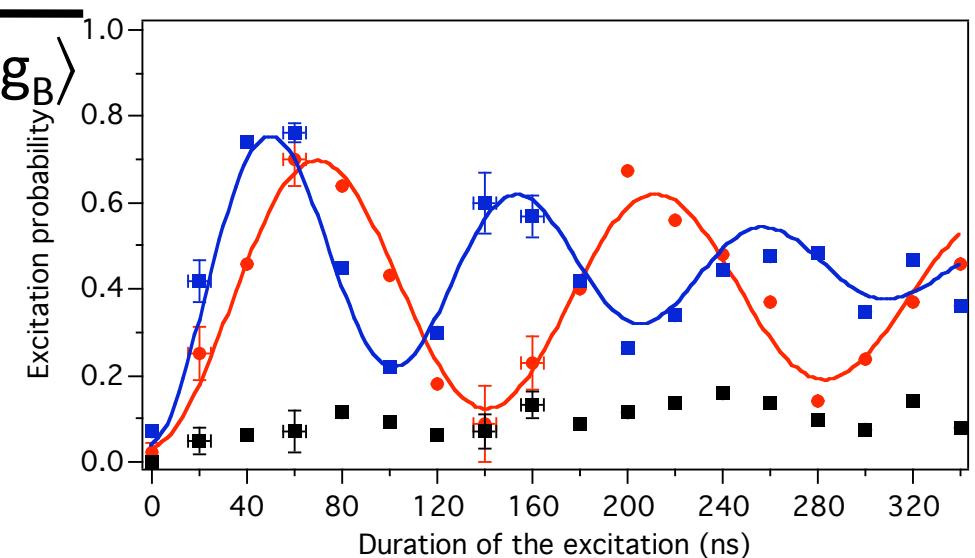
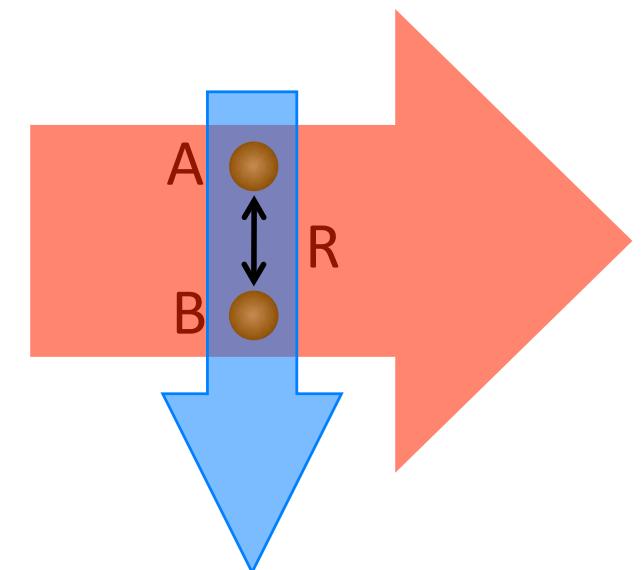
Recall (RW approx.):  $H_{\text{at-L}}^j = \frac{\hbar\Omega}{2} (|r\rangle\langle g| e^{i\mathbf{k}\cdot\mathbf{r}_j} + |g\rangle\langle r| e^{-i\mathbf{k}\cdot\mathbf{r}_j})$

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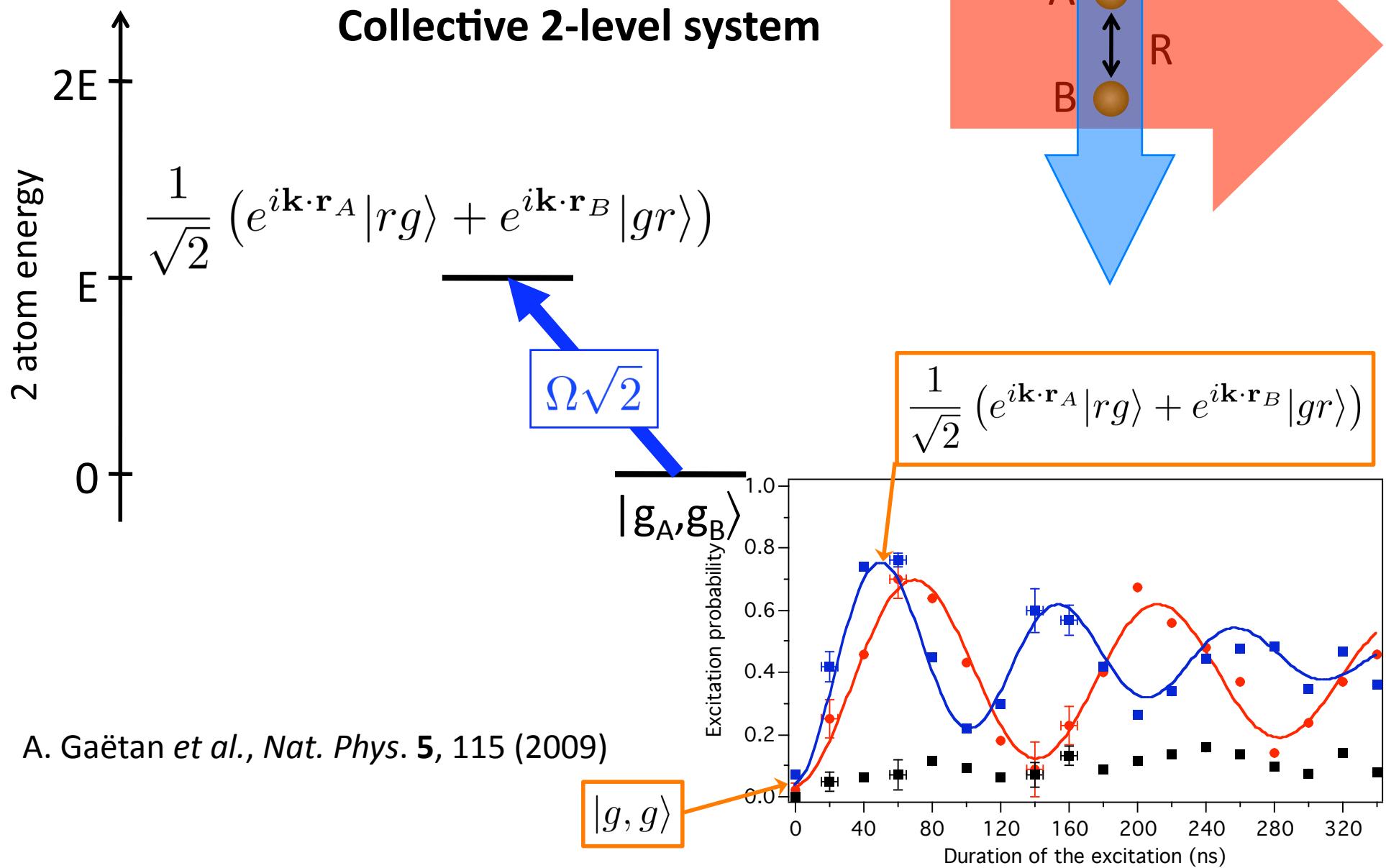
# Rydberg blockade: collective excitation (IO Palaiseau)



