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Strain and disorder effects on inhomogeneous states of manganites

Carmine Antonio Perroni

Coherentia-INFM and Universita' di Napoli "Federico II" (Italy)

currently at IFF, Forschungszentrum Juelich (Germany)

Coworkers:

V. Cataudella V. Marigliano Ramaglia G. De Filippis

Coherentia-INFM and Universita' di Napoli "Federico II" (Italy)



- Introduction to manganite properties
- Inhomogeneities and role of disorder and strain
- Model for manganite films: effects of thickness, strain and el-ph coupling
- Role of disorder at low temperatures
- Phase separation in films of $La_{1-x}Ca_xMnO_3$ and $La_{1-x}Sr_xMnO_3$
- Preliminar analysis about STM experimental data on LSMO films: features of multiphase modulation?

Manganites: $Re_{1-x}A_xMnO_3$ **Re-Rare earth** (La or Nd trivalent) A- Alkali (Sr or Ca divalent) Mixed valence compounds: Mn^{4+} instead of Mn^{3+} depending on hole doping x Compounds known for their colossal magnetoresistance properties CMR features in correspondence with metal-insulator ferro-para transition Perovskite structure Conductive -La, Sr channels Mn^{3+} 3d orbitals 10Dq $J_{l} > 1 \text{ eV}$ -Mn Core states t_{2g} (a)

Half-metallic systems with strong Hund's rule coupling and double-exchange mechanism

Classification of manganites



III) Small bandwidth (large el-ph coupling and disorder strength)

Inhomogeneities in manganites

As the Curie temperature is approached from below, diffuse scattering from lattice polarons

STM data in LCMO film for x=0.3 just below the critical temperature



Mechanisms of phase separation

Role of disorder (mainly electron mechanism)

Numerical studies through MonteCarlo simulations on double-exchange models characterized by first order transitions. Quenched disorder is important to smear the transition between competing states and induce clustered states crucial for CMR -> Prediction of characteristic temperature T^* (E. Dagotto "Nanoscale phase separation", Springer Verlag, 2003)

Role of (intrinsic) strain (mainly lattice mechanism)

Strong coupling between electronic and elastic degrees of freedom gives longrange strains. Natural mechanism for self-organized inhomogeneities over nanoand micrometre scale (K.H. Ahn et al., Nature 2004). **Prediction of Multiscale phase modulations** at high T with nanometer-scale polaron correlations(Cheong et al.)

Hamiltonian for manganites

Single orbital model

$$\begin{split} H &= -t \sum_{\langle i,j \rangle} \left(\frac{S_0^{i,j} + 1/2}{2S + 1} \right) c_i^{\dagger} c_j - \mu \sum_i c_i^{\dagger} c_i \\ &+ \omega_0 \sum_i a_i^{\dagger} a_i + g \omega_0 \sum_i c_i^{\dagger} c_i (a_i + a_i^{\dagger}) \\ &+ \epsilon \sum_{\langle i,j \rangle}^i \vec{S}_i \cdot \vec{S}_j - g_s \mu_B \sum_i \vec{h}_{ext} \cdot \vec{S}_i \end{split}$$

- Double-exchange term for strong Hund's rule coupling, with $(\vec{S}_0^{i,j})^2 = (\vec{S}_i + \vec{S}_j + \vec{s})^2$.
- Electron-phonon term.
- Super-exchange term.
- External magnetic field h_{ext} .

Mean-field phase diagram (bulk)



Films: confinement along the growth axis C.A. Perroni et al., PRB **68**, 224424 (2003)

•Qualitative agreement with experimental phase diagram for 0 < x < 0.5

•Evidence of phase separation



Manganite Films

Size effects



The M-I transition as a function of x for different bandwidth and couplings: the transition temperature decreases with film size



Strain effects due to mismatch with substrate



Role of size, strain and el-ph effects



 $2x10^{4}$

0.04

0.08

0,06

0.12

Т

0.14

0.16

0.18

0,2

0,1

Reduction of critical temperature, magnetization and increase of the resistivity up to

dead layer

0,22

Role of interfaces and dead-layer



You can take into account the surface term between ferro and antiferro layers when the thickness of the film is small.

With decreasing the size, the thickness of Antiferro layer increases. With small sizes, only the antiferro layer can become stable.

C.A. Perroni et al., to be submitted

Interesting issue within a collaboration with experimental groups

Low T transport: signature of disorder effects

Different mechanisms for resistivity: $ho \propto T^{lpha}$

 α =2: Fermi liquid or spin-flip (prohibited in truly half-metallic samples) mechanism (I. Mannari, 1959) (single crystals)

 α =2.5: finite density of states of the minority spins at Fermi energy due to disorder inducing spin-flip scattering (X. Wang et al. PRL 1999)

 α =3: anomalous single magnon scattering at finite temperatures through non rigid behavior of the bands due to spin fluctuations (T. Akimoto et al. PRL 2000).

Relevance of single-magnon scattering

 $\rho_{LT} = \rho_0 + AT^{\alpha}$

 $\rho_{\rm 0}$ measures the strength of the disorder

S. Mercone, C.A. Perroni et al., PRB **71,** 64415 (2005) ¹³

Analysis on many optimized films grown with different techiques

Correlation between residual resistivity and the exponent α seen in experimental data



 α =3 dominant as effective bandwidth decreases. With increasing disorder, the energies available to delocalized states get reduced.

Phase-separation in LCMO and LSMO films Experimental resistivity $\rho_{EX}(T)$



S. Mercone, C.A. Perroni et al., PRB **71,** 64415 (2005) ¹⁵



For LSMO bulk (single crystal) there is no metal-insulator transition ¹⁶

STM analysis in LSMO films with x=0.3 at room and nitrogen liquid temperature in collaboration with Coherentia –INFM STM group (R. Di Capua, M. Salluzzo, R. Vaglio)



Nanoscale correlations at room temperature

No inhomogeneities at nitrogen liquid temperature!

Features compatible with multiphase modulation

In addition to stress fields, amounts of disorder can tune the behavior of inhomogeneity and consequently CMR, see J. Burgy et al., PRL **92**, 97202 (2004).

R. Di Capua, C.A. Perroni, et al. to be submitted

Conclusions

•The role of thickness, strain and el-ph interaction are relevant in affecting the properties of manganite thin films and their tendencies toward "phase separation" regime.

•The role of disorder is important in influencing the transport properties.

•Open questions:: Nanometric and/or micrometric scale inhomogeneity. Feasibility of multiphase modulation?

•Perspectives: STM experiments and analysis at temperatures larger than MI transition temperature.