

MESOSCOPIC PHYSICS WITH LIGHT IN COLD ATOMS

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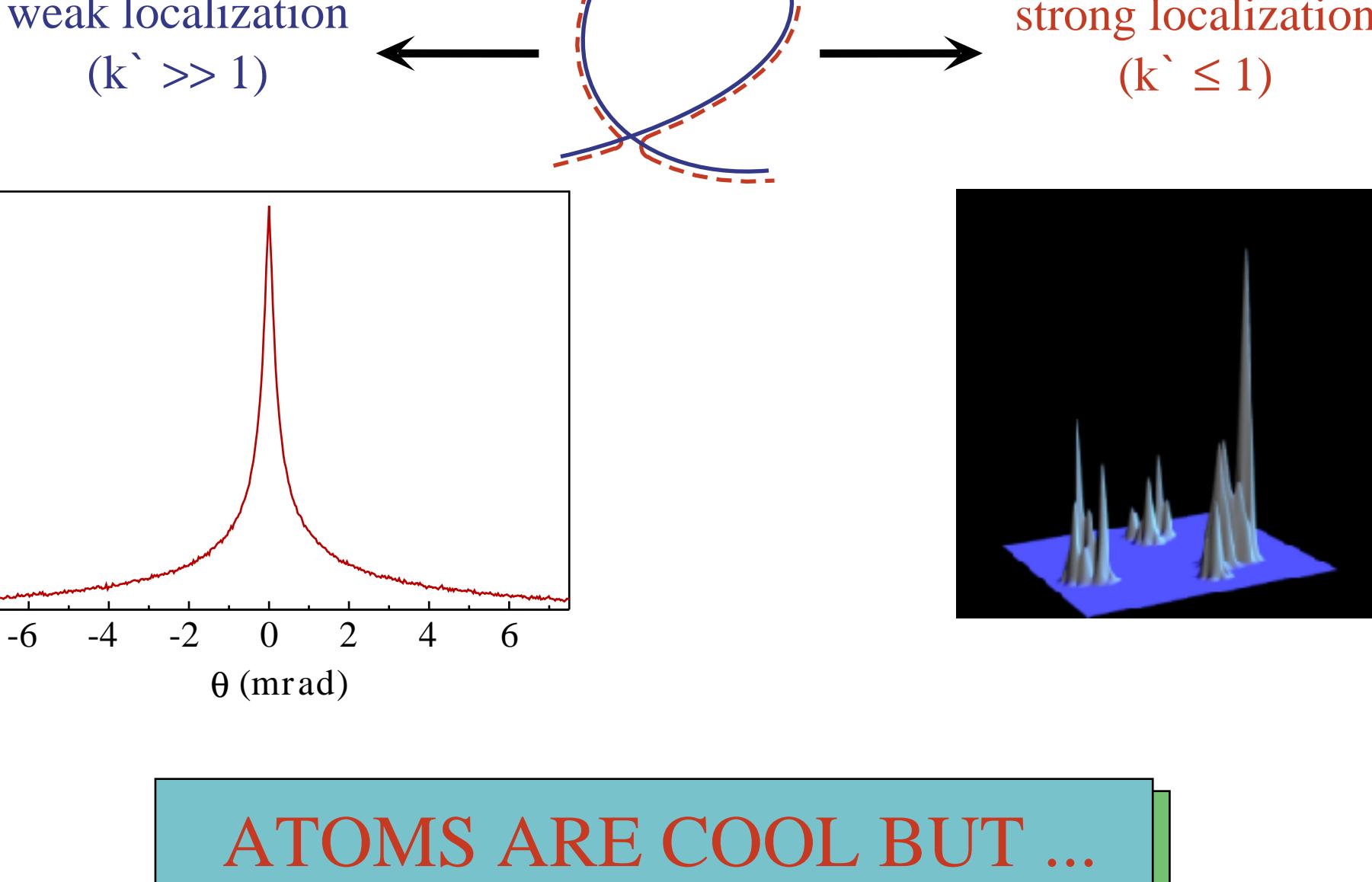
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MESOSCOPIC PHYSICS