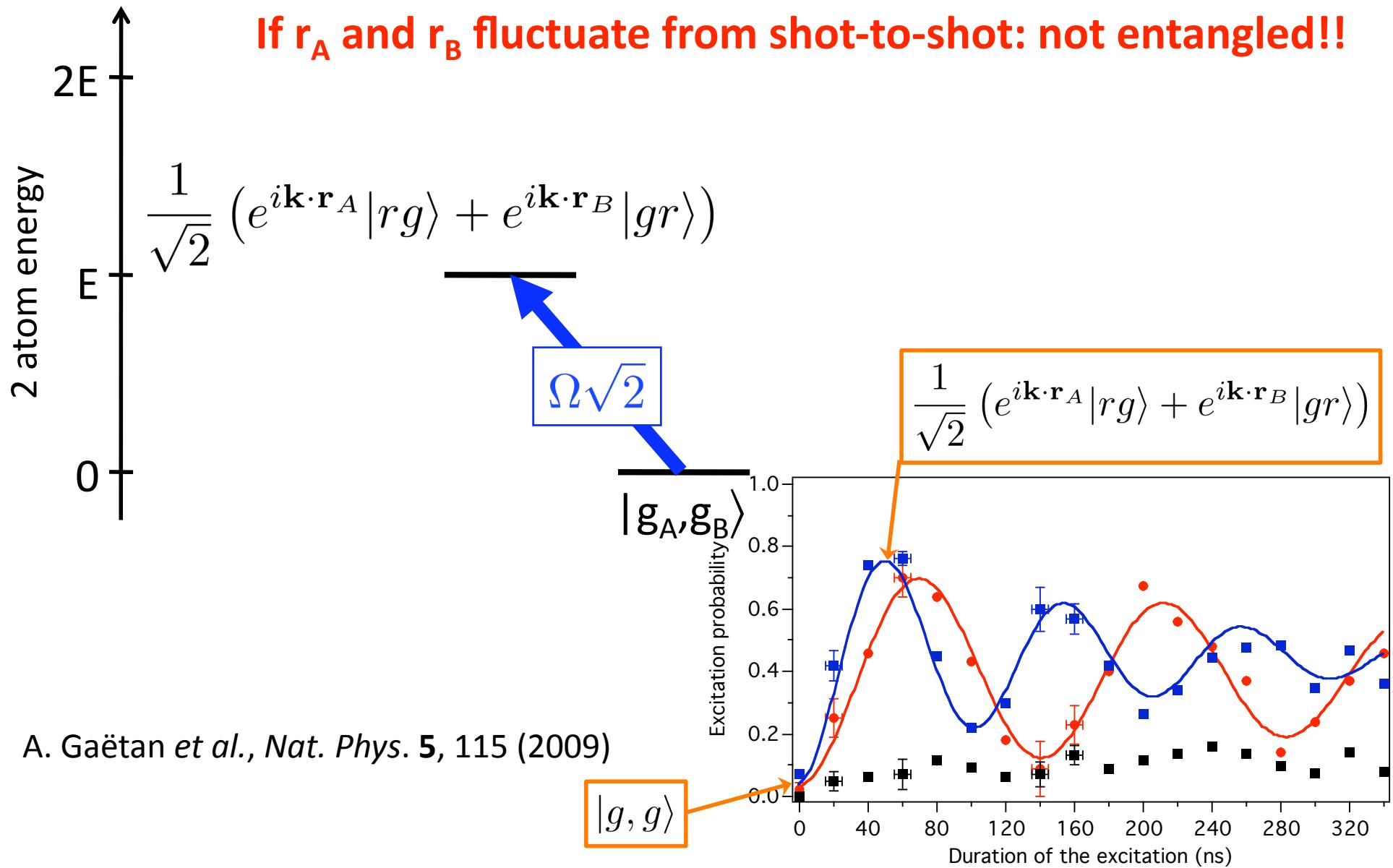
A. Gaëtan *et al.*, Nat. Phys. 5, 115 (2009)



# Rydberg blockade: collective excitation (IO Palaiseau)

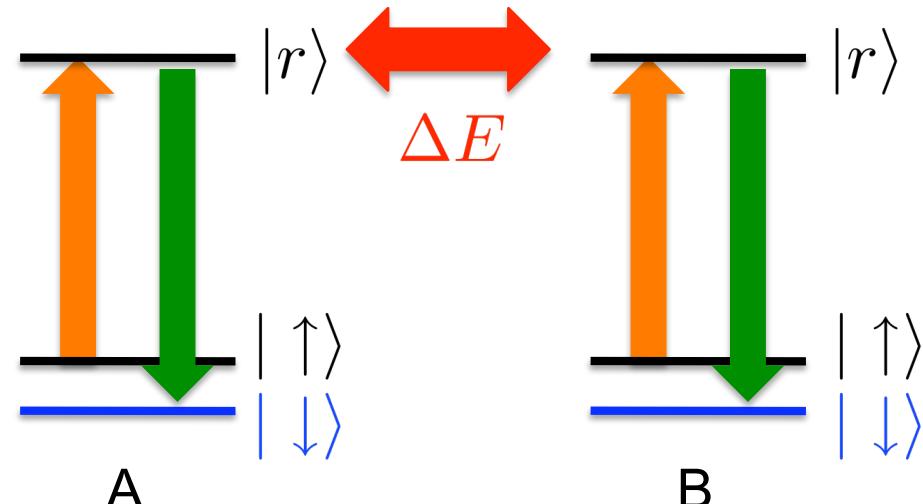
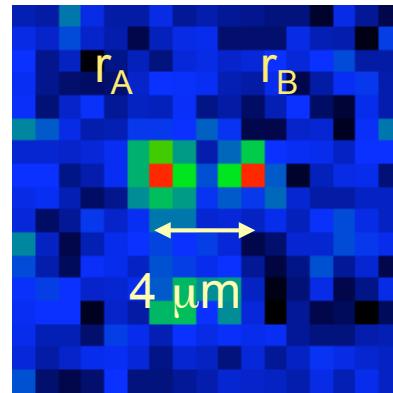


