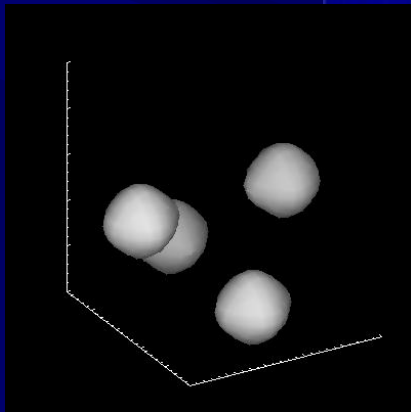
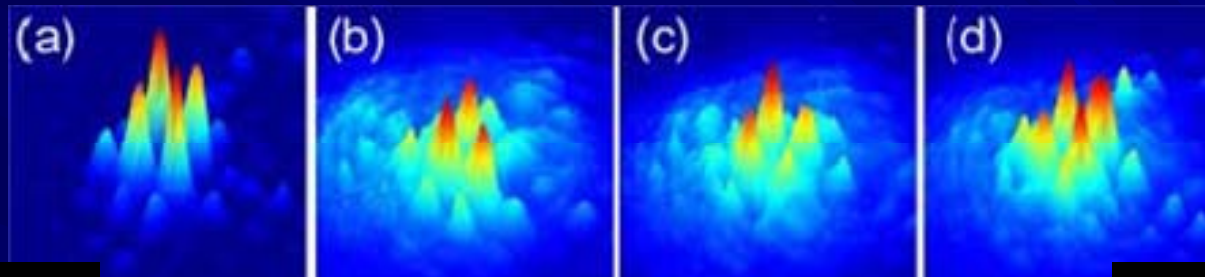
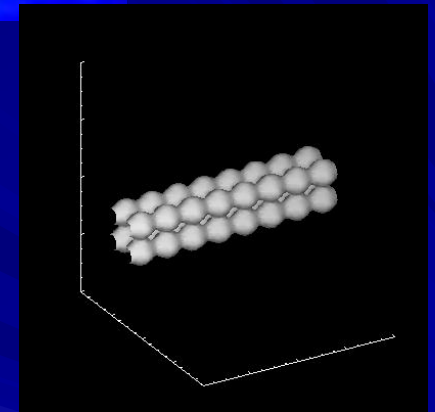


Nonlinear dynamics and solitons in optical lattices



Yuri Kivshar



Nonlinear Physics Centre
Australian National University
Canberra, Australia

Collaborators: “BEC group”

Elena Ostrovskaya, a key researcher

Tristram Alexander, a postdoc

Pearl Louis, a PhD student

Beata Dąbrowska, a PhD student



Our evolution: 1993-2004

Nonlinear Physics Center

2004

1993

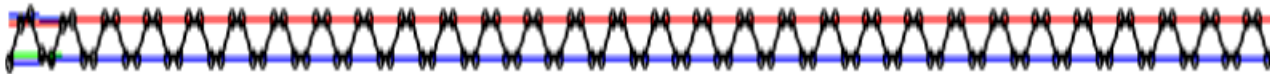
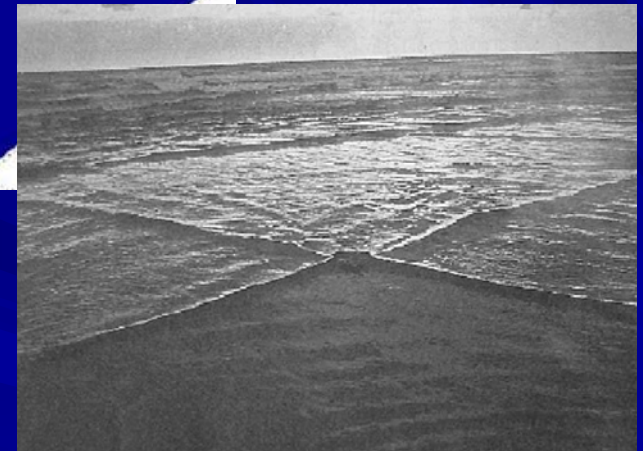
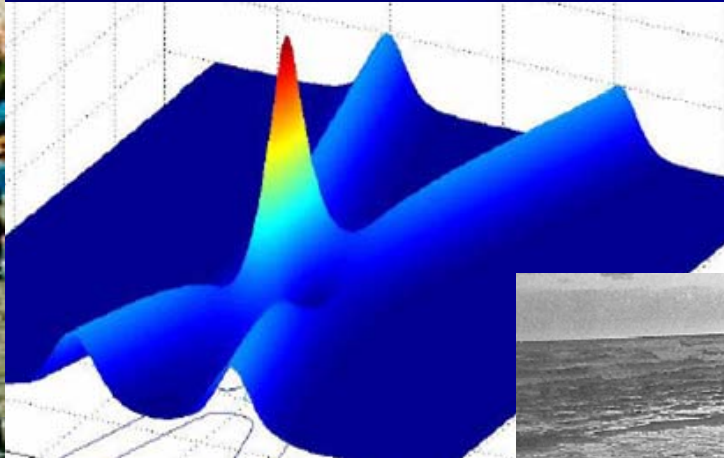


(Only “BEC people” of the Center)

Outline

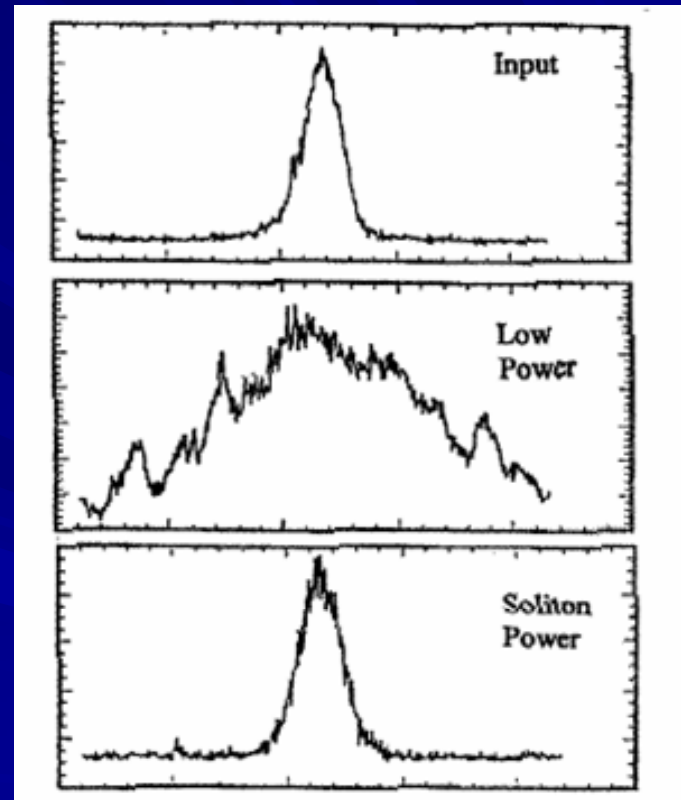
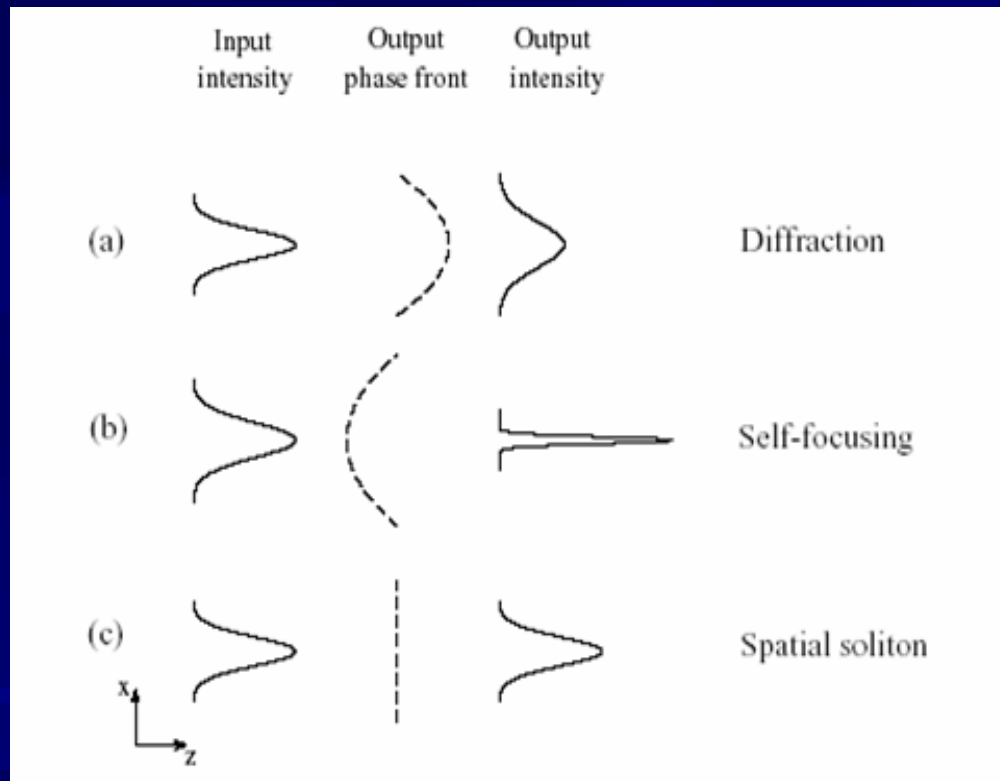
- Solitons: optics vs. BEC
- Bloch waves and band-gap spectrum
- Gap solitons in 1D lattices
- Mobility of 1D gap solitons
- 2D optical lattices and 2D gap solitons
- 2D discrete vortices
- Experimental observations in optics
- 3D discrete vortices

Solitons



Optics: Self-focusing and solitons

experiment



Nonlinear Kerr medium:

$$n^2 = n_0^2 + n_2 I$$

Matter waves in the mean-field picture

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi + V_{\text{Trap}}(\mathbf{r})\psi + g\psi|\psi|^2$$

- Striking similarity with coherent light waves
- “Kerr-type” nonlinearity due to atom-atom interactions
- Main difference: atoms have mass!
- Usually a trapping potential is included

Nonlinear matter waves

- **Nonlinearity** grows with the number of atoms N
(nonlinearity is always there)
- Repulsive: **dark** solitons and **vortices**
- Attractive: **bright** matter-wave solitons
- Can be tuned via **Feshbach resonances**
 - “Focusing”: attractive interactions ($g < 0$)
 - “Defocusing”: repulsive interactions ($g > 0$)
- **Optical lattices** create tunable Bragg gratings

Solitons in nonlinear atom optics

Macroscopic wave function: the mean field theory

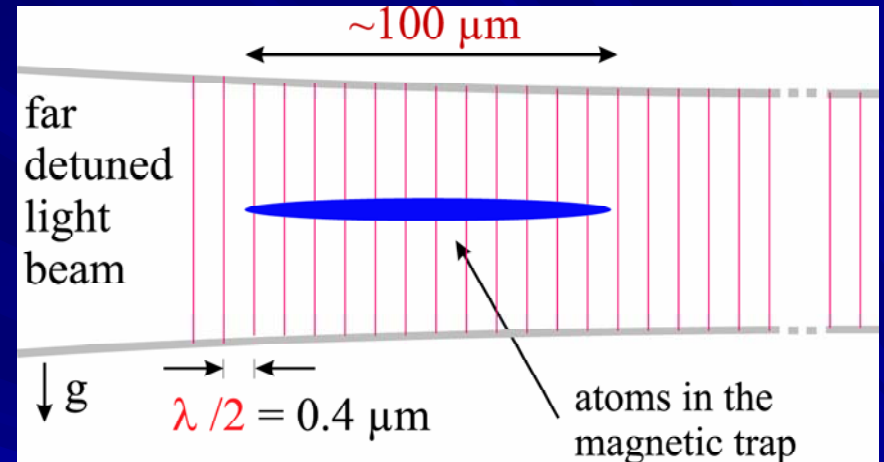
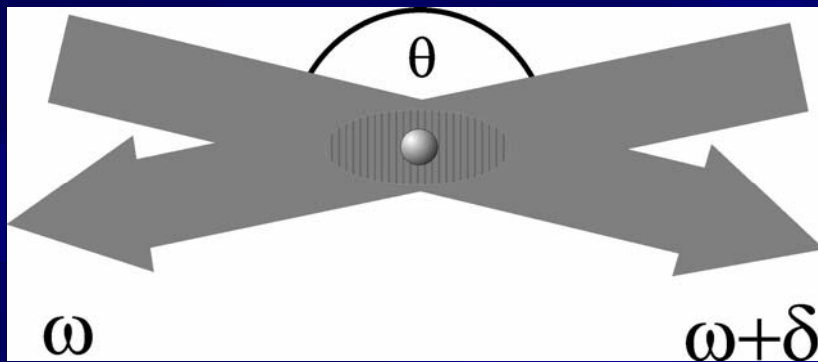
The Gross-Pitaevsky equation vs. the NLS equation

- Vortices and vortex solitons (1999)
- Dark solitons (1999-2000)
- Bright solitons (2002)
- Matter-wave gap solitons (2004)

Solitons in periodic structures

BEC in 1D optical lattices

$$V_{\text{Latt}} = V_0 \sin^2(k_L x + \delta t)$$



$$V_0 \sim I; \quad k_L \equiv \frac{\pi}{d} = \frac{\omega}{c} \cos \theta; \quad v = \frac{\delta}{2k_L}$$

$$\lambda_D \equiv \frac{2\pi\hbar}{p_x} \gg d \quad \Rightarrow \quad \frac{p_x^2}{2m} \equiv E_k \leq E_R \equiv \frac{\hbar^2 k_L^2}{2m}$$