INTERFERENCE CORRECTIONS TO TRANSPORT OF WAVES IN DISORDERED MEDIA



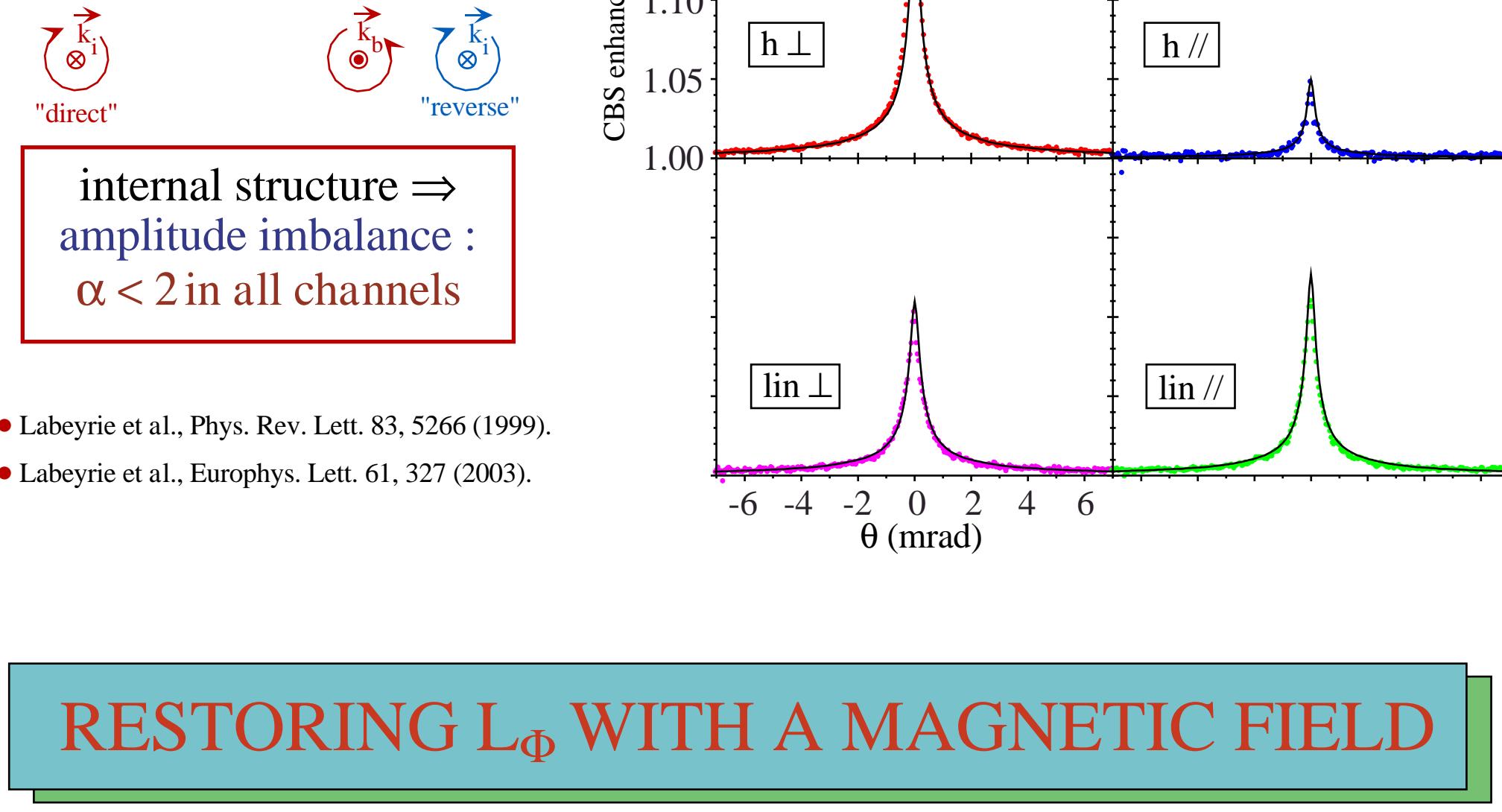
ATOMS ARE COOL BUT ...

SEVERAL MECHANISMS LIMIT L_ϕ

internal structure of ground state
atomic motion

saturation of atomic transition (inelastic scattering)

INTERNAL STRUCTURE OF GROUND STATE



• Labeyrie et al., Phys. Rev. Lett. 83, 5266 (1999).

• Labeyrie et al., Europhys. Lett. 61, 327 (2003).

RESTORING L_ϕ WITH A MAGNETIC FIELD

"usual" situations :

B breaks time-reversal symmetry
 $\Rightarrow L_\phi$ decreases

our case :

L_ϕ increases $\propto B$!!!

• Sigwarth et al., Phys. Rev. Lett. 93, 143906 (2004).

DYNAMICAL BREAKDOWN : PRELIMINARY CONCLUSIONS

• sharp decrease of L_ϕ when v_{rms} increases : $L_\phi / \gamma \approx (kv_{rms})^{-2/3}$

\Rightarrow mesoscopic criterion for optical thickness b :

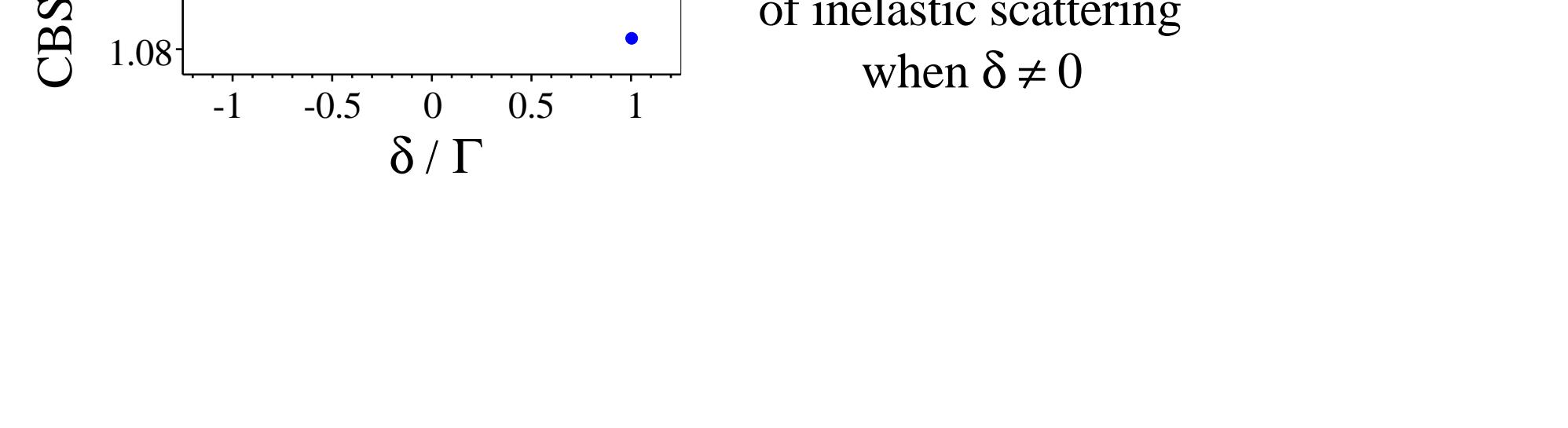
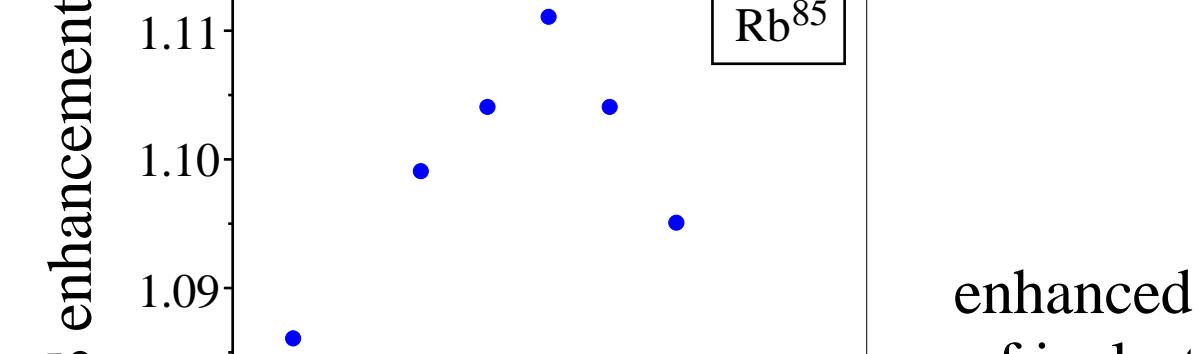
$$\frac{kv_{rms}}{\Gamma} < 1/b^2$$