# Rydberg blockade: collective excitation (IO Palaiseau)



## Entangled of two atoms using the Rydberg blockade

Collective excitation version



$$|\uparrow, \uparrow\rangle \xrightarrow{\text{orange}} \frac{1}{\sqrt{2}} (e^{i\mathbf{k} \cdot \mathbf{r}_A} |r, \uparrow\rangle + e^{i\mathbf{k} \cdot \mathbf{r}_B} |\uparrow, r\rangle) \quad k = k_R + k_B$$

$$\xrightarrow{\text{green}} \frac{1}{\sqrt{2}} (e^{i(\mathbf{k} \cdot \mathbf{r}_A - \mathbf{k}' \cdot \mathbf{r}'_A)} |\downarrow, \uparrow\rangle + e^{i(\mathbf{k} \cdot \mathbf{r}_B - \mathbf{k}' \cdot \mathbf{r}'_B)} |\uparrow, \downarrow\rangle)$$

$$\Rightarrow \frac{1}{\sqrt{2}} (|\downarrow, \uparrow\rangle + e^{i\phi} |\uparrow, \downarrow\rangle) \quad \text{with} \quad \phi = \mathbf{k} \cdot (\delta\mathbf{r}_A - \delta\mathbf{r}_B) \quad (k \approx k')$$

If atomic motion frozen  $\Rightarrow \delta r_A \approx \delta r_B \approx 0$

## Analyzing entanglement

$$|\psi_+\rangle = \frac{1}{\sqrt{2}}(|\downarrow, \uparrow\rangle + |\uparrow, \downarrow\rangle)$$

Measure the density matrix

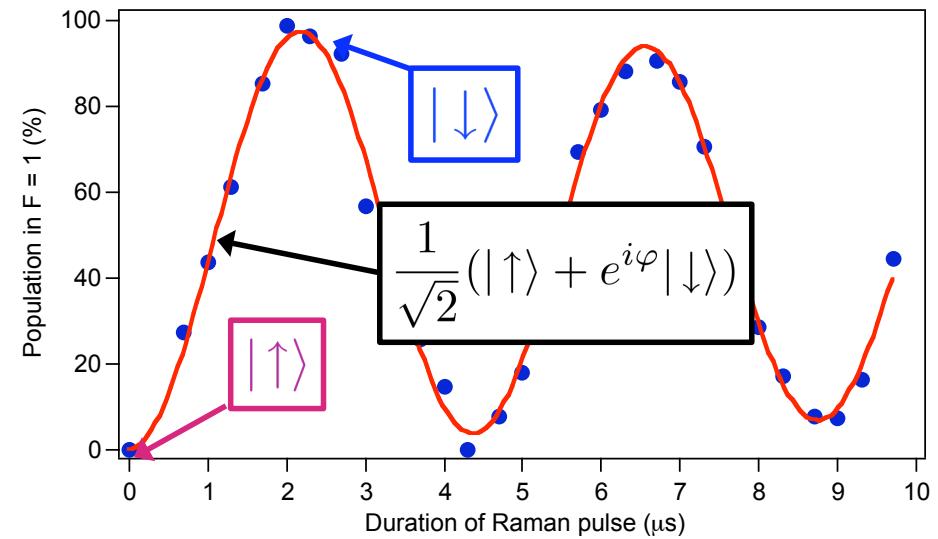
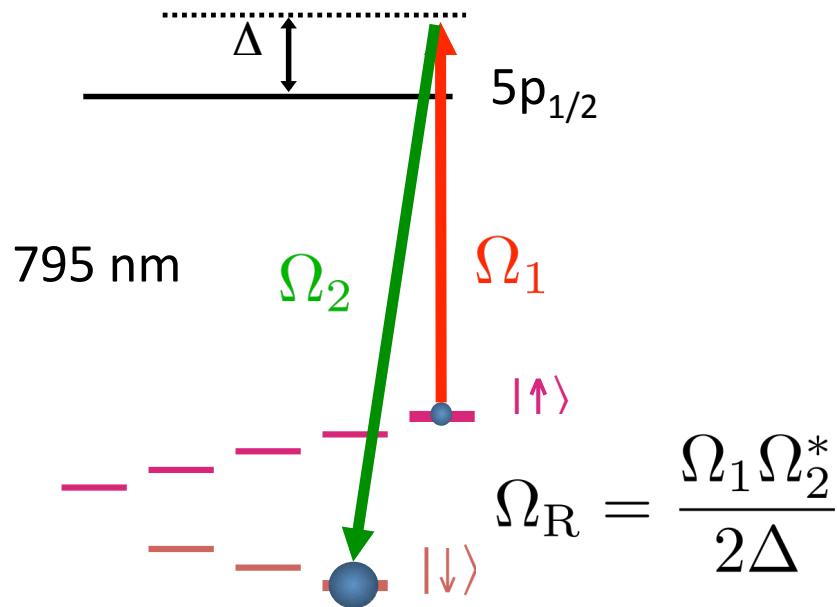
$$\hat{\rho} = \begin{pmatrix} P_{\downarrow\downarrow} & a & b & c \\ a^* & P_{\downarrow\uparrow} & \rho_{\uparrow\downarrow, \downarrow\uparrow} & d \\ b^* & \rho_{\uparrow\downarrow, \downarrow\uparrow}^* & P_{\uparrow\downarrow} & e \\ c^* & d^* & e^* & P_{\uparrow\uparrow} \end{pmatrix} |_{\downarrow\downarrow}, |\downarrow\uparrow\rangle, |\uparrow\downarrow\rangle, |\uparrow\uparrow\rangle$$

Extract the fidelity

$$F = \langle \psi_+ | \hat{\rho} | \psi_+ \rangle$$

$$F = \frac{1}{2}(P_{\downarrow\uparrow} + P_{\uparrow\downarrow} + 2\Re(\rho_{\downarrow\uparrow, \uparrow\downarrow}))$$

## Raman rotation: check coherence of the superposition



Global rotation on the two atoms:  $R_{A,B}(\theta, \varphi) = R_A(\theta, \varphi) \otimes R_B(\theta, \varphi)$

