

1. The set  $\{|\phi_1\rangle, |\phi_2\rangle\}$  is an orthonormal basis set for a 2-dimensional Hilbert space. The states  $|W\rangle$ ,  $|\theta_1\rangle$  and  $|\theta_2\rangle$  are defined as

$$\begin{aligned} |W\rangle &= \alpha |\phi_1\rangle + i\alpha |\phi_2\rangle, \\ |\theta_1\rangle &= \frac{i}{\sqrt{3}} |\phi_1\rangle - \frac{\sqrt{2}}{\sqrt{3}} |\phi_2\rangle, \quad |\theta_2\rangle = \frac{1}{\sqrt{2}} |\phi_1\rangle + \frac{i}{\sqrt{2}} |\phi_2\rangle. \end{aligned}$$

- (a) The state  $|W\rangle$  is normalized. The real and imaginary parts of the constant  $\alpha$  are equal. Find  $\alpha$ .

[6 marks]

- (b) Using the representation

$$|\phi_1\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad |\phi_2\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix},$$

express the operator  $\hat{M} = |\theta_1\rangle\langle\theta_2|$  as a matrix.

[5 marks]

- (c) Find out and explain whether or not  $\hat{M} = |\theta_1\rangle\langle\theta_2|$  is hermitian. Explain whether the operator  $\hat{M}$  can represent a physical observable.

[6 marks]

- (d) The observable  $B$  is represented by the operator  $\hat{B} = |\theta_2\rangle\langle\theta_2|$ . Calculate the uncertainty of the observable  $B$  in the state  $|\phi_1\rangle$ .

[10 marks]

(e) If  $B$  is measured, what are the possible answers that can be obtained?

**[8 marks]**

(f) If the system Hamiltonian is

$$\hat{H} = F |\phi_1\rangle\langle\phi_1| + G |\phi_2\rangle\langle\phi_2| = \begin{pmatrix} F & 0 \\ 0 & G \end{pmatrix}$$

and the system is in the state  $|\theta_1\rangle$  at time  $t = 0$ , find the state at any later time  $t$ . Here  $F$  and  $G$  are real constants.

**[15 marks]**

2. Consider a particle of mass  $m$  in one dimension, subject to the potential

$$V(x) = -\lambda\delta(x).$$

Here  $\lambda$  is a positive constant and  $\delta(x)$  is the Dirac delta function.

We will consider stationary solutions of the time-independent Schroedinger equation.

- (a) Explain why the derivative of the stationary solution,  $\psi'(x)$ , need not be continuous at  $x = 0$ .

**[3 marks]**

- (b) Derive the boundary conditions

$$\psi'_R(0) - \psi'_L(0) = -\frac{2m\lambda}{\hbar^2}\psi_L(0) = -\frac{2m\lambda}{\hbar^2}\psi_R(0)$$

where  $\psi_L(x)$  and  $\psi_R(x)$  are the wavefunctions on the left half-line and right half-line respectively.

Hint: you might consider integrating the Schroedinger equation from  $-\epsilon$  to  $+\epsilon$ , and then taking the limit  $\epsilon \rightarrow 0$ .

**[12 marks]**

- (c) This system has a single bound state at some negative energy,  $E < 0$ . This eigenfunction has the form

$$\psi(x) = \begin{cases} A_1 e^{\alpha x} + A_2 e^{-\alpha x} & \text{for } x < 0 \\ A_3 e^{\alpha x} + A_4 e^{-\alpha x} & \text{for } x > 0 \end{cases}$$

where  $\alpha$  is a positive constant.

Explain which of these terms should be dropped, and why.

Derive the relationship between  $\alpha$  and  $E$ .

**[10 marks]**

- (d) Use the boundary conditions and normalization to determine  $\alpha$  and the normalization constants  $A_i$ , in terms of  $\lambda$ . Assume  $A_1$  to be real and positive.

[15 marks]

- (e) Calculate the expectation value of momentum in the bound state.

[10 marks]

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SAMPLE ANSWERS

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1. Question 1.

