
Quantum dynamics of a photodetection process

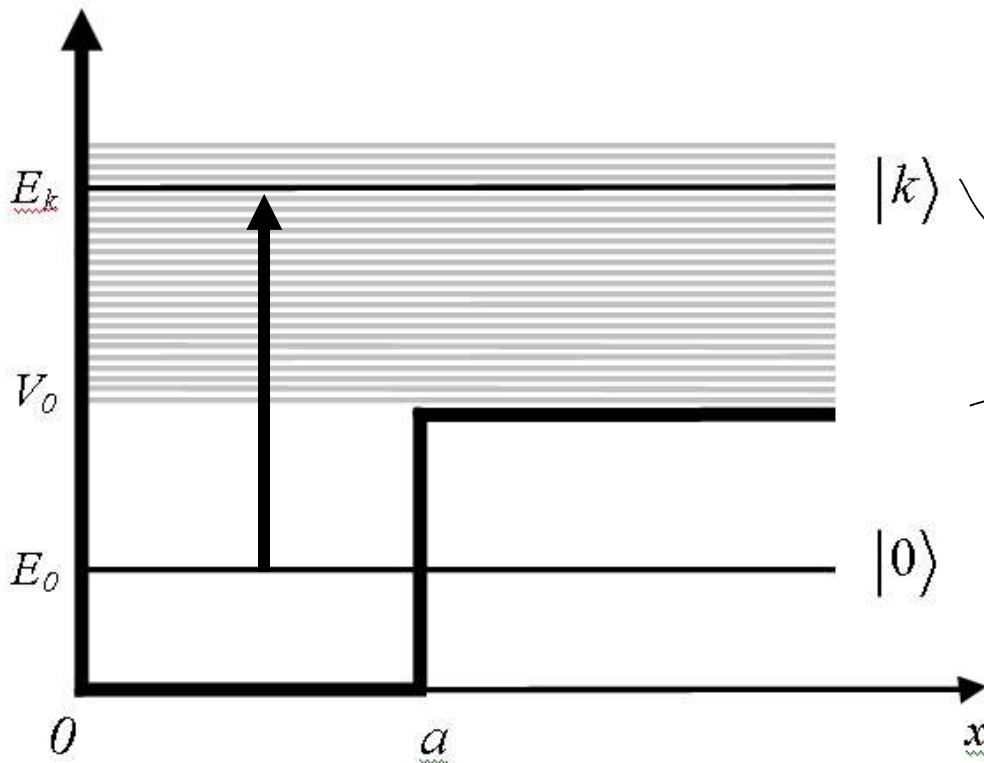
Andrew A. Ignatenko
(andi_blr@mail.ru)

B. I. Stepanov Institute of physics
National Academy of Sciences of Belarus

Motivation

- A central problem of quantum measurement theory
 - “Single-photon to single electron” transformation is a main process underlying in modern photodetection technique
 - Fundamental interest in investigation of matter-light interaction
-

Simple Model



E_k – excited energy level
 E_0 – ground level
 a – width of potential well
 V_0 – well height

Interaction constant

$$|g(\omega)|^2 = \frac{1}{\hbar^2} \frac{\hbar\omega_0}{2\varepsilon_0 V} |eX_{0k}|^2$$

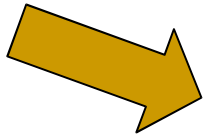
Matrix element

$$X_{0k} = \frac{-2C_1 C_2 ((k_1 a)(k'_1 a) - \xi_0^2 \sin(k_1 a) \sin(k'_1 a))}{a^2 (k_1^2 - k_1'^2)^2}$$

Wave vectors

$$k_1' = \sqrt{\frac{2mE}{\hbar^2}}, k_2' = \sqrt{\frac{2m(E - V_0)}{\hbar^2}}, k_1 = \sqrt{\frac{2mE_1}{\hbar^2}}, k_2 = \sqrt{\frac{2m(V_0 - E_1)}{\hbar^2}}$$

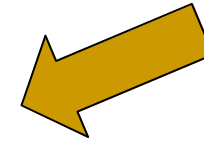
Wave function of free photoelectron



$$\Psi_k(x) = \sqrt{\frac{2}{L-a}} \sin(k_2' + \delta), \quad \delta = \arctg\left(\frac{k_2'}{k_1}, \operatorname{tg}(k_1' a)\right) - k_2' a$$