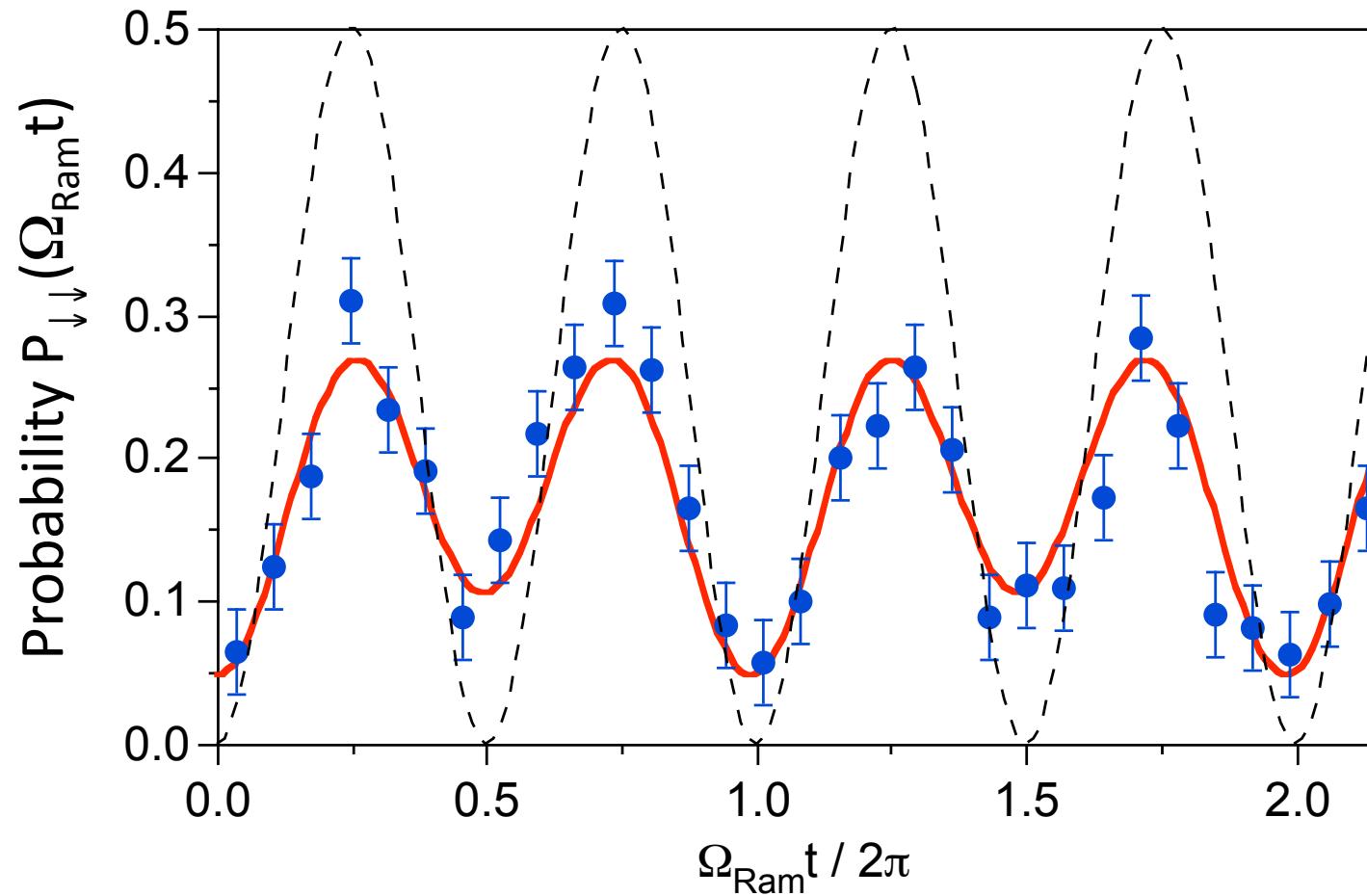
$$|\psi_+\rangle = \frac{1}{\sqrt{2}}(|\downarrow, \uparrow\rangle + |\uparrow, \downarrow\rangle) \quad \longleftrightarrow \quad \frac{1}{\sqrt{2}}(|\downarrow, \downarrow\rangle - |\uparrow, \uparrow\rangle)$$

$2 \times \Omega_{\text{Raman}}$

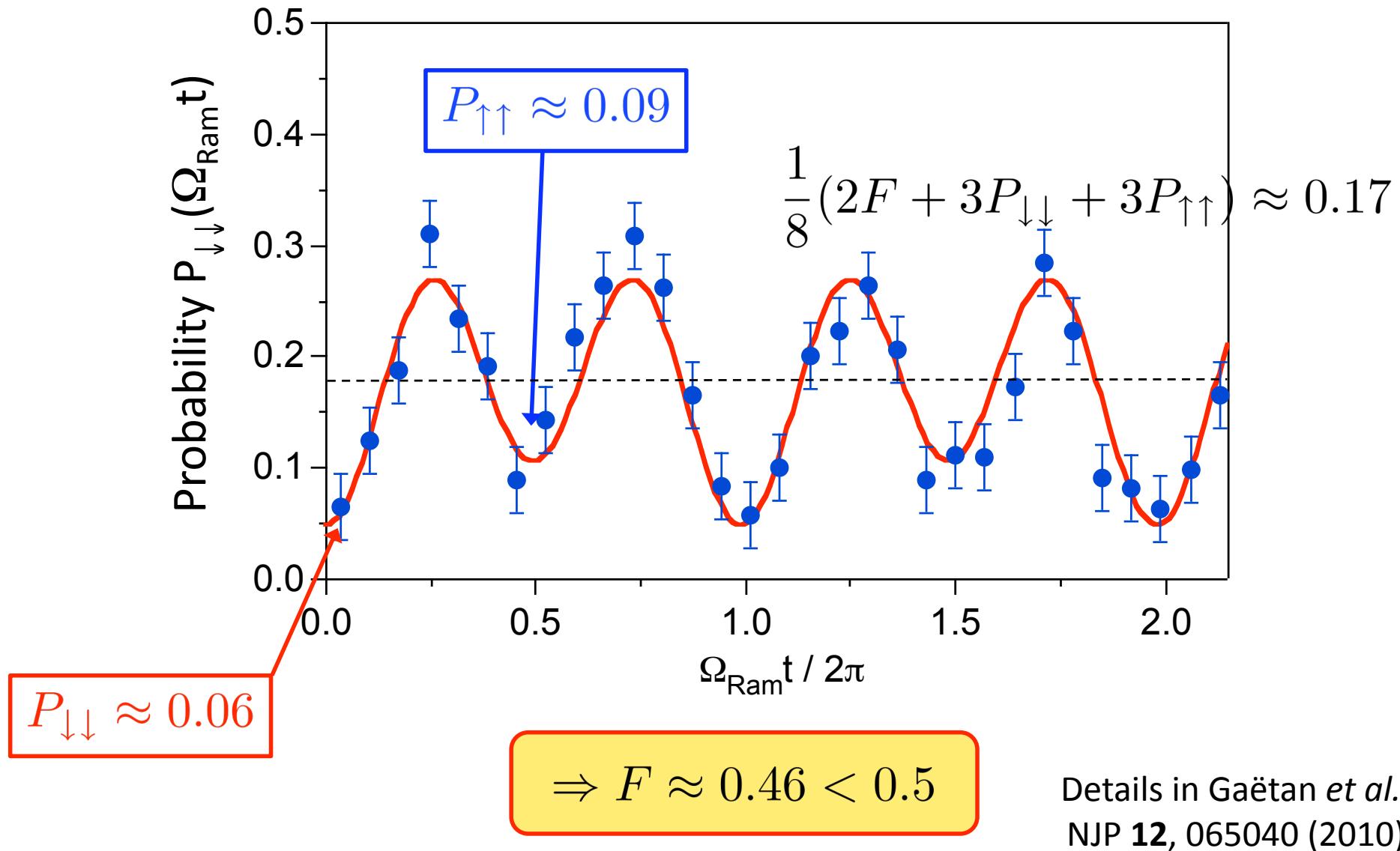
$$\Rightarrow \rho(\theta, \varphi) = R_{A,B}^{-1}(\theta, \varphi) \cdot \rho \cdot R_{A,B}(\theta, \varphi)$$