difficult to fulfill for $b \gg 1$

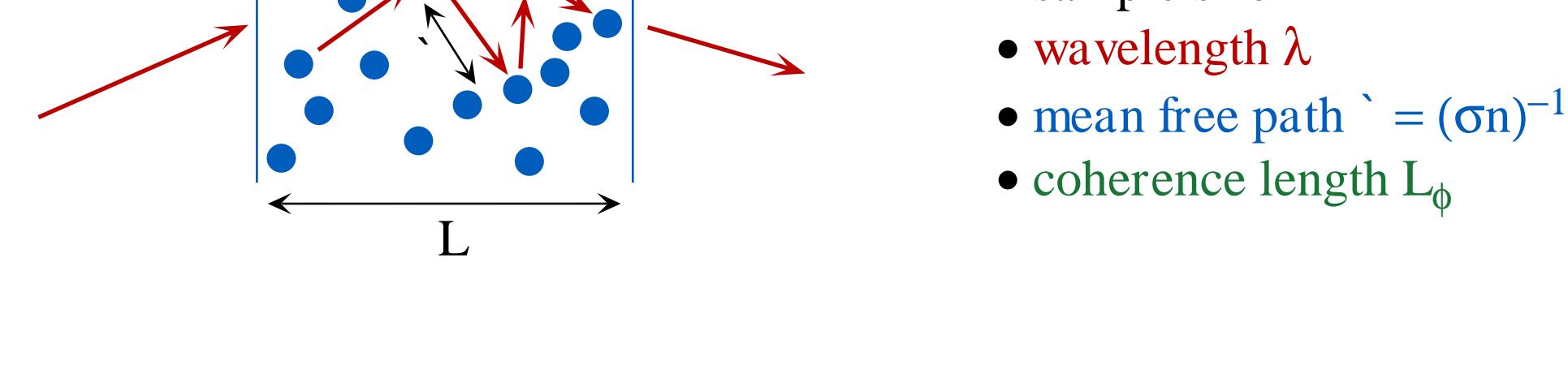
• partial restoration of CBS when δ increases



ROLE OF SCATTERED LIGHT SPECTRUM



COHERENCE LENGTH



"mesoscopic criterion" :

$L < L_\phi$

multiple scattering coherent transport

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WHY COLD ATOMS ?

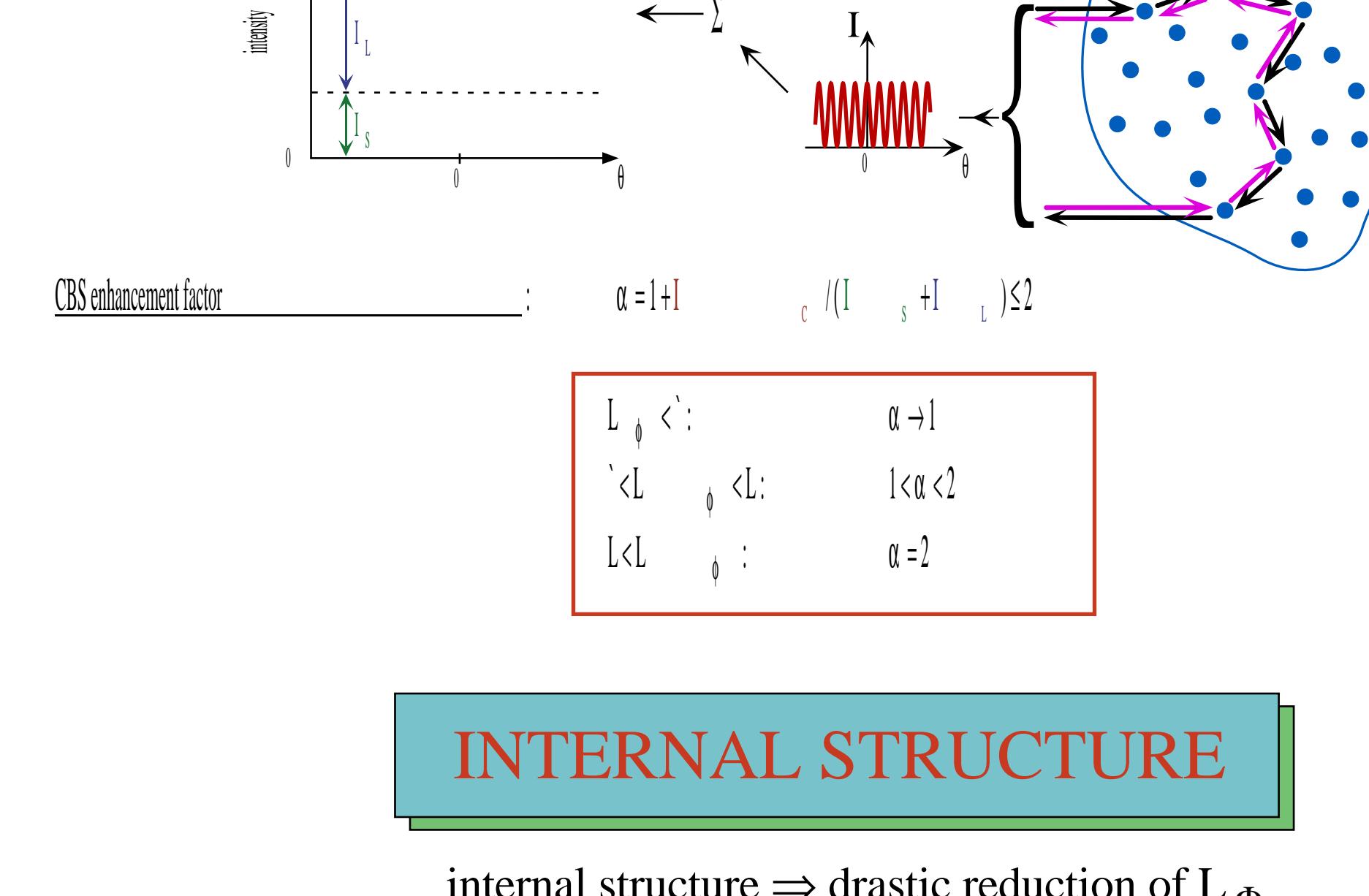
• "clean" & controllable system :