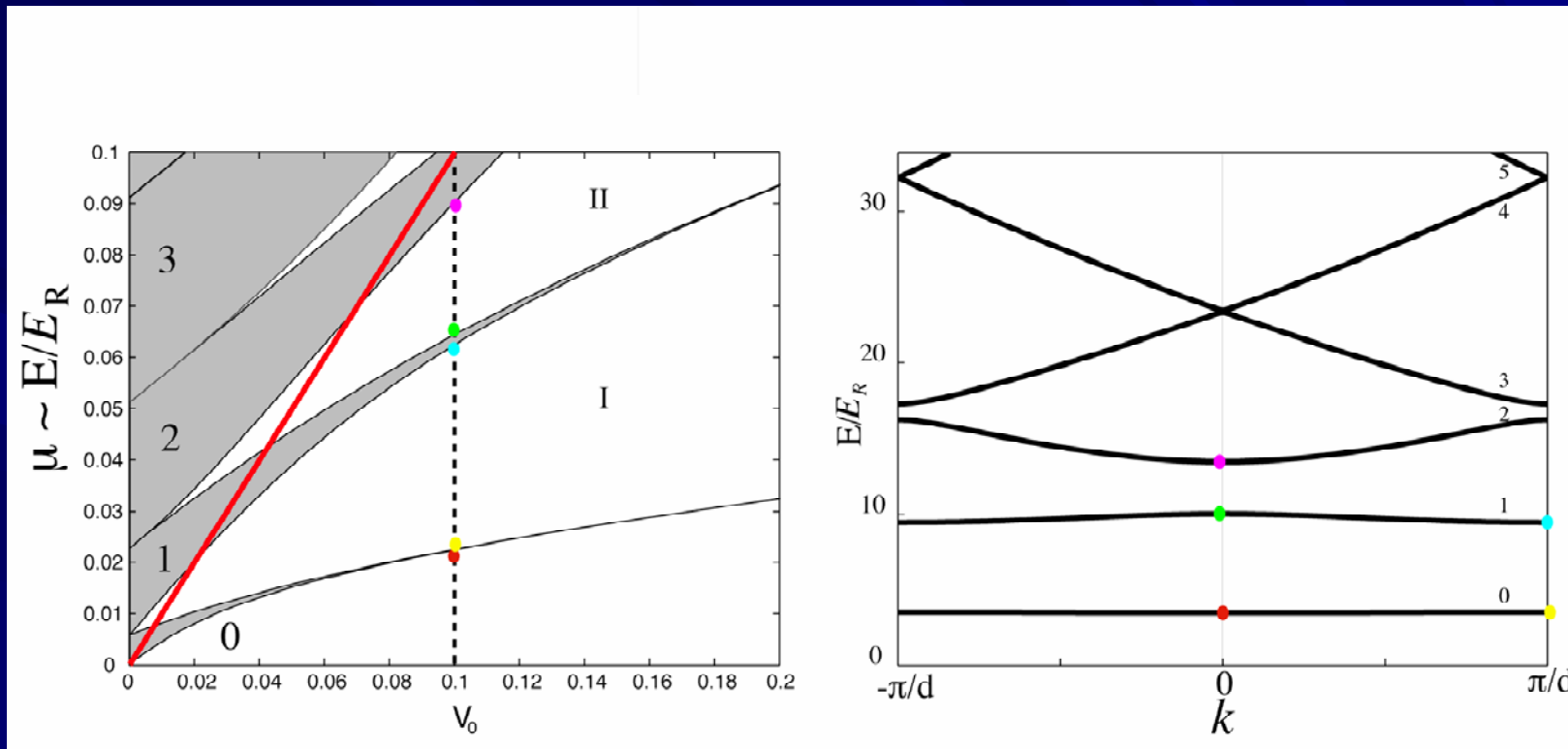
Atomic band-gap structures

BEC in optical lattices -

nonlinear band-gap structures for matter waves

- ➡ Unprecedented control over parameters
- ➡ Nonlinearity that can be turned off or reversed
- ➡ Degree of discreteness can be varied
- ➡ Dimensionality can be changed
- ➡ Effective dynamical tunability

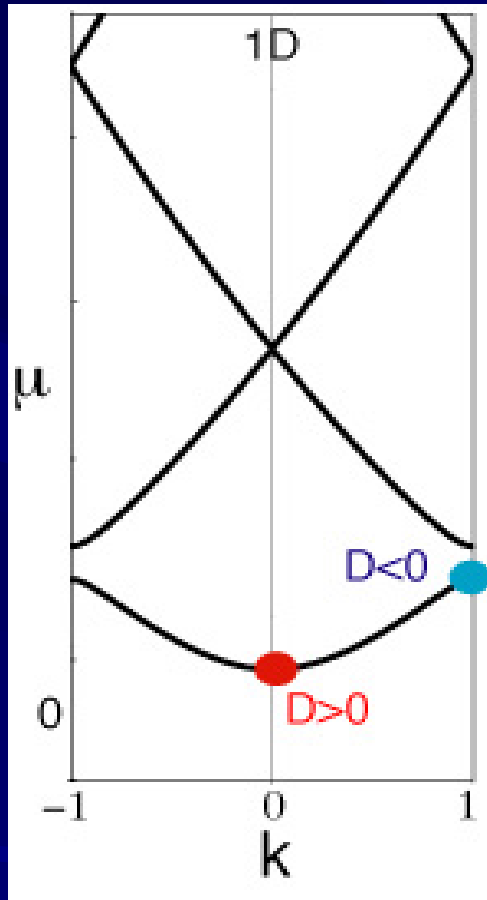
Band-gap spectrum



$$\psi(x, t) = \psi(x) \exp(-i\mu t)$$

$$\frac{1}{2} \frac{d^2 \psi}{dx^2} - \{V_0 \sin^2(k_L x) - \mu\} \psi = 0$$

Gap solitons



Dispersion relation in a free space:

$$E(p) = \frac{p^2}{2m}$$

Dispersion for BEC in a lattice is defined by

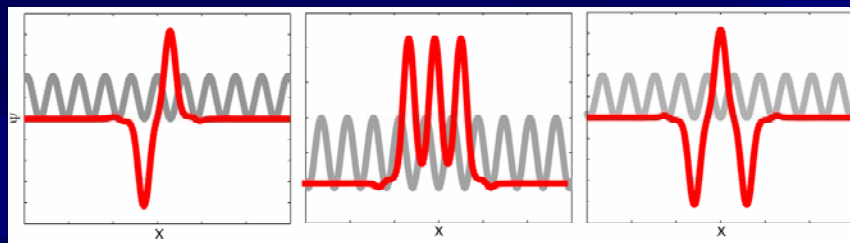
$$E(k) \equiv \mu(k)$$

$$D = \frac{\partial^2 \mu}{\partial k^2} \begin{cases} > 0 \text{ at the band's centre} \\ < 0 \text{ at the band's edge} \\ \text{(negative "mass")} \end{cases}$$

Gap solitons - balance between “anomalous” dispersion and positive interaction (repulsion) at the upper band edge

Gap solitons in 1D structures

Repulsive interactions!

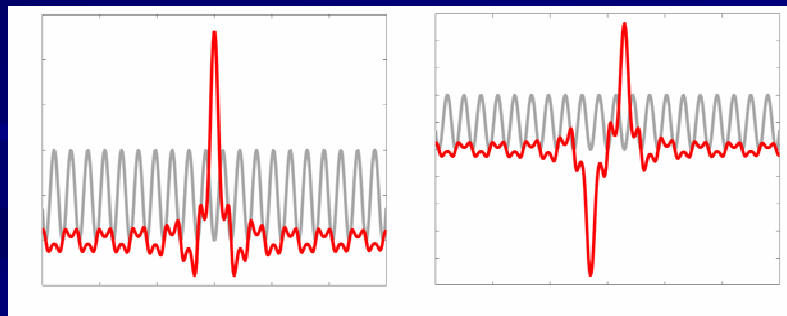


B

C

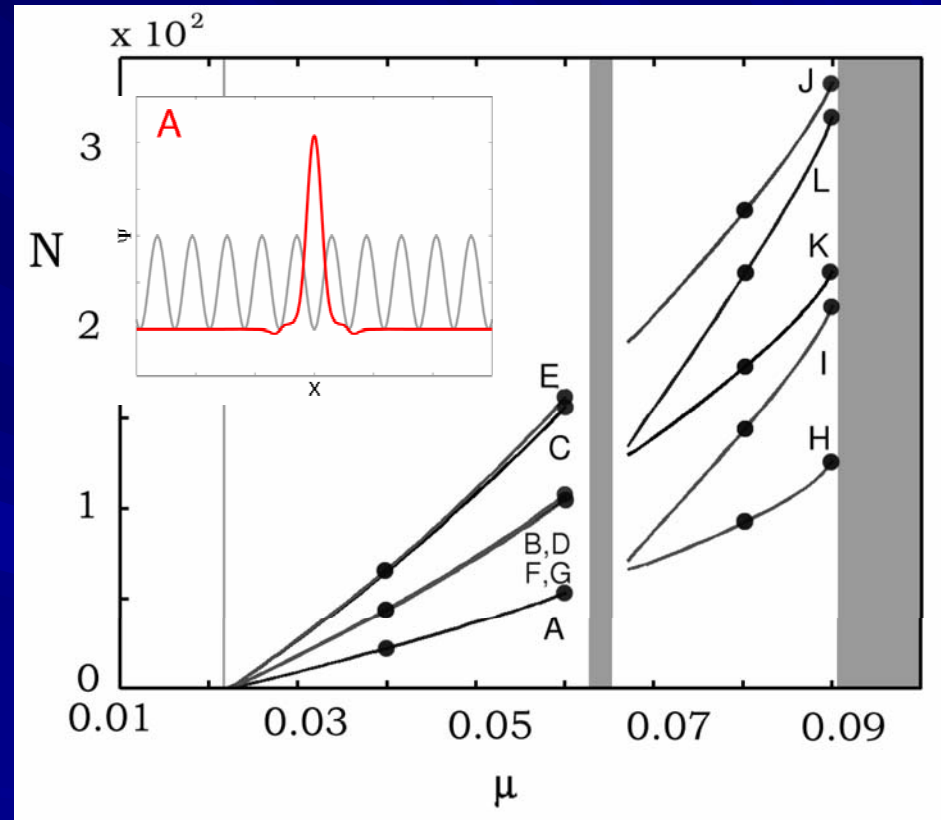
E

Inside the band:



A

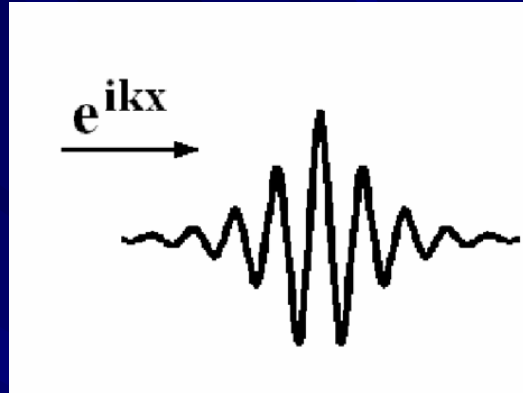
B



P. Louis et al., Phys. Rev. A **67**, 013602 (2003)

www.rsphysse.anu.edu.au/nonlinear

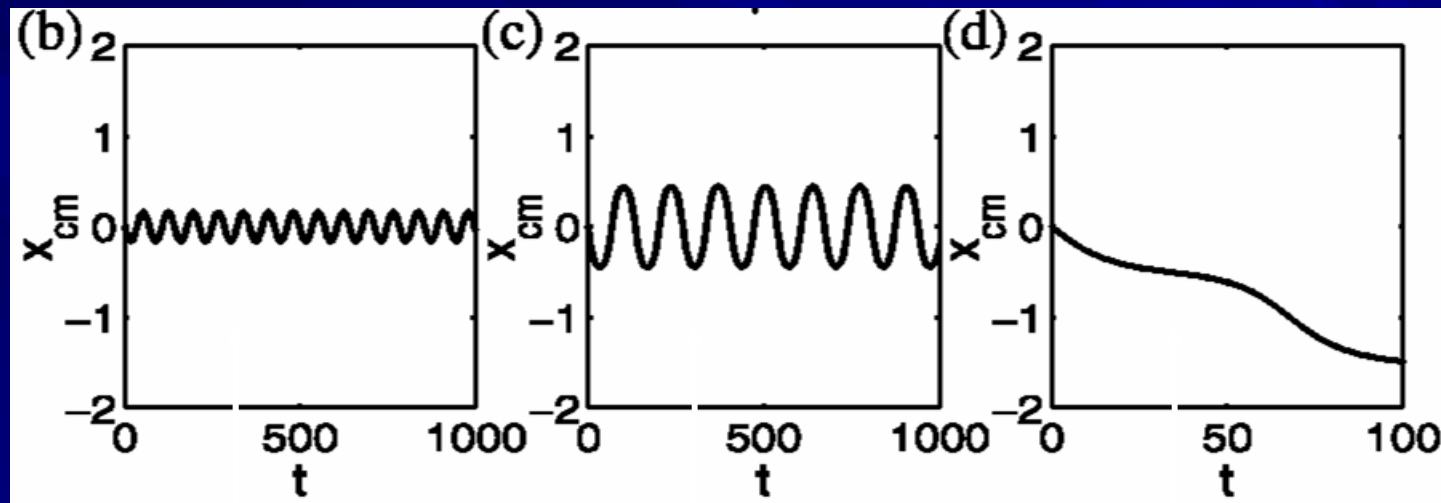
Soliton mobility



$$\mu = 7.7$$

$$k_{PN} = 1.2 \times 10^{-2}$$

(free motion)



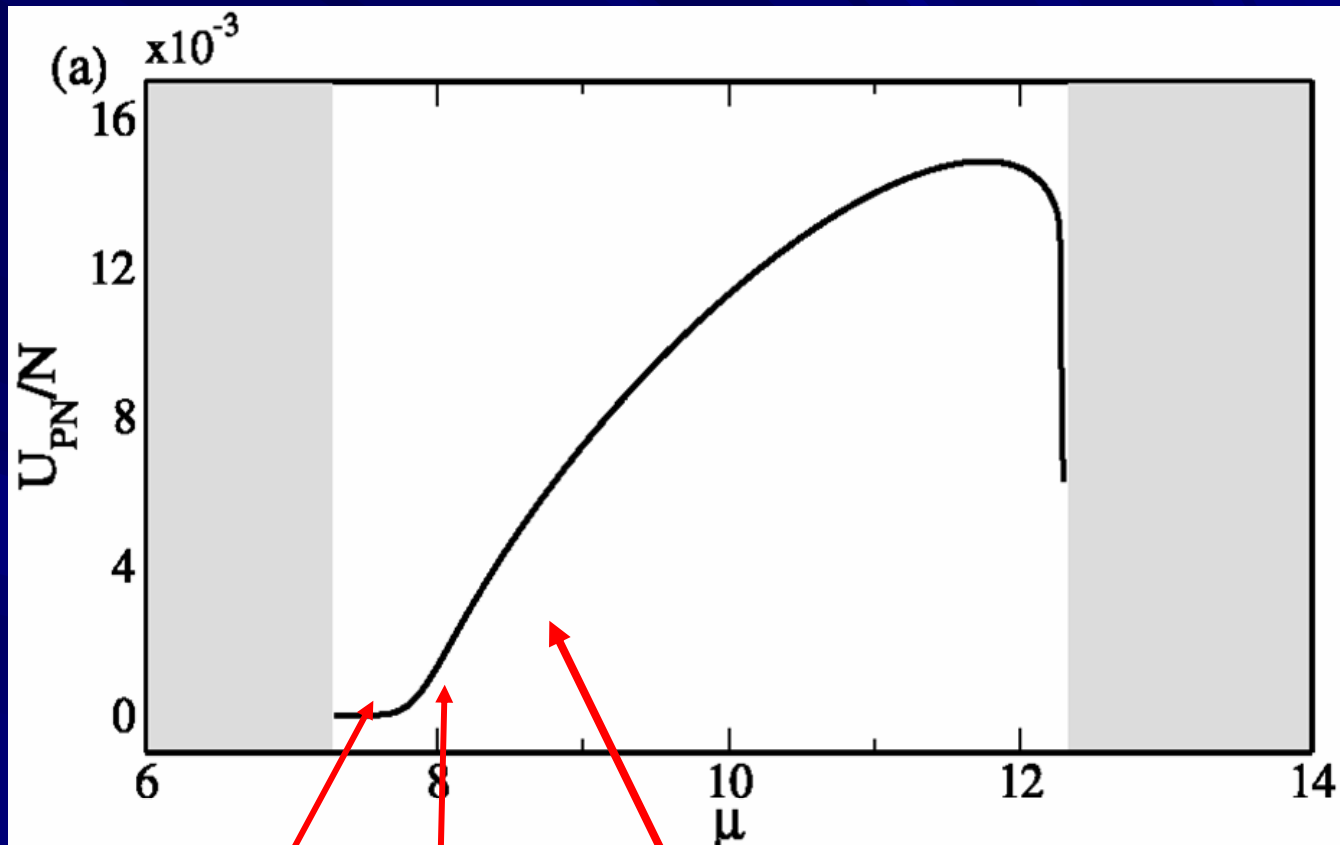
$$k < k_{PN}$$

$$k = k_{PN}$$

$$k > k_{PN}$$

(oscillations within one lattice site)