Rotation  $\hat{\rho} \rightarrow \hat{\rho}_{\text{rot}}$   $\Rightarrow$  transforms **coherence** into **population**

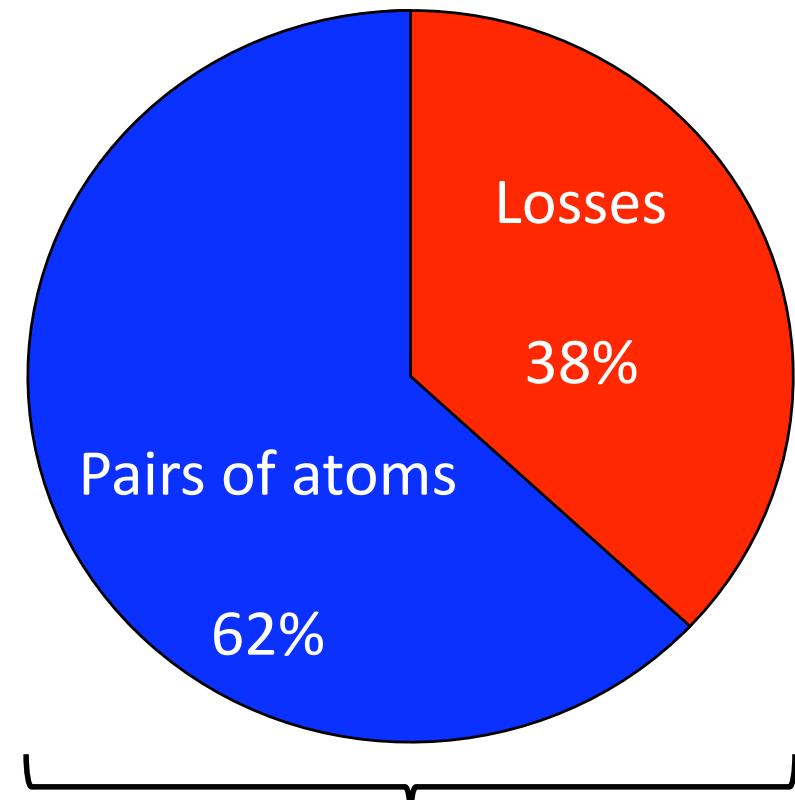
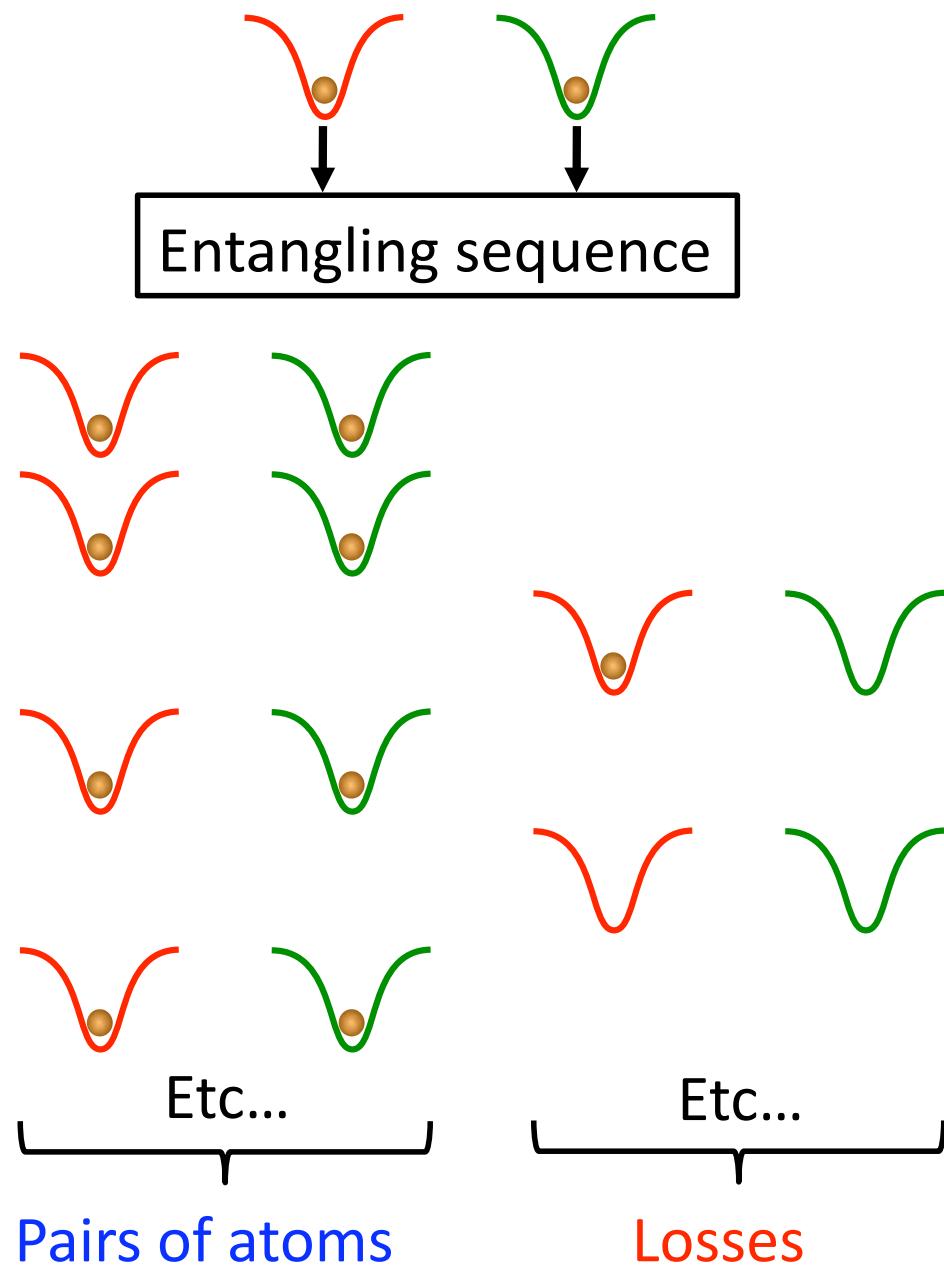
## Measurement of entanglement