=> well-known and well-controlled scatterers
=> no absorption
=> tunable δ)

• new physics & new regimes :

=> highly resonant and monodisperse samples
=> internal structure
=> feedback :
mechanical effects
optical pumping
entanglement
=> non-linearities

COHERENT BACKSCATTERING : A TOOL TO PROBE L_ϕ



setup :

collimation and spatial filtering

beam dump

lens (f=30mm)

MOT

chopper wheel

angular resolution : $\Delta\theta_{res} < 0.1$ mrad

CBS SETUP

timing :

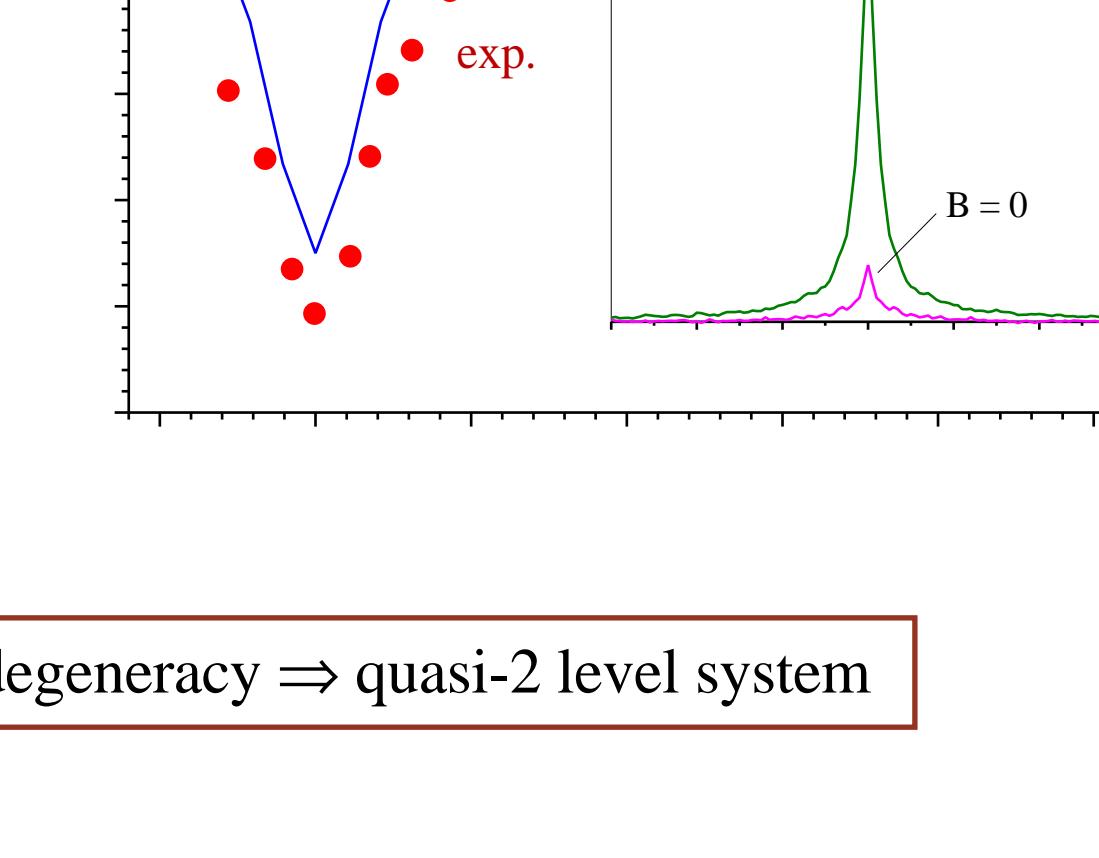
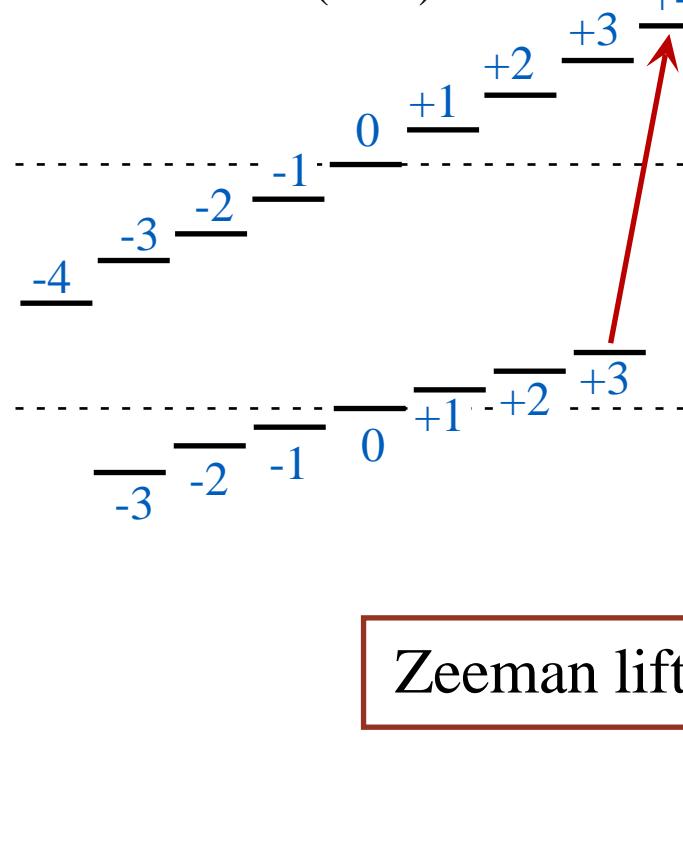
detection