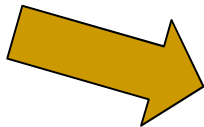
$C_0(t)$ is a probability of finding an electron in the well

$$\frac{dC_0(t)}{dt} = \int_0^\infty \int_\omega |g(\omega)|^2 \rho(\omega) e^{-i\omega(t-\tau)} C_0(t) d\tau d\omega$$



$$\rho(\omega) = \frac{\hbar \xi_0^2 (L-a)}{2\pi V_0 k_2' a^2}$$

Laplace transformation of $C_0(t)$



$$C_0^{[s]} = \frac{\hbar}{V_0} \frac{1}{s_v + ir_\omega - g_n^2 R(s_v)}$$

$$s_v = \frac{s\hbar}{V_0}$$

Wave packet of free photoelectron

$$\Psi_{ph}(x, t) = \sum_k C_k(t) \Psi_k(x)$$

$$C_k(t) = -ig_k \int_0^t e^{-i\omega_k(t-\tau)} C_0(\tau) d\tau$$

probability of finding the electron at the k-th level of continuum

Approximation of wave packet using Frenel integrals

$$\tilde{\Psi}(x, t) = A(t) e^{-\gamma t_x} S(\gamma t_x b(t))$$

$$S(z) = \int_0^z \sin\left(\frac{\pi t^2}{2}\right) dt$$

Frenel Integral