The set  $\{|\phi_1\rangle, |\phi_2\rangle\}$  is an orthonormal basis set for a 2-dimensional Hilbert space. The states  $|W\rangle$ ,  $|\theta_1\rangle$  and  $|\theta_2\rangle$  are defined as

$$\begin{aligned} |W\rangle &= \alpha |\phi_1\rangle + i\alpha |\phi_2\rangle, \\ |\theta_1\rangle &= \frac{i}{\sqrt{3}} |\phi_1\rangle - \frac{\sqrt{2}}{\sqrt{3}} |\phi_2\rangle, \quad |\theta_2\rangle = \frac{1}{\sqrt{2}} |\phi_1\rangle + \frac{i}{\sqrt{2}} |\phi_2\rangle. \end{aligned}$$

(a) Question 1(a)

The state  $|W\rangle$  is normalized. The real and imaginary parts of the constant  $\alpha$  are equal. Find  $\alpha$ .

**[6 marks]**

**[Sample Answer:]**

$$|\alpha|^2 + |i\alpha|^2 = 1 \quad \implies |\alpha|^2 = \frac{1}{2} \quad \implies |\alpha| = \frac{1}{\sqrt{2}}$$

Since the real and imaginary parts of  $\alpha$  are equal, we can write

$$\alpha = b + ib$$

where  $b$  is a real number. Then  $|\alpha|^2 = b^2 + b^2 = 2b^2$ . Hence

$$2b^2 = \frac{1}{2} \quad \implies b = \pm \frac{1}{2}$$

Therefore we have

$$\text{either } \alpha = \frac{1}{2} + \frac{i}{2} \quad \text{or} \quad \alpha = -\frac{1}{2} - \frac{i}{2}$$

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(b) Question 1(b)

Using the representation

$$|\phi_1\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad |\phi_2\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix},$$

express the operator  $\hat{M} = |\theta_1\rangle\langle\theta_2|$  as a matrix.**[5 marks]****[Sample Answer:]**

$$|\theta_1\rangle = \frac{i}{\sqrt{3}}|\phi_1\rangle - \frac{\sqrt{2}}{\sqrt{3}}|\phi_2\rangle = \begin{pmatrix} \frac{i}{\sqrt{3}} \\ -\frac{\sqrt{2}}{\sqrt{3}} \end{pmatrix}$$

$$|\theta_2\rangle = \frac{1}{\sqrt{2}}|\phi_1\rangle + \frac{i}{\sqrt{2}}|\phi_2\rangle = \begin{pmatrix} 1/\sqrt{2} \\ i/\sqrt{2} \end{pmatrix}$$

$$\begin{aligned} \hat{M} = |\theta_1\rangle\langle\theta_2| &= \begin{pmatrix} \frac{i}{\sqrt{3}} \\ -\frac{\sqrt{2}}{\sqrt{3}} \end{pmatrix} (1/\sqrt{2} \quad -i/\sqrt{2}) = \begin{pmatrix} \frac{i}{\sqrt{6}} & +\frac{1}{\sqrt{6}} \\ -\frac{1}{\sqrt{3}} & \frac{i}{\sqrt{3}} \end{pmatrix} \\ &= \frac{1}{\sqrt{6}} \begin{pmatrix} i & +1 \\ -\sqrt{2} & i\sqrt{2} \end{pmatrix} \end{aligned}$$

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(c) Question 1(c)

Find out and explain whether or not  $\hat{M} = |\theta_1\rangle\langle\theta_2|$  is hermitian.Explain whether the operator  $\hat{M}$  can represent a physical observable.**[6 marks]****[Sample Answer:]**

$$\hat{M}^\dagger = \begin{pmatrix} -\frac{i}{\sqrt{6}} & -\frac{1}{\sqrt{3}} \\ +\frac{1}{\sqrt{6}} & -\frac{i}{\sqrt{3}} \end{pmatrix} \neq \hat{M}$$

Hence not hermitian.

Physical observables are represented by hermitian operators (measured values need to be real). Hence  $\hat{M}$  cannot represent a physical observable.

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(d) Question 1(d)

The observable  $B$  is represented by the operator  $\hat{B} = |\theta_2\rangle\langle\theta_2|$ . Calculate the uncertainty of the observable  $B$  in the state  $|\phi_1\rangle$ .