20 ms

probe

time

RESTORING L_ϕ WITH A MAGNETIC FIELD



Zeeman lifting of degeneracy \Rightarrow quasi-2 level system

INTERNAL STRUCTURE

internal structure \Rightarrow drastic reduction of L_ϕ

J = 3 \rightarrow J' = 4

• C. A. Müller et al., Phys. Rev. A 64, 053804 (2001).

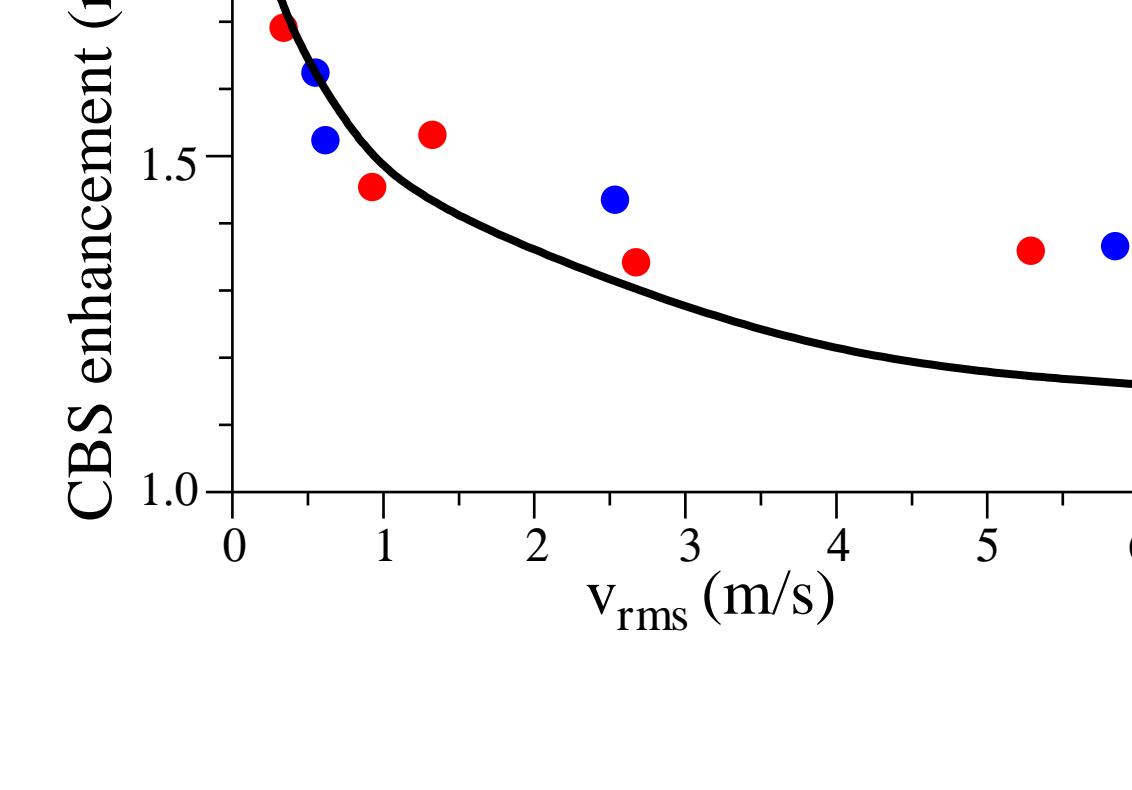
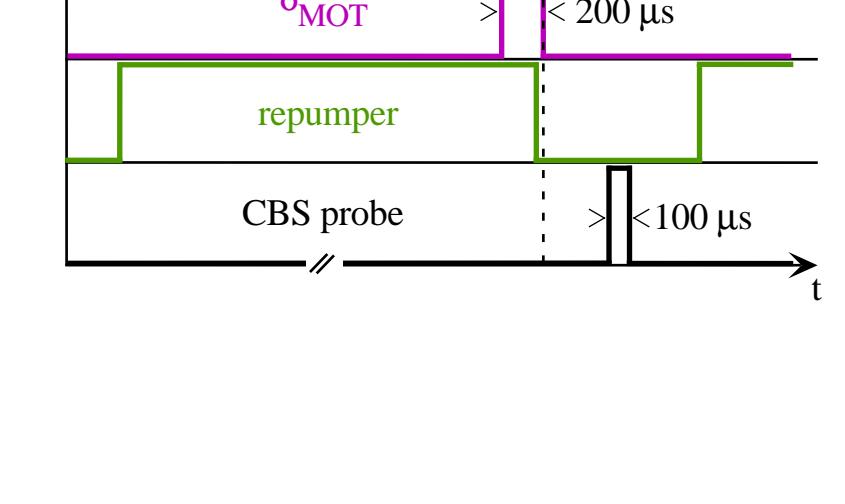
J = 3 \rightarrow J' = 4 $\Rightarrow L_\phi \approx \gamma / 3$

ATOMIC MOTION

motion of scatterers \Rightarrow random phase

• A. A. Golubentsev, Sov. JETP 59, 26 (1984).