## Measurement of entanglement



## Atom losses during the entangling sequence



$$F_{\text{pairs}} = F_{\text{tot}} / 0.62$$

$$F_{\text{pairs}} = 0.75 \pm 0.07$$

## Partial state reconstruction and error budget

**State prepared:**  $\rho_{exp} = 0.46|\psi_+\rangle\langle\psi_+| + \rho_{junk}$

**After loss correction:**  $\rho'_{exp} = 0.75|\psi_+\rangle\langle\psi_+| + \rho'_{junk}$

$$\rho'_{exp} = \begin{pmatrix} 0.1 & a & b & c \\ a^* & 0.38 & 0.38 & d \\ b^* & 0.38 & 0.38 & e \\ c^* & d^* & e^* & 0.12 \end{pmatrix}_{|\downarrow\downarrow\rangle, |\downarrow\uparrow\rangle, |\uparrow\downarrow\rangle, |\uparrow\uparrow\rangle}$$

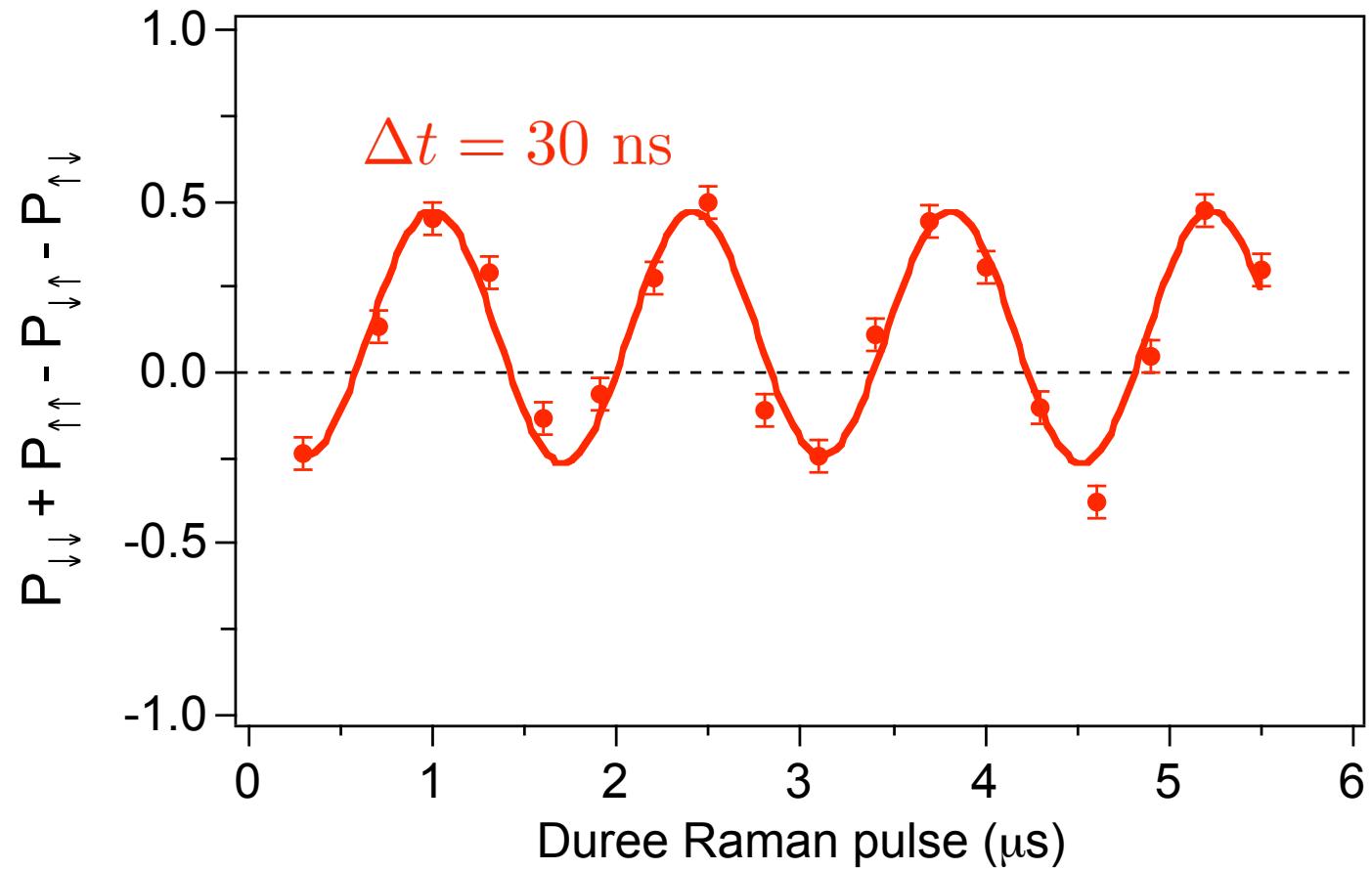
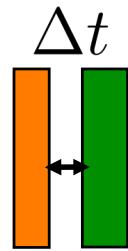
**Error budget**  $F_{pairs} = 0.75$ :