**[10 marks]**

**[Sample Answer:]**

$$|\theta_2\rangle = \frac{1}{\sqrt{2}}|\phi_1\rangle + \frac{i}{\sqrt{2}}|\phi_2\rangle = \begin{pmatrix} 1/\sqrt{2} \\ i/\sqrt{2} \end{pmatrix}$$

$$\begin{aligned} \hat{B} = |\theta_2\rangle\langle\theta_2| &= \begin{pmatrix} 1/\sqrt{2} \\ i/\sqrt{2} \end{pmatrix} \begin{pmatrix} 1/\sqrt{2} & -i/\sqrt{2} \end{pmatrix} = \begin{pmatrix} \frac{1}{2} & -\frac{i}{2} \\ \frac{i}{2} & \frac{1}{2} \end{pmatrix} \\ &= \frac{1}{2} \begin{pmatrix} 1 & -i \\ i & 1 \end{pmatrix} \end{aligned}$$

$$\hat{B}^2 = \frac{1}{4} \begin{pmatrix} 1 & -i \\ i & 1 \end{pmatrix} \begin{pmatrix} 1 & -i \\ i & 1 \end{pmatrix} = \frac{1}{4} \begin{pmatrix} 2 & -2i \\ 2i & 2 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 & -i \\ i & 1 \end{pmatrix} = \hat{B}$$

In the state  $|\phi_1\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ ,

$$\langle B \rangle = \langle \theta_1 | \hat{B} | \theta_1 \rangle = \begin{pmatrix} 1 & 0 \end{pmatrix} \hat{B} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = B_{11} = \frac{1}{2}$$

and

$$\langle B^2 \rangle = \langle B \rangle = \frac{1}{2}$$

so that

$$\langle B^2 \rangle - \langle B \rangle^2 = \frac{1}{2} - \frac{1}{4} = \frac{1}{4}$$

The uncertainty is

$$\Delta B = \sqrt{\langle B^2 \rangle - \langle B \rangle^2} = \frac{1}{2}$$

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(e) Question 1(e)

If  $B$  is measured, what are the possible answers that can be obtained?

[8 marks]

[Sample Answer:]

The eigenvalues of the  $B$  matrix are the possible results of a measurement. Calculation shows the eigenvalues to be 0 and 1. These are the possible values.

The eigenvalue calculation might go as follows:

$$\begin{aligned} \begin{vmatrix} \frac{1}{2} - \lambda & \frac{i}{2} \\ -\frac{i}{2} & \frac{1}{2} - \lambda \end{vmatrix} = 0 & \implies \left(\frac{1}{2} - \lambda\right)^2 - \frac{i}{2} \left(-\frac{i}{2}\right) = 0 \\ & \implies \frac{1}{4} + \lambda^2 - \lambda - \frac{1}{4} = 0 \\ & \implies \lambda^2 - \lambda = 0 \implies \lambda = 0, 1 \end{aligned}$$

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(f) Question 1(f)

If the system Hamiltonian is

$$\hat{H} = F |\phi_1\rangle\langle\phi_1| + G |\phi_2\rangle\langle\phi_2| = \begin{pmatrix} F & 0 \\ 0 & G \end{pmatrix}$$

and the system is in the state  $|\theta_1\rangle$  at time  $t = 0$ , find the state at any later time  $t$ . Here  $F$  and  $G$  are real constants.



[15 marks]

**[Sample Answer:]**

Let's denote the system wavefunction components as  $u(t)$  and  $v(t)$ , i.e., the wavefunction is

$$|\psi(t)\rangle = \begin{pmatrix} u(t) \\ v(t) \end{pmatrix}$$

Using the time-dependent Schrodinger equation,  $\hat{H}|\psi(t)\rangle = i\hbar\frac{\partial}{\partial t}|\psi(t)\rangle$ , we obtain

$$\begin{pmatrix} F & 0 \\ 0 & G \end{pmatrix} \begin{pmatrix} u(t) \\ v(t) \end{pmatrix} = i\hbar\frac{\partial}{\partial t} \begin{pmatrix} u(t) \\ v(t) \end{pmatrix} \implies \begin{cases} u'(t) = -\frac{iF}{\hbar}u(t) \\ v'(t) = -\frac{iG}{\hbar}v(t) \end{cases}$$

which can be solved to give

$$u(t) = u(0) \exp\left[-\frac{iF}{\hbar}t\right]; \quad v(t) = v(0) \exp\left[-\frac{iG}{\hbar}t\right]$$