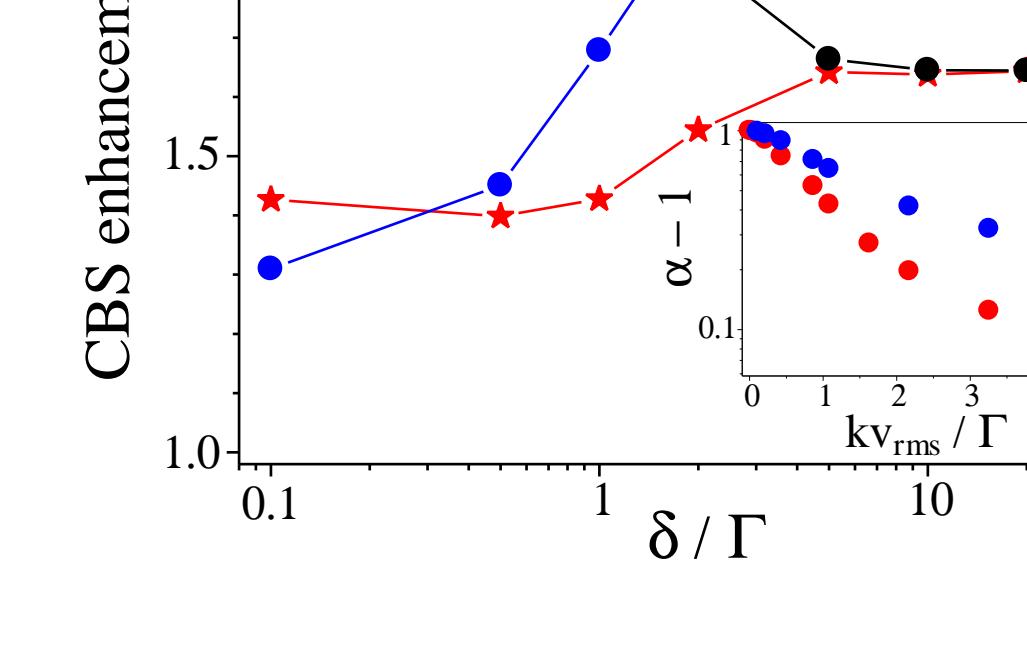
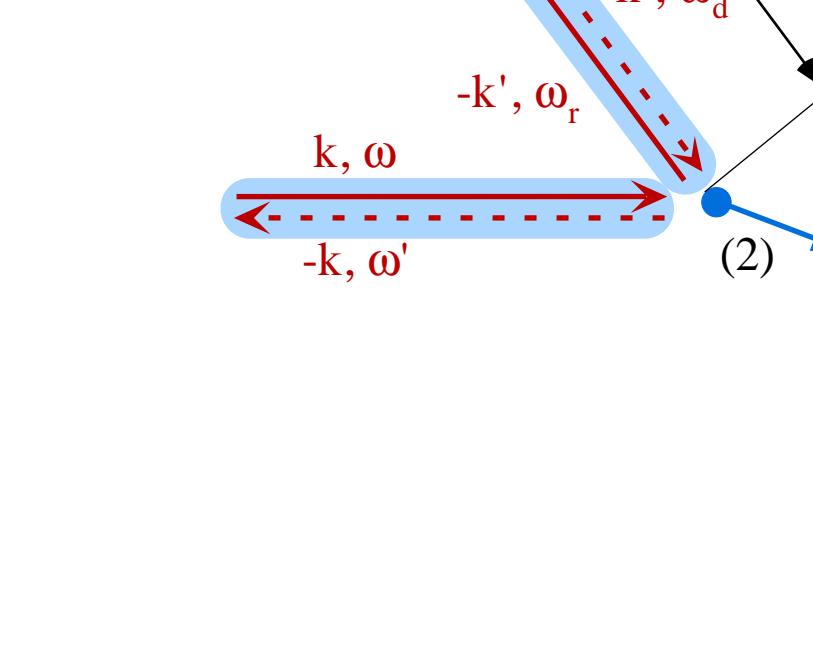
experimental observation of "dynamical breakdown" : heating by intense near-resonant optical molasse