**Imperfect blockade** **0.12**

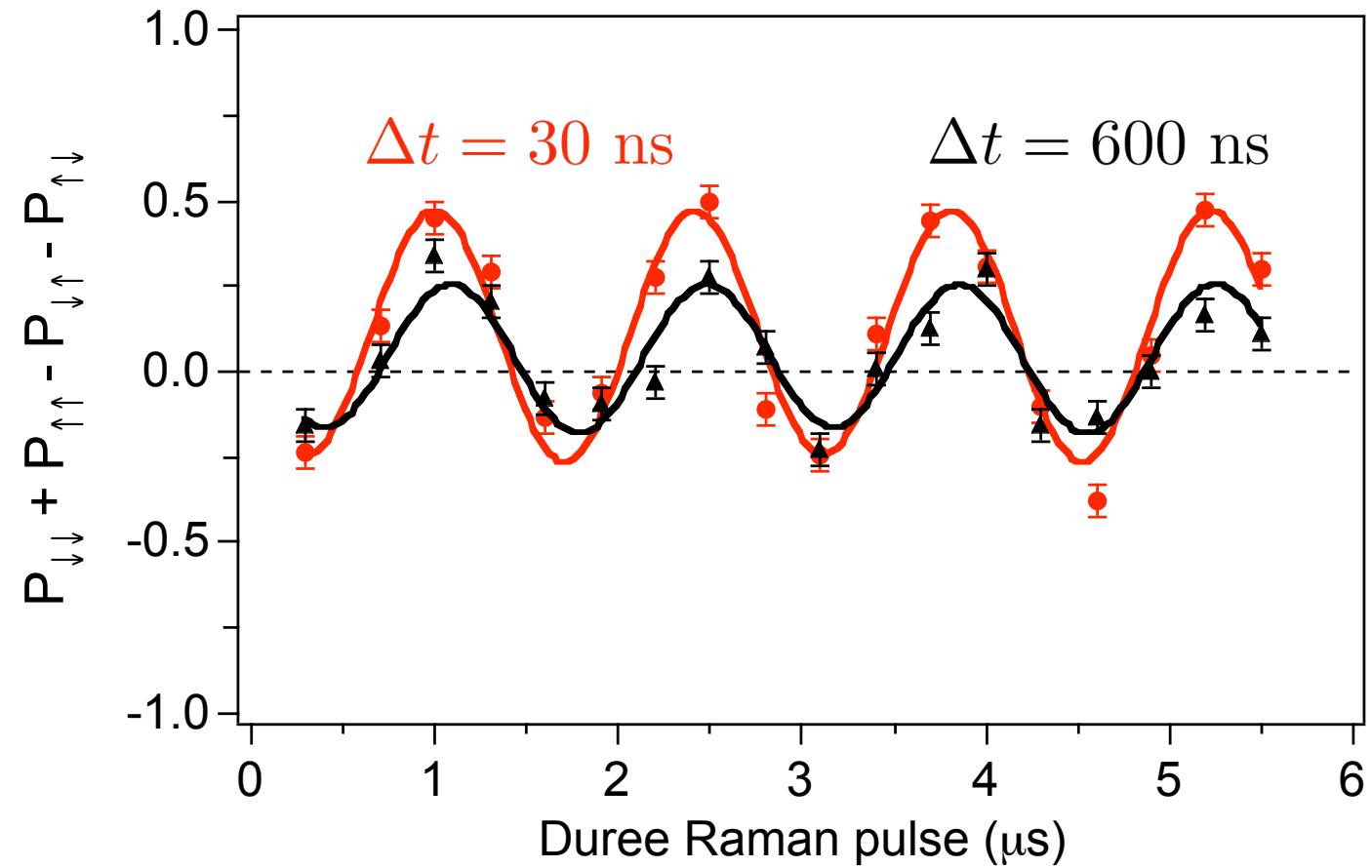
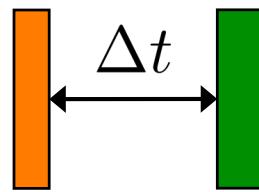
Technical (laser stability, sp. emission) 0.1

Motion of the atoms 0.03

## Influence of the motion of the atoms: how much frozen?



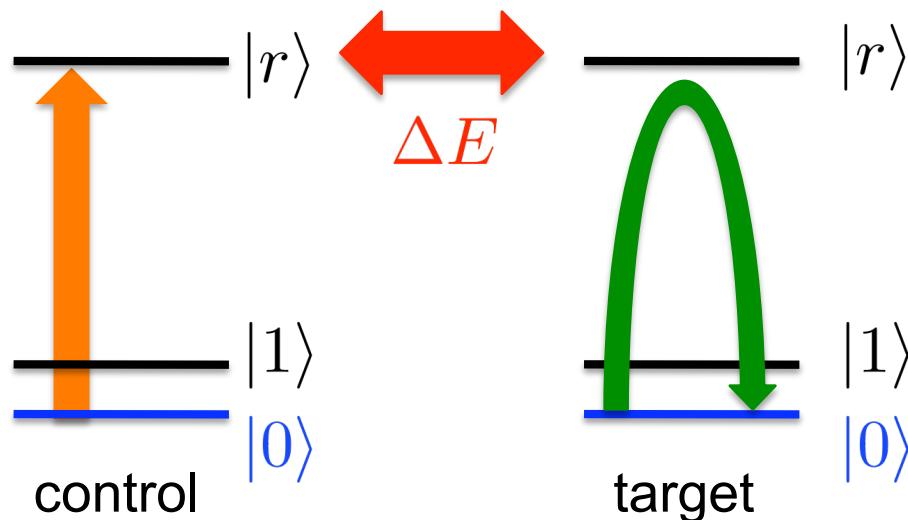
## Influence of the motion of the atoms: how much frozen?



Coherence  $\Re(\hat{\rho}_{\downarrow\uparrow,\uparrow\downarrow})$  reduced by 0.5

Theory:  $\langle e^{ik(v_A - v_B)\Delta t} \rangle = e^{-\frac{k_B T}{M} k^2 \Delta t^2} \approx 0.45$

# Quantum gate using Rydberg blockade



D. Jacksch *et al.*, PRL 85, 2208 (2000)

Sequence:

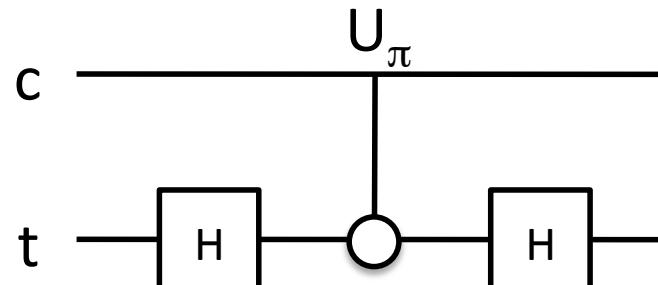
$$\pi_A - 2\pi_B - \pi_A$$

Table of truth:

$$\begin{array}{l}
 11 \rightarrow 11 \rightarrow 11 \rightarrow 11 \\
 10 \rightarrow 1r \rightarrow 1r \rightarrow -10 \\
 01 \rightarrow r0 \rightarrow r0 \rightarrow -01 \\
 00 \rightarrow r0 \xrightarrow{\text{red}} r0 \rightarrow -00
 \end{array}$$

**Blockade**

From  $\pi$ -gate to CNOT



$$\begin{array}{l}
 11 \rightarrow 10 \\
 10 \rightarrow 11 \\
 01 \rightarrow 01 \\
 00 \rightarrow 00
 \end{array}$$