The state at  $t = 0$  is  $|\theta_1\rangle = \frac{i}{\sqrt{3}}|\phi_1\rangle - \frac{\sqrt{2}}{\sqrt{3}}|\phi_2\rangle$ , i.e.,

$$\begin{pmatrix} u(0) \\ v(0) \end{pmatrix} = \begin{pmatrix} i/\sqrt{3} \\ -\sqrt{2}/\sqrt{3} \end{pmatrix}$$

Therefore

$$u(t) = \frac{i}{\sqrt{3}} \exp\left[-\frac{iF}{\hbar}t\right]; \quad v(t) = -\frac{\sqrt{2}}{\sqrt{3}} \exp\left[-\frac{iG}{\hbar}t\right]$$

so that the wavefunction at time  $t$  is

$$|\psi(t)\rangle = \begin{pmatrix} \frac{i}{\sqrt{3}} \exp\left[-\frac{iF}{\hbar}t\right] \\ -\frac{\sqrt{2}}{\sqrt{3}} \exp\left[-\frac{iG}{\hbar}t\right] \end{pmatrix} = \frac{i}{\sqrt{3}} e^{-iFt/\hbar} |\phi_1\rangle - \frac{\sqrt{2}}{\sqrt{3}} e^{-iGt/\hbar} |\phi_2\rangle$$

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2. Question 2.

Consider a particle of mass  $m$  in one dimension, subject to the potential

$$V(x) = -\lambda\delta(x).$$

Here  $\lambda$  is a positive constant and  $\delta(x)$  is the Dirac delta function.

We will consider stationary solutions of the time-independent Schroedinger equation.

(a) Question 2(a)

Explain why the derivative of the stationary solution,  $\psi'(x)$ , need not be continuous at  $x = 0$ .

**[3 marks]**

**[Sample Answer:]**

The derivative needs to be continuous if the potential is everywhere finite. The potential here is  $-\infty$  at  $x = 0$ , i.e., not finite everywhere. Hence the derivative  $\psi'(x)$  does not need to be continuous.

$$\psi'_R(0) - \psi'_L(0) = -\frac{2m\lambda}{\hbar^2}\psi_L(0) = -\frac{2m\lambda}{\hbar^2}\psi_R(0)$$

(b) Question 2(b)

Derive the boundary conditions

$$\psi'_R(0) - \psi'_L(0) = -\frac{2m\lambda}{\hbar^2}\psi_L(0) = -\frac{2m\lambda}{\hbar^2}\psi_R(0)$$

where  $\psi_L(x)$  and  $\psi_R(x)$  are the wavefunctions on the left half-line and right half-line respectively.

Hint: you might consider integrating the Schroedinger equation from  $-\epsilon$  to  $+\epsilon$ , and then taking the limit  $\epsilon \rightarrow 0$ .

**[12 marks]**

**[Sample Answer:]**

The second equality follows from the continuity of the wavefunction  $\psi(x)$ , which is continuous although its derivative is not:  $\psi_L(0) =$

$\psi_R(0) = \psi(0)$ . To obtain the equation for the derivative discontinuity, we need the Schroedinger equation.

The Schroedinger equation is

$$-\frac{\hbar^2}{2m}\psi''(x) - \lambda\delta(x)\psi(x) = E\psi(x)$$

$$\implies \psi''(x) = -\frac{2m\lambda}{\hbar^2}\delta(x)\psi(x) - \frac{2mE}{\hbar^2}\psi(x)$$

Integrating both sides from  $x = -\epsilon$  to  $x = +\epsilon$ , one obtains

$$\int_{-\epsilon}^{\epsilon} \psi''(x)dx = -\frac{2m\lambda}{\hbar^2} \int_{-\epsilon}^{\epsilon} \delta(x)\psi(x)dx - \frac{2mE}{\hbar^2} \int_{-\epsilon}^{\epsilon} \psi(x)dx$$

The first integral is equal to  $\psi'(\epsilon) - \psi'(-\epsilon)$  because  $\frac{d}{dx}\psi'(x) = \psi''(x)$ . (Fundamental Theorem of Calculus.)