DYNAMICAL BREAKDOWN OF CBS

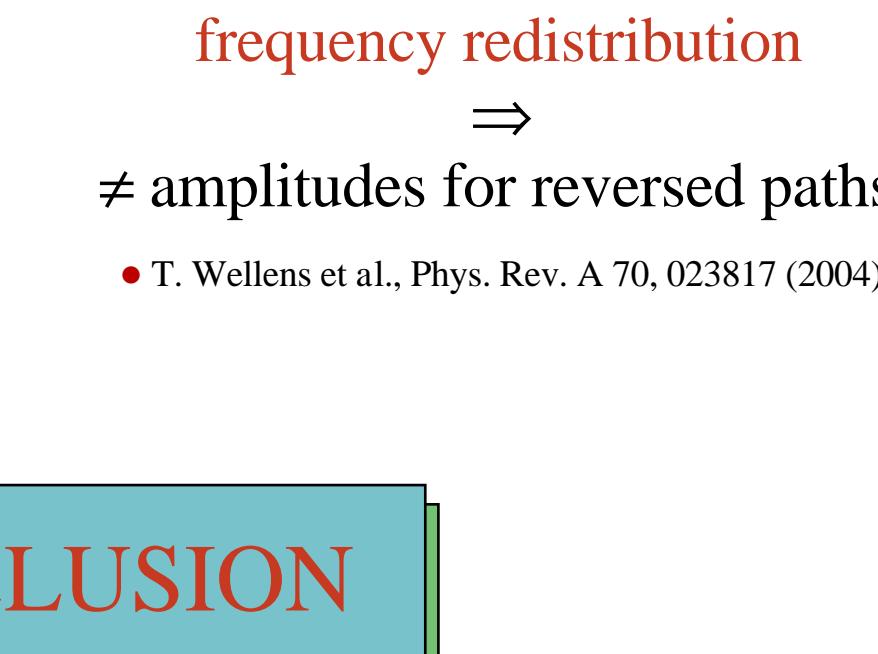
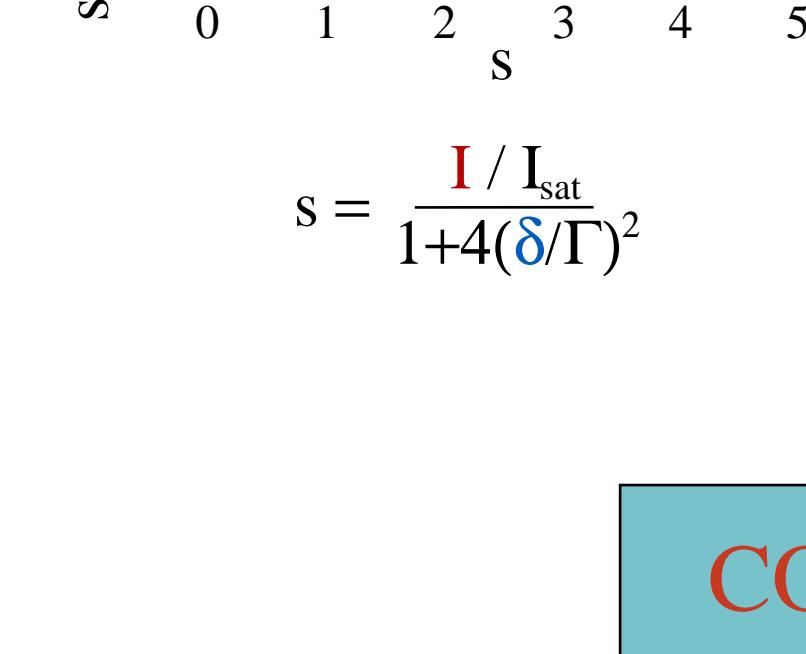
scattering phase : $\langle \Delta\phi_s \rangle = 4 \frac{kv_{rms}}{\Gamma} \frac{1}{1+4(\delta/\Gamma)^2}$

propagation phase : $\langle \Delta\phi_p \rangle = -2 \frac{kv_{rms}}{\Gamma} \frac{1-4(\delta/\Gamma)^2}{1+4(\delta/\Gamma)^2}$



