## REALITY OF THE QUANTUM NORMAL FORM GENERATED BY PT-SYMMETRIC PERTURBATIONS

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ABSTRACT. Consider the classical Hamiltonian  $\mathcal{H}_{\varepsilon}(\xi, x) = \omega_1 \xi_1 + \ldots + \omega_l \xi_l + \varepsilon \mathcal{V}(\xi, x), \varepsilon \in \mathbb{R}$ , defined on the phase space  $\mathbb{R}^l \times \mathbb{T}^l$ , where  $\mathcal{V}$  is a periodic potential with real Fourier coefficients  $\mathcal{V}_k$  satisfying the condition:  $\mathcal{V}_{2k} = \mathcal{V}_{-2k}$  and  $\mathcal{V}_{2k+1} = -\mathcal{V}_{-2k-1}$ , for all  $k \in \mathbb{Z}$ . Its (Weyl) quantization yields the non self-adjoint *PT*-symmetric operator  $H_{\varepsilon} = -i\hbar[\omega_1 \frac{\partial}{\partial x_1} + \ldots + \omega_l \frac{\partial}{\partial x_l}] + \varepsilon V$  acting on  $L^2(\mathbb{T}^l)$ . We consider a class of perturbations  $\mathcal{V}$  for which the quantum normal form is uniformly convergent with respect to  $\hbar$ , and prove that the *PT*-symmetry implies the reality of

 $\varepsilon V$  acting on  $L^2(\mathbb{T}^l)$ . We consider a class of perturbations  $\mathcal{V}$  for which the quantum normal form is uniformly convergent with respect to  $\hbar$ , and prove that the *PT*-symmetry implies the reality of its symbol. This entails the explicit contruction of a similarity transformation mapping  $H_{\varepsilon}$  into a self-adjoint operator, as well as of an exact quantization formula for its real spectrum. For  $\hbar = 0$ the construction reduces to a classical complex canonical transformation mapping the complex Hamiltonian  $\mathcal{H}_{\varepsilon}(\xi, x)$  into a real one.