$$t_x = t - \frac{x}{v}$$

$$\gamma = 2\pi |g(\omega)|^2 \rho(\omega)$$

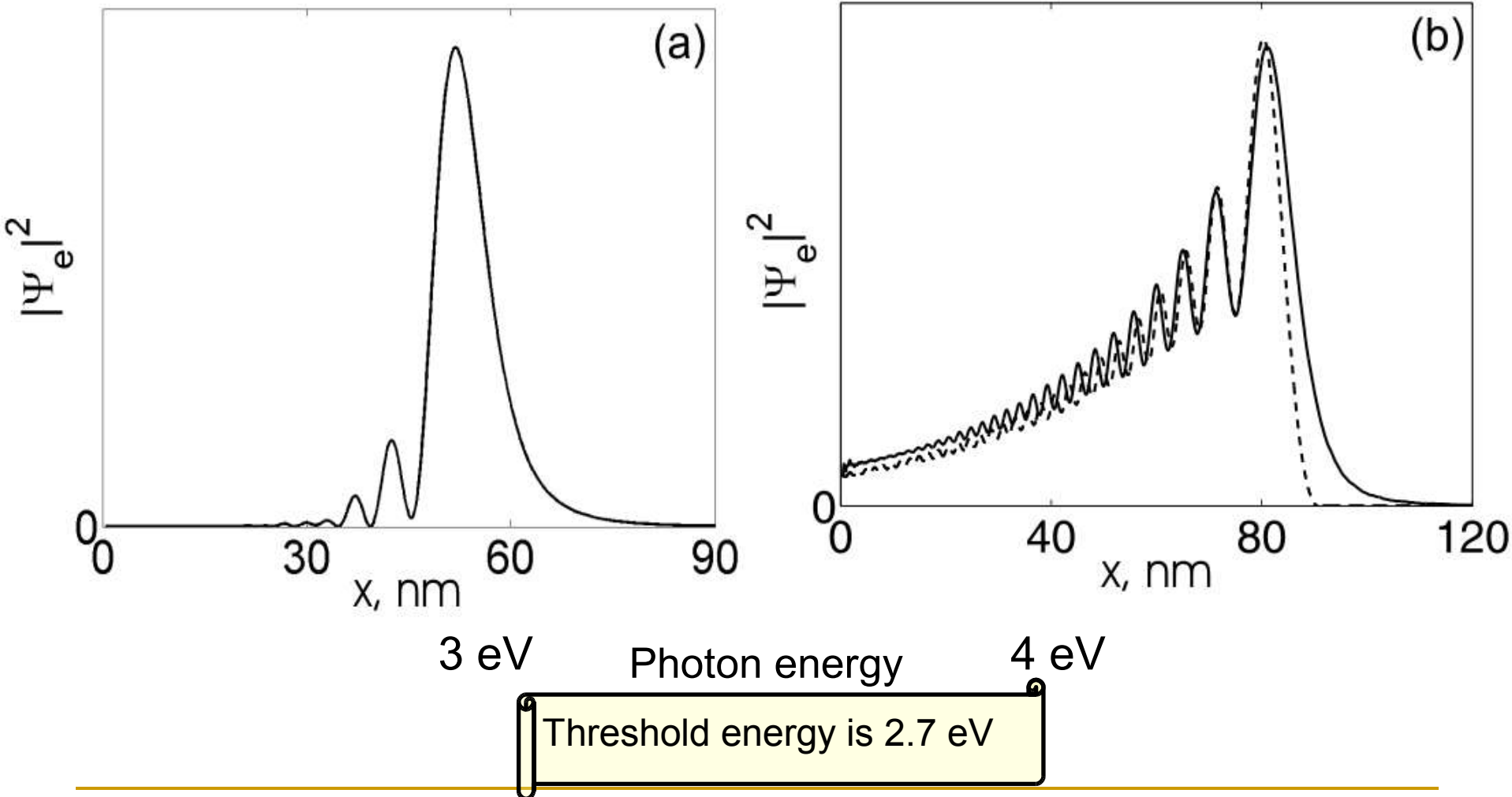
decay

$$A(t) = e^{-\alpha t}$$

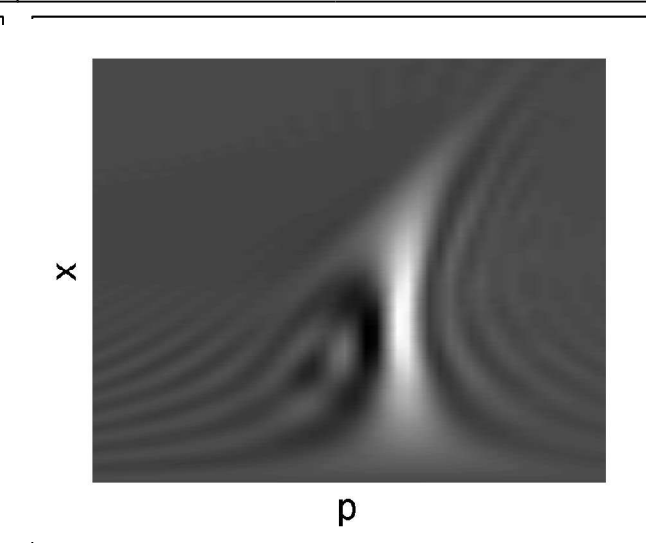
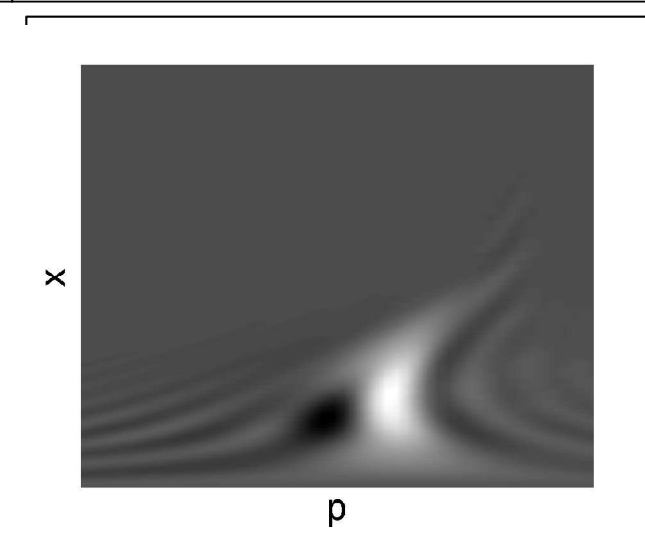
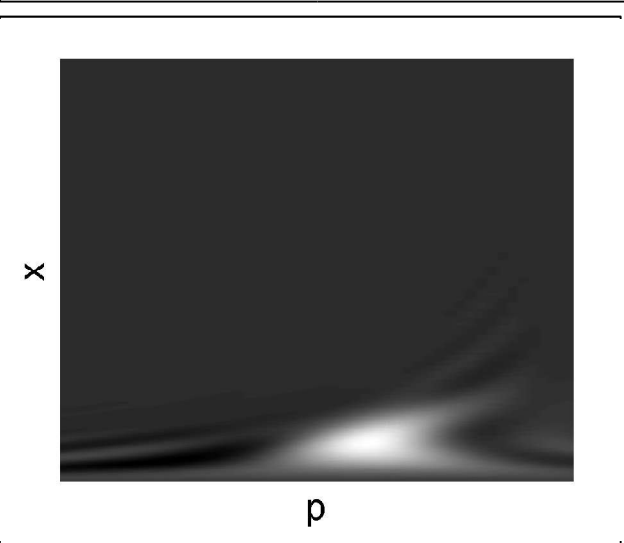
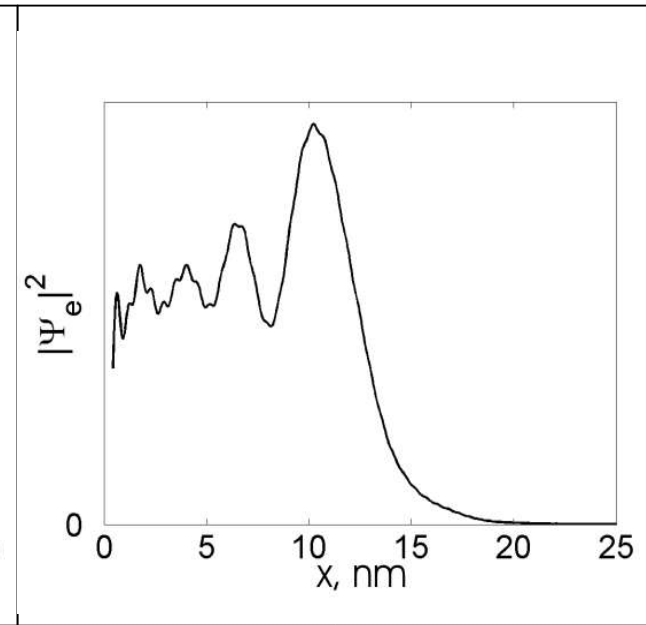
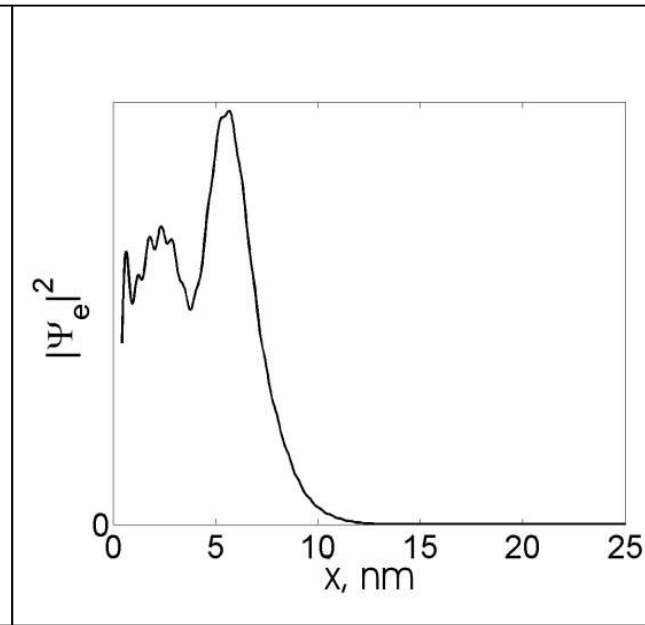
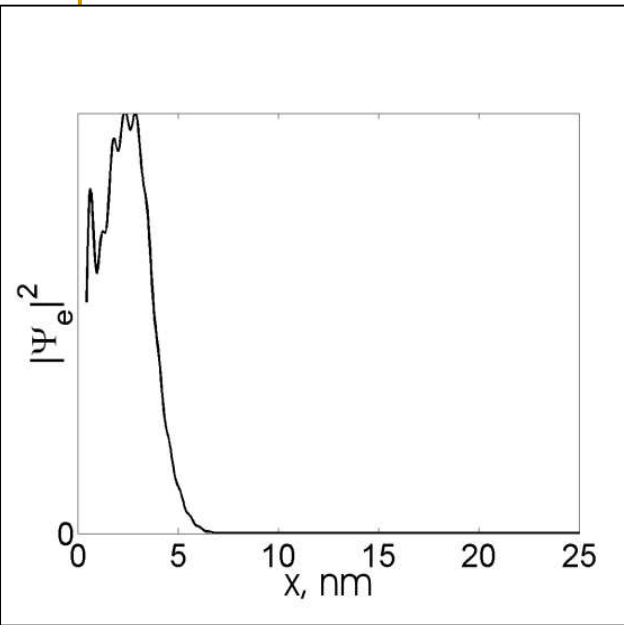
$$b(t) = b_1 e^{-\gamma b_2 t}$$