# The CNOT gate at uni. Wisconsin (1)

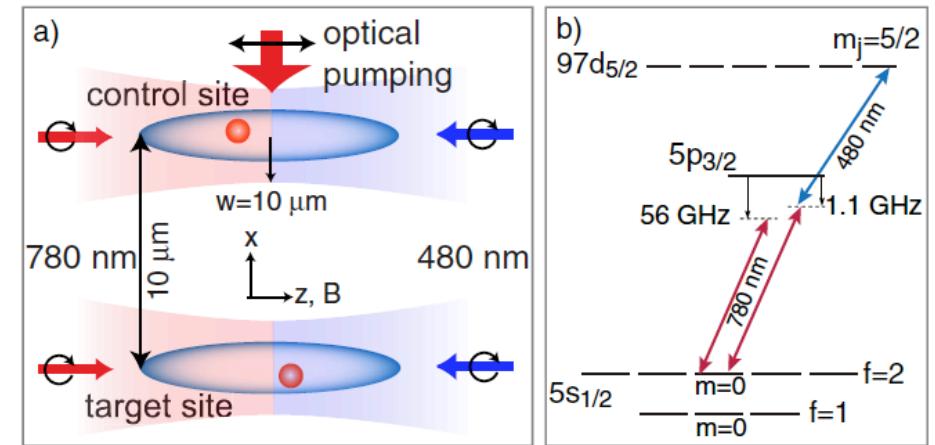
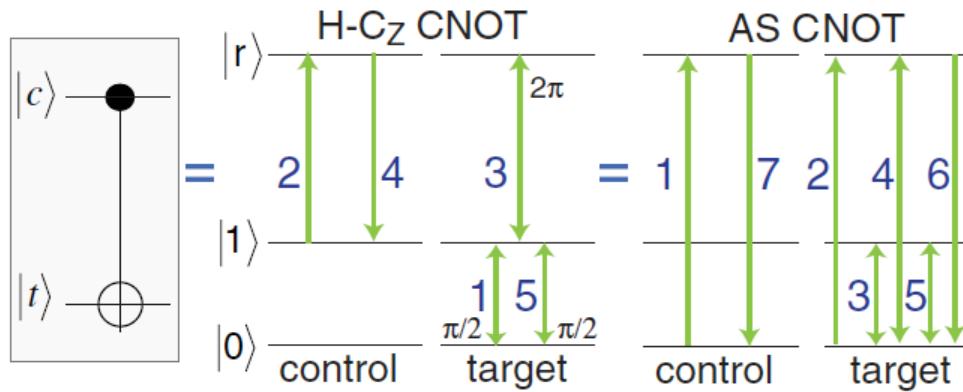
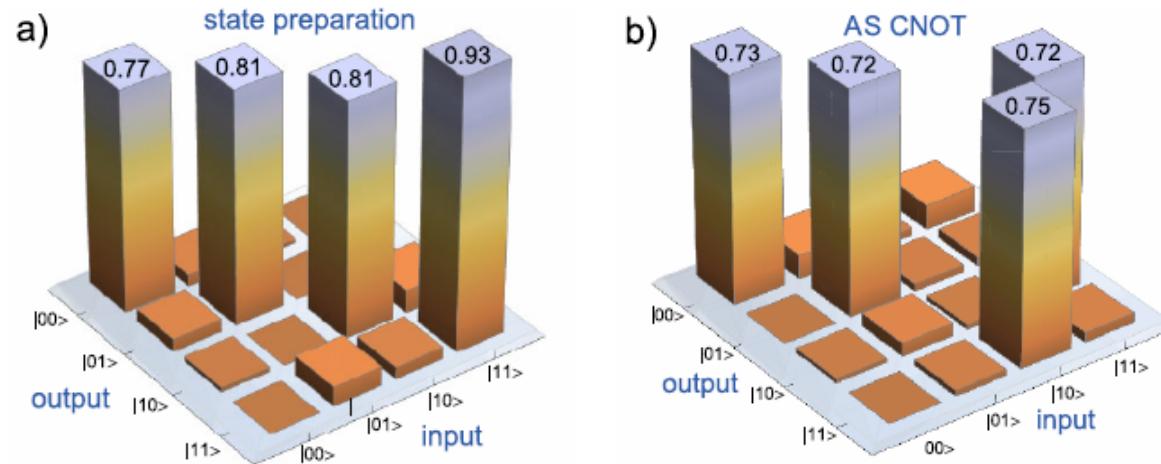


Table of truth: check blockade



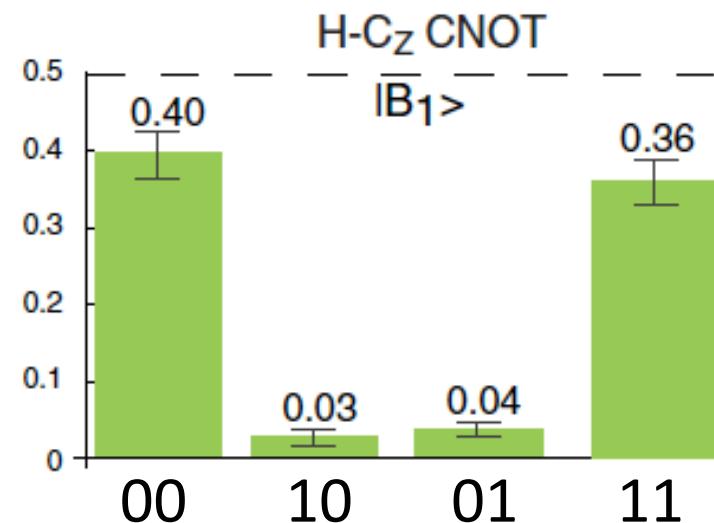
Isenhower *et al.*, PRL 104, 010503 (2010)

## The CNOT gate at uni. Wisconsin (2)

Prepare entangled states

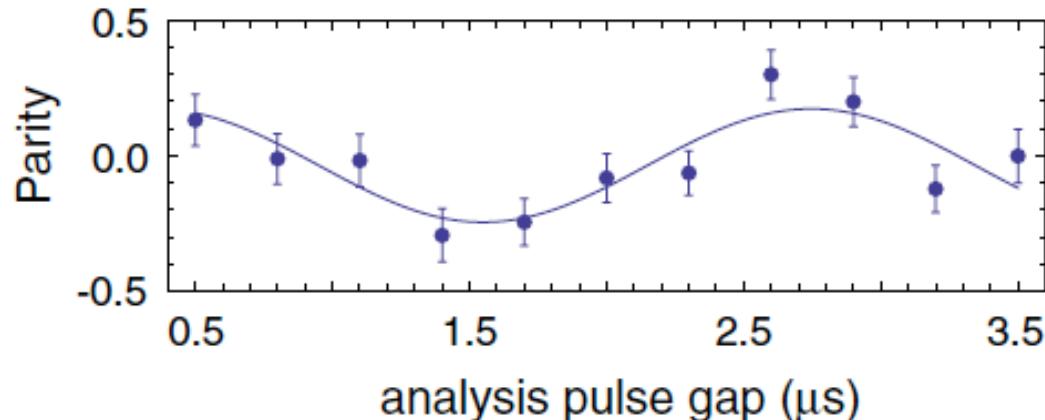
Prepare  $(|0\rangle + |1\rangle)|0\rangle$

gate



Check coherence: global Raman rotation

$$\Pi = P_{00} + P_{11} - P_{01} - P_{10}$$



Amplitude  
 $\Rightarrow \Re(\rho_{0011})$

$$\mathcal{F} = \frac{1}{2}(P_{00} + P_{11}) + \Re(\rho_{0011})$$

$F = 0.58$   
(with loss correction)

See also: Zhang *et al.*, PRA **82**, 030306(R) (2010)

## Conclusion on the gate and entanglement

Proof-of-principle: OK

But needs to improve fidelity!!

And careful with the phases...