

Micromanipulation of ultra-cold atoms close to conducting surfaces

S. Wildermuth¹, P. Krüger¹, S. Hofferberth¹, L.M. Andersson¹, S. Groth^{1,2},
E. Haller¹, D. Gallego¹, I. Bar-Joseph², J. Schmiedmayer¹

¹Physikalisches Institut, Universität Heidelberg, Philosophenweg 12, 69120 Heidelberg, Germany

²Department of Condensed Matter Physics, The Weizmann Institute of Science, Rehovot, Israel

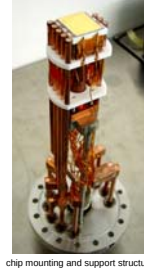
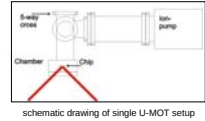
homepage: www.atomchip.org



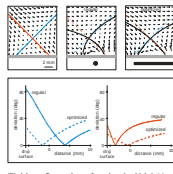
⁸⁷Rb-setup

Apparatus:

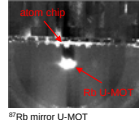
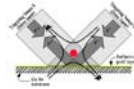
- single chamber vacuum system ($p < 10^{-11}$ mbar), operated with pulsed Rb-dispensers
- optimized U-MOT, collects $> 3 \cdot 10^8$ atoms in 20s of loading
- magnetic trap with $\sim 10^8$ atoms
- BEC in Cu-structure trap: $\sim 10^5$ atoms
- precooled thermal atoms ($\sim 10 \mu\text{K}$) can be transferred to various traps and guides on the chip
- BEC in different chip wire traps
- trapping and cooling in both $|F=2, m_F=2\rangle$ and $|F=1, m_F=-1\rangle$ states



Main loading technique: Optimized mirror U-MOT

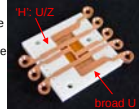


The quadrupole field originates from a specially designed U-shaped wire underneath the chip in combination with a homogenous bias field. The shape of the U is optimized to create a large volume quadrupole field (no bulky quadrupole coils necessary).

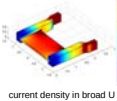


Multi-layer copper structure underneath the chip to be able to create a sizeable BEC:

- broad U-wire for large-scale MOT
- different Z-wires for evaporative cooling

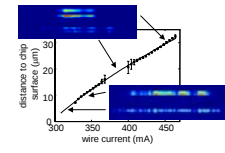


Optimized U-MOT for experiments with ultracold atoms near surfaces, S. Wildermuth, P. Krüger, C. Becker, M. Brajic, S. Haupt, A. Kasper, R. Folman, J. Schmiedmayer, PRA 69, 030901(R) (2004)

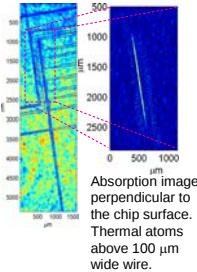
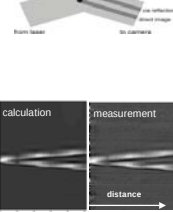


Imaging near surfaces

Horizontal absorption imaging with slightly tilted laser beam allows to measure distance to surface.



Two images of BEC overlap for $d < 7 \mu\text{m}$. Distance can be extrapolated taking the finite size of the wire into account.

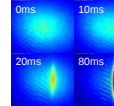


Condensate propagation in a 1d guide

trap geometry



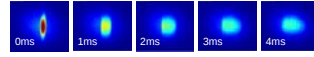
BEC formation



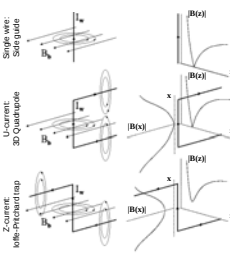
- thermal atoms just above T_c in a 1d trap
- at 0ms electric dimple is suddenly created
- 1d BEC forms slowly
- BEC becomes larger and 3d
- equilibrium reached after $\sim 80\text{ms}$

BEC propagating along 1d guide

- BEC is formed in 3d dimple
- off-center release to elongating (1d) trap by sudden ($< 1\mu\text{s}$) switching off of dimple voltage
- background gas temperature $\sim T_c$
- movement of edge at $\sim 5\text{cm/s}$
- thermalization into equilibrium state (centered cloud at T_c) on slower time scale ($\sim 100\text{ms}$)



Main building block: magnetic wire traps



$V(r) = -\mu_B \mathbf{B}(r) \cdot \mathbf{m}_F$; in adiabatic approximation: m_F is constant of motion; only important parameter is $|\mathbf{B}(r)|$

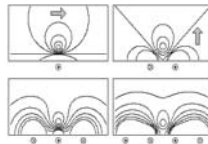
Basic building block: wire traps, derived from combining the magnetic field from current carrying wire(-structures) with homogeneous bias fields.