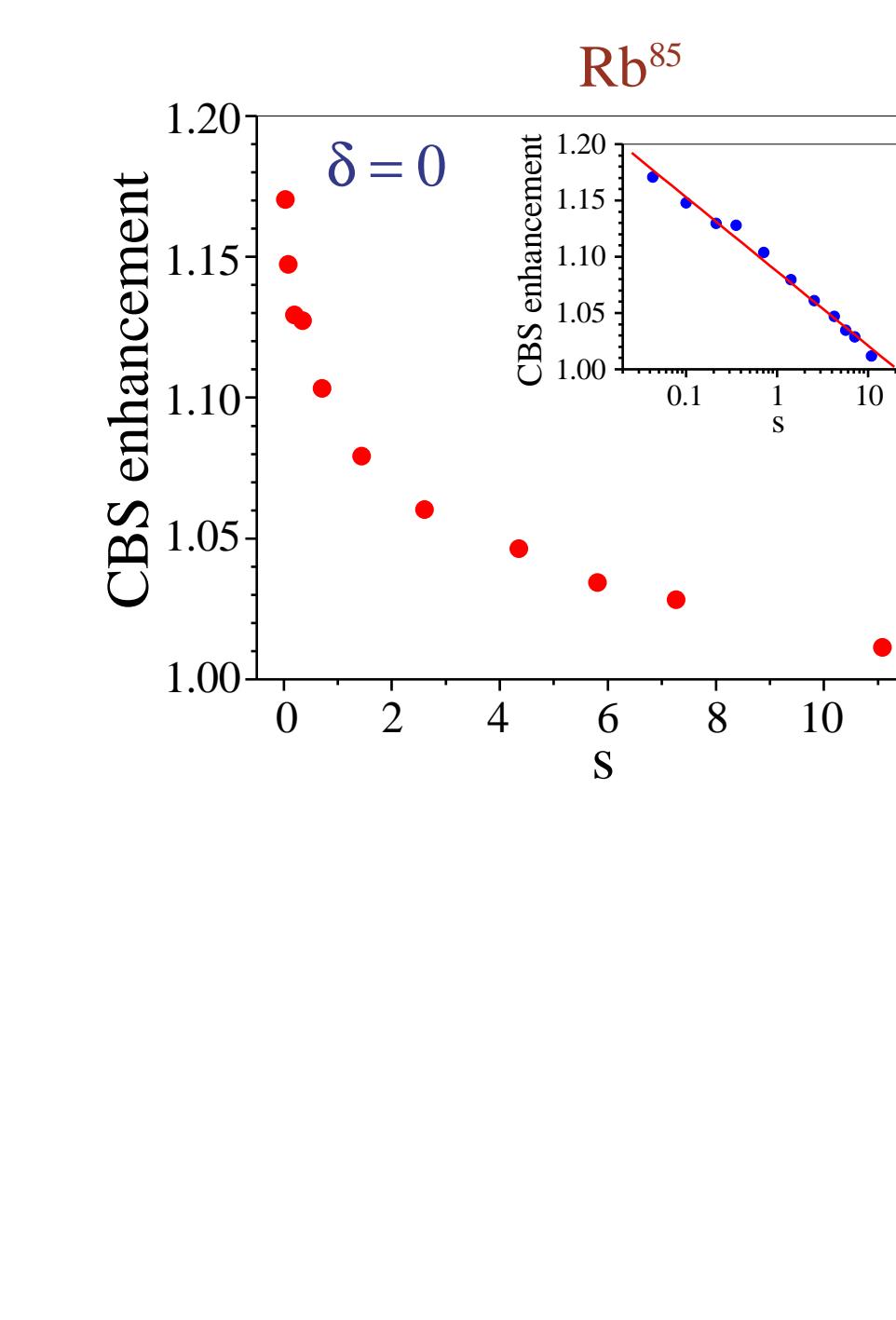
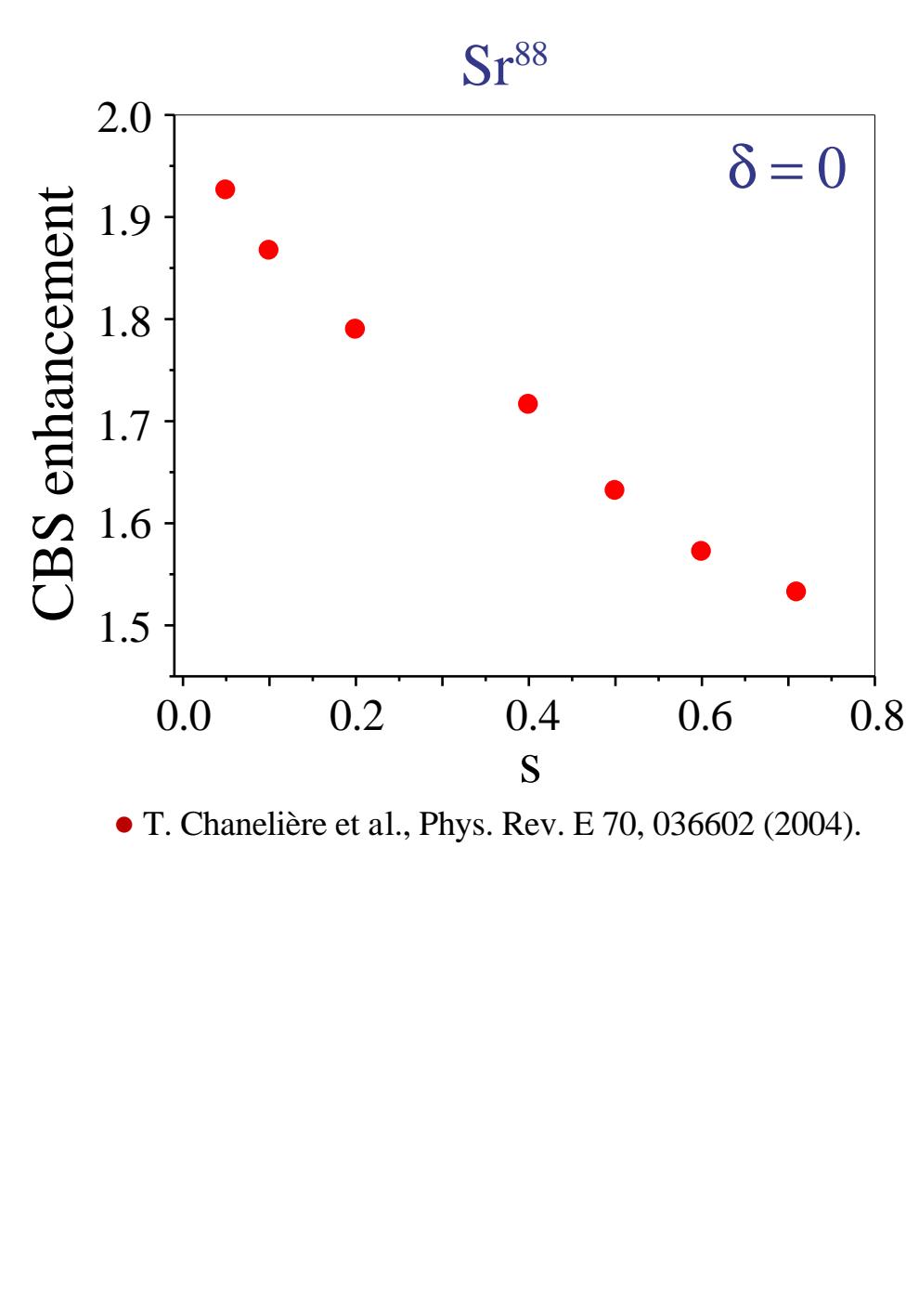
SATURATION OF THE ATOMIC TRANSITION

saturation \Rightarrow inelastic scattering



\neq amplitudes for reversed paths

• T. Wellens et al., Phys. Rev. A 70, 023817 (2004).



CONCLUSION

• identification and demonstration of several coherence-reducing mechanisms :

\Rightarrow internal structure

\Rightarrow temperature

\Rightarrow saturation

\Rightarrow internal structure :

apply B field (or use $J = 0 \rightarrow J' = 1$)

\Rightarrow temperature :

cool down ... but severe limitation for long scattering paths !

\Rightarrow saturation :

use $s \ll 1$

possible problems in random laser

and strong localization