Peierls-Nabarro potential

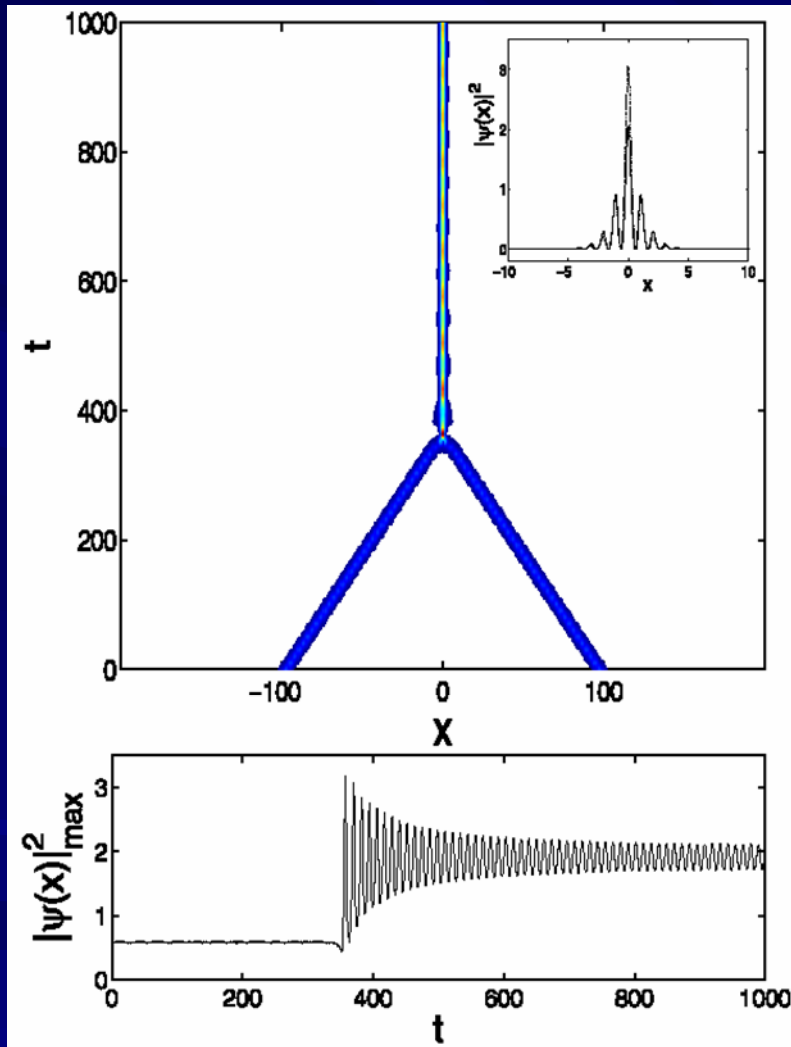


$\mu < 7.5$
lattice-free
solitons

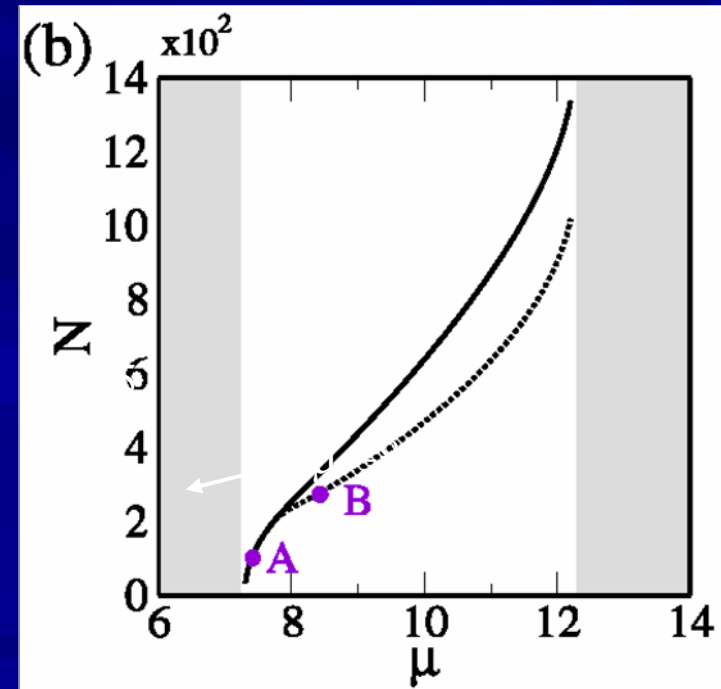
$7.5 < \mu < 8.2$
discrete NLS
solitons

$\mu > 8.2$
centre of mass
oscillations

Soliton fusion

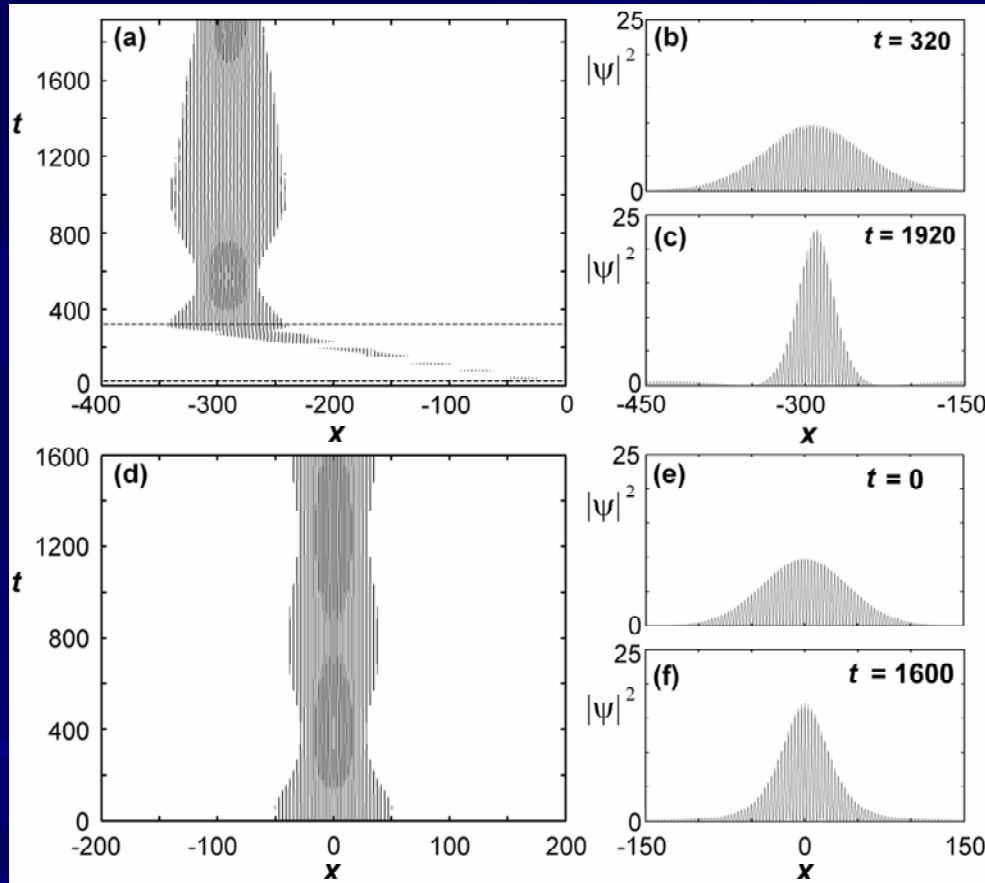


discrete NLS $\mu \geq 7.5$
strongly inelastic scattering



170 $\rightarrow$ 270 atoms

Generation of gap solitons



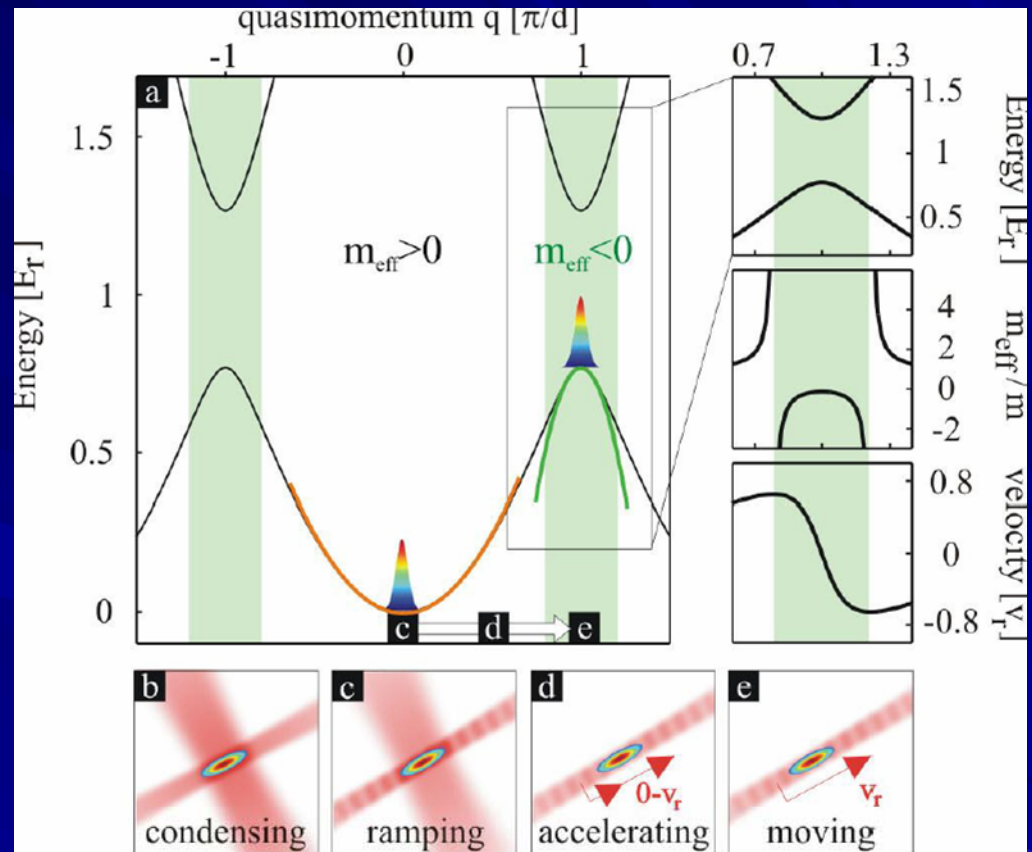
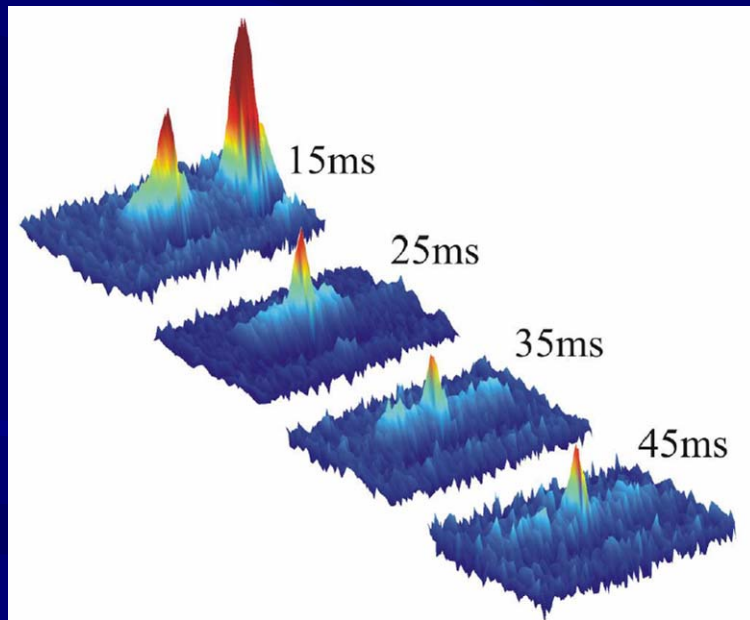
Top half

1. BEC loaded at $k=0$ by slowly ramping up the lattice.
2. Lattice slowly accelerated to $v = v_{\text{recoil}}$. Adiabatically moves BEC to the gap edge.
3. Wavepacket evolved at the gap-edge in a lattice moving with velocity $v = v_{\text{recoil}}$.

Bottom half

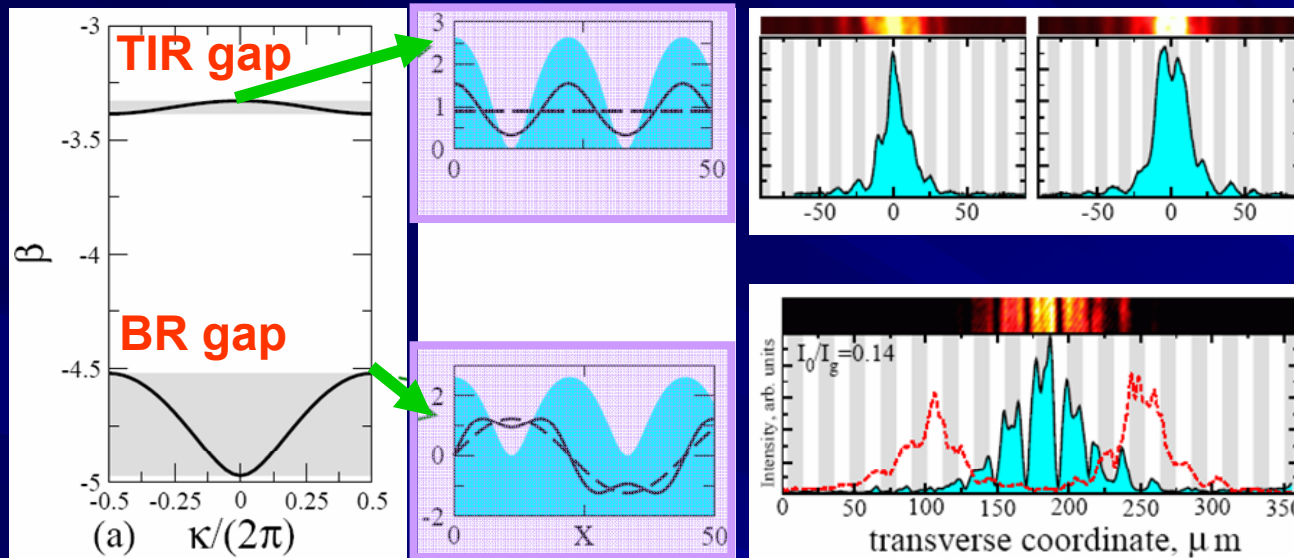
Generates the correct internal structure using interference rather than acceleration.