These functions are responsible for decreasing the amplitude of the packet and widening of the structure accordingly

Wavepacket of the free photoelectron and its approximation



Wigner Function of the free electron



3D wave packet

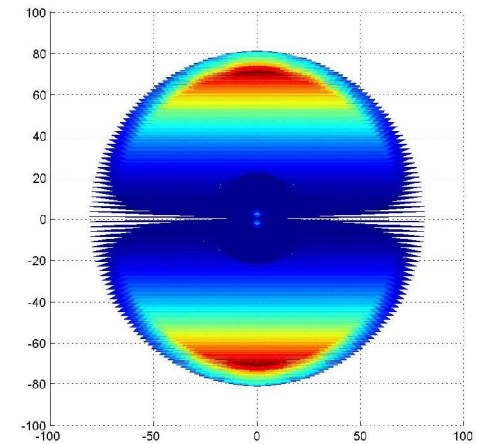
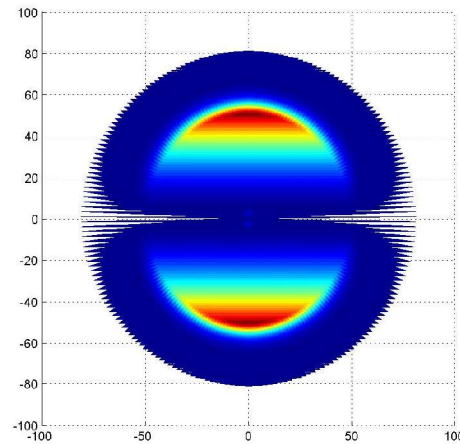
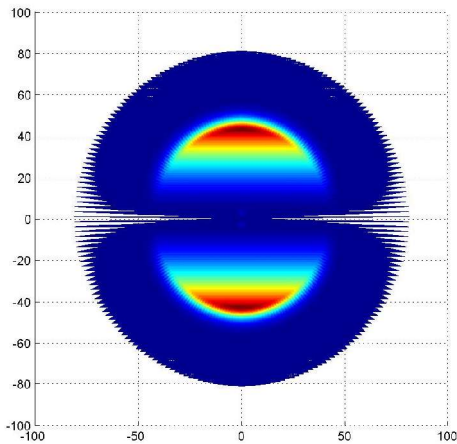
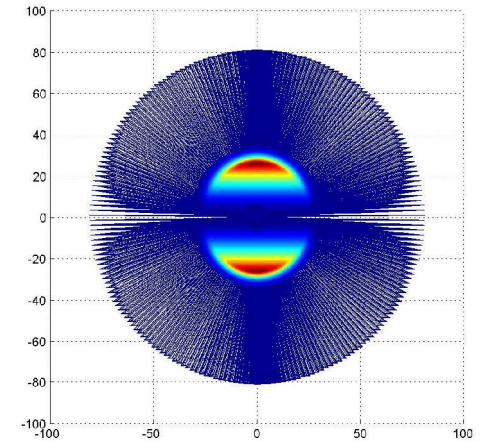
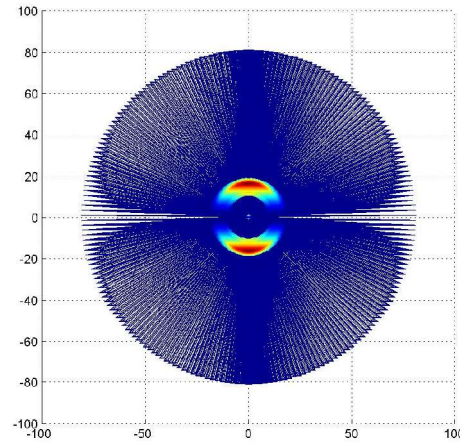
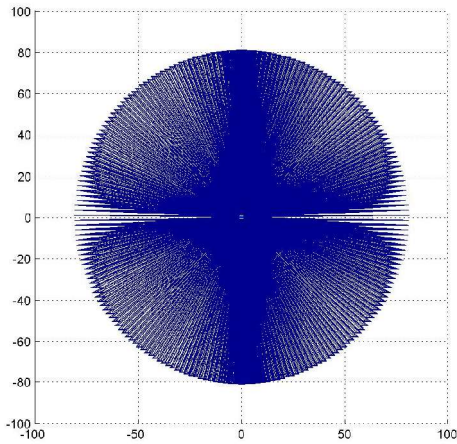
$$\Psi_{3D} = R(r) Y_{lm}(\theta, \varphi)$$

$$R(r) = \frac{C}{\sqrt{r}} \left\{ J_{3/2}(k_2' r) - \text{ctg}(k_2' L) J_{-3/2}(k_2' r) \right\}$$

l, m	0, 0	1, 0	1, ± 1
Y_{lm}	$\frac{1}{\sqrt{4\pi}}$	$i\sqrt{\frac{3}{\pi}} \cos\theta$	$\mp i\sqrt{\frac{3}{8\pi}} \sin\theta e^{i\varphi}$

l is equal to 0
for initial state
of the electron

3D – model of the wavepacket



Conclusion

- Temporal dynamics of the free electron wavepacket has been investigated and finite time of its creation was found
 - The oscillations of the trailing edge of the wavepacket were found
 - Wigner function showed that photoelectron is in a squeezed state
 - 3D-model of wavepacket has been built
-