The second integral is  $\psi(0)$  due to the definition of the Dirac delta function.

Since  $\psi(x)$  is continuous, the last integral can be approximated for very small  $\epsilon$  as  $\psi(0)2\epsilon$ , which is an infinitesimal quantity.

Thus we have

$$\psi'(\epsilon) - \psi'(-\epsilon) = -\frac{2m\lambda}{\hbar^2}\psi(0) - \frac{4mE}{\hbar^2}\psi(0)\epsilon$$

In the limit  $\epsilon \rightarrow 0$ , the last term vanishes, and also

$$\lim_{\epsilon \rightarrow 0} \psi'(\epsilon) = \psi'_R(0), \quad \lim_{\epsilon \rightarrow 0} \psi'(-\epsilon) = \psi'_L(0).$$

Hence we obtain the boundary condition

$$\psi'_R(0) - \psi'_L(0) = -\frac{2m\lambda}{\hbar^2}\psi(0)$$

or

$$\psi'_R(0) - \psi'_L(0) = -\frac{2m\lambda}{\hbar^2}\psi_L(0) = -\frac{2m\lambda}{\hbar^2}\psi_R(0)$$

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(c) Question 2(c)

This system has a single bound state at some negative energy,  $E < 0$ . This eigenfunction has the form

$$\psi(x) = \begin{cases} A_1 e^{\alpha x} + A_2 e^{-\alpha x} & \text{for } x < 0 \\ A_3 e^{\alpha x} + A_4 e^{-\alpha x} & \text{for } x > 0 \end{cases}$$

where  $\alpha$  is a positive constant.

Explain which of these terms should be dropped, and why.

Derive the relationship between  $\alpha$  and  $E$ .

[10 marks]

**[Sample Answer:]**

For  $x < 0$ , the term  $e^{-\alpha x}$  blows up at large negative  $x$ , so that the wavefunction becomes impossible to normalize. Hence for normalizability, the term  $A_2 e^{-\alpha x}$  should be dropped.

For the corresponding reason, the term  $e^{+\alpha x}$  is not allowed in the  $x > 0$  half-line. Hence the term  $A_3 e^{\alpha x}$  should be dropped as well. The eigenfunction has the form

$$\psi(x) = \begin{cases} A_1 e^{\alpha x} & \text{for } x < 0 \\ A_4 e^{-\alpha x} & \text{for } x > 0 \end{cases}$$

To determine the constant  $\alpha$  in terms of the energy, we need to consider the Schrodinger equation either for  $x < 0$  or for  $x > 0$ , away from the problematic point  $x = 0$ .

Considering  $x < 0$ , the time-independent SE is

$$-\frac{\hbar^2}{2m} \psi''(x) = E \psi(x)$$

Using  $\psi(x) = A_1 e^{\alpha x}$  gives

$$-\frac{\hbar^2}{2m} A_1 (\alpha^2) e^{\alpha x} = E A_1 e^{\alpha x} \quad \implies \quad \alpha = \sqrt{\frac{2m(-E)}{\hbar^2}}$$

Since  $E$  is negative,  $\alpha$  is a real constant.

Can also express  $E$  in terms of  $\alpha$ :

$$E = -\frac{\hbar^2 \alpha^2}{2m}$$

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(d) Question 2(d)

Use the boundary conditions and normalization to determine  $\alpha$  and the normalization constants  $A_i$ , in terms of  $\lambda$ . Assume  $A_1$  to be real and positive.

[15 marks]

**[Sample Answer:]**

The wavefunction is  $\psi_L(x) = A_1 e^{\alpha x}$  on the left and  $\psi_R(x) = A_4 e^{-\alpha x}$  on the right.