Experiment: Gap Solitons in BEC



Oberthaler's group, 2004

Spatial gap solitons in optics



Discrete solitons:
Single beam excitation

Gap solitons:
two beams

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20 AUGUST 2004

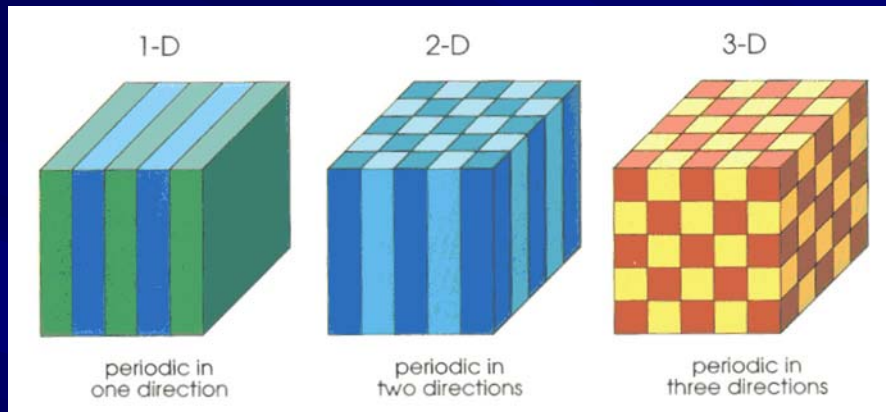
Controlled Generation and Steering of Spatial Gap Solitons

Dragomir Neshev, Andrey A. Sukhorukov, Brendan Hanna, Wieslaw Krolikowski, and Yuri S. Kivshar

*Nonlinear Physics Centre and Laser Physics Centre, Centre for Ultra-high bandwidth Devices for Optical Systems (CUDOS),
Research School of Physical Sciences and Engineering, Australian National University, Canberra, ACT 0200, Australia*

2D optical lattices

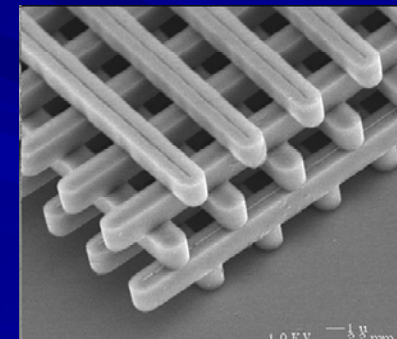
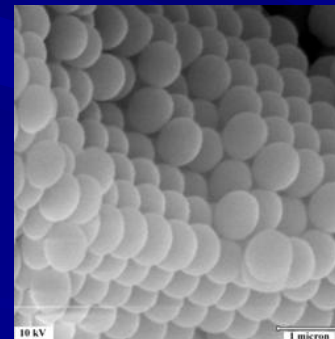
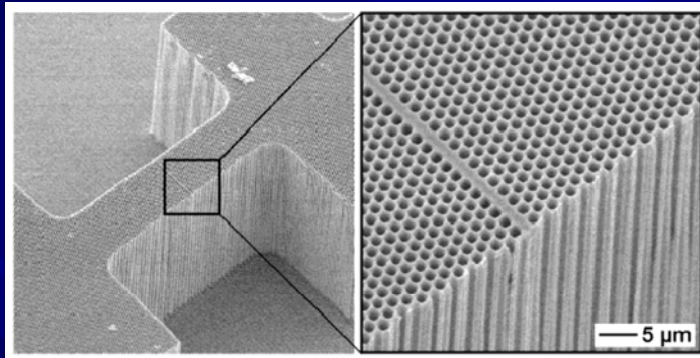
Photonic band-gap structures



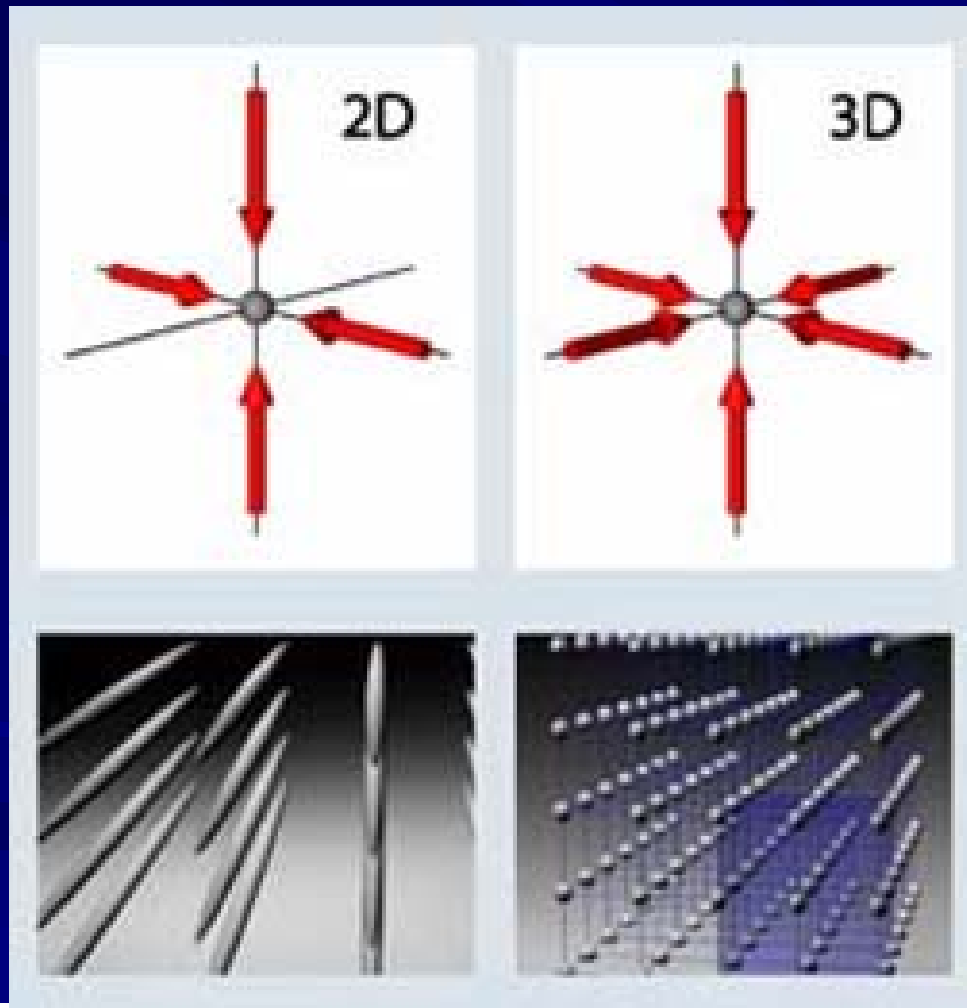
PRL 58 (1987):

- *S. John*

- *E. Yablonovitch*



Optical lattices in higher dimensions



Munich BEC group

2D: phase coherence
M. Greiner et al.
PRL 87, 160405 (2001)

3D: coherence break-up
M. Greiner et al.
Nature 415, 39 (2002)

BEC in a 2D optical lattice

$$i\frac{\partial\psi}{\partial t} + \nabla^2 \psi - V_{\text{Latt}}(\mathbf{r})\psi - \sigma\psi|\psi|^2 = 0$$

$$V_{\text{Latt}} = V_0 \left[\sin^2(k_L x) + \sin^2(k_L y) + 2\mathbf{e}_1 \cdot \mathbf{e}_2 \cos \theta \sin(k_L x) \sin(k_L y) \right]$$

“Square” lattice and Repulsive nonlinearity $\sigma > 0$

Stationary states of noninteracting condensate - Bloch waves:

$$\psi(\mathbf{r}, t) = \varphi(\mathbf{r}) \exp(-i\mu t); \quad \varphi(\mathbf{r}) = u_{\mathbf{k}}(\mathbf{r}) \exp(i\mathbf{k}\mathbf{r})$$

$$\mathbf{k} \in \text{BZ}; \quad u_{\mathbf{k}}(\mathbf{r}) = u_{\mathbf{k}}(\mathbf{r} + \mathbf{d})$$

2D band-gap spectrum

$$i\frac{\partial\psi}{\partial t} + \nabla^2\psi - V_L(\mathbf{r})\psi - \psi|\psi|^2 = 0$$

$$V_{\text{Latt}} = V_0 [\sin^2(x) + \sin^2(y)]$$

Stationary states:

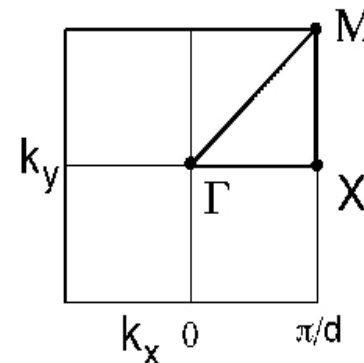
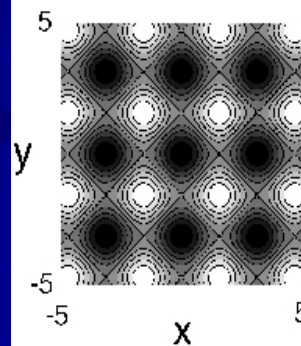
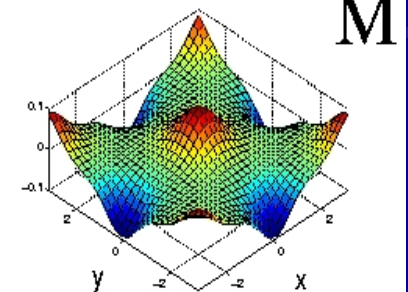
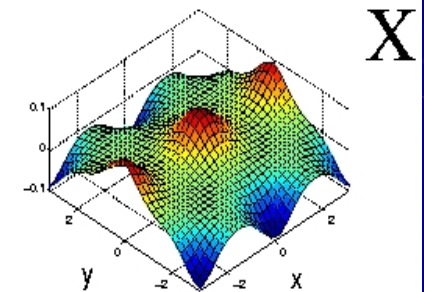
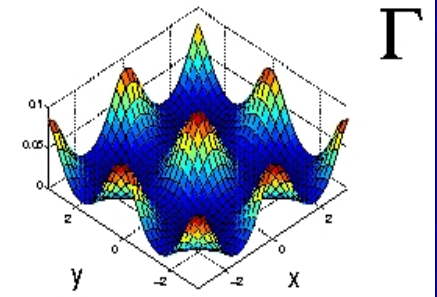
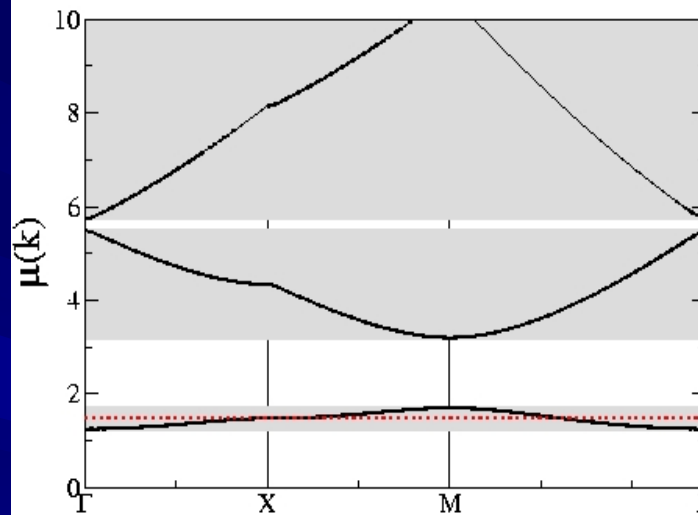
$$\psi(\mathbf{r}, t) = \varphi(\mathbf{r}) \exp(-i\mu t);$$

Bloch states:

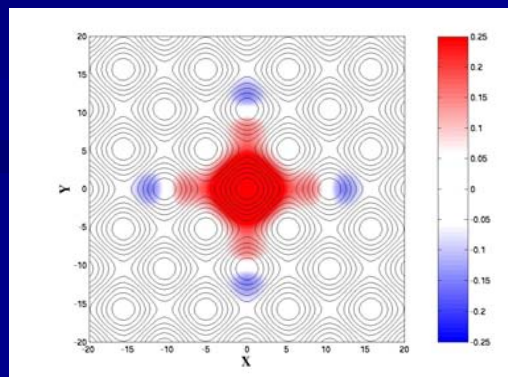
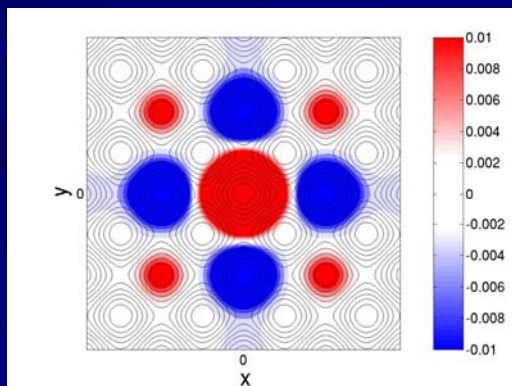
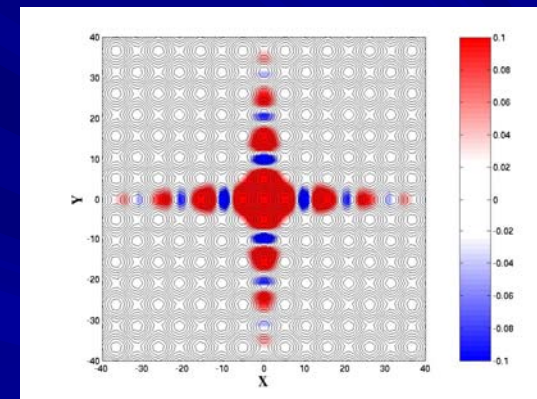
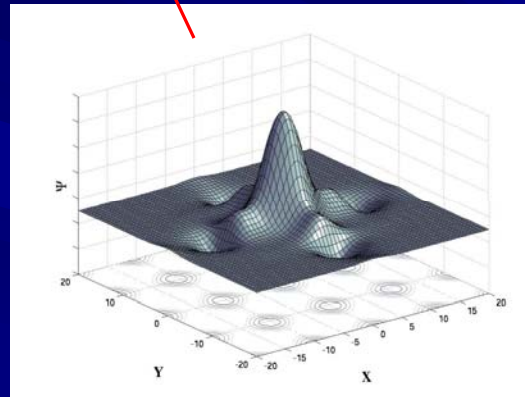
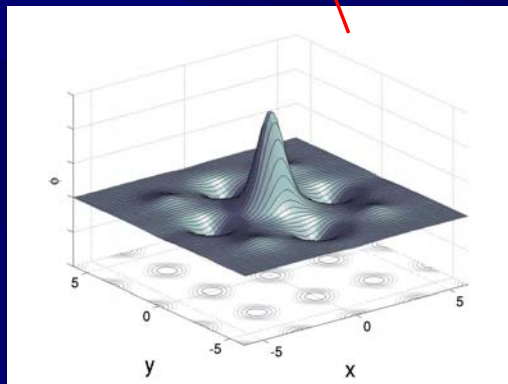
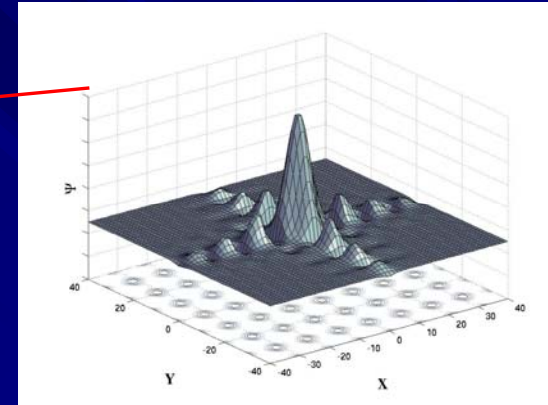
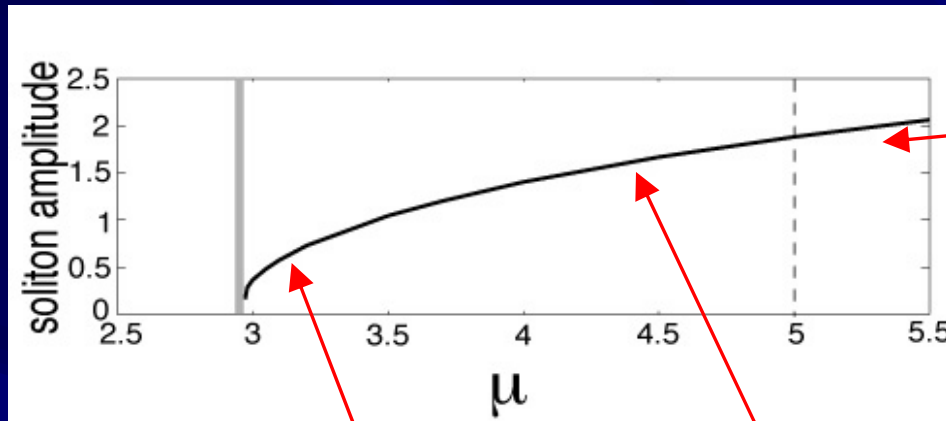
$$\varphi(\mathbf{r}) = u_{\mathbf{k}}(\mathbf{r}) \exp(i\mathbf{k}\mathbf{r})$$

$$\mathbf{k} \in \text{BZ};$$

$$u_{\mathbf{k}}(\mathbf{r}) = u_{\mathbf{k}}(\mathbf{r} + \mathbf{d})$$



Gap solitons in 2D lattices

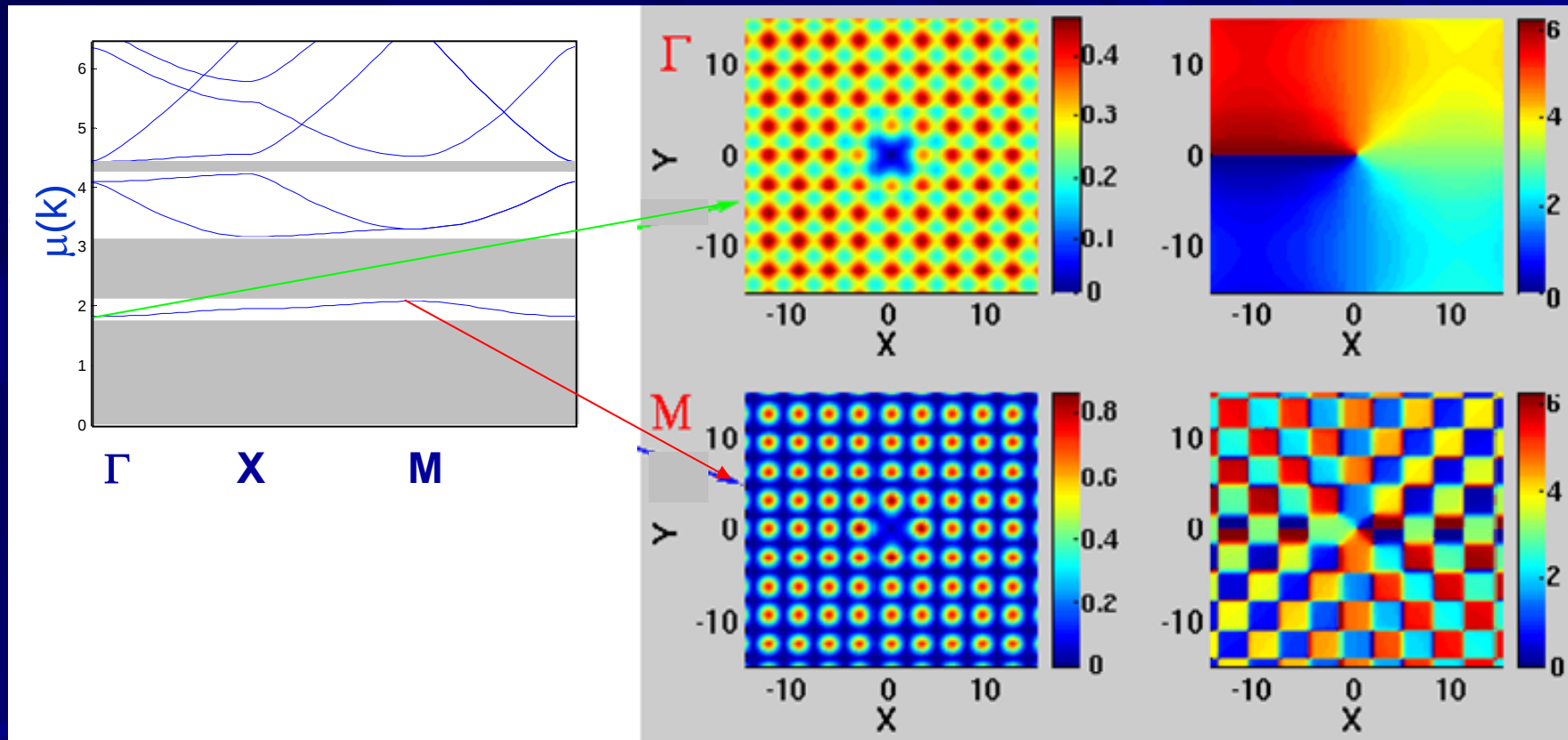


Ostrovskaya et al, PRL
90, 160407 (2003)

www.rsphysse.anu.edu.au/nonlinear

2D optical lattices and discrete vortices

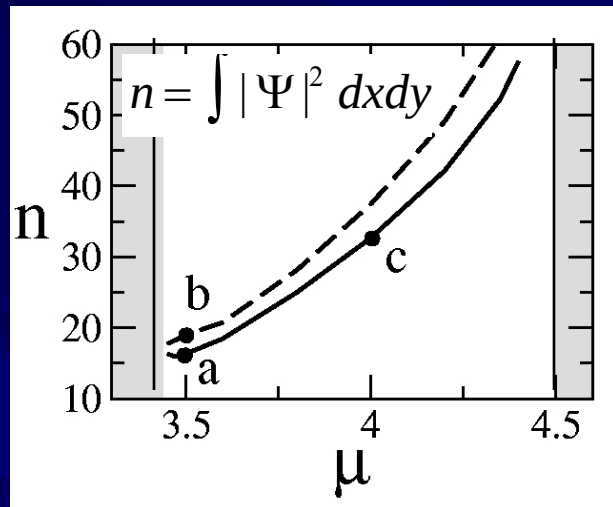
Imprinting a phase: BEC vortex in-band



Bloch vortex at $k=0$: $\psi(r, \theta) \sim A \exp(im\theta) \exp(-a^2 r^2) u_{1, k_\Gamma}(r, \theta)$

Bloch vortex at $k=k_M$: $\psi(r, \theta) \sim A \exp(im\theta) \exp(-a^2 r^2) u_{1, k_M}(r, \theta)$

Stationary in-gap states: gap vortices



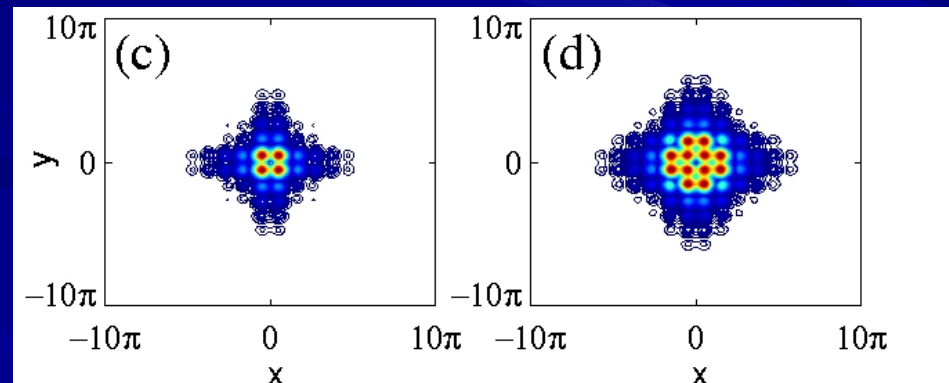
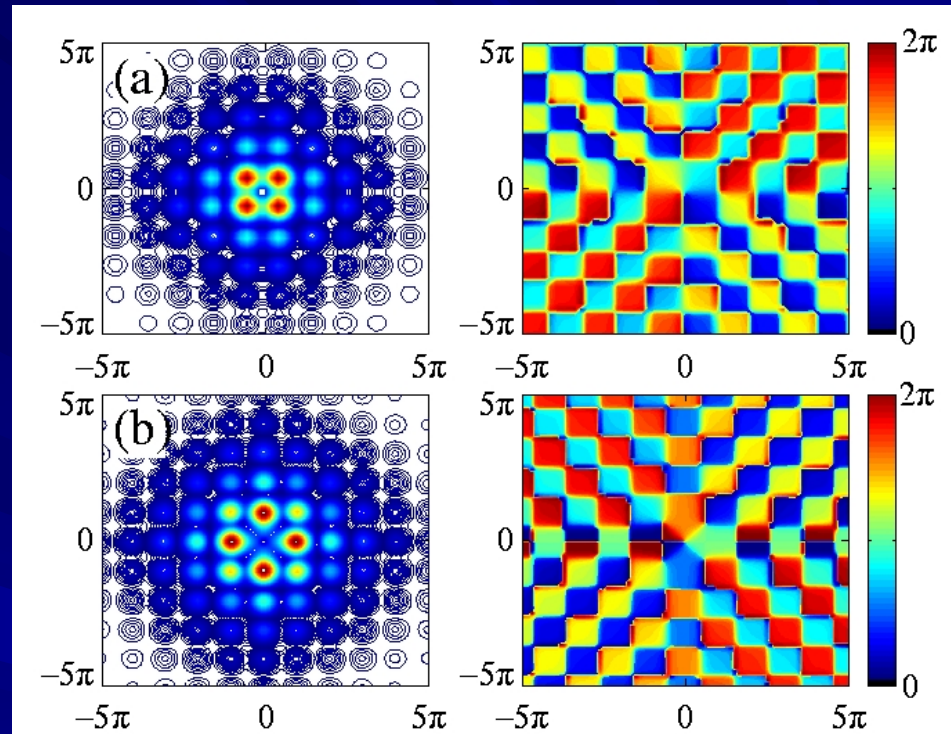
two symmetry types:

(a,c) off-site vortex

(b) on-site vortex

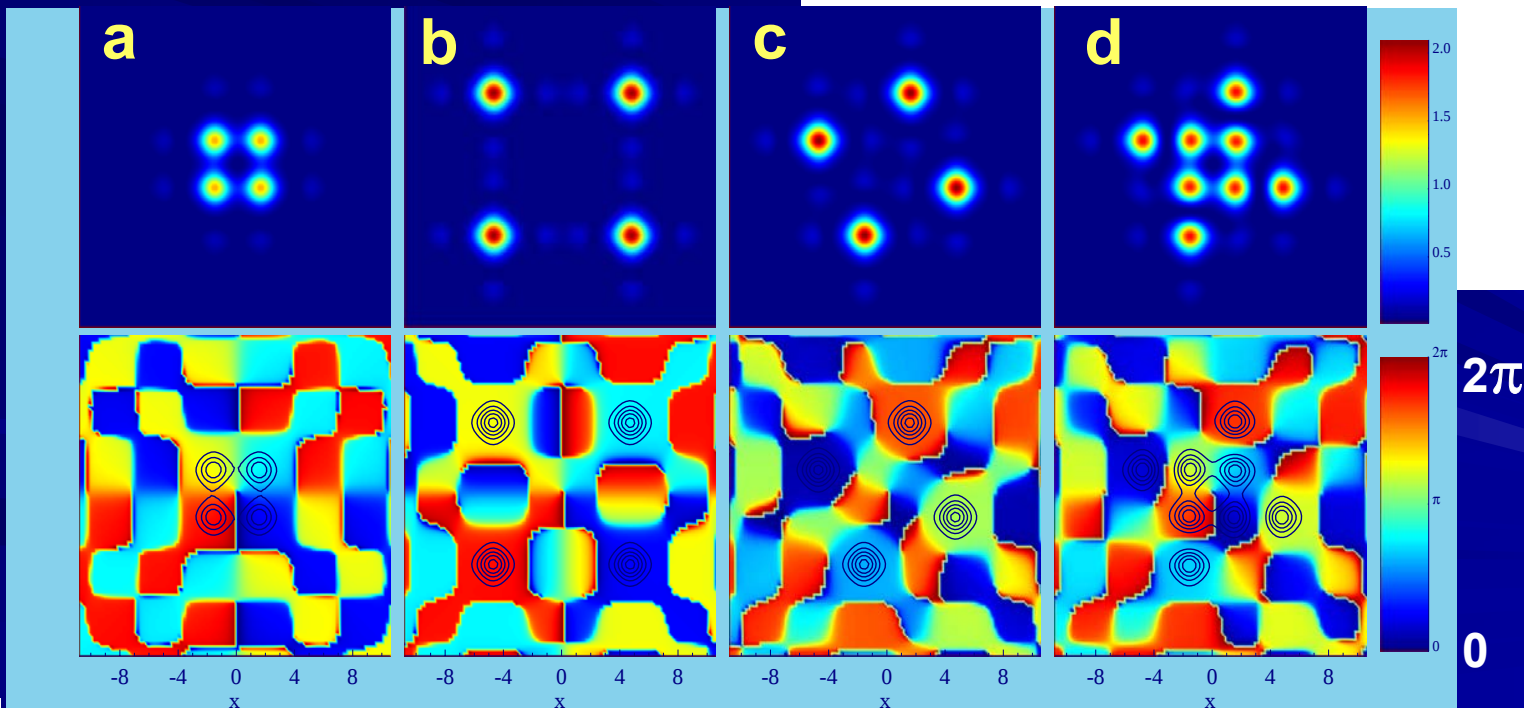
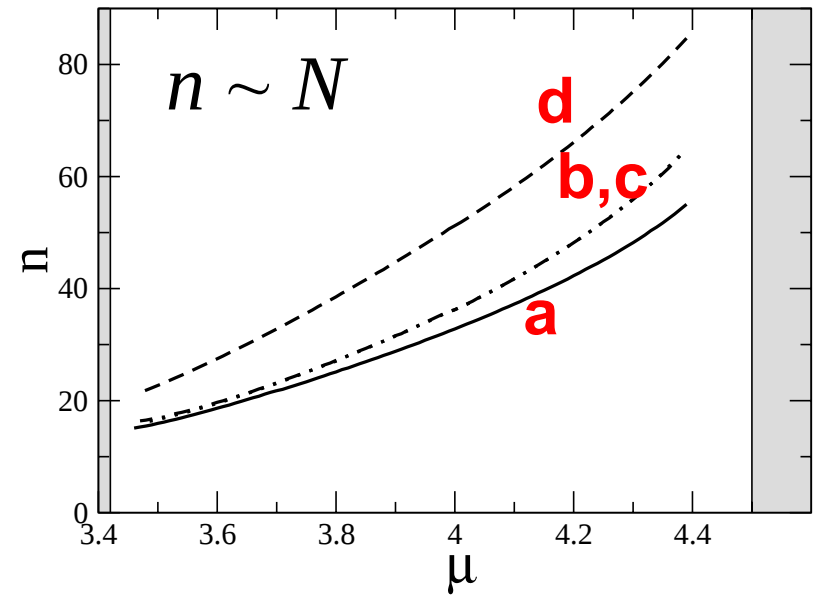
and different widths:

(d) wide off-site vortex



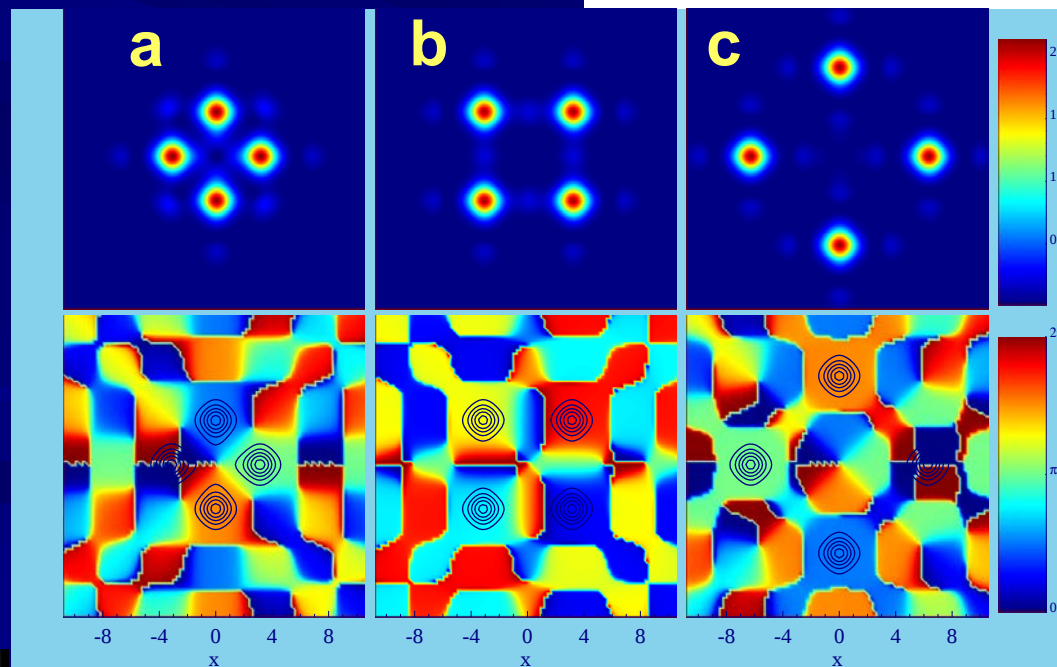
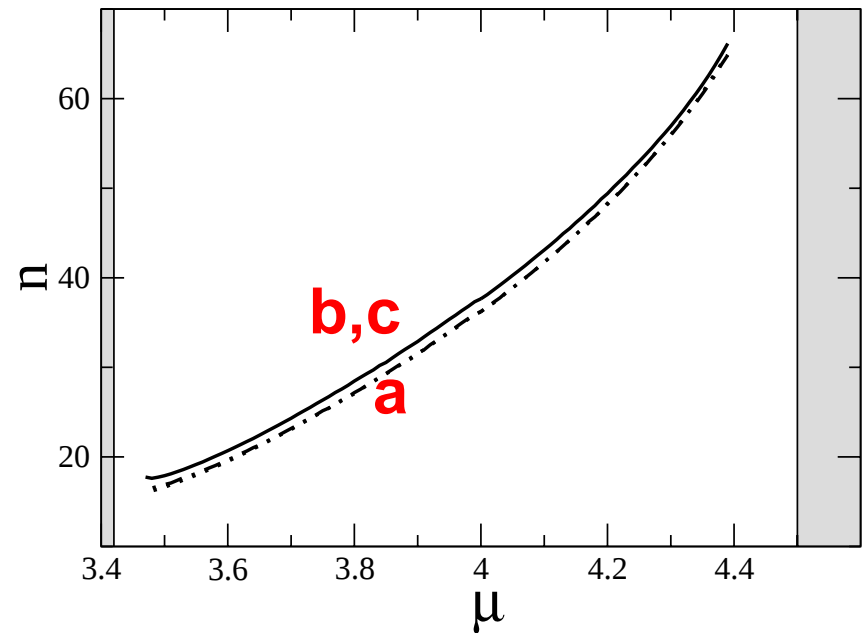
Off-site vortices

- Off-site vortices of different radii
- Almost degenerate in atom number



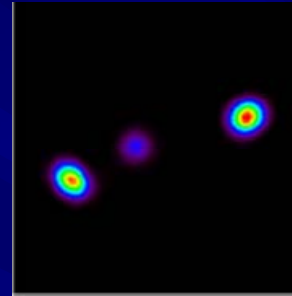
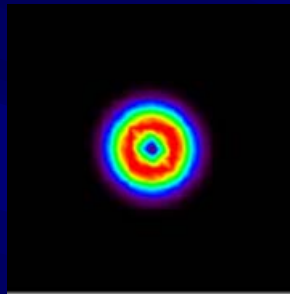
On-site vortices

- On-site vortex families are even closer in atom numbers

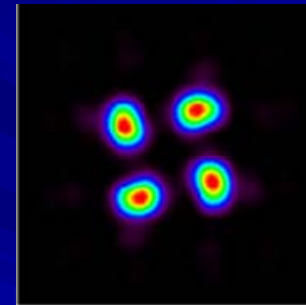


Discrete vortices in nonlinear optics

Why use a lattice?



Unstable,
no lattice



Stable on a lattice

Periodic potential leads
to vortex stability

Discrete optical vortex: observation

Optically induced 2D grating in a photorefractive crystal

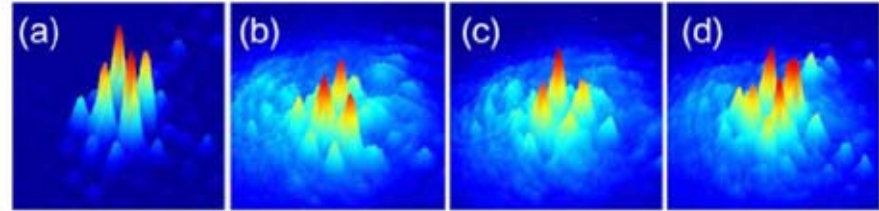
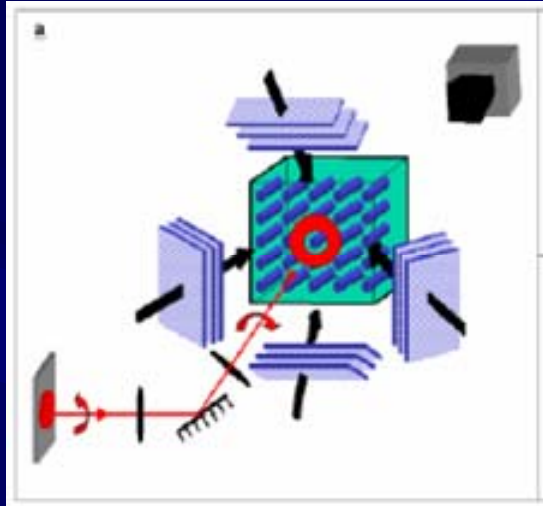
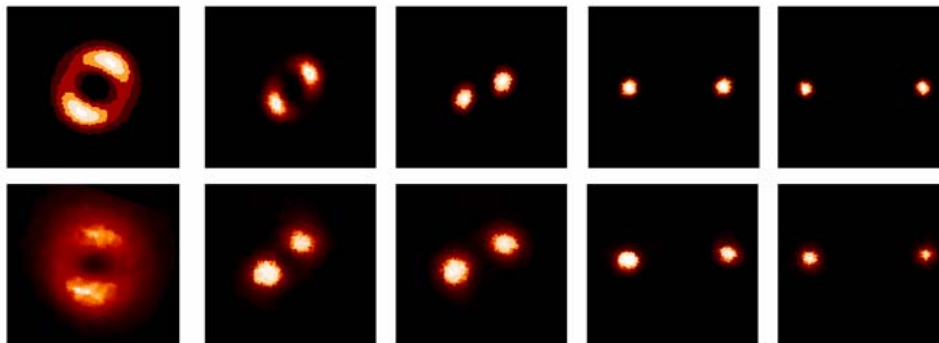


FIG. 5 (color online). Phase structure measurements. (a) Discrete vortex soliton for a lattice period of $20 \mu\text{m}$. (b)–(d) Interferograms of the vortex soliton with a weak broad coherent beam, whose relative phase is changed in steps of $\pi/2$.

Neshev et al & Fleischer et al, PRL 92 (2004)



**Without a lattice
vortex rings decay!**

Theory + experiment
Tikhonenko et al, 1995

Energy balance condition

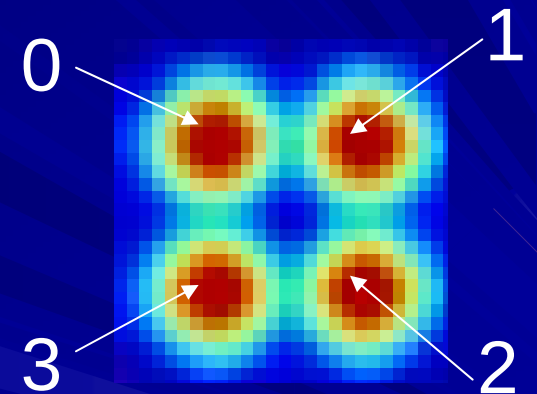
- We examine the existence properties of discrete vortices by solving the energy balance conditions for coupled modes:

N: number of lobes

Inter-lobe coupling strength

$$\sum_{m=0}^{N-1} f_{nm} \sin(\varphi_m - \varphi_n) = 0.$$

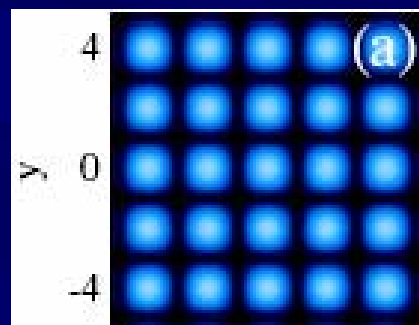
Phase of lobe



- Energy into and out of a lobe must be equal

Vortex lattice solitons

Different types of the vortex symmetries

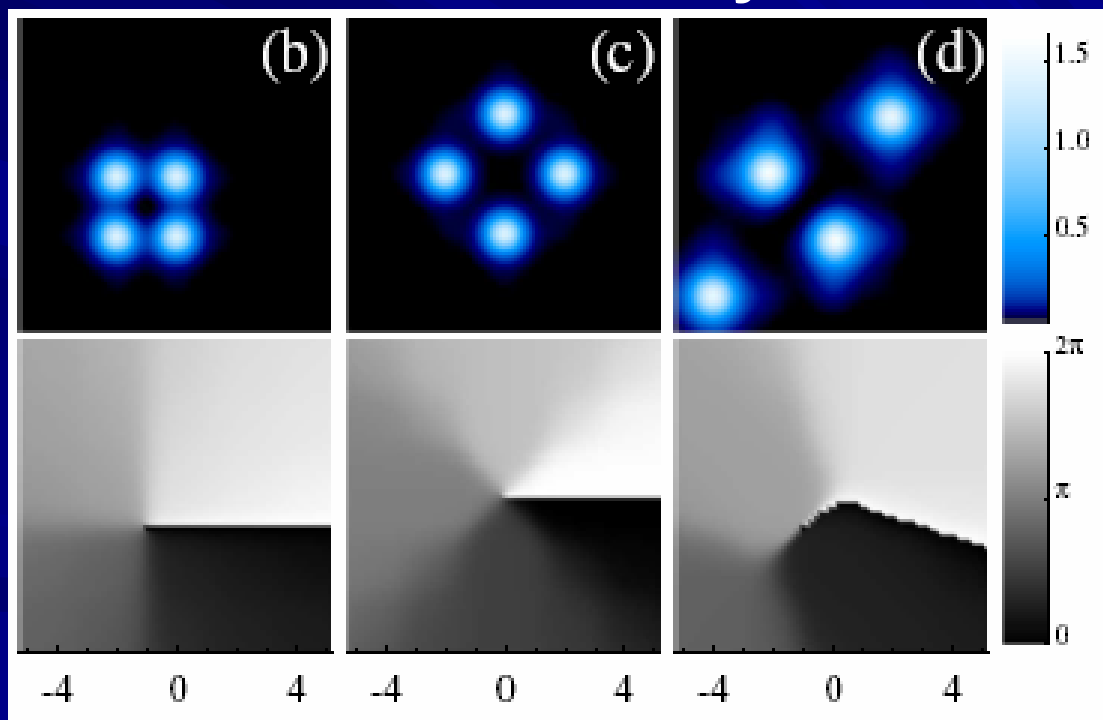


lattice

Off-site

On-site

asymmetric



Two experimental observations

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26 MARCH 2004

Observation of Discrete Vortex Solitons in Optically Induced Photonic Lattices

Dragomir N. Neshev, Tristram J. Alexander, Elena A. Ostrovskaya, and Yuri S. Kivshar

*Nonlinear Physics Group, Research School of Physical Sciences and Engineering, Australian National University,
Canberra ACT 0200, Australia*

Hector Martin, Igor Makasyuk, and Zhigang Chen

*Department of Physics and Astronomy, San Francisco State University, San Francisco, California 94132, USA
TEDA College, Nankai University, Tianjin, China*

(Received 4 September 2003; published 25 March 2004)

We report on the first experimental observation of *discrete vortex solitons* in two-dimensional optically induced photonic lattices. We demonstrate strong stabilization of an optical vortex by the lattice in a self-focusing nonlinear medium and study the generation of the discrete vortices from a broad class of singular beams.