Single wire + bias field: 2D quadrupole along the wire: side guide
U-shaped wire + bias field: 2D quadrupole trap
Z-shaped wire + bias field: 3D Ioffe-Pritchard trap

The bias fields need not be applied from the outside; extra wires can suffice. This allows even more freedom when designing a device.

Magnetic traps on atom chips can be combined with:

- electric fields
- dipole traps
- other quantum-optics techniques like cavities, etc.



Magnetic surface disorder potentials

Thermal clouds and BECs can be brought close ($d \sim 3 \mu\text{m}$) to the wire surface to probe surface disorder potentials.

Thermal clouds at $T = 1 \mu\text{K}$ do not fragment down to distances of $3 \mu\text{m}$.

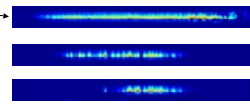
Boltzmann distribution

$$n \sim \exp(V / k_B T)$$

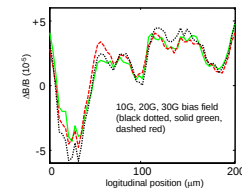
gives upper limit on potential roughness of $\Delta B/B < 2 \times 10^{-4}$.

Disorder potential can be reconstructed from density profiles of 1d BECs:

$$V(x) = -2\hbar\omega_{\perp} a_{\text{scat}} n_{1d}(x)$$

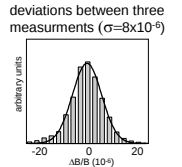


scaling of disorder potential with current in the wire



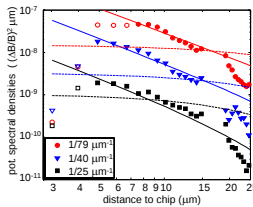
Measurement at same position and constant distance ($10 \mu\text{m}$) to the wire but three different wire currents / bias fields.

- we find a linear current scaling
- disorder potential stemming from irregular current flow in wire
- excludes electric patch effects at a scale of 10^{-13}eV ($d \sim 5 \mu\text{m}$)

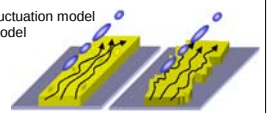


variation of disorder potential with distance to the wire

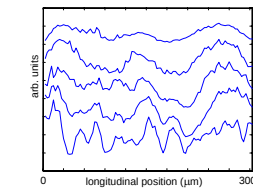
potential spectral densities ($100 \mu\text{m}$ wide wire)



dashed curves: wire edge fluctuation model
solid curves: local current model



In our case, disorder potential is dominated by local deviations from a straight current flow. Sensitivity $\Delta B/B = 10^{-5}$ corresponds to a resolution of 10^{-5}rad .

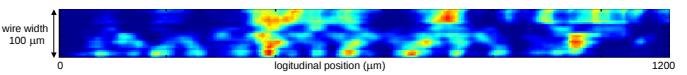


Disorder potential reconstructed from density profiles of 1d BECs. Curves are shifted by equal amount from each other. Heights above the chip are indicated on the right.

- at $d = 10 \mu\text{m}$ we find rms $\Delta B/B = 3 \times 10^{-5}$ ($k > 1/200 \mu\text{m}^{-1}$)
- fluctuation spectrum scales as $\sim 1/k^2$

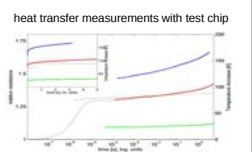
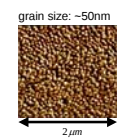
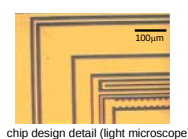
variation of disorder potential with transverse position

2d potential landscape at constant height of $10 \mu\text{m}$ above a $100 \mu\text{m}$ wide wire.



Chip fabrication

- semiconductor (Si, GaAs) and sapphire substrates
- lithographically patterned gold layer ($5 \mu\text{m}$ thickness, resolution $< 1 \mu\text{m}$)
- current densities up to 10^7A/cm^2 possible
- voltages of more than 300V ($10 \mu\text{m}$ gaps)



Atom Chips: Fabrication and Thermal Properties, S. Groth, P. Krüger, S. Wildermuth, R. Folman, T. Fernholz, D. Mahalu, I. Bar-Joseph, J. Schmiedmayer, APL in press, cond-mat/0404141