Continuity of the wavefunction:

$$\psi_L(0) = \psi_R(0) \quad \implies \quad A_1 = A_4$$

Normalization:

$$\begin{aligned} \int_{-\infty}^0 |\psi_L(x)|^2 dx + \int_0^{\infty} |\psi_R(x)|^2 dx &= 1 \\ \implies |A_1|^2 \int_{-\infty}^0 e^{2\alpha x} dx + |A_1|^2 \int_0^{\infty} e^{-2\alpha x} dx &= 1 \\ \implies |A_1|^2 \frac{1}{2\alpha} + |A_1|^2 \frac{1}{2\alpha} &= 1 \\ \implies |A_1|^2 = \alpha = \sqrt{-\frac{2mE}{\hbar^2}} \end{aligned}$$

Discontinuity of  $\psi'$ :

$$\begin{aligned} -\alpha A_1 - \alpha A_1 &= -\frac{2m\lambda}{\hbar^2} A_1 \quad \begin{cases} \text{using } \psi'_L(0) = A_1(\alpha)e^0 = \alpha A_1 \\ \text{and } \psi'_R(0) = A_4(-\alpha)e^0 = -\alpha A_1 \end{cases} \\ \implies \alpha &= \frac{m\lambda}{\hbar^2} \quad \implies \quad A_1 = \sqrt{\alpha} = \sqrt{\frac{m\lambda}{\hbar^2}} \end{aligned}$$

(We could also calculate the energy,  $E = -\frac{\hbar^2 \alpha^2}{2m} = -\frac{m\lambda^2}{2\hbar^2}$ , if we were asked to.)

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(e) Question 2(e)

Calculate the expectation value of momentum in the bound state.

**[10 marks]****[Sample Answer:]**

The wavefunction of the bound state is  $\psi_L(x) = A_1 e^{\alpha x}$  on the left and  $\psi_R(x) = A_1 e^{-\alpha x}$  on the right. We will continue to treat  $A_1$  as a real positive number, so the wavefunction is everywhere real.

The momentum operator is  $\hat{p} = -i\hbar \frac{\partial}{\partial x}$ ; hence the expectation value of momentum in the state  $\psi(x)$  is

$$\begin{aligned} \langle p \rangle &= \int_{-\infty}^{\infty} dx \psi^*(x) \hat{p} \psi(x) = \int_{-\infty}^{\infty} dx \psi^*(x) \left( -i\hbar \frac{\partial}{\partial x} \right) \psi(x) \\ &= -i\hbar \int_{-\infty}^{\infty} dx \psi(x) \psi'(x) \end{aligned}$$

We've used  $\psi^*(x) = \psi(x)$  since the wavefunction is real.

Now we can use the expression for the wavefunction to calculate  $\int dx \psi(x) \psi'(x)$ . The problem is that there are different expressions for the negative half-line and positive half-line. Hence we split up the integral into two parts:

$$\begin{aligned} \int_{-\infty}^{\infty} dx \psi(x) \psi'(x) &= \int_{-\infty}^0 dx \psi(x) \psi'(x) + \int_0^{\infty} dx \psi(x) \psi'(x) \\ &= \int_{-\infty}^0 dx \psi_L(x) \psi'_L(x) + \int_0^{\infty} dx \psi_R(x) \psi'_R(x) \\ &= \int_{-\infty}^0 dx (A_1 e^{\alpha x}) (A_1 \alpha e^{\alpha x}) + \int_0^{\infty} dx (A_1 e^{-\alpha x}) (-A_1 \alpha e^{-\alpha x}) \\ &= A_1^2 \alpha \int_{-\infty}^0 dx e^{2\alpha x} - A_1^2 \alpha \int_0^{\infty} dx e^{-2\alpha x} \end{aligned}$$

The two integrals are equal, this can be seen by evaluating both or by a variable transformation:

$$\begin{aligned} \int_{x=-\infty}^{x=0} dx e^{2\alpha x} &= \int_{u=\infty}^{u=0} d(-u) e^{2\alpha(-u)} = \int_{u=\infty}^{u=0} d(-u) e^{2\alpha(-u)} \\ &= - \int_{\infty}^0 du e^{-2\alpha u} + \int_0^{\infty} du e^{-2\alpha u} = \int_0^{\infty} dx e^{-2\alpha x} \end{aligned}$$

Hence the two integrals cancel, and we have

$$\langle p \rangle = -i\hbar \int_{-\infty}^{\infty} dx \psi(x) \psi'(x) = -i\hbar A_1^2 \alpha \times 0 = 0$$

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