VOLUME 92, NUMBER 12

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26 MARCH 2004

Observation of Vortex-Ring “Discrete” Solitons in 2D Photonic Lattices

Jason W. Fleischer,^{1,2} Guy Bartal,¹ Oren Cohen,¹ Ofer Manela,¹ Mordechai Segev,¹
Jared Hudock,² and Demetrios N. Christodoulides²

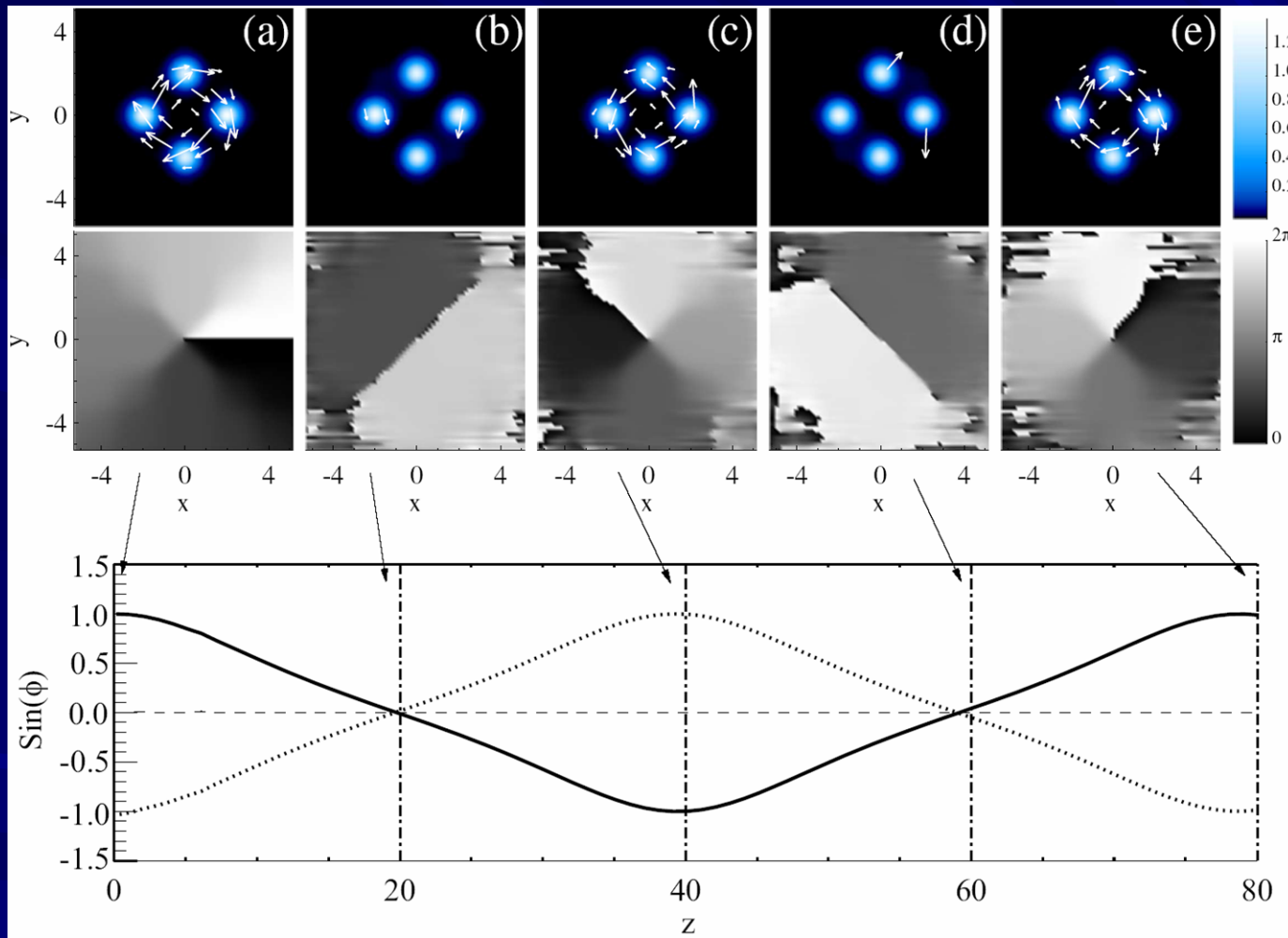
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We present the experimental observation of both on-site and off-site vortex-ring solitons of unity topological charge in a nonlinear photonic lattice, along with a theoretical study of their propagation dynamics and stability.

Charge flipping of on-site vortex



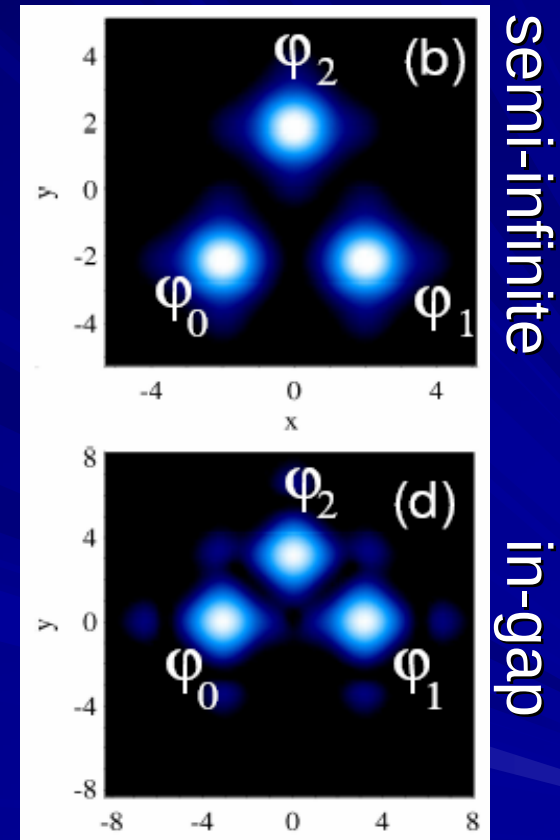
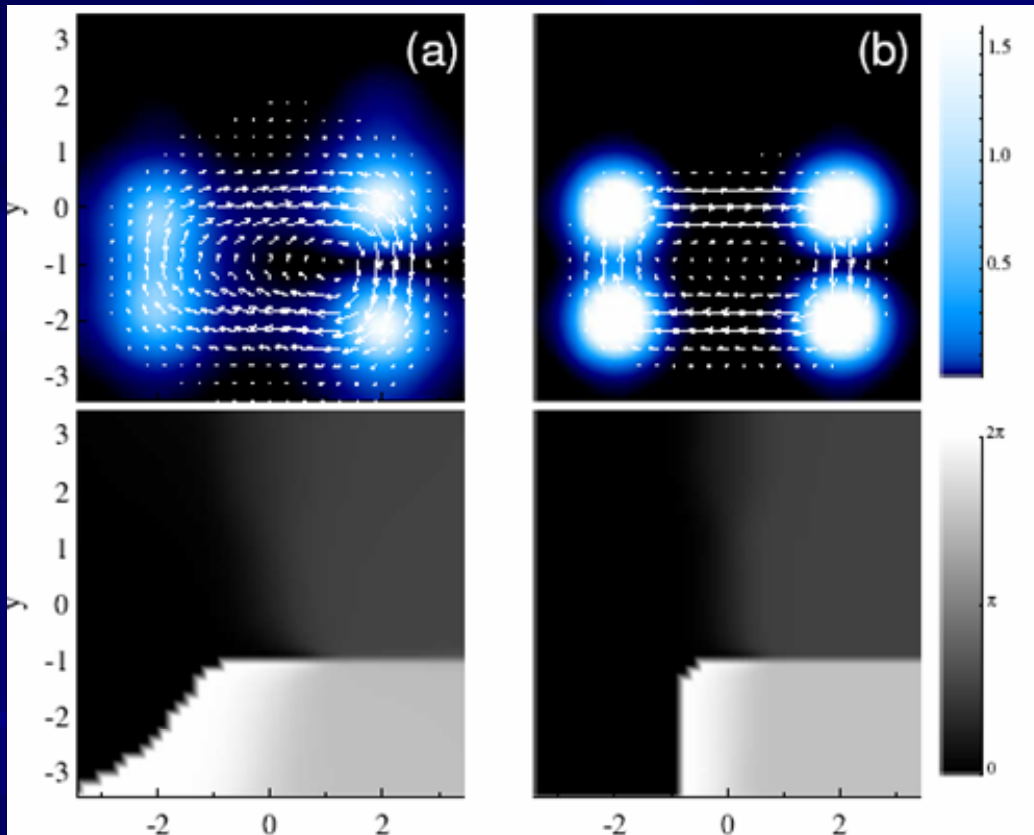
A small, but finite, perturbation leads to a charge reversal and a new effect:

periodic flipping of the vortex charge

Alexander, Sukhorukov, & Kivshar, PRL **93**, 063901 (2004)

Asymmetric vortices

Nonlinearity allows for the symmetry breaking



- ❖ Symmetric rectangular vortex does not exist
- ❖ Diagonal coupling - weaker than nearest-neighbor coupling

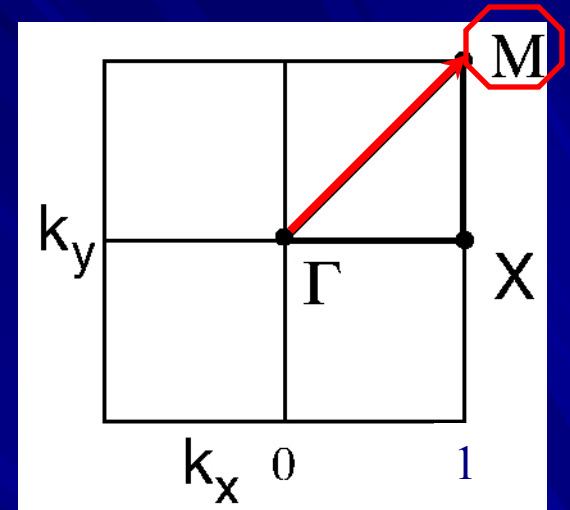
How to generate a gap vortex?

Gap-state preparation requirements:

- Negative effective diffraction regime
- Adiabatic drive to the Brillouin zone's edge
- Phase ramp imprint

$$D_{x,y} < 0$$

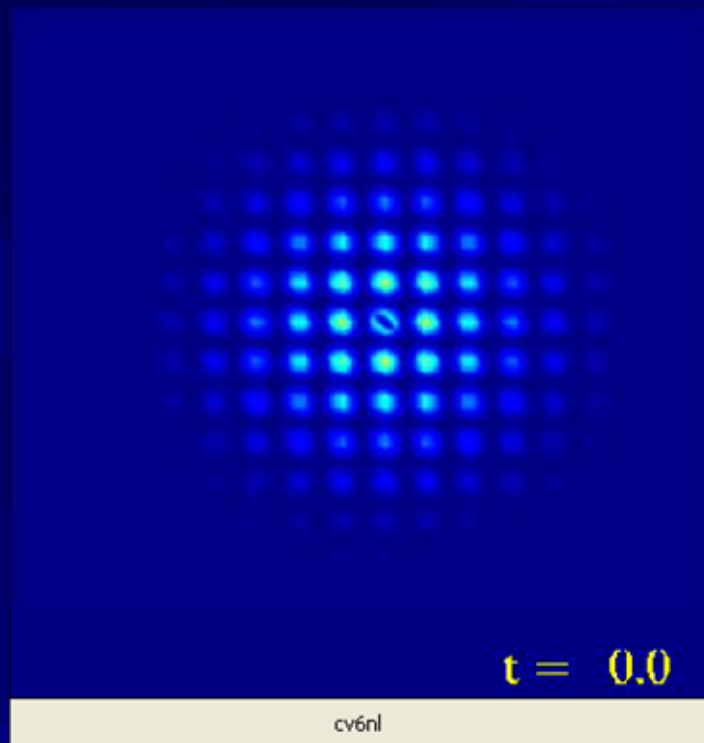
$$t \gg 1/\Delta\mu_{gap}$$



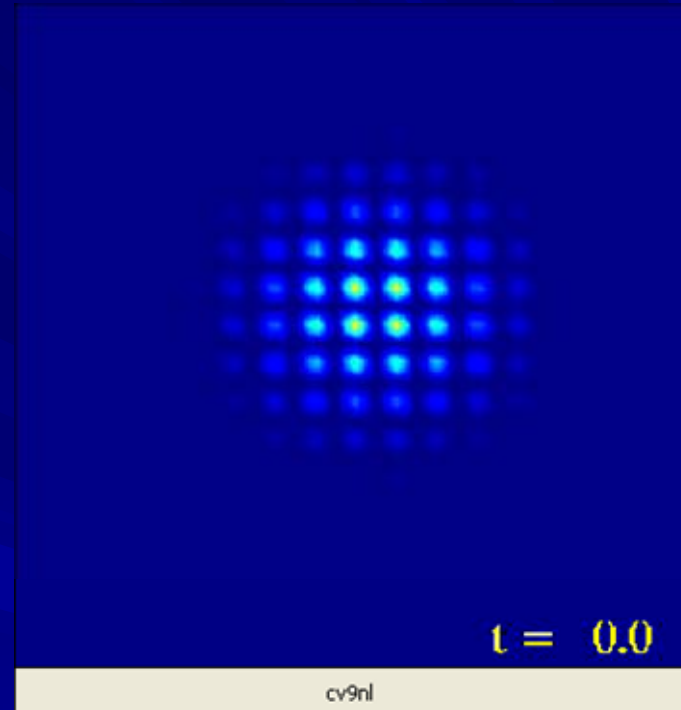
Simulation procedure:

- Start with a broad BEC wavepacket **at the correct band edge** (with a nontrivial Bloch-wave phase)
- Imprint a charge one vortex phase
- Let evolve in time

Generation dynamics



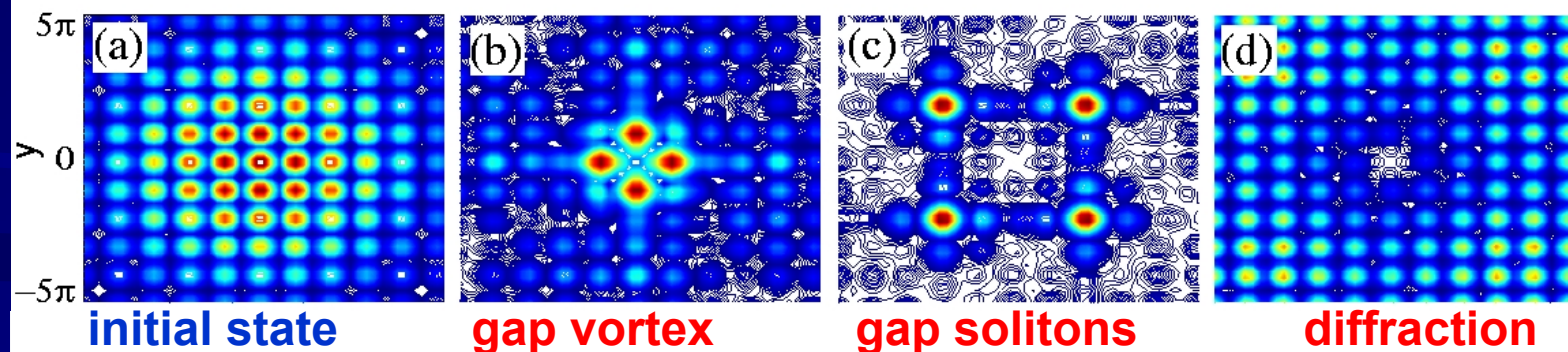
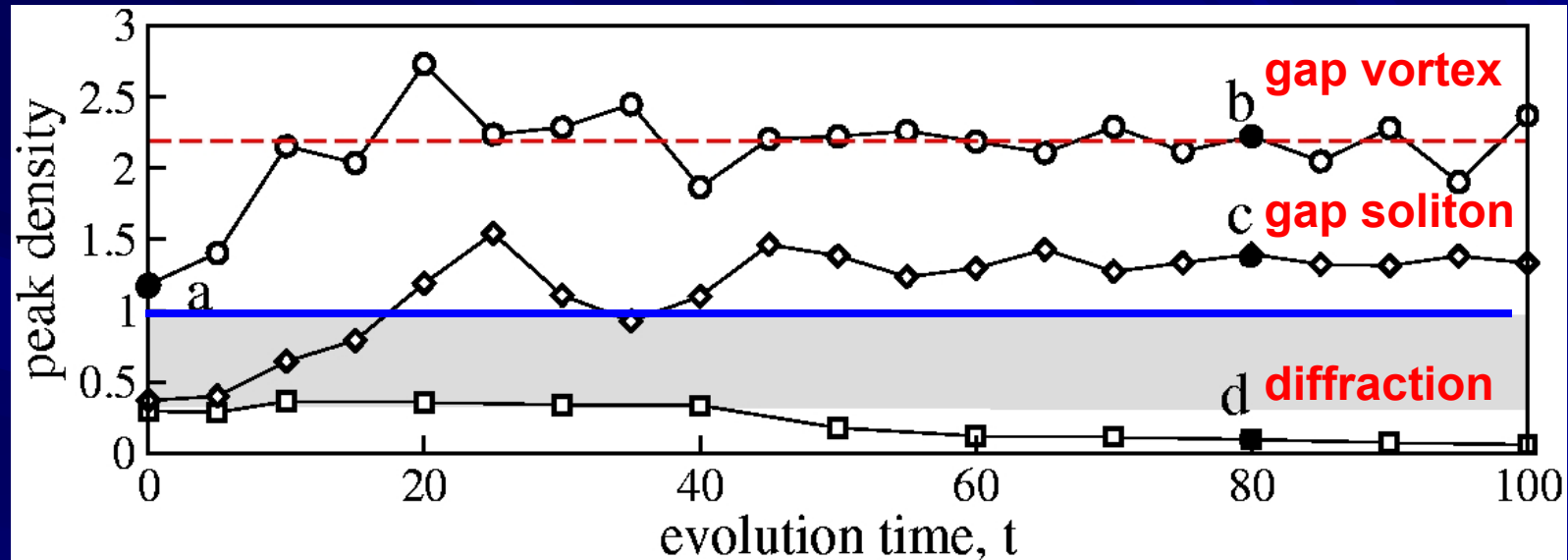
Discrete diffraction



Gap vortex

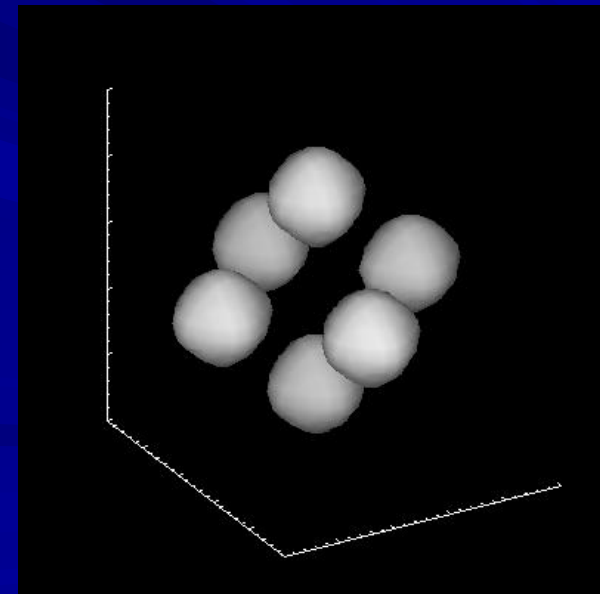
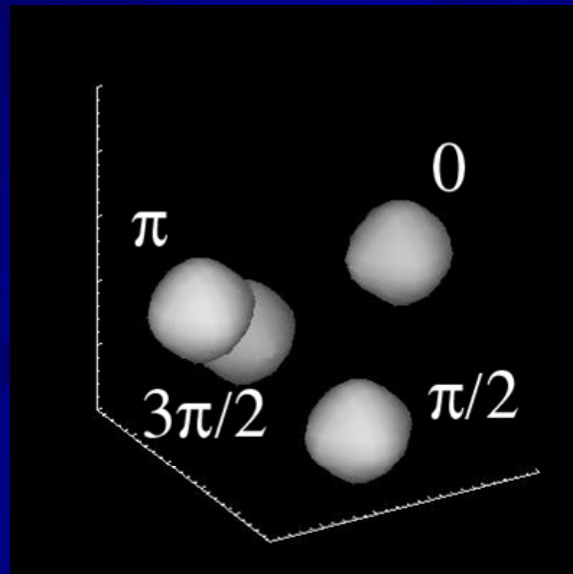
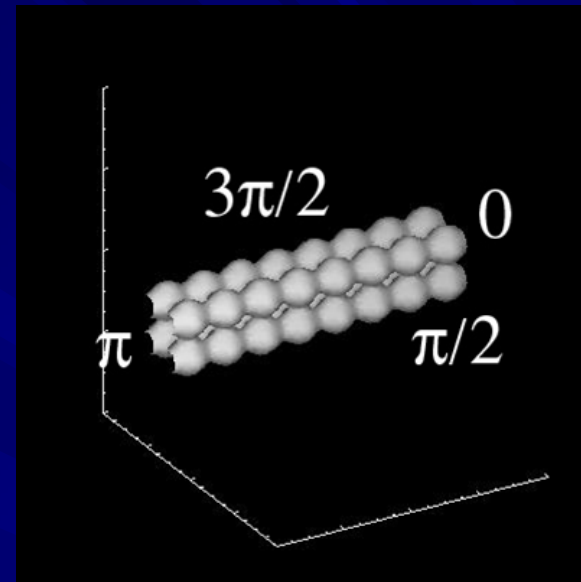
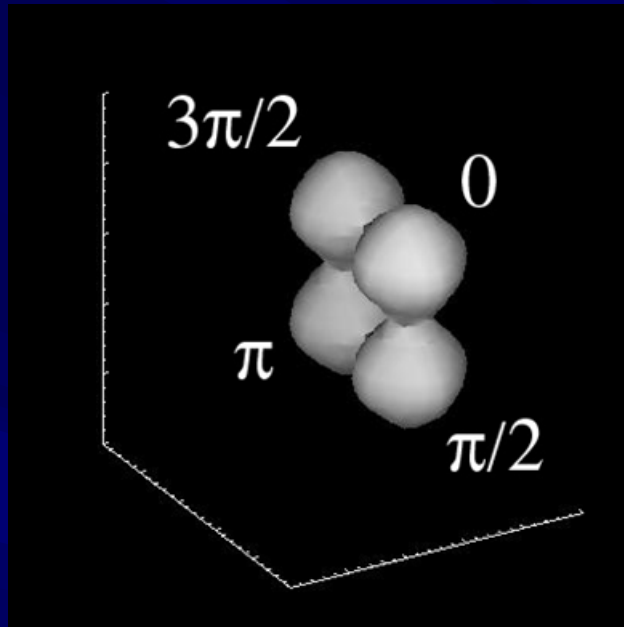
Possible outcomes

- Necessary: **atom number** above vortex state threshold
- Sufficient: **peak density** above threshold

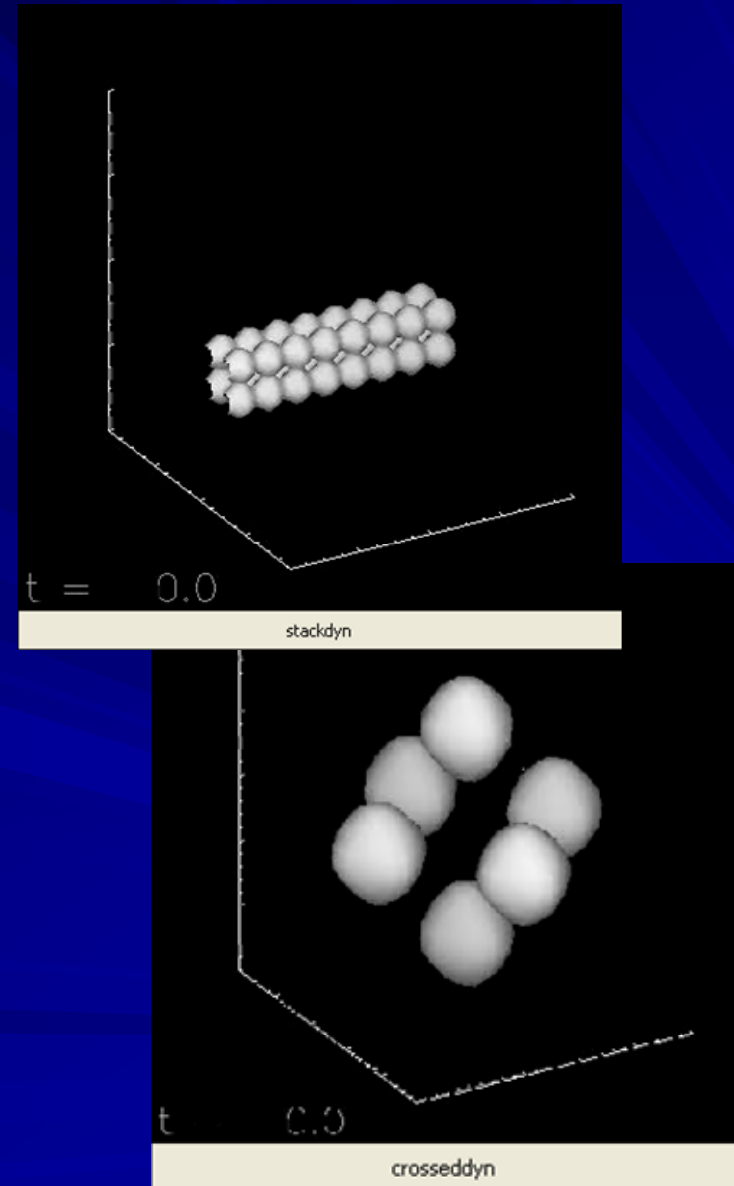
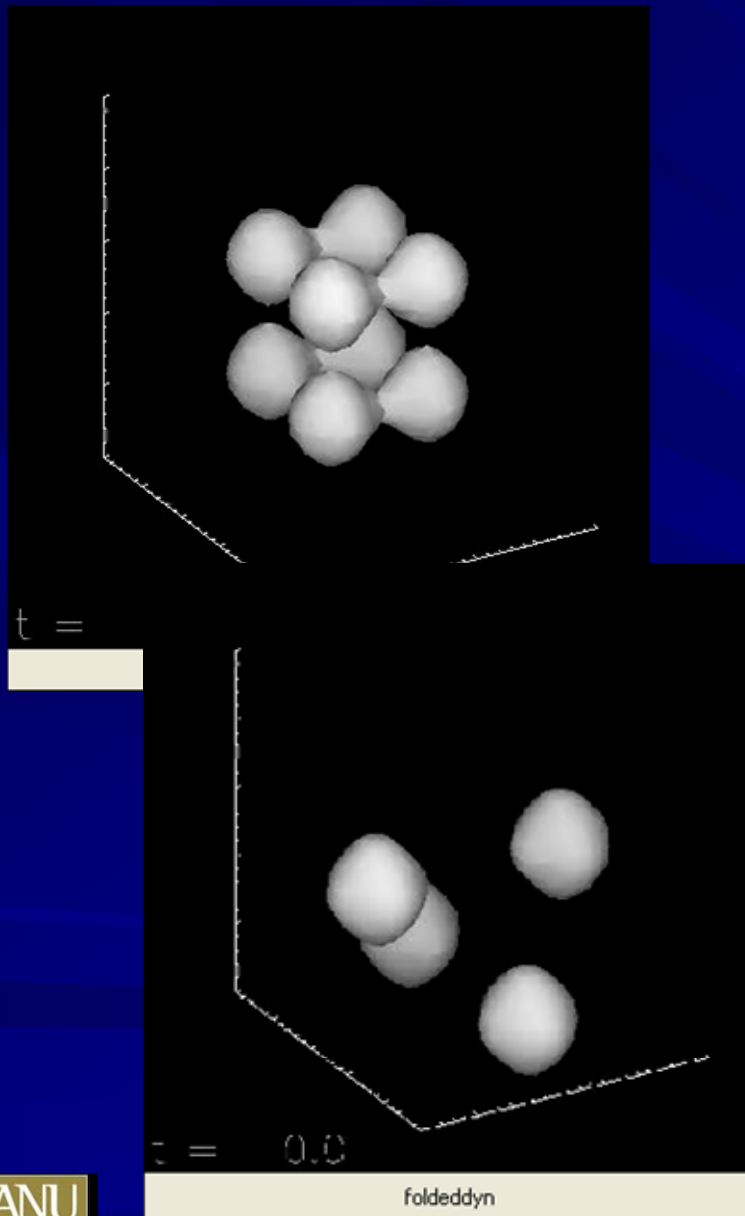


3D discrete vortices

3D discrete vortices



Dynamics of 3D discrete vortices



Conclusions

- **Nonlinear optics provides a solid background for the physics of nonlinear matter waves and solitons**
- **Much to learn from the field of spatial optical solitons**
- **Novel features: a strong trap; 2D vs. 3D; strong repulsion**

Last 2-3 years show a strong influence of both the fields on each other

Novel directions for matter waves in lattices

Attractive condensate in a lattice (1D vs 2D & 3D)

- > In-gap nonlinear effects (e.g. gap vortex)
- > Atomic bandgap engineering - BEC in superlattices
- > Spinor, multi-component, and atom-molecular BEC

